CSE291 Topics in Computer Graphics Mesh Animation

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Today

- Review introduction to FEM
- Elastic bodies in 3D

Finite Element Analysis (FEA)

- Class of techniques to solve PDEs
- FEM recipe
 - 0. Strong formulation
 - 1. Weak formulation
 - 2. Trial solutions and test functions
 - 3. Galerkin approximation
 - 4. Matrix formulation

1D example

- One dimensional bar, unit length
- Subject to distributed load (forces) f(t)
- · Boundary conditions
- Find longitudinal displacement u(t)



Strong (differential) formulation

• Constitutive equation (1D Poisson problem)

$$\frac{d^2u(t)}{dt^2} + f(t) = 0$$

Strong (differential) formulation

• Constitutive equation (1D Poisson problem)

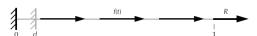
$$\frac{d^2u(t)}{dt^2} + f(t) = 0$$

- · Boundary conditions
 - Prescribed displacement at the beginning (essential or geometric boundary condition)

$$u(0) = d$$

- Concentrated force at the end (*natural* or *force* boundary condition)

$$u'(1) = R$$



Weak formulation

• Constitutive equation holds "in average"

$$-\int_{0}^{1} v(t) \left(\frac{d^{2}u(t)}{dt^{2}} + f(t) \right) dt = 0 \quad (\#)$$

where v(t) is a test function

Above needs to be true for arbitrary test function

Boundary conditions

ullet Restrict $trial\ solutions\ u(t)\$ to satisfy $geometric\$ boundary condition

$$u(t) \in \mathcal{U} = \{u(t)|u(0) = d\}$$

• Restrict test function v(t) to satisfy homogeneous boundary conditions

$$v(t) \in \mathcal{V} = \{v(t)|v(0)=0\}$$

Weak formulation

- Integrate (#) by parts
- Principle of *virtual displacements* or principle *of virtual work*

$$\int_0^1 \frac{du}{dt} \frac{dv}{dt} dt = \int_0^1 v f dt + Rv(1)$$

Internal work

External work

• Body is in equilibrium

Solving the weak formulation

• Find $u(t) \in \mathcal{U}$ such that for all $v(t) \in \mathcal{V}$

$$\int_0^1 \frac{du}{dt} \frac{dv}{dt} dt = \int_0^1 v f dt + Rv(1)$$

Galerkin approximation

 Restrict test functions and trial solutions to finite dimensional function spaces

$$v^h(t) = \sum_{i=1}^n \hat{v}_i N_i(t) \in \mathcal{V}^h \subset \mathcal{V}$$

$$u^h(t) = \sum_{i=1}^n \hat{w}_i N_i(t) + dN_0(t) \in \mathcal{U}^h \subset \mathcal{U}$$

$$N_i(0) = 0, i = 1, \dots, n$$
 and $N_0(0) = 1$

 $N_i(t)$ basis (shape) functions

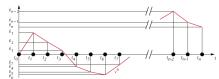
 \hat{v}_i, \hat{w}_i scalar weights

Galerkin approximation

• Shape functions usually have local support (hence the term *finite element*)



ullet Example member of \mathcal{V}^h



Galerkin approximation

• Plug test, trial functions into weak form

$$\sum_{j=1}^{n} \hat{w}_{j} a(N_{i}, N_{j}) = (f, N_{i}) + N_{i}(1)R - da(N_{i}, N_{0})$$
for $i = 1, \dots, n$

where

$$a(N_i, N_j) = \int_0^1 \frac{dN_i}{dt} \frac{dN_j}{dt} dt$$

$$(f, N_i) = \int_0^1 N_i f dt$$

Matrix formulation

• Define

$$K_{ij} = a(N_i, N_j)$$

$$F_i = (f, N_i) + N_i(1)R - da(N_i, N_0)$$

· Matrix form

$$\mathbf{K}\mathbf{w} = \mathbf{F}$$

Stiffness matrix ${f K}$ Force vector ${f F}$

Matrix formulation

Problem statement

- Given stiffness matrix and force vector, find parameters of displacement function
- \bullet I.e., solve $\mathbf{K}\mathbf{w}=\mathbf{F}$ for \mathbf{w}
- Stiffness matrix is positive definite, sparse
- Use standard techniques, e.g., conjugate gradient solvers

Summary

Basic "FEA recipe" (one of many)

- Derive weak formulation of PDE
- Multiply PDE with test function
 - Integrate, apply integration by parts
- Specify constraints on test and trial functions to fulfill boundary conditions
- Galerkin approximation
 - Discretize test and trial functions
 - Plug into weak formulation
 - Derive matrix equation

Questions?

