

SUPPLEMENTARY MATERIAL

Resonant Dipole-Dipole Induced Mechanical Rotation of a Diamond Crystal

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I. NV⁻ CENTER THEORY

A. NV spin hamiltonian

The hamiltonian of the electronic spin of the negatively charged nitrogen-vacancy center in its ground state can be written as :

$$\hat{\mathcal{H}}_s = \hbar D S_z^2 + \hbar \gamma_e \mathbf{B} \cdot \hat{\mathbf{S}},$$

where $D = (2\pi)2.87$ GHz is the crystal field splitting originating from spin-spin interactions, and $\gamma_e = 28\text{GHz/T}$ is the electron gyromagnetic ratio. We neglect contributions from the strain and local electric field in the hamiltonian since we are working with magnetic fields on the order of 100G. We also neglect the hyperfine interaction with the nuclear spin of the ^{14}N atom since we are working with ensembles with a typical inhomogeneous broadening of 5 MHz. The \mathbf{z} axis of the S_z operator here is the axis formed by the nitrogen atom and the vacancy in the body frame.

B. Diamond crystalline axes and degeneracy conditions

There are four possible crystalline axis for the NV centers (so-called “classes” of NV) which are depicted in Fig. 1 b) and correspond to the crystalline directions $[111]$, $[1\bar{1}\bar{1}]$, $[\bar{1}1\bar{1}]$ and $[\bar{1}\bar{1}1]$.

The magnetic field is represented in the diamond basis in Fig. 1 a) where the usual polar and azimuthal angles θ and ϕ are defined with respect to the \mathbf{z} ($[001]$) direction.

For some orientations of the magnetic field, the projection of the magnetic field on two or more NV axes will be identical, and therefore the energy level of the corresponding classes will be the same. These degeneracies are represented in Fig 1 c), where the dashed lines are the locus of the crystalline planes orthogonal to the $[110]$ directions (6 planes in total). When the magnetic field belongs to these planes, we observe a degeneracy between two classes of NVs, as can be seen in the Fig. 4 of the main paper.

The plain lines are the loci of the planes orthogonal to the $[100]$ directions (3 planes total). When the magnetic field angle lies in these planes, every class is resonant with another class,

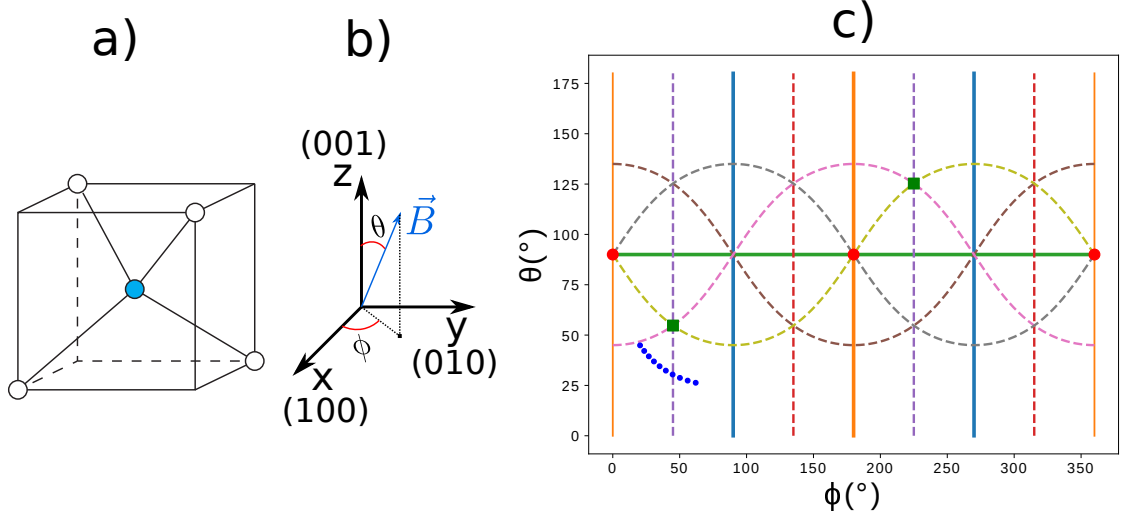


FIG. 1: a) Representation of the four crystalline axes of the diamond b) Representation of the magnetic field in the diamond crystalline basis c) Representation of the crystalline planes in the (θ, ϕ) basis. The $[110]_{\perp}$ family of planes are shown using dashed lines, $[100]_{\perp}$ using plain lines. The $[100]$ direction is marked by red circles, the $[111]$ direction by green squares. The path of the magnetic field in the experiment of Fig. 3 of the main text is shown with blue dots.

as can be seen in Fig. 5. The red circles correspond to the $[100]$ directions, for which the four classes of NV are degenerate. The green squares correspond to the $[111]$ direction where one class is aligned with the magnetic field, and the three others are degenerate. Finally the blue dots correspond to the path followed by the magnetic field in the experiment presented in Fig 3 of the main text, where we can see that a plane from the $[110]_{\perp}$ family is being crossed.

