

Chapter 1

NV-NV cross-relaxations : the fluctuator model

In this chapter,...

1.1 Experimental observation of NV-NV cross-relaxation

Before we discuss the theoretical complications related to NV-NV cross-relaxations (CR), let us first show the unambiguous experimental proof of the presence of NV-NV CR.

1.1.1 NV-NV CR between nonequivalent NV centers

NV-NV CR was first observed more than thirty years ago [2, 3]. The first observations were between non-equivalent NV centers, meaning that the two NV centers involved in the dipole-dipole coupling were not polarized equally.

This scenario can happen for instance when two NV centers from different classes see a different transverse (and longitudinal) magnetic field. Fig 1.1 illustrates this in the case where the magnetic field is parallel with one of the four classes: $\mathbf{B} \parallel [111]$

When $\mathbf{B} \parallel [111]$, as we discussed in the last chapter, one class sees no transverse field and is therefore always polarized. The three other (equivalent) classes on the other hand get more and more depolarized as the magnetic field amplitude is increased.

It turns out that there is a co-resonance at $B=592$ G between the class parallel to \mathbf{B} and the three other classes. This co-resonance is represented in Fig. 1.1-b) by an orange circle.

Fig. 1.1-c), which we already saw in the last chapter, shows the change in PL of an ensemble of NV centers as \mathbf{B} is scanned along the $[111]$ axis. For $B=592$ G, we can see a drop in PL also circled in orange. This drop is

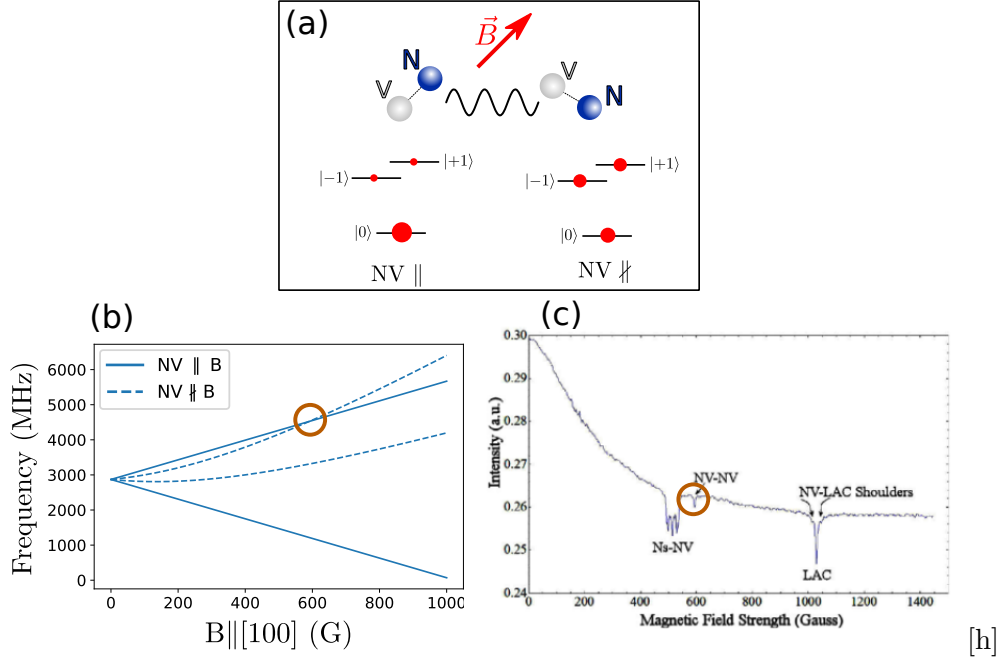


Figure 1.1: label. Taken from [1]

the result of the CR between NV centers from the class parallel to \mathbf{B} and NV centers from the three other classes.

While a difference in polarization is needed to observe a population transfer, it is not enough to explain the drop in PL. This PL drop is due to the difference in brightness between the $|0\rangle$ and $|+1\rangle$ from the different classes involved. The PL contrast between the two states is higher for the class parallel to \mathbf{B} than for the three other ones, which is another consequence of the transverse field.