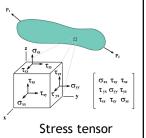
Linear elastic bodies

Physically-based model

- Stress, body forces
- Equilibrium conditions
- Deformations and displacements
- Strair
- Stress-strain relationship, constitutive equations

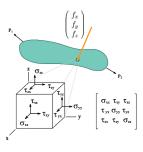
Stress

- Forces exerted by a piece of material onto its environment
- Measured as force per unit area
- Normal, shear stress
- Stress tensor
- Symmetric!



Body forces

- External force applied to a piece of material
- Force per volume



Equilibrium conditions

• Independent of material properties

$$\begin{pmatrix} \frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \tau_{xy}}{\partial y} + \frac{\partial \tau_{xz}}{\partial z} + f_x \\ \frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} + \frac{\partial \tau_{yz}}{\partial z} + f_y \\ \frac{\partial \tau_{xz}}{\partial x} + \frac{\partial \tau_{yz}}{\partial y} + \frac{\partial \sigma_{zz}}{\partial z} + f_z \end{pmatrix} = 0$$

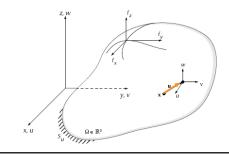
where f_x, f_y, f_z are body forces (force per volume)

· Holds for fluids too

Questions?

Deformations and displacments

- Deformation represented by displacement field $\mathbf{u}(\mathbf{x})$



Strain

- Geometric deformation
- Green tensor, or Cauchy's infinitesimal strain tensor

$$\varepsilon_{ij} = \frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right)$$

where u is the diplacement field

- "Symmetric part" of Jacobian of u
- Also called "kinematic equation"

Strain

• Matrix form

$$\varepsilon = \begin{bmatrix} \varepsilon_{xx} \\ \varepsilon_{yy} \\ \varepsilon_{zz} \\ \varepsilon_{xy} \\ \varepsilon_{yz} \\ \varepsilon_{zx} \end{bmatrix} = \begin{bmatrix} \frac{\partial}{\partial x} & 0 & 0 \\ 0 & \frac{\partial}{\partial y} & 0 \\ 0 & 0 & \frac{\partial}{\partial z} \\ \frac{\partial}{\partial y} & \frac{\partial}{\partial x} & 0 \\ 0 & \frac{\partial}{\partial z} & \frac{\partial}{\partial y} \end{bmatrix} \begin{bmatrix} u \\ v \\ w \end{bmatrix}$$

Stress-strain relationship

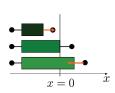
- Given geometric deformation of a piece of a body, what are strains (forces)?
- Stress-strain relationship captures physical material properties

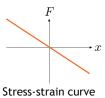
Linear-elastic bodies

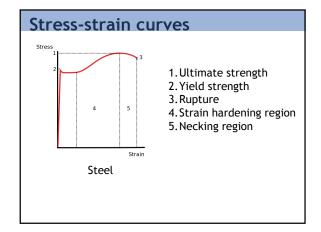
• Hooke's law in 1D (springs)

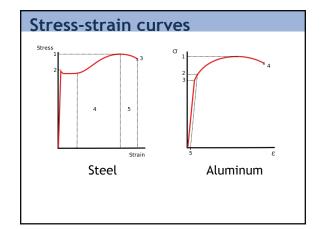
$$F = -kx$$

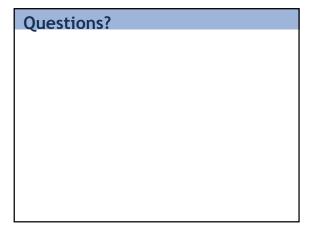
force F, spring constant k, displacement from rest length \boldsymbol{x}











