

Chapter 1

NV-NV cross-relaxations : the fluctuator model

In this chapter,...

1.1 Experimental observation of NV-NV cross-relaxation

Before we discuss the theoretical complications related to NV-NV cross-relaxations (CR), let us first show the unambiguous experimental proof of the presence of NV-NV CR.

1.1.1 NV-NV CR between nonequivalent NV centers

NV-NV CR was first observed more than thirty years ago [2, 3]. The first observations were between non-equivalent NV centers, meaning that the two NV centers involved in the dipole-dipole coupling were not polarized equally.

This scenario can happen for instance when two NV centers from different classes see a different transverse (and longitudinal) magnetic field. Fig 1.1 illustrates this in the case where the magnetic field is parallel with one of the four classes: $\mathbf{B} \parallel [111]$

When $\mathbf{B} \parallel [111]$, as we discussed in the last chapter, one class sees no transverse field and is therefore always polarized. The three other (equivalent) classes on the other hand get more and more depolarized as the magnetic field amplitude is increased.

It turns out that there is a co-resonance at $B=592$ G between the class parallel to \mathbf{B} and the three other classes. This co-resonance is represented in Fig. 1.1-b) by an orange circle.

Fig. 1.1-c), which we already saw in the last chapter, shows the change in PL of an ensemble of NV centers as \mathbf{B} is scanned along the $[111]$ axis. For $B=592$ G, we can see a drop in PL also circled in orange. This drop is

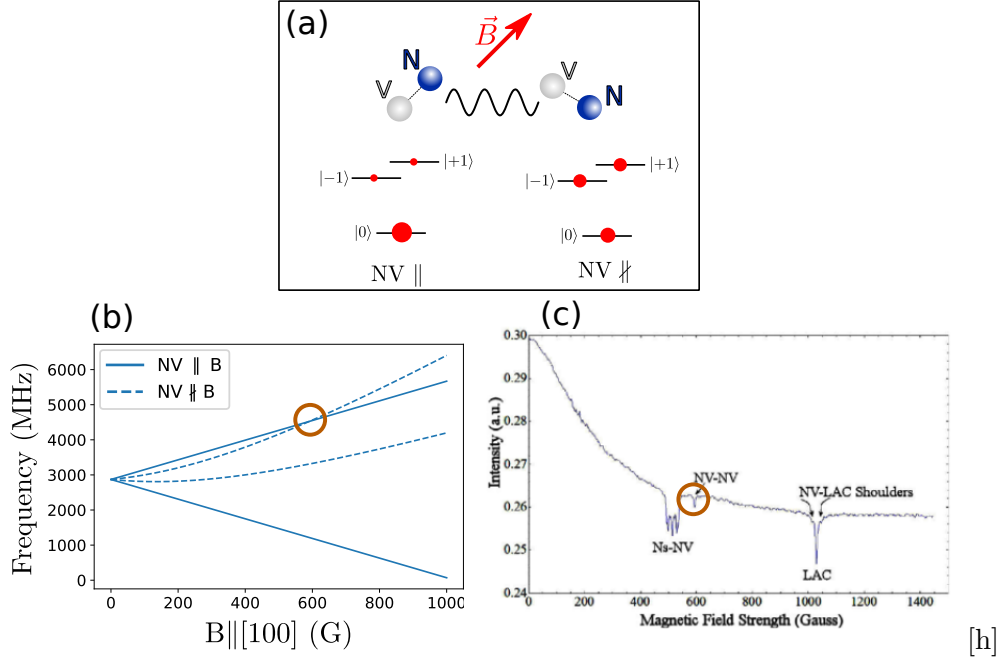


Figure 1.1: label. Taken from [1]

the result of the CR between NV centers from the class parallel to \mathbf{B} and NV centers from the three other classes.

While a difference in polarization is needed to observe a population transfer, it is not enough to explain the drop in PL. This PL drop is due to the difference in brightness between the $|0\rangle$ and $|+1\rangle$ from the different classes involved. The PL contrast between the two states is higher for the class parallel to \mathbf{B} than for the three other ones, which is another consequence of the transverse field.

The co-resonance between two different classes of NV centers with different magnetic field projection is a relatively rare event: it can only occur for the $|0\rangle \rightarrow |+1\rangle$ transition, and only for magnetic fields greater than 592 G.

1.2 Microscopic origin of the fluctuators

1.2.1 Tunneling

1.2.2 Modulation of the J-coupling by the phonons

”Electron Spin-Lattice Relaxation in Phosphorus-Doped Silicon”

”On the interaction of nuclear spins in a crystalline lattice”

”Concentration- and Compensation-Dependent Spin-Lattice Relaxation in n-Type Silicon”

Attention quand même, ce process dépend e la température si je dis pas de bêtises, c'est p-e pas ça chez nous. Nan ok, ça modifie γ mais pas le T1 directement. P-e que c'est mesurable ?

Attention bis : C'est l'échange interaction (le terme en δ) qui est modifié, visiblement le couplage dipolaire est trop faible pour expliquer ça (papier de Pines et "Spin Resonance of Impurity Atoms in Silicon")

Bibliography

- [1] Seiji Armstrong et al. “NV–NV electron–electron spin and NV–NS electron–electron and electron–nuclear spin interaction in diamond”. In: *Physics Procedia* 3.4 (2010), pp. 1569–1575.
- [2] K Holliday et al. “Optical hole-bleaching by level anti-crossing and cross relaxation in the NV centre in diamond”. In: *Journal of Physics: Condensed Matter* 1.39 (1989), p. 7093.
- [3] Eric van Oort and Max Glasbeek. “Cross-relaxation dynamics of optically excited N-V centers in diamond”. In: *Physical Review B* 40.10 (1989), p. 6509.