The co-resonance between two different classes of NV centers with different magnetic field projection is a relatively rare event: it can only occur for the $|0\rangle \rightarrow |+1\rangle$ transition, and only for magnetic fields greater than 592 G.

1.1.2 NV-NV CR between equivalent NV centers

More recently, experiments [4, 5, 6] on dense NV ensemble ($[NV] > 1$ ppm) or at low temperature have shown that there was CR even between equivalent NV centers.

I refer by equivalent NV centers to NV centers with the same spin Hamiltonian. This means either NV center from the same class, or NV centers from different classes with the same projection of the magnetic field along their axis. Equivalent NV centers are by nature resonant with each other, and therefore susceptible to CR.

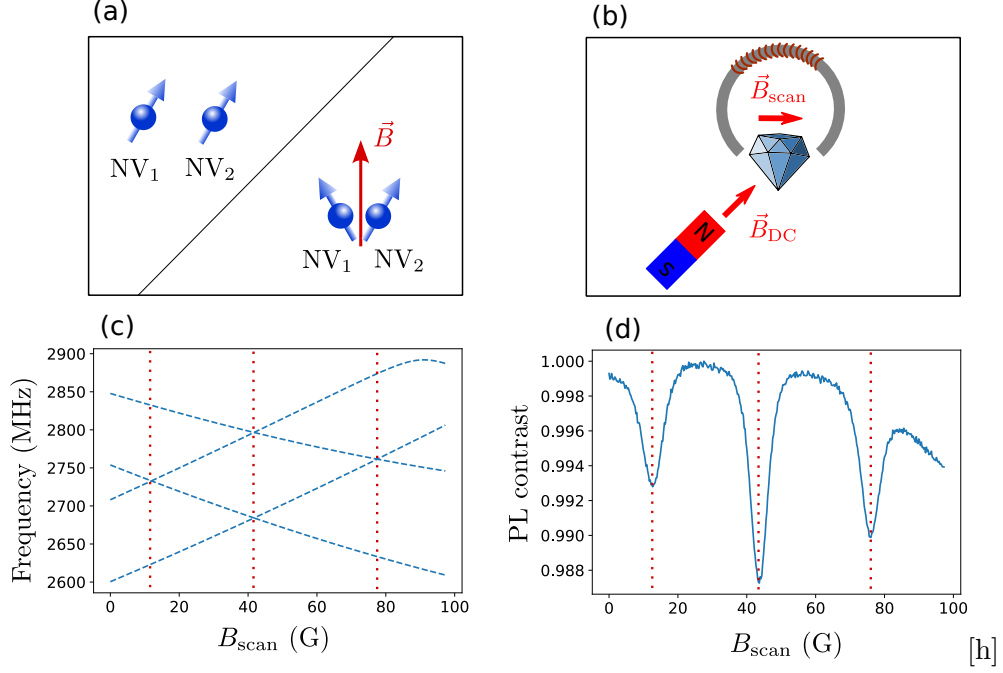


Figure 1.2: CR between equivalent NV centers for sample ADM-15-1. (a) Representation of equivalent NV centers: either two NVs from the same class, or two classes with the same projected magnetic field. (b) Magnetic setup used for the experiment: a permanent magnet is used to apply a bias magnetic field, and an electromagnet is used to add a variable magnetic field. (c) Simulation of the $|0\rangle \rightarrow |-1\rangle$ transition for the 4 classes of NV centers as a function of the scanned magnetic field. The transitions were computed based on several ODMR spectra. The red dotted lines correspond to inter-class resonance (d) Change in the PL of the NV centers ensemble as a function of the scanned magnetic field. The inter-class resonances are reported here.

Fig. 1.2 shows an experimental observation of equivalent NV-NV Cr. The key to observe NV-NV Cr is to bring some NV centers in and out of resonance with other NV centers, which we can do with the different NV classes. To do so, we need to add an initial bias magnetic field in addition to the scanned magnetic field, as represented in Fig. 1.2-b).