Linear elastic bodies

- Generalization of Hooke's law to 3D bodies
- · Isotropic materials

$$\sigma = \begin{bmatrix} \sigma_{xx} \\ \sigma_{yy} \\ \sigma_{zz} \\ \tau_{zy} \\ \tau_{zz} \\ \tau_{yz} \\ \tau_{zz} \end{bmatrix} = \frac{E}{d} \begin{bmatrix} (1-\nu) & \nu & \nu & 0 & 0 & 0 & 0 \\ \nu & (1-\nu) & \nu & 0 & 0 & 0 & 0 \\ \nu & \nu & (1-\nu) & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & (1-2\nu)/2 & 0 & 0 \\ 0 & 0 & 0 & 0 & (1-2\nu)/2 & 0 \\ 0 & 0 & 0 & 0 & 0 & (1-2\nu)/2 \end{bmatrix} \begin{bmatrix} \varepsilon_{xx} \\ \varepsilon_{yy} \\ \varepsilon_{xz} \\ \varepsilon_{xy} \\ \varepsilon_{zy} \\ \varepsilon_{zz} \\ \varepsilon_{zz} \end{bmatrix}$$

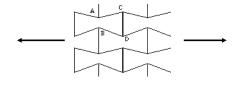
$$d = (1 + \nu)(1 - 2\nu)$$

Linear elastic bodies

- E Young's modulus of elasiticity, slope of stress-strain curve
- ν Poisson's ratio
- Measure of how much thinner the material gets in one direction as you pull in the other
- Indicates change in volume
- 0.5 corresponds to no change in volume, i.e., incompressible

Negative Poisson's ratio

• http://silver.neep.wisc.edu/~lakes/Poisson.html



Constitutive equations

Strong formulation

• Equilibrium conditions

$$\mathcal{A}(\mathbf{u}) = \begin{pmatrix} A_1 \\ A_2 \\ A_3 \end{pmatrix} = \begin{pmatrix} \frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \tau_{xy}}{\partial y} + \frac{\partial \tau_{xz}}{\partial z} + f_x \\ \frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} + \frac{\partial \tau_{yz}}{\partial z} + f_y \\ \frac{\partial \tau_{xx}}{\partial x} + \frac{\partial \tau_{yz}}{\partial y} + \frac{\partial \sigma_{zz}}{\partial z} + f_z \end{pmatrix}$$

and linear stress-strain relationship

• Linear elastic bodies: linear displacementstrain, linear strain-stress, linear constitutive equations

Questions?

FEA recipe

- Derive weak formulation of PDE
 - Multiply PDE with test function
 - Integrate, apply integration by parts
 - Specify constraints on test and trial functions to fulfill boundary conditions
- Galerkin approximation
 - Discretize test and trial functions
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Derivation of weak form

• Multiplication of constitutive equations by test function $\mathbf{v} \equiv \delta \mathbf{u} = [\delta u, \delta v, \delta w]^T$ (also called *virtual displacements*)

$$\begin{split} & \int_{\Omega} \mathbf{v}^{T} \mathcal{A}(\mathbf{u}) d\Omega = \\ & - \int_{\Omega} \delta u \left(\frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \tau_{xy}}{\partial y} + \frac{\partial \tau_{xz}}{\partial z} + f_{x} \right) + \delta v(A_{2}) + \delta w(A_{3}) d\Omega \end{split}$$

- Integration by parts (Green's formula)
- · No details here

Weak form

• Weak form, or principle of virtual work

$$\int_{\Omega} \delta \varepsilon^T \sigma d\Omega - \int_{\Omega} \delta \mathbf{u}^T \mathbf{f} d\Omega = 0$$

- Virtual strain $\delta \varepsilon \in \mathbf{R}^{6 \times 1}$
- Strain vector $\sigma \in \mathbf{R}^{6 \times 1}$