II. EXPERIMENTAL DETAILS

A. Experimental setup

The experimental setup illustrated in Fig.2 is similar to the one used in [1] with the addition of a permanent magnet and an electromagnetic (EM) coil in order to perform magnetic field scans. The diamond sample is typically illuminated with 1mW of 532 nm laser light, focused by an objective with a numerical aperture of 0.5. An acousto-optic modulator (AOM) is used to switch on and off the 532nm laser and to finely tune its

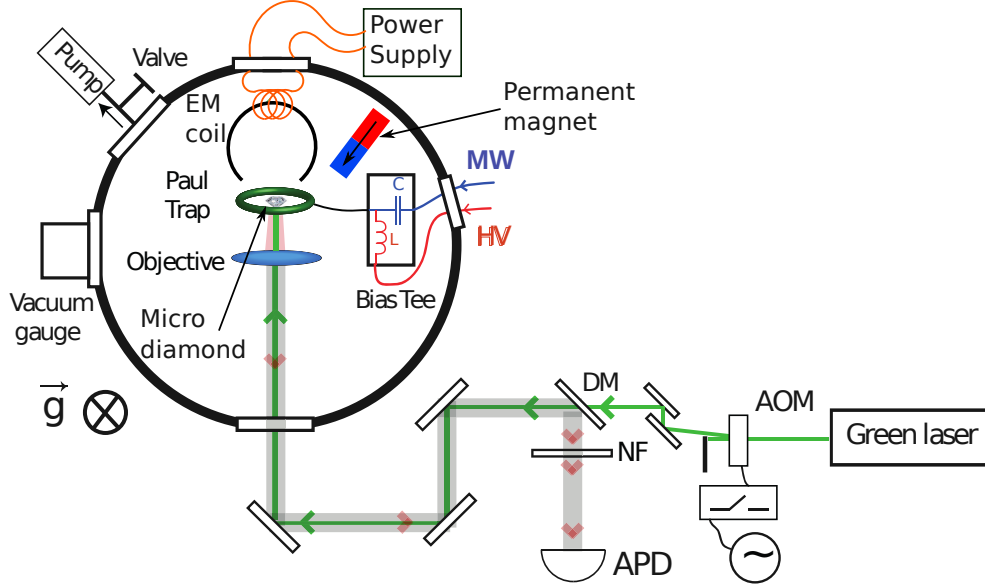


FIG. 2: Illustration of the experimental setup.

power. The photo-luminescence (PL) is collected by the objective, separated from the excitation light using a dichroic mirror (DM) and a 532nm notch filter (NF), and detected using a multimode-fiber single-photon avalanche photo-detector (APD) (SPCM-AQRH-15 from Perkin Elmer). Typically, from the heavily doped samples that we use, we can detect PL photons at a rate of 1 MHz after attenuating the signal by a factor 100 with neutral density filters. The Paul trap is a pseudo-ring with a diameter of approximately $200 \mu\text{m}$, as can be seen in [2]. It acts both as trap through the high voltage (HV) and as a microwave (MW) antenna.

The magnetic field generated by the (homemade) EM coil is controlled by a programmable power supply (Rhode & Schwartz NGE 103) performing current ramps. While the levitating setup is located in a vacuum chamber, all the experiments presented in this article are performed at atmospheric pressure.

B. T_1 Measurement

As shown in the Fig. 2-c) of the main text, the spin lifetime of the NV centers is modified in the presence of cross-relaxation with other classes of NV centers. Here we present the protocol for removing the effects of charge state transfer in the dark, which masks the PL signal decay induced solely by spin depolarization. The protocol described in Fig. 3 consists