Fig. 1.2-c) shows the transition frequencies of the $|0\rangle \rightarrow |-1\rangle$ transition for the 4 classes of NV centers on sample ADM-15-1, an HPHT micro-diamond with $[NV] \sim 3$ ppm. There are three magnetic field values $B_{scan} \sim 12, 42$ and 77 G for which there is an inter-class resonance. Each of these resonances correspond to a "geometric" resonance where the B field is in a symmetry plane between the two classes, meaning that the two resonant NV classes are indeed equivalent.

Finally, Fig. 1.2-d) shows the PL contrast as the electromagnet field is scanned. There are 3 very clear dips when $B_{\text{scan}} \sim 12, 42$ and 77 G which coincide with inter-class resonances. We will see later that this PL dip is associated with a decrease of the NV centers spin T_1 time from both classes.

This observation, and similar ones done by many groups, seem to indicate that there are CR between equivalent NV centers. This is however incompatible with our previous assumptions that equivalent NV centers are equally polarized and bright. To understand this phenomenon, we will need to introduce new hypotheses.

1.2 NV inhomogeneity and the fluctuator model

1.2.1 CR in an inhomogeneous NV bath

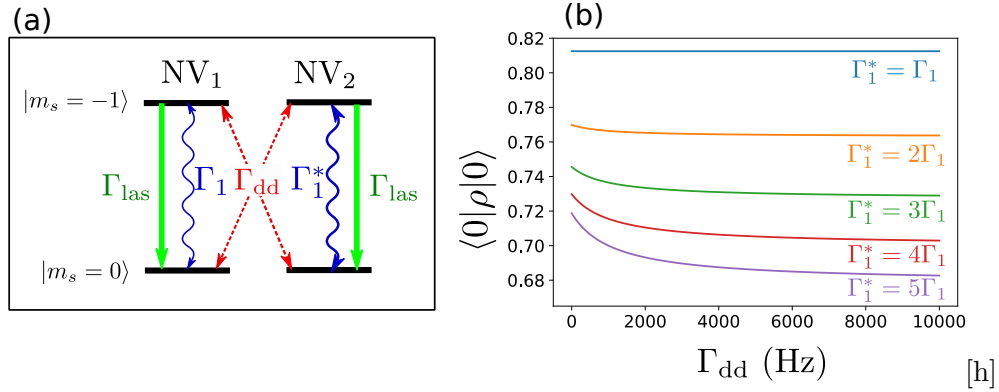


Figure 1.3: label

A likely explanation to the observed NV-NV CR is that the NV centers are not all equivalents.

Indeed, if we assume that, for instance, the relaxation rate $\Gamma_1 = 1/T_1$ is not strictly the same for each NV centers, but instead follows a certain distribution $\rho(\Gamma_1)$, then CR between NV centers with different Γ_1 can explain the observations in Fig. 1.2.

Fig. 1.3 illustrates this case. We are assuming here a coupling between two NV centers with respective relaxation rate Γ_1 and Γ_1^* . Fig. 1.2-a) illustrates the model I chose for this system: Γ_{las} represent the optical pumping rate from the $|+1\rangle$ to the $|0\rangle$ state, and Γ_{dd} the flip-flop rate between the two spins. I consider here only the incoherent dynamics of the population which I model with various rates. I also only consider the $|0\rangle$ and $|+1\rangle$ states as the results would be the same for the $|0\rangle$ and $|-1\rangle$ states.

By finding the steady state of these rate equations, I can compute the final population in the $|0\rangle$ for both spins, which I will assume be proportional

to the total PL. Fig. 1.3-b) shows this $|0\rangle$ population as a function of the flip-flop rate Γ_{dd} for various values of Γ_1^* .

We can see that when $\Gamma_1^* = \Gamma_1$, as expected, the PL is not modified by the coupling of the two spins regardless of the coupling strength. When $\Gamma_1^* > \Gamma_1$ however, not only is the starting PL slightly lower since NV_2 is less polarized, but the PL now decreases as we increase the coupling strength. This effect, which is the origin of the CR contrast, is stronger when the difference between Γ_1 and Γ_1^* is higher.