Weak form

· Virtual strain

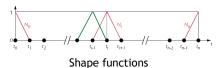
$$\delta \varepsilon = \left(\frac{\partial \delta u}{\partial x}, \frac{\partial \delta v}{\partial y}, \frac{\partial \delta w}{\partial z}, \frac{\partial \delta u}{\partial y} + \frac{\partial \delta v}{\partial x}, \ldots\right)^T = \mathcal{S} \delta \mathbf{u}$$

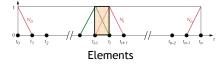
Matrix form

$$\mathcal{S}\delta\mathbf{u} = \begin{pmatrix} \frac{\partial}{\partial x} & 0 & 0\\ 0 & \frac{\partial}{\partial y} & 0\\ 0 & 0 & \frac{\partial}{\partial z}\\ \frac{\partial}{\partial y} & \frac{\partial}{\partial z} & 0\\ 0 & \frac{\partial}{\partial z} & \frac{\partial}{\partial y}\\ \frac{\partial}{\partial z} & 0 & \frac{\partial}{\partial z} \end{pmatrix} \begin{pmatrix} \delta u\\ \delta v\\ \delta w \end{pmatrix}$$

Discretization

Elements vs. shape functions





Discretization

- · Galerkin approximation
- · Element point of view
- Displacement in each element is

$$\mathbf{u}^{(m)}(x,y,z) = \sum_{i=0}^{k-1} \hat{\mathbf{u}}_i^{(m)} N_i^{(m)}(x,y,z)$$

where $\hat{\mathbf{u}}_i^{(m)} \in \mathbf{R}^3$ number of shape functions in element k and (x,y,z) lies within element m

Discretization

Maintain map relating element indices and global indices

$$(m,i) \rightarrow j, \quad j = 0, \dots n-1$$

where total number of shape functions is

- Note this mapping is many-to-one
- Correspondence

$$\hat{\mathbf{u}}_i^{(m)} \equiv \hat{\mathbf{u}}_j$$

$$N_i^{(m)} \equiv N_j$$

Questions?

Discretization

Matrix form

$$\mathbf{u}^{(m)}(x,y,z) = \mathbf{H}^{(m)}(x,y,z)\hat{\mathbf{u}}^{(m)}$$

where

$$\mathbf{H}^{(m)}(x,y,z) = \begin{pmatrix} N_0^{(m)}, \dots, N_{k-1}^{(m)} & 0 & 0 \\ 0 & N_0^{(m)}, \dots, N_{k-1}^{(m)} & 0 \\ 0 & 0 & N_m^{(m)}, \dots, N_{k-1}^{(m)} \end{pmatrix} \in \mathbf{R}^{3 \times 3k}$$

$$\hat{\mathbf{u}}^{(m)} = \left(\hat{u}_{x_0}^{(m)}, \dots, \hat{u}_{x_{k-1}}^{(m)}, \hat{u}_{y_0}^{(m)}, \dots, \hat{u}_{y_{k-1}}^{(m)}, \hat{u}_{z_0}^{(m)}, \dots, \hat{u}_{z_{k-1}}^{(m)}\right)^T \in \mathbf{R}^{3k}$$

Discretization

• Strain

$$\varepsilon^{(m)}(x,y,z) = \mathcal{S}\mathbf{u}^{(m)}(x,y,z) = \mathbf{B}^{(m)}(x,y,z)\hat{\mathbf{u}}^{(m)}$$