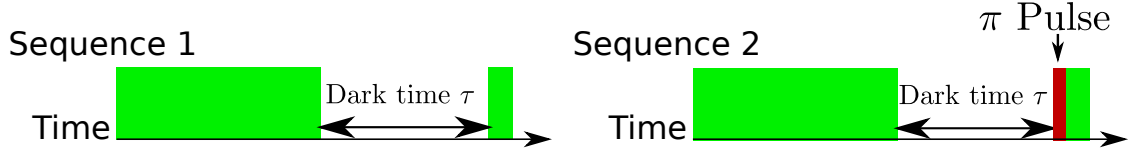


FIG. 3: T_1 measurement protocol. Green bars represent laser excitation, red bar represent resonant microwave π pulse.

in using two sequences. In the first sequence the spins are initially polarized in their $m_s = |0\rangle$ state through a 1 ms green laser excitation pulse and then left to evolve in the dark for a variable dark time τ . The spin state is finally read out using a 10 μs laser pulse (shorter than the polarization time of the spins).

The second sequence uses the same parameters (polarization time, dark time and readout time) than the first sequence, but uses an extra resonant microwave π pulse tuned to a transition of one of the four classes of NV^- right before the readout pulse. The latter sequence brings the population of the $m_s = |0\rangle$ state to the $m_s = |\pm 1\rangle$ state of one class of NV centers.

Writing S_1 the signal obtained by the first sequence and S_2 the one obtained by the second sequence, and then measuring the difference $S_1 - S_2$ between the two, we can measure the spin state population of a single class, and at the same time remove unwanted contributions to the photoluminescence, such as charge state transfer in the dark (which would be common to the two signals and therefore subtracted). In order to avoid low frequency noises or drift, we alternate both sequence when performing the measurement.

C. Magnetic field calibration

A permanent magnet is placed a few cm away from the diamond sample in order to apply a uniform magnetic field to the NV centers together with an electro-magnet.

To calibrate the magnetic field magnitude B , and its orientation θ with respect to the NV axis, we record Optically Detected Magnetic Resonance (ODMR) spectra and use two orientations to determine both the angle and the magnetic field amplitude.

D. Spin-mechanical detection

High sensitivity of the spin-torque is achieved by using a speckle pattern produced by the rough surface of the micro-diamond under coherent illumination. When the particle is stably levitating, at the particle image plane, we then focus a small area of the speckle image onto an optical fibre and detect the photons transmitted through the fibre with the APD₁. The detected signal is then highly sensitive to the particle position and orientation.

For the spin-torque measurements presented in Fig. 4-a), the microwave detuning is scanned in 2 MHz steps with a duration of 10 ms per points. During those 10ms, the diamond orientation has enough time to reach its equilibrium position and the spin torque effect can be observed. The average count-rate is about 1 Mega-counts/s.

E. Signal drift for levitating particles

Measurements on levitating diamonds have to be relatively short (few minutes at most) because of a slow drift on the particle orientation which changes the detection location on the specular reflection off the diamond surface. The most likely origin of this drift is the loss of charges of the diamond due to photoionization by the laser, which changes the trapping conditions over time.

III. PRINCIPLE OF THE MECHANICAL DETECTION

A. Origin of the magnetic torque

The magnetic torque responsible for the motion of the diamond fundamentally comes from the crystalline anisotropy of the NV centers and from the transverse field B_{\perp} responsible for mixing the eigenstates in the stationary state. We will start by considering the torque from a single NV center. We simplify the problem by assuming that the motion is in the $x - z$ plane so the magnetic part of the hamiltonian reads $H_B = \hbar\gamma_e B(\hat{S}_x \sin \theta + \hat{S}_z \cos \theta)$ where θ is the angle between the B field and NV center quantization axis. We thus obtain the spin torque

$$\hat{\tau}_s = -\partial\hat{H}/\partial\theta = \hbar\gamma_e B(-\cos\theta\hat{S}_x + \sin\theta\hat{S}_z). \quad (1)$$

The mean value of the torque in terms of the reduced density matrix elements ρ_{ij} in the basis of the S_z eigenstates $(-1, 0, 1)$ in the body fixed frame is

$$\langle \hat{\tau}_s \rangle = \hbar \gamma_e B (\rho_{-1-1} - \rho_{11}) \sin \theta - \hbar \frac{\gamma_e B}{\sqrt{2}} S \cos \theta \quad (2)$$

where we introduced $S = \rho_{0,1} + \rho_{1,0} + \rho_{0,-1} + \rho_{-1,0}$. The reduced density matrix elements are traced over a bath consisting of laser photons used to polarize the NV at a rate γ_{las} , phonons or spin-fluctuators acting on the spin populations at a rate $R_1 = 1/T_1$ and P_1 or nuclear spins dephasing the electronic spin at a rate $1/T_2^*$. In the limit $\gamma_e B \ll D$, and $\gamma_{\text{las}} \gg \Gamma_1$ the laser still polarizes the spins in the ground state so that $\rho_{00} \gg \rho_{11}, \rho_{-1-1}$. Assuming that the pure dephasing T_2^* is much shorter than the sum of the population relaxation $T_1/2$ and the laser induced repolarization γ_{las} the equations of motion for the coherences are

