

Explorations into the Finite Element Method

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Main test case:

$$-u''(x) + q(x)u(x) = f(x)$$

In order to use the weak form of the ODE we must first employ the "Method of Weighted Residuals". The residual for our case is simply:

$$R(x) = -u''(x) + q(x)u(x) - f(x)$$

Thus to complete the method of weighted residuals we have

$$\int w(x)R(x) = \int w(x) \left(-u''(x) + q(x)u(x) - f(x) \right)$$

Note that through integration by parts:

$$\int -w(x) \frac{d^2 u}{dx^2} = \int \frac{dw}{dx} \frac{du}{dx} dx - \int w(x) \frac{du}{dx} d\Gamma$$

where the $d\Gamma$ integral refers to the boundary condition.

Thus plugging in we get,

$$\int \frac{dw}{dx} \frac{du}{dx} + q(x)w(x)u(x)dx = \int f(x)w(x)dx + \int w(x) \frac{du}{dx} d\Gamma$$

Now we must make the leap and approximate our solution $u(x)$ as

$$u_{app} = \sum_{j=1}^n U_j S_j(x)$$

where $S_j(x) = \delta_{ij}$ Substituting this in we get,

$$\int \frac{dw}{dx} \sum_{j=1}^n U_j S_j(x) + q(x)w(x) \sum_{j=1}^n U_j S_j(x)dx = \int f(x)w(x)dx + \int w(x)SV d\Gamma$$

where SV refers to the boundary values.

$$\sum_{j=1}^n \left(\int \frac{dw}{dx} \frac{dS_j(x)}{dx} + q(x)w(x)S_j(x)dx \right) U_j = \int f(x)w(x)dx + \int w(x)SV d\Gamma$$

In other notation,

$$[K]\{U\} = \{F\} + \{B\}$$

where K is the stiffness matrix. Since we are using the Galerkin method $w(x) = S_i$

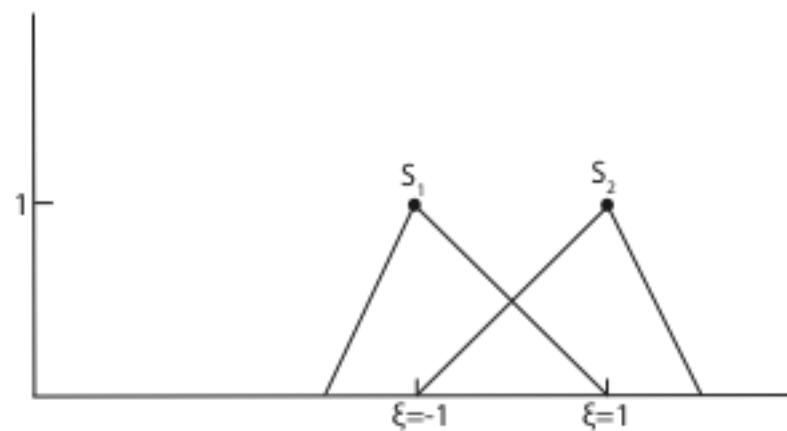
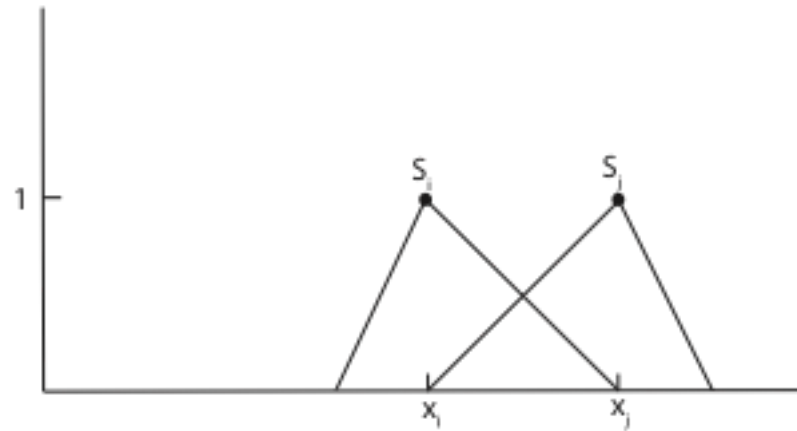
$$K_{ij} = \int \frac{dS_i}{dx} \frac{dS_j}{dx} + q(x)S_iS_j dx$$

$$F_i = \int f(x)S_i dx$$

$$B \rightarrow \text{Boundary}$$

For our $S_j(x) = \delta_{ij}$ we examine one element to see our change of coordinates.