1.2.2 Presentation of the fluctuator model

Choi et al. in [6] take this approach a step further by separating the NV centers into two groups: "normal" NV centers with a phonon-limited lifetime outside of dipole-dipole interaction, and "fluctuators" which are NV centers with an intrinsic, extremely fast relaxation mechanism. We are assuming that the fluctuator relaxation rate, γ_f is much greater than the optical pumping rate Γ_{las} . The fluctuators are therefore unpolarized and effectively act as dark spins, similarly to what we discussed in the last chapter.

We should note here that, while this simplifications seem a bit extreme, the fluctuator model is more than a toy model. There are actually good evidence of the presence of these dark NV centers in dense NV ensemble, as will be discussed below.

We will consider the fluctuators act as a Markovian bath, meaning that the fluctuator density matrix will always read $\rho = \frac{1}{3}I$, regardless of its interaction with NV centers. With this assumption, we can compute the modification caused by the fluctuators on the rest of the NV lifetimes.

1.2.3 Single NV center coupled to a single fluctuator

We will follow here the notations and calculation steps in [6]. To compute the depolarization induced by the fluctuator bath the NV centers, we should first consider the interaction between a single NV center and a fluctuator. Since we assume the fluctuators to be always depolarized, this step is similar to the coupling of an NV center to a dark spin done in the last chapter.

We will first decompose the dipole-dipole Hamiltonian in a radial and angular part:

$$\mathcal{H}_{\text{dd}} \approx -\frac{J_0}{r^3} [(g + ih)(|0, +1\rangle \langle +1, 0| + |0, -1\rangle \langle -1, 0| + qS_z^1 S_z^2] + h.c., \quad (1.1)$$

where the expression of J_0 , g , h and q is given in appendix [REF]. g , h and q are dimensionless factors that are function of the relative orientation of the two dipoles.

We then introduce the dimensionless number η defined as:

$$\eta^2 = \frac{1}{3}(|g|^2 + |h|^2) \frac{4\gamma_f^2}{(\omega_f - \omega_{\text{NV}})^2 + 4\gamma_f^2}, \quad (1.2)$$

where $\gamma_f = 1/T_1^f$ is the fluctuator decay rate. This number η encapsulate both the geometry of the problem (through g and h) and the resonance condition between the NV center and the fluctuator.

Finally, the additional depolarization rate induced by the fluctuator on the NV center reads:

$$\gamma_s(\mathbf{r}) = \left(\frac{J_0}{r^3} \right) \frac{\eta^2}{\gamma_f}. \quad (1.3)$$

Compared to eq. [REF], we have replaced here Γ_2^* by γ_f . This is because we assume that the pure dephasing, supposedly equal for the NV center and the fluctuator, is smaller than the depolarization rate γ_f . We will further discuss this hypothesis later.

1.2.4 Ensemble of NV centers coupled to a bath of fluctuators

We now need to compute the average depolarization caused by the ensemble of fluctuators on the ensemble of NV centers.

We consider each fluctuator as an independant depolarization source, therefore the depolarization on a single NV center reads:

$$\gamma_{\text{NV}} = \sum_i \gamma_s(\mathbf{r}_i). \quad (1.4)$$

We then want to compute the distribution $\rho(\gamma_{\text{NV}})$ defined as:

$$\rho(\gamma_{\text{NV}}) = \int d\{r_i\} \rho(\{r_i\}) \delta \left(\sum_i \gamma_s(\mathbf{r}_i) - \gamma_{\text{NV}} \right). \quad (1.5)$$

To do so, we need to determine the distribution of the fluctuators positions $\{r_i\}$. Assuming that they are homogeneously distributed in the bulk of the material, we find:

$$\rho(\gamma_{\text{NV}}) = \frac{e^{-1/(4\gamma_{\text{NV}}T)}}{\sqrt{4\pi\gamma_{\text{NV}}^3T}}, \quad (1.6)$$

where the time constant T was introduced and is defined as:

$$\frac{1}{T} = \left(\frac{4\pi n_f J_0 \bar{\eta}}{3} \right)^2 \frac{\pi}{\gamma_f}, \quad (1.7)$$

with n the fluctuator density and $\bar{\eta}$ the averaged value of $|\eta|$: $\bar{\eta} = \int \text{Prob}(\eta) |\eta| d\eta$.

