

# Jacobi theta functions from a Gaussian-cosecant convolution

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## 0.1 Background: The Faddeeva, or plasma dispersion, function

A function of some interest is the convolution of a Gaussian with a simple pole,

$$Z(\zeta) \equiv \frac{1}{\sqrt{\pi}} \int_{-\infty}^{\infty} \frac{e^{-z^2}}{z - \zeta} dz, \quad \text{Im}(\zeta) > 0 \quad (1)$$

defined initially in the upper half  $\zeta$ -plane. The function is defined and rescaled in various ways, and goes by names such as the Faddeeva function, Kramp function, or plasma dispersion function. Typically  $Z(\zeta)$  is defined in the upper half-plane and continued into the lower half. Continued in this way, a convenient form valid throughout the complex  $\zeta$ -plane is

$$Z(\zeta) = i\sqrt{\pi}e^{-\zeta^2}(1 + \text{erf}(i\zeta)). \quad (2)$$

## 0.2 A convolution with resonance at $z = n \in \mathbb{Z}$

The cosecant function  $\pi \csc(\pi z)$  is locally a simple pole at each integer, due to its Laurent series

$$\pi \csc(\pi z) = \sum_{n=-\infty}^{\infty} \frac{z(-1)^n}{z^2 - n^2} = \sum_{n=-\infty}^{\infty} \frac{(-1)^n}{z - n}, \quad (3)$$

$$\implies \lim_{z \rightarrow m} \pi \csc(\pi z) \rightarrow \frac{(-1)^m}{z - m}. \quad (4)$$

Now consider the convolution of a Gaussian with the cosecant function,

$$\mathcal{J}(\zeta) \equiv \frac{1}{\sqrt{\pi}} \int_{-\infty}^{\infty} \pi \csc(\pi(z - \zeta)) e^{-z^2} dz, \quad \text{Im}(\zeta) > 0. \quad (5)$$

### 0.2.1 Representation of the convolution as a series of Faddeeva functions

**Theorem 1** *One representation of the integral is*

$$\mathcal{J}(\zeta) = \sum_{n=-\infty}^{\infty} (-1)^n Z(\zeta - n) \quad (6)$$

*as a series in terms of, essentially, Faddeeva functions.*

**Proof:** *Apply the expansion in Eq. 3 and integrate term-by-term.*

An issue with this series representation is that in order to resolve the behavior at large real argument, terms are needed up to  $|n| \geq |\text{Re}(\zeta)|$ .

### 0.2.2 Representation of the convolution as a Fourier-like series

**Theorem 2** *An alternative representation of the integral is the Fourier-like series*

$$\mathcal{J}(\zeta) = 2\pi i \sum_{n=0}^{\infty} e^{2\pi i(n+\frac{1}{2})\zeta} e^{-(\pi(n+\frac{1}{2}))^2} \quad (7)$$

**Proof:** *To find an alternative representation, expand the cosecant in geometric series*

$$\pi \csc(\pi(z - \zeta)) = 2\pi i \frac{e^{-\pi i(z-\zeta)}}{1 - e^{-2\pi i(z-\zeta)}} \quad (8)$$

$$= 2\pi i \sum_{n=0}^{\infty} e^{-2\pi i(z-\zeta)(n+\frac{1}{2})}, \quad (9)$$

also called  $q$ -series in terms of the nome  $q = \exp(-2\pi i(z - \zeta))$ , convergent for  $\text{Im}(\zeta) > 0$ . Then  $\mathcal{J}(\zeta)$  may be integrated term-by-term as a series of Gaussian integrals

$$\mathcal{J} = 2\pi i \sum_{n=0}^{\infty} e^{2\pi i(n+\frac{1}{2})\zeta} \frac{1}{\sqrt{\pi}} \int_{-\infty}^{\infty} e^{-z^2 - 2\pi i z(n+\frac{1}{2})} dz \quad (10)$$

$$= 2\pi i \sum_{n=0}^{\infty} e^{2\pi i(n+\frac{1}{2})\zeta} e^{-(\pi(n+\frac{1}{2}))^2} \quad (11)$$

$$= 2\pi i e^{-\zeta^2} \sum_{n=0}^{\infty} e^{-(i\zeta - \pi(n+\frac{1}{2}))^2} \quad (12)$$

The last equality above is another representation similar in form to that for  $Z(\zeta)$  as the product of a Gaussian and an error function. However, the form as written in the theorem is more obviously convergent for  $\text{Im}(\zeta) > 0$ . It is interesting to observe that  $\mathcal{J}(\zeta)$  is “half” of a  $\vartheta$ -function.

### 0.2.3 Continuation of $\mathcal{J}(\zeta)$ into the lower half-plane

Now consider the extension of  $\mathcal{J}(\zeta)$  defined in Eq. 5 into the lower half-plane. The series of Eq. 7 in fact gives the desired function.