Here is our initial element using the "hats" - kronecker delta. The first image is a pre-transformation image and the final image is post-transformation



Hence $S_1 = \frac{1}{2}(1 - \xi)$ and $S_2 = \frac{1}{2}(1 + \xi)$ which correlates to a transformation of

$$x = \frac{h^e}{2}\xi + \frac{x_1^e + x_2^e}{2}$$

where $h^e = x_2^e - x_1^e$

Therefore, with our coordinate change, we can rewrite our stiffness matrix noting that $\frac{dS}{dx} = \frac{dS}{d\xi} \frac{d\xi}{dx}$

$$K_{ij} = \int_{-1}^1 \left(\left(\frac{dS_i}{d\xi} \frac{d\xi}{dx} \right) \left(\frac{dS_j}{d\xi} \frac{d\xi}{dx} \right) + q S_i S_j \right) \frac{dx}{d\xi} d\xi$$

It is worth noting that

$$\frac{dx}{d\xi} = \frac{h^e}{2} = J^e$$

Thus we have,

$$K_{ij} = \int_{-1}^1 \left(\left(\frac{dS_i}{d\xi} \frac{1}{J^e} \right) \left(\frac{dS_j}{d\xi} \frac{1}{J^e} \right) + q S_i S_j \right) J^e d\xi$$

And

$$F_I^e = \int_{-1}^1 S_i f(\xi) J^e d\xi$$

Since we only have 2 S equations we need to calculate only $K_{11}, K_{12}, K_{21}, K_{22}$ and F_1, F_2 and then assemble the matrix based on them, created a tridiagonal matrix!

Implementation of Code on a mini test case

In [1]:

```
import numpy as np
import scipy.misc
import scipy as sp
import matplotlib.pyplot as plt
%matplotlib inline
from FEM import Stiff,F
```

Our ODE is

$$-u'' + 0 * u = 1$$

which is simplistic since we allow q to be zero and f is a constant.

Our ODE has the solution $u = x(1 - x)$

In [2]:

```
n_elements = 100
s_elements = 2
start = 0
end = 1
q = 0 #q function
```

And now for our equations of the stiff matrix and transformations:

In [3]:

```
def S_i(x):
    return (1/2)*(1-x)

def S_j(x):
    return (1/2)*(1+x)

def K(S_i,S_j,dS_i,dS_j,q,J):
    return ((dS_i*(1/J))*(dS_j*(1/J))+q*S_i*S_j)*J

def f(x):
    return 1.
```

In [4]:

```
prac_space = np.linspace(start,end,n_elements)
prac_k = Stiff(K,S_i,S_j,q,n_elements,s_elements,start,end)
prac_k
```

Out[4]:

```
array([[ 200., -100.,   0., ...,   0.,   0.,   0.],
       [-100.,  200., -100., ...,   0.,   0.,   0.],
       [   0., -100.,  200., ...,   0.,   0.,   0.],
       ...,
       [   0.,   0.,   0., ..., 200., -100.,   0.],
       [   0.,   0.,   0., ..., -100.,  200., -100.],
       [   0.,   0.,   0., ...,   0., -100.,  200.]])
```

In [5]:

```
prac_f = F(f,S_i,S_j,n_elements,s_elements,start,end)
prac_f[0] *= 2.
prac_f[-1] *= 2.
```

In [6]:

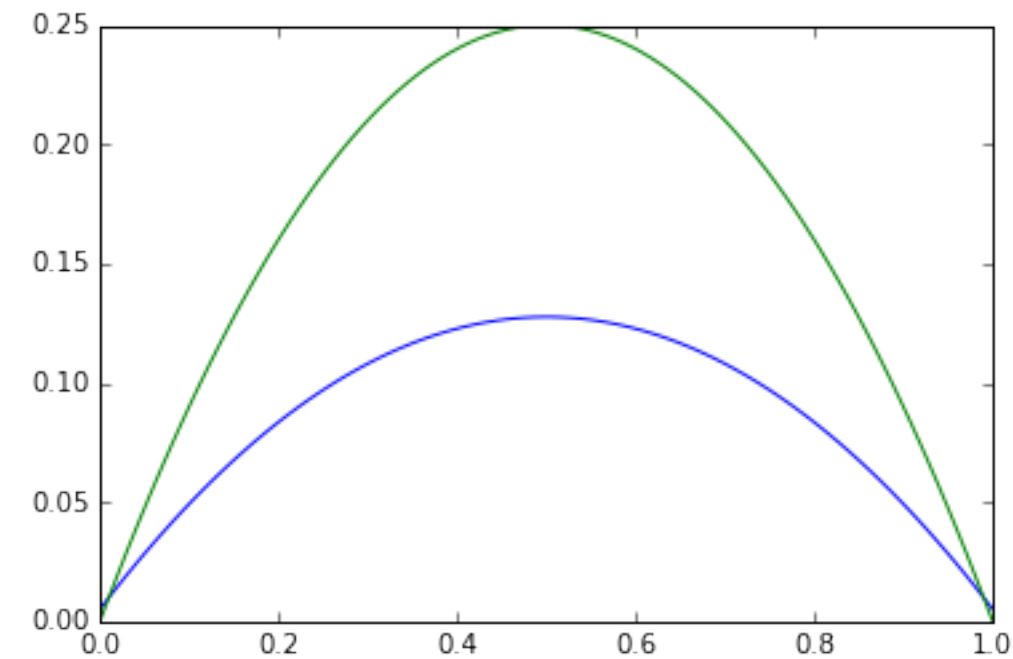
```
exact = lambda x: x*(1-x)
exact_val = np.zeros(n_elements)
for i in range(n_elements):
    exact_val[i] = exact(prac_space[i])
```

In [7]:

```
plt.plot(prac_space, sp.linalg.solve(prac_k, prac_f))
plt.plot(prac_space, exact_val)
```

Out[7]:

[<matplotlib.lines.Line2D at 0x1060f58d0>]



Test Case:

For a test case I will be solving the equation

$$-y'' + y = \sin(ax), \quad t \in (0, \pi)$$

both analytically and numerically in order to asses errors.

Numerical Solving

In [8]:

```
n_elements = 100
s_elements = 2
start = 0
end = 2*np.pi
q = 1
a = 1
B_init = np.zeros(n_elements)
B_start = np.sin(a*start)
B_end = np.sin(a*end)
B_init[0] = B_start
B_init[-1] = B_end

def S_i(x):
    return (1/2)*(1-x)

def S_j(x):
    return (1/2)*(1+x)

def K(S_i,S_j,dS_i,dS_j,q,J):
    return ((dS_i/J)*(dS_j/J)+q*S_i*S_j)*J

def f(x):
    return np.sin(a*x)
```

In [9]:

```
spacing = np.linspace(start,end,n_elements)
```

In [10]:

```
K_matrix = Stiff(K,S_i,S_j,q,n_elements,s_elements,start,end);
F_matrix = F(f,S_i,S_j,n_elements,s_elements,start,end);
B_matrix = B_init
```

Solving the matrix via my software

In [11]:

```
from FEM import LUalgobanded, solveUnitLowerTriangularbanded, solveUnitUpperTriangu
```

Since we know the boundary conditions, we can nix the initial and final row/column. So keep in mind that with our example we let the boundary conditions be such that the $x_{initial} = 0$ and $x_{final} = 0$.

In [12]:

```
L_stiff,U_stiff = LUalgobanded(K_matrix,1)
```

In [13]:

```
C_stiff = solveUnitLowerTriangularbanded(L_stiff,F_matrix,1)
```

In [14]:

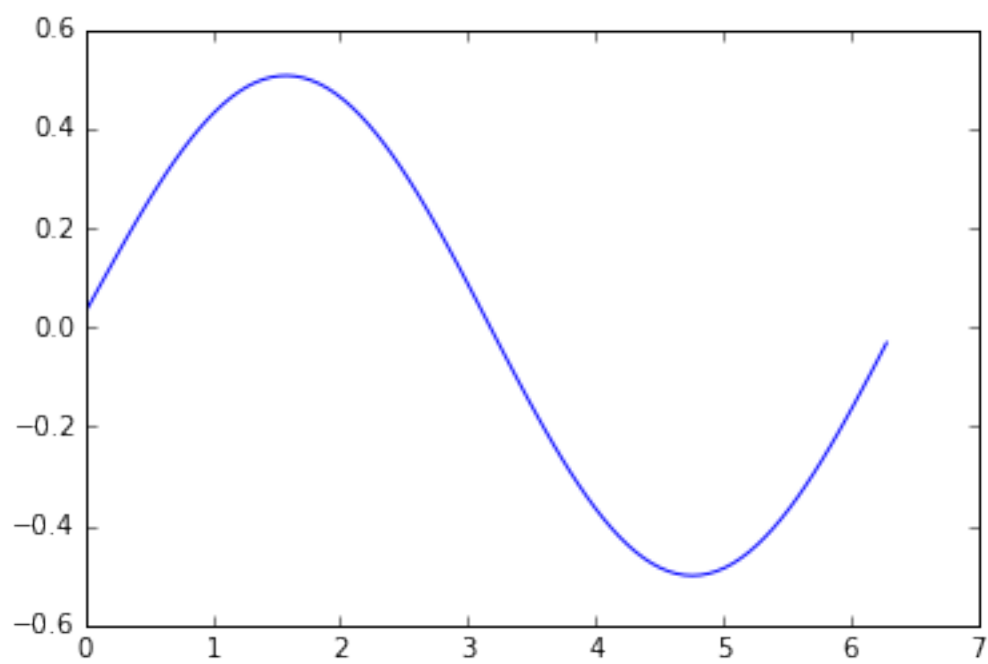
```
Sol_stiff = np.zeros(n_elements)  
Sol_stiff = solveUnitUpperTriangularbanded(U_stiff,C_stiff,1)
```

In [15]:

```
plt.plot(spacing,Sol_stiff)
```

Out[15]:

[<matplotlib.lines.Line2D at 0x106d81208>]



Solve with Scipy.linalg.solve

In [16]:

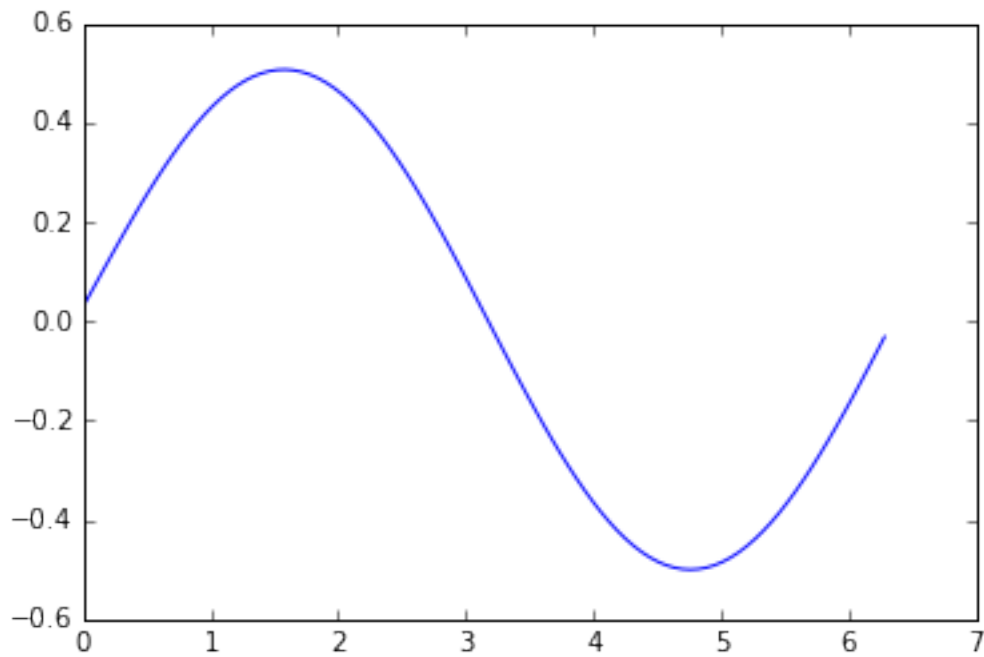
```
Scipy_sol = sp.linalg.solve(K_matrix,F_matrix+B_matrix)
```

In [17]:

```
plt.plot(spacing,Scipy_sol)
```

Out[17]:

```
[<matplotlib.lines.Line2D at 0x106f6ce48>]
```



Exact Solution

Let's quickly find the exact solution.

$$y'' + y = \sin(ax)$$

So my guess is $y = A\sin(ax) \rightarrow y'' = -Aa^2\sin(ax)$ Hence plugging in we get,

$$Aa^2\sin(ax) + A\sin(ax) = \sin(x)$$

$$A = \frac{1}{1+a^2}$$

Therefore,

$$y = \frac{1}{1+a^2}\sin(ax)$$

In [18]:

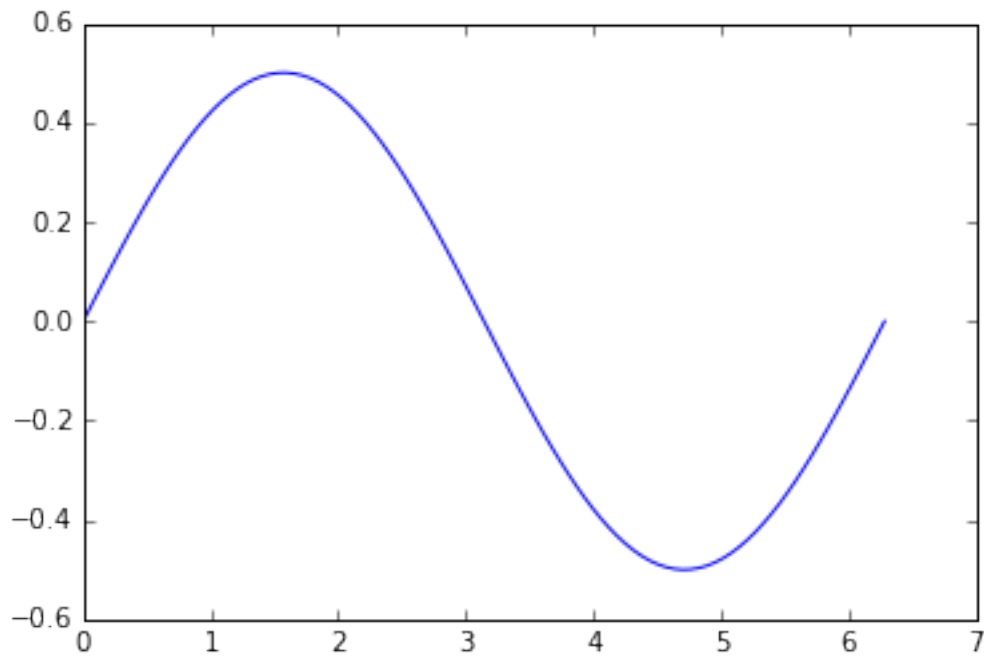
```
exact = lambda x: (1/(1+a**2))*np.sin(a*x)
exact_val = np.zeros(n_elements)
for i in range(n_elements):
    exact_val[i] = exact(spacing[i])
```


In [19]:

```
plt.plot(spacing,exact_val)
```

Out[19]:

```
[<matplotlib.lines.Line2D at 0x107035d30>]
```



Errors:

We are going to check the errors comparing the exact solution to the scipy solving software and my banded algo software.

In [20]:

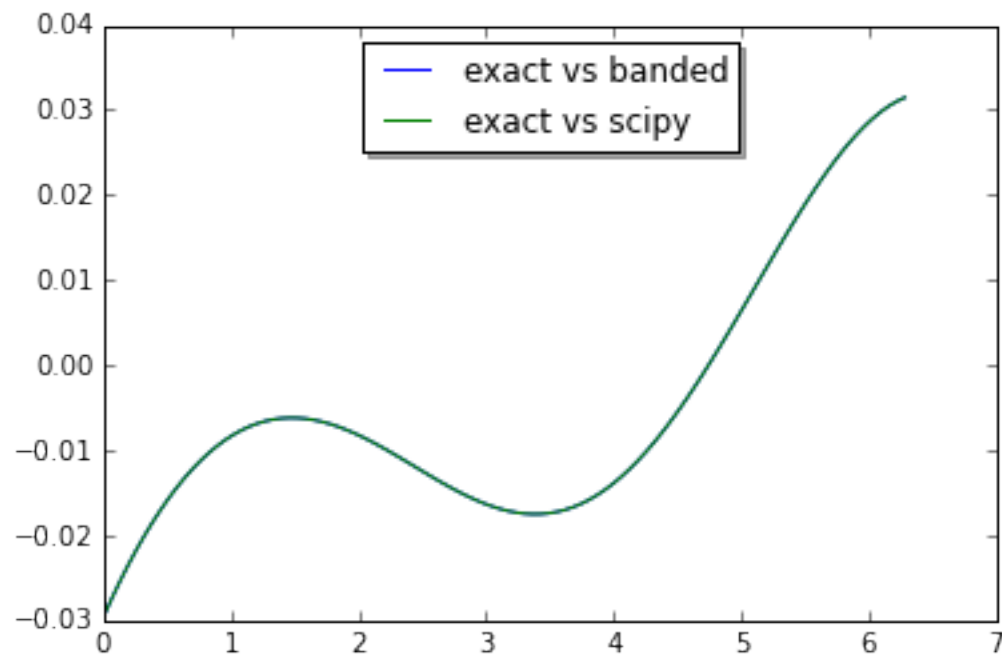
```
exact_vs_banded_error_max = sp.linalg.norm(exact_val-Sol_stiff,np.inf)
exact_vs_scipy_error_max = sp.linalg.norm(exact_val-Scipy_sol,np.inf)
print("The maximum error between the exact solution and my banded algorithm solution is",exact_vs_banded_error_max)
print("The maximum error between the exact solution and the scipy algorithm solution is",exact_vs_scipy_error_max)
```