Finally, the polarization dynamics from the ensemble of NV centers can be computed:

$$P(t) = \int_0^\infty \rho(\gamma_{\text{NV}}) e^{-\gamma_{\text{NV}}t} d\gamma = e^{-\sqrt{t/T}}. \quad (1.8)$$

In conclusion, the fluctuators causes a depolarization on the NV ensembles with a timescale T given by 1.7, and the dynamics of this depolarization is that of a stretched exponential, unlike the phonon-induced depolarization which is purely exponential.

It should be noted that the stretched exponential nature of the decay is not proper to the fluctuator model. [7] came to the same conclusion by analyzing the resonant coupling of NV centers with a P1 bath. The stretched exponential is a result from localized noise sources, randomly distributed in the bulk (3D).

1.3 Validation of the fluctuator model

We will now show experimental results validating the fluctuator model.

1.3.1 The stretched exponential lifetime

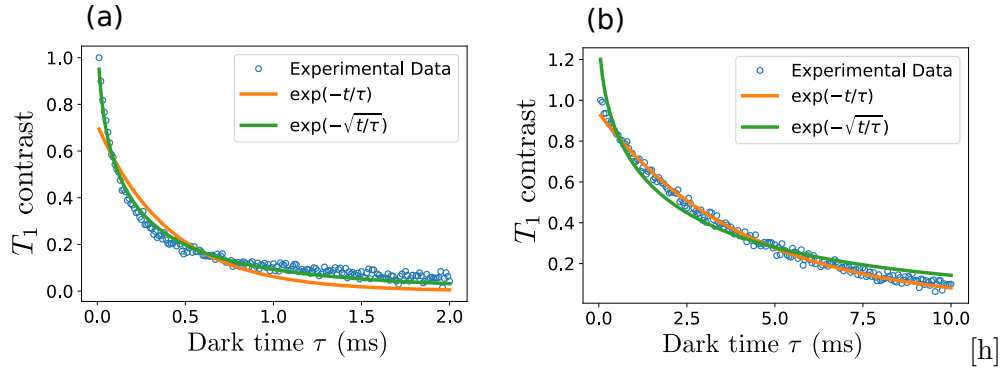


Figure 1.4: T_1 measurement following the protocol described in [REF]. The best exponential and stretched exponential fits are given each time. a) Sample ADM-15-2 with $B=0$. The optimal T_1 values are $T_1^{\text{stretch}} = 150 \mu\text{s}$ and $T_1^{\text{exp}} = 410 \mu\text{s}$. b) Sample CVD-pink with $B \neq 0$ and a single class probed. The optimal T_1 values are $T_1^{\text{stretch}} = 1.9 \text{ ms}$ and $T_1^{\text{exp}} = 4.1 \text{ ms}$

T_1 measurements with dense ensemble of NV centers and when many classes are resonant with each other (typically when $B=0$) indeed tend to show stretched exponential profiles.

Fig. ?? shows an example of that: in Fig. ??-a) the T_1 measurement comes from sample ADM-15-2 at $B=0$. The profile is clearly better fitted by a stretched exponential than a regular exponential, and the timescale found $T_1^{\text{stretch}} = 150 \mu\text{s}$ is much shorter than the expected phonon limited lifetime of a few ms.

On the other hand, Fig. ??-b) shows a T_1 measurement coming from sample CVD-pink for $B \neq 0$. This time the profile is more exponential, and

the value $T_1^{\text{exp}} = 4.1$ ms is coherent with a phonon-limited lifetime. Even though CVD-pink shows some NV-NV CR feature, the spin dynamics in this sample is still dominated by the phonons, which explains the exponential profile.