$$\frac{\partial \rho_{01}}{\partial t} = -\frac{1}{2T_2^*} \rho_{01} - i \frac{\gamma_e B}{\sqrt{2}} \sin \theta - i \rho_{01} D + \mathcal{O}\left(\frac{(\gamma_e B)^2}{D}\right) \quad (3)$$

$$\frac{\partial \rho_{0-1}}{\partial t} = -\frac{1}{2T_2^*} \rho_{0-1} - i \frac{\gamma_e B}{\sqrt{2}} \sin \theta - i \rho_{0-1} D + \mathcal{O}\left(\frac{(\gamma_e B)^2}{D}\right). \quad (4)$$

The motional dynamics is very slow compared to the zero-field and magnetic field splittings. The latter are also larger than the decoherence rate $1/T_2^*$ so we can adiabatically eliminate the coherences and find

$$\rho_{01} = \rho_{10} \approx -\frac{\gamma_e B \sin \theta}{\sqrt{2} D} \quad \text{and} \quad \rho_{0-1} = \rho_{-10} \approx -\frac{\gamma_e B \sin \theta}{\sqrt{2} D} \quad (5)$$

Since

$$\rho_{11} - \rho_{-1-1} = \mathcal{O}\left(\frac{\gamma_e B}{D}\right). \quad (6)$$

Re-injecting these expressions in the expression for the mean torque, we get

$$\langle \hat{\tau}_s \rangle = \frac{\hbar (\gamma_e B)^2}{D} \sin 2\theta + \mathcal{O}\left(\frac{(\gamma_e B)^3}{D^2}\right). \quad (7)$$

It is in fact the gradient of the energy in the ground state taken at the angle θ .

Indeed, supposing that $\gamma B \ll D$, so that H_B can be treated as a perturbation to the spin-spin hamiltonian, the perturbed energy ϵ_0 of $|0\rangle$ is

$$\epsilon_g = \sum_{m_s=\pm 1} \frac{|\langle 0 | H_B | \pm 1 \rangle|^2}{-\epsilon_{\pm 1}^0} = -\hbar \frac{(\gamma_e B)^2}{D} \sin^2 \theta$$

The formula $-\partial\epsilon_g/\partial\theta$ then gives Eq. 7. It is the equation that is used in the core of the manuscript. In the approximate regime of the present study, the Feynman-Hellman theorem that relates the angular derivative of the mean energy to the torque in the pure state case, can in fact also be applied directly here.

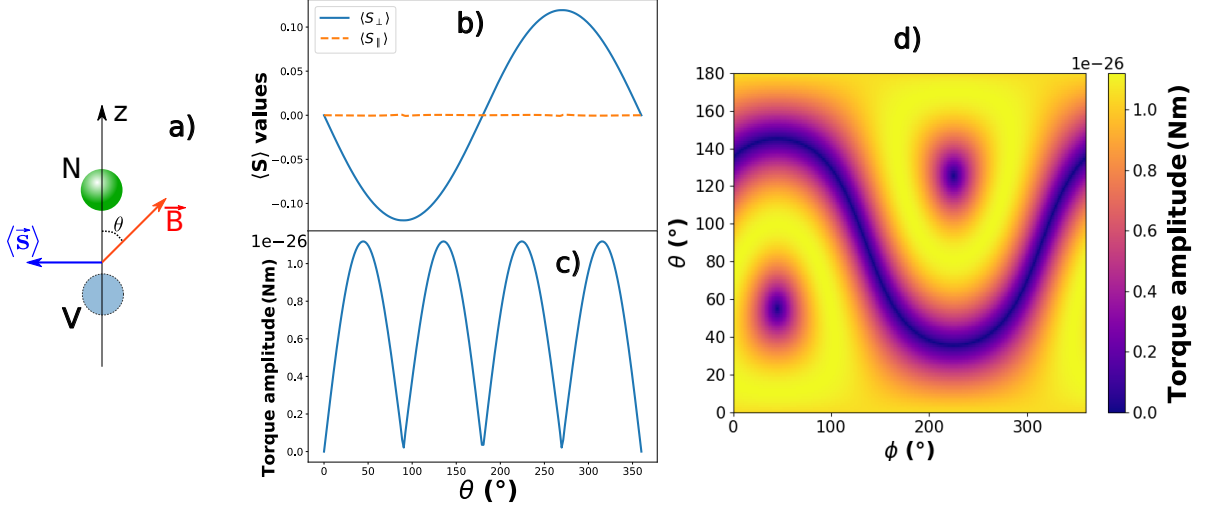


FIG. 4: **a)** Sketch showing an NV center aligned to the z direction in the presence of an external magnetic field \mathbf{B} at an angle θ with respect to the NV axis. The resulting spin vector $\langle \hat{S} \rangle$ of the NV center is shown by the blue arrow. **b)** Longitudinal (S_{\parallel} in dashed line) and transverse (S_{\perp} in plain line) components of the average value of the spin operator, in \hbar unit, as a function of θ and for a magnetic field amplitude $|\mathbf{B}| = 100$ G. **c)** Amplitude of the magnetic torque acting on a single spin as a function of θ for $|\mathbf{B}| = 100$ G. **d)** Amplitude of the same magnetic torque in the crystalline basis with θ and ϕ being the polar and azimuthal angle with respect to the $[100]$ direction.

Another way to estimate the torque is depicted in Fig 4. Under green excitation and in the presence of an external magnetic field, the spins will acquire a magnetization $\gamma_e \langle \hat{S} \rangle$ which, under the low magnetic fields (< 200 G) we are working at, will be oriented at an angle of 90° from the NV axis : $\langle \hat{S}_z \rangle \approx 0$ and $\langle \hat{S}_{\perp} \rangle \neq 0$. This magnetization vanishes when the magnetic field is aligned with the NV center since there is no longer a transverse field responsible for the mixing of the eigenstates.

The magnetization of the NV center is therefore not aligned with the magnetic field, except when the field is also at a 90° angle from the NV axis, which means that the magnetic torque $\mathbf{\Gamma} = \gamma_e \langle \hat{S} \rangle \times \mathbf{B}$ will be non-zero everywhere except when the field is aligned with the center, or in the plane normal to the direction of the center. We can describe each NV center

as a paramagnetic defect with the anisotropic magnetic susceptibility $\chi = \begin{pmatrix} -\chi_{\perp} & 0 & 0 \\ 0 & -\chi_{\perp} & 0 \\ 0 & 0 & 0 \end{pmatrix}$ in the $(\mathbf{x}, \mathbf{y}, \mathbf{z})$ basis where \mathbf{z} is the orientation of the NV center.

The amplitude of the torque with respect to the magnetic field orientation (at a B field amplitude of 100 G) is represented in 1D in Fig 4-c) where we can see a behavior very close to $|\sin(2\theta)|$, which is the formula (3) of the main text that we obtained through a perturbative approach. The same torque amplitude is represented in 2D in Fig 4 d). The two purple dots in the map correspond to the $[111]$ direction when the magnetic field is aligned with the centers. The curvy line corresponds to the $[111]_{\perp}$ plane. Importantly, the maximum torque value is $1 \cdot 10^{-26}$ N.m for a single spin.

B. Total spin torque and influence of cross-relaxation

Fig 3 a) in the main text, represents the same map as Fig 4-c) but including the four NV centers, one in each of the possible $[111]$ orientations. We can see that the maximum torque actually decreased to $3 \cdot 10^{-28}$ N.m even though we increased the number of NV centers by four. This is due to the directional averaging of the torque generated by the four centers. The torque per NV center is decreased by more than two orders of magnitude when taking the directional averaging into consideration. Fig 3 b) in the main text shows the same map, this time taking into account the modification of the spin lifetime due to cross-relaxations. The detailed model is presented in the section V. There are two things to note here :

1. The maximum torque has increased by an order of magnitude compared to the previous case. It reached up to $3 \cdot 10^{-27}$ N.m for four spins, so about 10^{-27} Nm per spin. Qualitatively, this is because cross-relaxation will lower the torque contribution of specific classes (the ones that get depolarized), meaning that the end result is closer to the single spin case (there is less directional averaging).
2. The change in magnetic torque is resonant, it happens only when different classes are brought to resonance. This can be seen by comparing Fig 3 b) of the main text to the $[110]_{\perp}$ planes that were drawn in Fig. 1. The change in the signal when scanning a magnetic field across a CR will be much sharper than the sinusoidal change in the

spin-torque.

C. Torque sensing with a levitating diamond

The way we experimentally measure spin-torques applied on the levitating diamond is by measuring the induced diamond orientational displacement from equilibrium. We model the trap as a pure harmonic potential, both for the center of mass and for the librational degrees of freedom of the diamond with trapping frequencies $\omega_t \approx (2\pi) \cdot 1$ kHz. Considering a single librational degree of freedom, we can write the torque exerted by the trap as $\Gamma_t = -K(\theta - \theta_{eq})$, where $K = I\omega_t^2$ is the stiffness of the trap, I being the moment of inertia of the diamond.

The application of an external torque Γ_{ext} to the diamond will therefore shift the angular equilibrium position in such a way that :

$$-K(\theta - \theta_{eq}) + \Gamma_{ext} = -K(\theta - \theta'_{eq}) \quad (8)$$

so that

$$\delta\theta = \theta'_{eq} - \theta_{eq} = \frac{\Gamma_{ext}}{K} = \frac{\Gamma_{ext}}{I\omega_t^2} \quad (9)$$

In our case, Γ_{ext} is the magnetic torque exerted by the NV^- spins on the diamond. We can write it $\Gamma_{ext} = N_{NV}\langle\Gamma_{1spin}\rangle$ where $\langle\Gamma_{1spin}\rangle = \gamma_e\langle\hat{\mathbf{S}}\rangle \times \mathbf{B} \approx 10^{-27}$ Nm is the expected magnetic torque applied by one spin.

By using the inertia moment formula of a sphere : $I = \frac{2}{5}mr^2$, we can then rewrite the angular displacement as

$$\delta\theta = \frac{\langle\Gamma_{1spin}\rangle n(NV^-)}{\frac{2}{5}m_C r^2 \omega_T^2} \approx 10^{-4} rad$$

where $n(NV^-) \approx 5 \cdot 10^{-6}$ (5 ppm) is the number of NV centers per atoms in the crystal, $m_C \approx 2 \cdot 10^{-26}$ kg is the average weight of a carbon atom (we assume that the bulk of the diamond weight comes from carbon atoms), $r = 7.5 \mu m$ is the typical radius of our diamonds and $\omega_T = 6.3 \cdot 10^3$ rad/s is the typical value of the trap angular frequency.

It should be noted that the main uncertainty comes here from the diamond size can change the expected result by an order of magnitude.

IV. CROSS-RELAXATION DETECTION FOR ANOTHER TYPE OF DEGENERACY

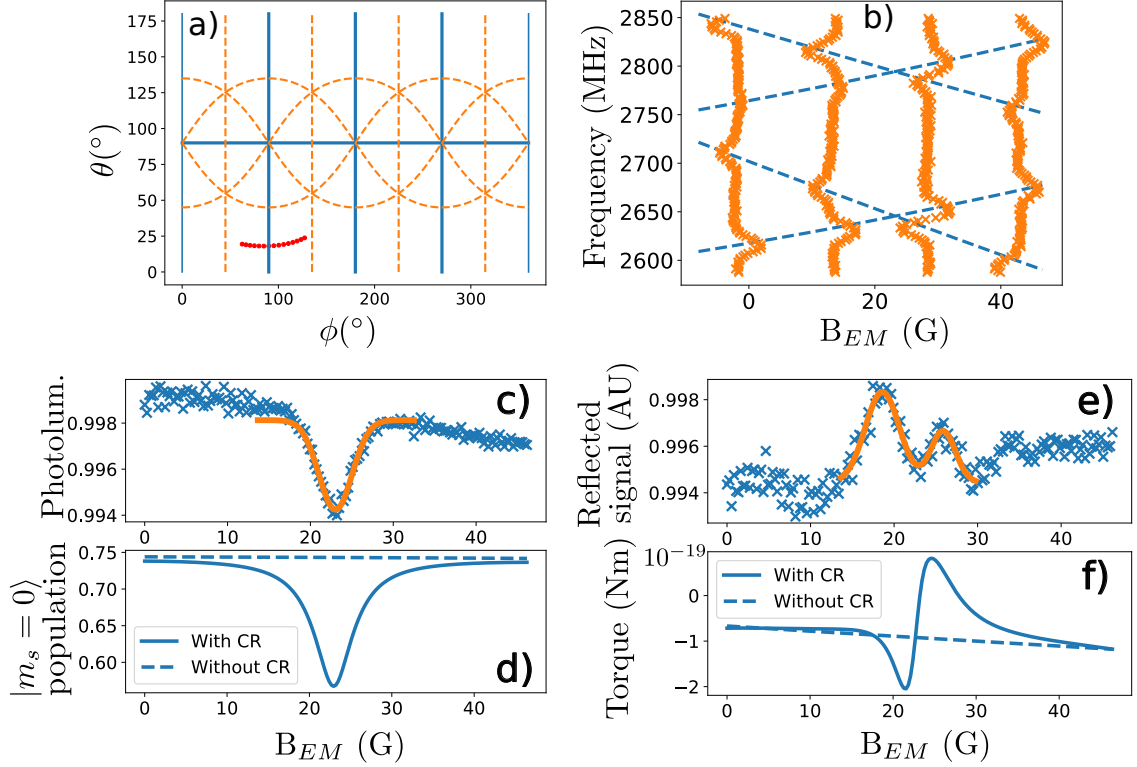


FIG. 5: Mechanical detection of a dipolar interaction when crossing a $[100]_{\perp}$ plane. **a)** Path of the magnetic field angle (red dots) in the (θ, ϕ) basis in the measurements, with respect to the $[100]$ direction. The $[100]_{\perp}$ family of planes are shown in plain blue lines. The $[110]_{\perp}$ family of planes is shown in orange dashed lines. **b)** ESR spectra measured at 4 different magnetic field values. Calculated evolution of the NV transition frequencies as a function of the electromagnet B field B_{EM} , for the four $|0\rangle \rightarrow |-1\rangle$ transitions (dashed lines). **c)** Photoluminescence of the NV centers as a function of the scanning magnetic field (blue crosses) and gaussian fit (orange line). **d)** Simulated population in $|m_s = 0\rangle$ for the stationary state with (plain) or without (dashed) taking into account the decrease of the T_1 induced by the cross-relaxations. **e)** Signal reflected off the diamond surface as a function of magnetic field amplitude B_{EM} (blue crosses). The orange line shows a double gaussian fit. **f)** Simulated torque applied by the spins on the diamond, with (plain) or without (dashed) taking into account the cross-relaxations.

Similarly to in Fig. 3 of the main text, we managed to mechanically detect other types of

mechanically induced dipolar couplings. Fig. 5b) shows a calculation of the frequencies of the $|0\rangle \rightarrow |-1\rangle$ transitions for all four classes of NV centers and the mechanically detected ESR spectra measured using the reflected laser for various magnetic field values. Unlike the experiment reported in the main text, this time all four classes of NV are resonant with another class at $B=23$ G. This indicates that we are crossing a $[100]_{\perp}$ plane instead of a $[110]_{\perp}$ one, as can be seen in Fig 5a).

Fig. 5c) shows the recorded photoluminescence of the NV centers during the magnetic field scan. As expected, a drop in the photoluminescence rate is observed when the degeneracy occurs. The PL drop is slightly more pronounced in this case compared to the experiments presented in the main text because all classes are depolarized here, instead of only two. This is well predicted by the numerical estimates shown in Fig. 5d).

Fig. 5e) shows the signal of the laser reflected off the diamond surface, proportional to the angular displacement. Here, there is a clear difference between the angular response and the experiment shown in the main text. Instead of a single Gaussian drop centered on the CR resonance, two bumps are observed on both sides of the resonance. Fig. 5f) shows the result of simulations, where we can see a dispersive profile with an almost zero torque at the resonance. The reason we do not observe a change of sign in the experiment (with two positive bumps instead of a positive and negative one) is likely to be because of the non-linearity of our detection : if the signal initially corresponds to a dark spot of the speckle, then a change in the motion of the diamond can only result in an increased signal.

Let us give a physical interpretation of the dispersive angular profile for this degeneracy condition. Here, the magnetic torque generated by the four classes of NV is not modified exactly on resonance, since all four classes are depolarized identically. Close to resonance, all classes will not be identically depolarized however : looking at Fig. 5b), we can see that the two classes which have a higher frequency are always slightly closer to each other than the two classes of lower frequency. This can be since by computing the slope in the evolution of the transitions frequencies, which are found to be smaller for the two upper classes. This effect results in more depolarization for these two classes, except when they are exactly on resonance. This interpretation explains the overall shape of the torque and thus of the angular response. A quantitative analysis would require knowledge about the directions of the 4-NV directions with respect to the three main directions of the trapped diamond axes as well as a calibration of the sensitivity of the speckle detection method for

the three corresponding angular modes.

V. SIMULATION DETAILS

In this part we will discuss the method used to simulate the average torque as well as the population in the $|m_s = 0\rangle$ state. Numerical solving of the master equation was performed using the Quantum Toolbox in Python (QuTiP) [4] [5].

In order to describe the dynamics of our spin ensemble, we introduce a incoherent optical pumping through the jump operators $\mathcal{L}_+ = \Gamma_l |0\rangle \langle +1|$ and $\mathcal{L}_- = \Gamma_l |0\rangle \langle -1|$, where $\Gamma_l \approx (2\pi)10$ kHz is the laser induced polarizing rate to the ground state. We also introduce the usual T_1 jump operators $\mathcal{L}_i^j = \frac{1}{T_1} |i\rangle \langle j|$ where $|i, j\rangle = |0, \pm 1\rangle$.

In order to describe the T_1 modification induced by the cross-relaxations, we use a phenomenological model where each class has its own T_1^i ($i \in \{1, 2, 3, 4\}$) that depends on the energy levels of the other classes with the formula :

$$\frac{1}{T_1^i} = \frac{1}{T_1^0} + \sum_{j \neq i} \frac{1}{T_1^{dd}} e^{-\frac{(\nu_i - \nu_j)^2}{2(\sigma^{dd})^2}},$$

where ν_i and ν_j are the transition frequencies of the classes i and j (we are arbitrarily considering the $|0\rangle \rightarrow |-1\rangle$ transition here, since the resonance condition is the same for both transitions at the magnetic fields we are working at. This is not always true for magnetic fields greater than 592 G[6]). σ^{dd} is the width of the CR features. We measured them to be similar to the inhomogeneous broadening given by the dipolar coupling to P_1 centers, *i.e.* ≈ 6 MHz.

$T_1^0 = 1.03$ ms and $T_1^{dd} = 0.38$ ms were chosen to match the T_1 measurements presented in Fig. 1 of the main text. We only focus on the T_1 without degeneracy and the one with a single degeneracy since our experiments will not have more than two resonant classes at once. Our model is probably not suited to deal with triple or quadruple resonances.

Finally, according to previous measurements performed in [7], only the $|0\rangle \langle \pm 1|$ and $|\pm 1\rangle \langle 0|$ (corresponding to a single quantum exchange in the dipole-dipole interaction) operators are modified by the cross-relaxations. Using this model, we can numerically solve the master equation and get the density matrix in the stationary state ρ_s . With ρ_s we can directly obtain the $|m_s = 0\rangle$ population, corresponding to the experimentally measured photoluminescence.

With regards to the torque estimation, we use a semi-classical formula :

$$\mathbf{\Gamma} = N_0 \gamma_e \langle \hat{\mathbf{S}} \rangle \times \mathbf{B},$$

where $N_0 \approx 10^9$ is an estimate of the number of spins in our sample based on the average size and NV density of our diamonds, γ_e is the gyromagnetic ratio of the electron and $\langle \hat{\mathbf{S}} \rangle = \text{Tr}(\rho_s \hat{\mathbf{S}})$ is the averaged spin vector in the stationary state, averaged again over the four possible orientations of NV. This formula assumes that the spin dynamics is faster than the dynamics of the motion of the diamond, which is the case in our experiments.

In our plots in Fig. 5 and Fig. 3 of the main text, we only represent one spatial component (e.g. Γ_x) of the torque, because the three components behave similarly.

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- [1] T. Delord, P. Huillery, L. Schwab, L. Nicolas, L. Lecordier, and G. Hétet, Phys. Rev. Lett. **121**, 053602 (2018).
 - [2] T. Delord, Ph.D. thesis, École Normale Supérieure (2019).
 - [3] T. Delord, P. Huillery, L. Nicolas, and G. Hétet, Nature **580**, 56 (2020).
 - [4] J. R. Johansson, P. D. Nation, and F. Nori, Computer Physics Communications **183**, 1760 (2012).
 - [5] J. Johansson, Comp. Phys. Comm **184**, 1234 (????).
 - [6] E. van Oort and M. Glasbeek, Phys. Rev. B **40**, 6509 (1989), ISSN 0163-1829, number: 10, URL <https://link.aps.org/doi/10.1103/PhysRevB.40.6509>.
 - [7] J. Choi, S. Choi, G. Kucsko, P. C. Maurer, B. J. Shields, H. Sumiya, S. Onoda, J. Isoya, E. Demler, F. Jelezko, et al., Phys. Rev. Lett. **118**, 093601 (2017), ISSN 0031-9007, 1079-7114, number: 9, URL <https://link.aps.org/doi/10.1103/PhysRevLett.118.093601>.