**Theorem 3** *The series given by Eq. 7 is valid throughout the entire complex  $\zeta$ -plane as the continuation given by adding the sum of residues when  $\zeta$  crosses the real axis from the upper half-plane.*

**Proof:** Denote the function of Eq. 5 as  $\mathcal{J}^+(\zeta)$ . Using the geometric series for  $\text{Im}(\zeta) < 0$  gives

$$\mathcal{J}^-(\zeta) = -2\pi i \sum_{n=0}^{\infty} e^{-2\pi i(n+\frac{1}{2})\zeta} e^{-(\pi(n+\frac{1}{2}))^2}. \quad (13)$$

Then the jump across the real axis is given by

$$\Delta\mathcal{J} = \mathcal{J}^+ - \mathcal{J}^- = 2\pi i \sum_{n=0}^{\infty} e^{-(\pi(n+\frac{1}{2}))^2} (e^{2\pi i(n+\frac{1}{2})\zeta} + e^{-2\pi i(n+\frac{1}{2})\zeta}) \quad (14)$$

$$= 4\pi i \sum_{n=0}^{\infty} \cos\left(\pi(1+2n)\zeta\right) e^{-(\pi(n+\frac{1}{2}))^2} \quad (15)$$

$$= 2\pi i \vartheta_2(\pi\zeta, e^{-\pi^2}) \quad (16)$$

in terms of the Jacobi theta function  $\vartheta_2(z, q) = 2 \sum_{n=0}^{\infty} \cos((1+2n)z) q^{(n+\frac{1}{2})^2}$ .

On the other hand, the residue of the integrand at any simple pole  $z = \zeta + n$  is

$$\text{Res}\left(\pi \csc(\pi(z - \zeta)) e^{-\zeta^2}\right) \Big|_{z=\zeta+n} = (-1)^n e^{-(\zeta+n)^2} \quad (17)$$

$$\Rightarrow \sum \text{Res}\left(\pi \csc(\pi(z - \zeta)) e^{-\zeta^2}\right) = \sum_{n=-\infty}^{\infty} (-1)^n \frac{e^{-(\zeta+n)^2}}{\sqrt{\pi}} \quad (18)$$

$$= e^{i\pi\zeta} \sum_{n=-\infty}^{\infty} \frac{1}{\sqrt{\pi}} e^{-(\zeta+n)^2} e^{-i\pi(\zeta+n)} \quad (19)$$

Recall Poisson's summation formula relating a periodic summation to one of its Fourier transform,

$$\sum_{n=-\infty}^{\infty} f(x+n) = \sum_{n=-\infty}^{\infty} \hat{f}(n) e^{i2\pi nx}, \quad (20)$$

and observe that the Fourier transform of the summand  $f(x+n)$  in Eq. 19 is  $\hat{f} = e^{-(\pi(m+\frac{1}{2}))^2}$ , so

$$\sum_{n=-\infty}^{\infty} (-1)^n \frac{e^{-(\zeta+n)^2}}{\sqrt{\pi}} = \sum_{n=-\infty}^{\infty} e^{2\pi i(n+\frac{1}{2})\zeta} e^{-(\pi(n+\frac{1}{2}))^2} = \vartheta_2(\pi\zeta, e^{-\pi^2}). \quad (21)$$

Therefore,  $\Delta\mathcal{J} = 2\pi i \sum(\text{Res}(\mathcal{J}^+))$ . Adding the sum of residues back to  $\mathcal{J}^-$  in the lower half-plane results in the desired continued function,

$$\mathcal{J}(\zeta) = \mathcal{J}^- + 2\pi i \sum(\text{Res}(\mathcal{J}^+)) = \mathcal{J}^- + \mathcal{J}^+ - \mathcal{J}^- = \mathcal{J}^+, \quad \text{Im}(\zeta) < 0. \quad (22)$$

This proves the theorem.

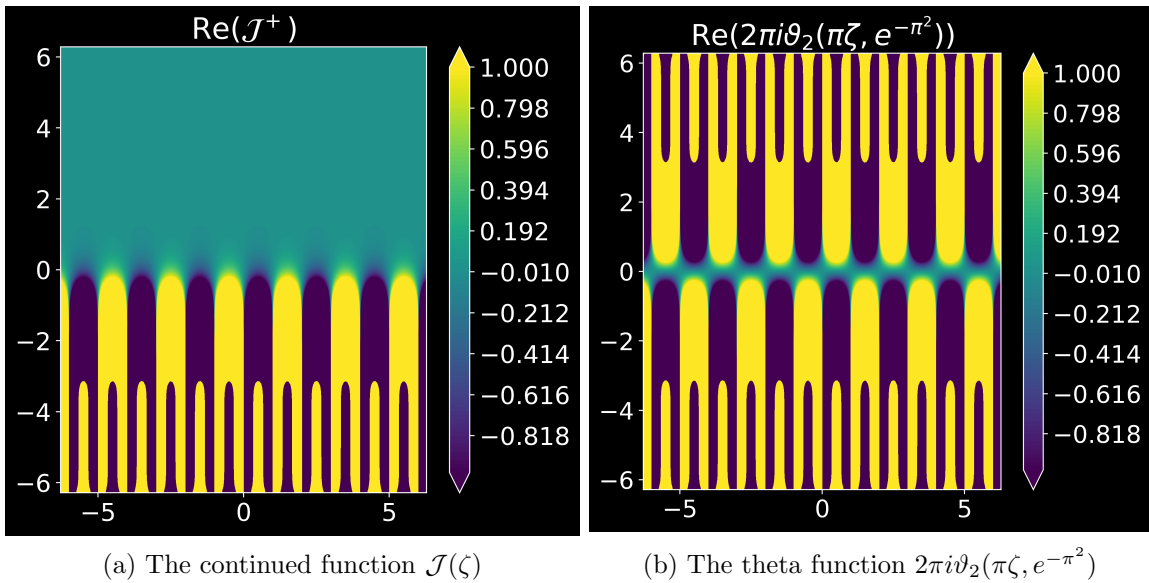


Figure 1: Comparison of the function considered (“half” a  $\vartheta$ -function) and a full theta function, given by the sum of residues of the integrand. Plots are on the complex  $\zeta$ -plane.