The maximum error between the exact solution and my banded algorithm solution is 0.031471016971

The maximum error between the exact solution and the scipy algorithm solution is 0.031471016971

In [21]:

```
exact_vs_banded_line = plt.plot(spacing,exact_val-Sol_stiff,label="exact vs banded")
exact_vs_scipy_line = plt.plot(spacing,exact_val-Scipy_sol,label="exact vs scipy")
legend = plt.legend(loc='upper center', shadow=True)
```



Hence we can say that my banded software is doing essentially the same as scipy's built in software. More importantly, my Finite Element Method code is functioning properly. For further test cases I will just be using scipy software to solve the matrix equation.

Next Main Test Form:

$$-u'' + p(x)u' + q(x)u = f(x)$$

To start solving lets first let $A(x) = e^{-\int p(x)dx}$.

Thus multiplying our test equation by A we get,

$$-A(x)u'' + A(x)p(x)u' + A(x)q(x)u = A(x)f(x)$$

Since $Au'' + A(x)p(x)u' = A(x)u'$ by the chain rule, we can fix up our equation to:

$$-(A(x)u')' + Q(x)u = F(x)$$

where $Q(x) = A(x)q(x)$ and $F(x) = A(x)f(x)$

Once again using the "Method of Weighted Residuals" in order to obtain the Weak Form of the ODE...

$$\int -w(x)(A(x)u')' + w(x)Q(x)u = \int w(x)F(x)$$

It is certainly worth noting that

$$\int -w(A(x)u')' = \int \frac{dw}{dx}A(x)u' dx - \int w(x)A(x)u' d\Gamma$$

where $d\Gamma$ just represents our boundary. Now everything proceeds like the previous example:

$$\int A(x) \frac{dw}{dx} \frac{du}{dx} + Q(x)w(x)u(x)dx = \int F(x)w(x)dx + \int A(x)w(x) \frac{du}{dx} d\Gamma$$

Now we must make the leap and approximate our solution $u(x)$ as

$$u_{app} = \sum_{j=1}^n U_j S_j(x)$$

where $S_j(x) = \delta_{ij}$ Substituting this in we get,

$$\int A(x) \frac{dw}{dx} \sum_{j=1}^n U_j S_j(x) + Q(x)w(x) \sum_{j=1}^n U_j S_j(x)dx = \int F(x)w(x)dx + \int A(x)w(x)SVd\Gamma$$

where SV refers to the boundary values.

$$\sum_{j=1}^n \left(\int A(x) \frac{dw}{dx} \frac{dS_j(x)}{dx} + Q(x)w(x)S_j(x)dx \right) U_j = \int F(x)w(x)dx + \int A(x)w(x)SVd\Gamma$$

In other notation,

$$[K]\{U\} = \{F\} + \{B\}$$

where K is the stiffness matrix. Since we are using the Galerkin method $w(x) = S_i$

$$K_{ij} = \int A(x) \frac{dS_i}{dx} \frac{dS_j}{dx} + Q(x)S_i S_j dx$$

$$F_i = \int F(x)S_i dx$$

$$B \rightarrow \text{Boundary}$$

Using the same hat functions as before the transformation follows in the same exact manner... Hence,

$$K_{ij} = \int_{-1}^1 \left(A(x) \left(\frac{dS_i}{d\xi} \frac{1}{J^e} \right) \left(\frac{dS_j}{d\xi} \frac{1}{J^e} \right) + Q(x)S_i S_j \right) J^e d\xi$$

And

$$F_I^e = \int_{-1}^1 S_i F(\xi) J^e d\xi$$

In []: