

The Dynamical Matrix

Recall: Hamiltonian: $\mathcal{H} = \sum_{\ell} \frac{P_{\ell}^2}{2M_{\ell}} + \frac{1}{2} \sum_{\ell, \ell'} \psi^{\circ}(\ell, \ell') u_{\ell} u_{\ell'}$

Classical eqn's of motion:

$$1. \frac{\partial \mathcal{H}}{\partial P_{\ell}} = \dot{R}_{\ell} = \ddot{u}_{\ell} \quad 2. \frac{\partial \mathcal{H}}{\partial R_{\ell}} = -\dot{P}_{\ell} = \frac{\partial \mathcal{H}}{\partial u_{\ell}}$$

$$1. \frac{\partial \mathcal{H}}{\partial P_{\ell}} = \ddot{u}_{\ell} \Rightarrow P_{\ell} = M_{\ell} \dot{u}_{\ell}$$

$$2. \frac{\partial \mathcal{H}}{\partial R_{\ell}} = -\dot{P}_{\ell} = \frac{\partial \mathcal{H}}{\partial u_{\ell}} \Rightarrow \frac{\partial \mathcal{H}}{\partial u_{\ell}} = \frac{1}{2} \frac{\partial}{\partial u_{\ell}} \left\{ \sum_{\ell''} \psi^{\circ}(\ell, \ell'') u_{\ell} u_{\ell''} \right\}$$

Note ℓ, ℓ', ℓ''
are just indices

$$\begin{aligned} &= \frac{1}{2} \left\{ \sum_{\ell'} \psi^{\circ}(\ell', \ell) u_{\ell'} + \sum_{\ell''} \psi^{\circ}(\ell, \ell'') u_{\ell''} \right\} \\ &= \sum_{\ell'} \psi^{\circ}(\ell, \ell') u_{\ell'} \end{aligned}$$

and: $\dot{P}_{\ell} = M_{\ell} \ddot{u}_{\ell}$

$$\therefore M_{\ell} \ddot{u}_{\ell} = - \sum_{\ell'} \psi^{\circ}(\ell, \ell') u_{\ell'}$$

OR $M_{\ell} \ddot{u}_{\ell} + \sum_{\ell'} \psi^{\circ}(\ell, \ell') u_{\ell'} = 0$

This is valid in general for an interaction potential energy function which is dependent on the ion positions.

$$M_e \ddot{u}_e + \sum_{e'} \psi^*(e, e') u_{e'} = 0$$

→ leads to set of N coupled diff. eqn's.

$$\text{e.g. } M_1 \ddot{u}_1 + \psi^*(1, 1) u_1 + \psi^*(1, 2) u_2 + \dots + \psi^*(1, N) u_N = 0$$

⋮

$$M_N \ddot{u}_N + \psi^*(N, 1) u_1 + \psi^*(N, 2) u_2 + \dots + \psi^*(N, N) u_N = 0$$

Assume solutions of the form:

$$u_e(t) = u_e^* e^{\pm i\omega t} \quad \xrightarrow{\text{amp. of oscillation}}$$

$$\Rightarrow -M_e \omega^2 u_e^* + \sum_{e'} \psi^*(e, e') u_{e'}^* = 0$$

or

$$\sum_{e'} \left[-M_e \omega^2 \delta_{ee'} + \psi^*(e, e') \right] u_{e'}^* = 0$$

$\underbrace{\quad}_{N \times N \text{ matrix.}} \equiv L$

$$\text{define: } -L_{ee'} \equiv -M_e \omega^2 \delta_{ee'} + \psi^*(e, e')$$

$$\Rightarrow L u^* = 0$$

$\begin{matrix} \nwarrow \\ N \times N \\ \text{matrix} \end{matrix}$

\curvearrowright vector of u_e^* 's

Recall: $\Psi^0(L, \ell') \equiv \left. \frac{\partial^2 \omega}{\partial R_{\ell'} \partial R_{\ell}} \right|_0 \rightarrow \text{real}$

$$\Psi^0(L, \ell') = \Psi^0(\ell', L)$$

$$L_{\ell, \ell'} = M_e \omega^2 f_{\ell, \ell'} - (\Psi^0(L, \ell'))$$

→ elements are real & symmetric

∴ Matrix L is Hermitian

Our set of equations: $\sum_{\ell'} L_{\ell, \ell'} \Psi^0_{\ell'} = 0$

has (nontrivial) solutions provided $|L| = 0$

→ Results in solution of N roots for ω^2

→ Since L is Hermitian, the roots i.e. ω^2 are real

→ while this corresponds to $2N$ soln for ω
 $(\pm \sqrt{\omega^2})$

$$\Psi_e(t) \propto e^{i\omega t} + e^{-i\omega t}$$

∴ still just N soln's. → 1D

For N degrees of freedom, N roots

i.e. 1D → N roots

2D → $2N$

3D → $3N$

→ each value of ω^2 represents a phonon frequency or normal mode of vibration

Assume solutions:

$$u^o(k) = A(k) \frac{e^{ik \cdot R^o}}{\sqrt{M_k}} \rightarrow \text{in 1D } R^o = a\ell.$$

↓ ↓

amplitude of mode k normalization

→ apply periodic boundary conditions: $u_\ell = u_{\ell+N}$

$$k = \frac{2\pi}{aN} m, \text{ m integer}$$

→ for 1D but
g.m. expressions for
 $2 \rightarrow 3$ D.