Where

$$\mathbf{B}^{(m)}(x,y,z) = \begin{pmatrix} \frac{\partial}{\partial z} [N_0^{(m)}, \dots, N_{k-1}^{(m)}] & 0 & 0 & 0 \\ 0 & \frac{\partial}{\partial y} [N_0^{(m)}, \dots, N_{k-1}^{(m)}] & 0 & 0 \\ 0 & 0 & \frac{\partial}{\partial z} [N_0^{(m)}, \dots, N_{k-1}^{(m)}] & 0 \\ 0 & 0 & \frac{\partial}{\partial z} [N_0^{(m)}, \dots, N_{k-1}^{(m)}] \\ \frac{\partial}{\partial y} [N_0^{(m)}, \dots, N_{k-1}^{(m)}] & \frac{\partial}{\partial z} [N_0^{(m)}, \dots, N_{k-1}^{(m)}] & \frac{\partial}{\partial y} [N_0^{(m)}, \dots, N_{k-1}^{(m)}] \\ 0 & \frac{\partial}{\partial z} [N_0^{(m)}, \dots, N_{k-1}^{(m)}] & 0 & \frac{\partial}{\partial z} [N_0^{(m)}, \dots, N_{k-1}^{(m)}] \\ \frac{\partial}{\partial z} [N_0^{(m)}, \dots, N_k^{(m)}] & 0 & \frac{\partial}{\partial z} [N_0^{(m)}, \dots, N_k^{(m)}] \end{pmatrix} \in \mathbf{R}^{6 \times 3k}$$

• Note: stress is

$$\sigma^{(m)}(x,y,z) = \mathbf{C}^{(m)}\varepsilon^{(m)}(x,y,z) = \mathbf{C}^{(m)}\mathbf{B}^{(m)}(x,y,z)\hat{\mathbf{u}}^{(m)}$$

Discretization

• Remember: weak form

$$\int_{\Omega} \delta \varepsilon^T \sigma d\Omega = \int_{\Omega} \delta \mathbf{u}^T \mathbf{f} d\Omega$$

- Use same per-element interpolation for displacement/strain and virtual displacement/strain
- Per-element perspective

$$\textstyle \sum\limits_{m} \int_{V^m} \delta \varepsilon^T \sigma d\Omega = \sum\limits_{m} \int_{V^{(m)}} \delta \mathbf{u}^T \mathbf{f} d\Omega$$

Discretization

• Galerkin approximation of weak form

$$\begin{split} \delta \hat{\mathbf{U}}^T \left[\sum_{m} \int_{V^{(m)}} \mathbf{B}_{*}^{(m)^T} \mathbf{C}_{*}^{(m)} \mathbf{B}_{*}^{(m)} dV^{(m)} \right] \hat{\mathbf{U}} &= \delta \hat{\mathbf{U}}^T \left[\sum_{m} \int_{V^{(m)}} \mathbf{H}_{*}^{(m)^T} \mathbf{f}^{(m)} dV^{(m)} \right] \\ \Rightarrow \left[\sum_{m} \int_{V^{(m)}} \mathbf{B}_{*}^{(m)^T} \mathbf{C}_{*}^{(m)} \mathbf{B}_{*}^{(m)} dV^{(m)} \right] \hat{\mathbf{U}} &= \left[\sum_{m} \int_{V^{(m)}} \mathbf{H}_{*}^{(m)^T} \mathbf{f}^{(m)} dV^{(m)} \right] \end{split}$$

where $\hat{\mathbf{U}} = {\{\hat{\mathbf{u}}_i\}, j = 1, ..., n}$

Assembly (map element indices to global indices)

$$\mathbf{B}_*^{(m)}[u,v] \equiv \mathbf{B}^{(m)}[s,t], \text{ where } (m,s) \to u, (m,t) \to v$$

Matrix formulation

• Global stiffness matrix

$$\mathbf{K} = \sum\limits_{m} \int_{V^{(m)}} \mathbf{B}_{*}^{(m)^{T}} \mathbf{C}_{*}^{(m)} \mathbf{B}_{*}^{(m)} dV^{(m)}$$

Global force vector

$$\mathbf{R} = \sum_{m} \int_{V^{(m)}} \mathbf{H}_{*}^{(m)T} \mathbf{f}^{(m)} dV^{(m)}$$

· Solve for nodal weights

$$\mathbf{K}\mathbf{\hat{U}} = \mathbf{R}$$

Questions?

Dynamic equation

$$\partial_j \sigma_{ij} + f_i = \rho \, \partial_{tt} u_i$$

• Note equilibrium equation is the same, just with zero right hand side

Viscoelastic and plastic materials

- Viscoelastic: material that has damping, some energy during deformation converted to heat
- Plastic: when stress exceeds threshold, material changes shape permanently

Next time

- "Pose space deformation", Lewis et al.
- Presentation by Iman Mostafavi