1.3.2 Fitting T_1 profile in the general case

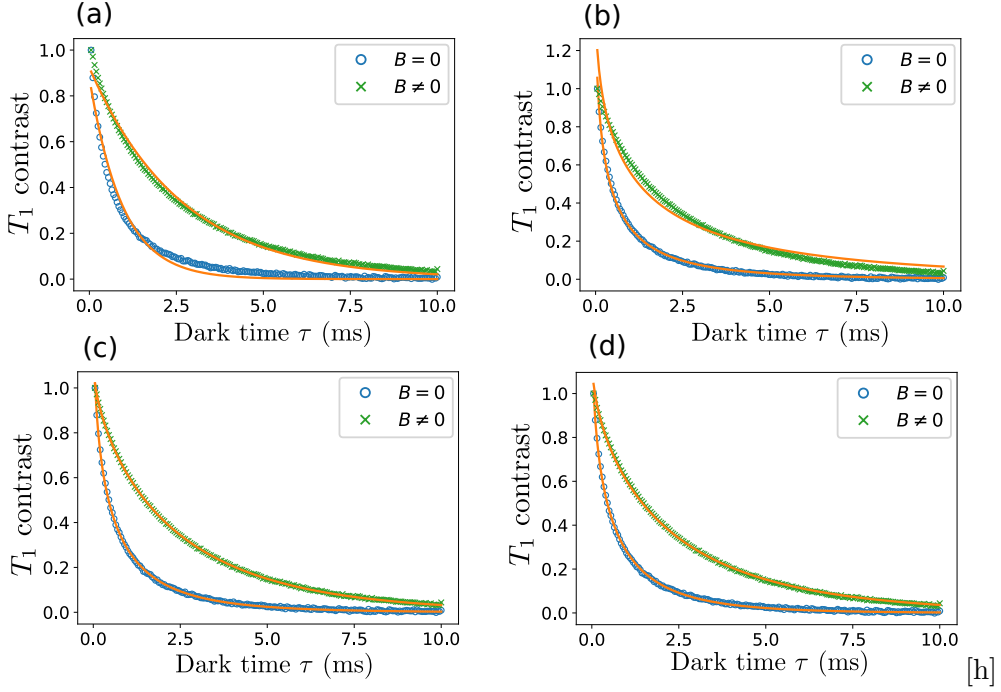


Figure 1.5: Different fitting procedure on two T_1 measurements done on sample ADM-150-1 for $B = 0$ and $B \neq 0$.

- a) Exponential fit $S(\tau) = \exp(-\tau/T_1^{\text{ph}})$.
- b) Stretched exponential fit $S(\tau) = \exp(-\sqrt{\tau}/T_1^{\text{dd}})$.
- c) Bi-exponential fit $S(\tau) = \exp(-\tau/T_1^{\text{ph}} - \sqrt{\tau}/T_1^{\text{dd}})$.
- d) Bi-exponential fit with fixed $T_1^{\text{ph}} = 5$ ms: $S(\tau) = \exp(-\tau/5 - \sqrt{\tau}/T_1^{\text{dd}})$.

For many samples however, the T_1 profile is neither fully exponential nor fully stretched exponential, because the relaxation time associated with CR, which we will call T_1^{dd} is of the same order as the phonon limited lifetime T_1^{ph} . Sometimes the same sample can go from mostly exponential to mostly stretched exponential depending on the number of NV-NV co-resonances,

Table 1.1: Fitting parameters on Fig. ??

Figure	$B = 0$		$B \neq 0$	
	T_1^{ph} (ms)	T_1^{dd} (ms)	T_1^{ph} (ms)	T_1^{dd} (ms)
Fig. a)	0.95	*	2.64	*
Fig. b)	*	0.32	*	1.03
Fig. c)	6.75	0.44	4.22	7.19
Fig. d)	5	0.51	5	4.65

Bold characters indicate parameters that were arbitrarily fixed.

forcing us to include both aspects if we want a unique fitting formula.

Fig. ?? shows four different fitting procedure on the two same measurements. The measurements here are two T_1 measurements done on the same sample ADM-150-1, following the protocol described in [REF], with and without an external magnetic field ($\sim 50G$). The magnetic field was strong enough to split the four classes and only one class was probed in the second case.

The fits used here are either purely exponential or purely stretched exponential, or a combination of both where the exponential lifetime T_1^{ph} was either left as a free parameter or fixed at a value $T_1^{\text{ph}} = 5$ ms. The values of the different fitting parameters used is reported in Table 1.1.

We can see that neither the purely exponential or stretched exponential fits can be satisfying for both measurements: the $B = 0$ curve is poorly fitted by the exponential fit and the $B \neq 0$ curve is purely fitted by the stretched exponential. The protocols that include both exponential and stretched lifetimes correctly fit both curves. In the case of Fig. ??-d), we arbitrarily fixed $T_1^{\text{ph}} = 5$ ms since this is the value we typically measure on samples with low NV density, and we expect that the phonon-limited exponential lifetime is not modified by the NV concentration.

Since the protocol where T_1^{ph} was fixed and the one where it wasn't both correctly fit our data, we decided to use the protocol where T_1^{ph} was fixed because there is one less free parameter in the fitting function, and the values we obtain on T_1^{dd} can be directly compared.

We should note however that the values of T_1^{dd} we obtain are not absolute. Indeed, in the case of Fig. ??, the value of T_1^{ph} can be somewhat arbitrarily fixed between 3 and 6 ms with satisfying fits, and the resulting values T_1^{dd} can vary by almost an order of magnitude in the case of $B \neq 0$. While this technique of measuring T_1^{dd} is relatively sensitive and reproducible, it is not very accurate.

Finally, another fitting procedure is shown in Fig. ?. This time

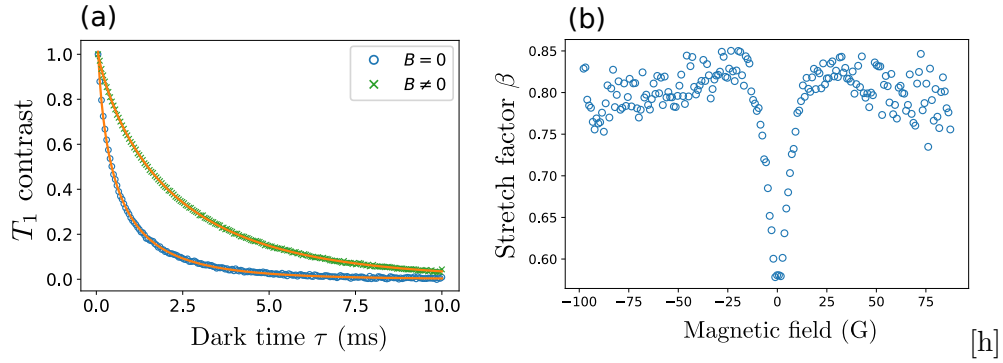


Figure 1.6: a) Same two measurements as Fig. ?? with the fitting formula $S(\tau) = \exp((- \tau/T_1)^\beta)$. The fitting parameters are $\beta = 0.58$ and $T_1 = 0.46$ ms for $B = 0$, and $\beta = 0.80$ and $T_1 = 2.15$ ms for $B \neq 0$. b) Optimal β parameter found for each T_1 measurement on the same sample as a function of the external magnetic field.

1.3.3 The fluctuators linewidth

1.3.4 Parallel with P-doped Si

”Electron Spin-Lattice Relaxation in Phosphorus-Doped Silicon”

”On the interaction of nuclear spins in a crystalline lattice”

”Concentration- and Compensation-Dependent Spin-Lattice Relaxation in n-Type Silicon”

1.3.5 Microscopic origin of the fluctuators

On va dire subsection plutot. Du coup ici discussion tunneling Vs Modulation of J Vs surface spins

Attention quand même, ce process dépend e la température si je dis pas de bêtises, c’est p-e pas ça chez nous. Nan ok, ça modifie gammaf mais pas le T1 directement. P-e que c’est mesurable ?

Attention bis : C’est l’échange interaction (le terme en delta) qui est modifié, visiblement le couplage dipolaire est trop faible pour expliquer ça (papier de Pines et ”Spin Resonance of Impurity Atoms in Silicon”)

1.4 Geometric conditions for NV-NV resonance

En gros la carte, les 5 ESR en SI du PRL (lol) et le 121 Vs 22

1.5 Limitations of the fluctuator model

La modification des T1, tjr plus petite que prévue (faire un tableau), et les T1 non-exponentiels (je crois que y'a des cas sur Ludo, bon courage mdr)
Parler aussi des alternatives au fluctuator model, ici ou dans la 2e section

Bibliography

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