Recall, there are N modes, i.e. # of m 's. By convention, limit m to:

$$m \in \left(-\frac{N}{2} + 1, \frac{N}{2} \right) \Rightarrow \text{total } N \text{ modes}$$

Verify Solutions

$$\sum_{R'} L_{\ell R'} u_{\ell'}^o = 0 = \sum_{\ell'} \left(M_{\ell} \omega^2(k) f_{\ell\ell'} - \psi_{\ell\ell'}^o \right) A(k) \frac{e^{ikR_{\ell'}}}{\sqrt{M_{\ell'}}}$$

$$= M_{\ell} \omega^2(k) A(k) \frac{e^{ikR_{\ell}}}{\sqrt{M_{\ell}}} + - \sum_{\ell'} \frac{e^{ikR_{\ell'}}}{\sqrt{M_{\ell'}}} A(k) \psi_{\ell\ell'}^o e^{\underline{ik(R_{\ell'} - R_{\ell})}}$$

$= 1 \text{ when } \ell' = \ell$

$$O = \sqrt{M_2} A(k) e^{ikR_e^0} \left[\omega^2(k) - \sum_{\ell'} \frac{\psi^0(\ell, \ell')}{\sqrt{M_\ell M_{\ell'}^0}} e^{ik(R_{\ell'}^0 - R_\ell^0)} \right]$$

ω^0 

$= 0$

$$\therefore \omega^2(k) = \sum_{\ell'} \frac{\psi^0(\ell, \ell')}{\sqrt{M_\ell M_{\ell'}^0}} e^{ik(R_{\ell'}^0 - R_\ell^0)}$$

$\equiv D(k) \rightarrow$ Dynamical Matrix.

$$\omega^2(k) = D(k)$$

$$D(k) = \sum_{\ell'} \frac{\psi^0(\ell, \ell')}{\sqrt{M_\ell M_{\ell'}^0}} e^{ik(R_{\ell'}^0 - R_\ell^0)}$$

Note that $D(k)$ does not depend on ℓ :

\rightarrow take $M_{\ell'} = M_\ell = M$

$\rightarrow R_e^0$ are lattice vectors $\therefore R_{\ell'}^0 - R_\ell^0 = R_{\ell'-\ell}^0$ is a lattice vector

$\rightarrow \psi^0(\ell, \ell')$ depends on separation of $\ell \neq \ell'$

let $\ell' - \ell \equiv m$

$$D(k) = \frac{1}{M} \sum_m e^{ikR_m^0} \psi^0(m)$$

$$D(k) = \frac{1}{M} \sum_m e^{ikR_m^0} \phi^0(m) = \omega^2(k)$$

mass geometric phase factor force constants

generalization of:

$$\omega^2 = \frac{k}{m}$$

One more note:

$$\sum_{\ell \neq \ell'} \phi^0(\ell, \ell') = 0 \quad \rightarrow \text{think Newton's 3rd Law.}$$

i.e. $\sum_m \phi^0(m) = 0$

$$\Rightarrow \phi^0(0) + \sum_{m \neq 0} \phi^0(m) = 0$$

\hookrightarrow separate self interaction

$$\therefore D(k) = \frac{1}{M} \sum_{m \neq 0} e^{ikR_m^0} \phi^0(m) + \frac{1}{M} \phi^0(0)$$

$$= - \sum_{m \neq 0} \phi^0(m)$$

$$D(k) = \frac{1}{M} \sum_{m \neq 0} (e^{ikR_m^0} - 1) \phi^0(m)$$

geometric factor depends only on lattice

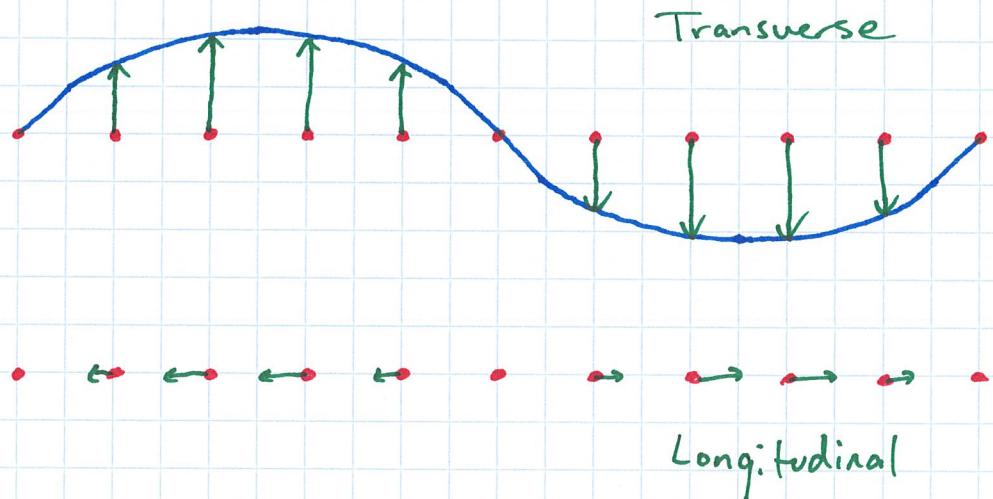
"self-force" components removed

Physical Interpretation of Solutions

$$u_k(k, t) = \frac{E(k)}{\sqrt{M}} e^{ikR^0} e^{-iw(k)t}$$

↳ displacement of ion k vibrating in mode k

- Normal modes are collective vibrations.



- sound waves are low frequency longitudinal & trans. phonons ($20\text{ Hz} - 20\text{ kHz}$)

$$\text{largest } \lambda? : k = \frac{2\pi}{\lambda} = \frac{2\pi m}{Na} \Rightarrow m=1 \Rightarrow \lambda = Na$$

$\underbrace{Na}_{\text{size of crystal}}$

$$\text{smallest } \lambda? : m = \frac{N}{2} \Rightarrow \lambda = 2a$$

↳ 2 unit cells