

Inference of inter-particle forces from stochastic trajectories

Cristina Melnic

26.04.2021

System dynamics

- Langevin equation

$$dx = \frac{1}{\gamma} F(x) dt + \sigma \xi(t) dt$$

- $p(x_0, t_0), p(x_i, t_i | x_{i-1}, t_{i-1}) - ?$
- Fluctuation-dissipation theorem $D\beta\gamma = 1$
-

$$dx = D\beta F(x) dt + \sigma \xi(t) dt \quad (1)$$

Numerically generated trajectories

- Euler-Maruyama scheme
- Potentials

$$V_{HO}(x) = \frac{kx^2}{2} \quad (2)$$

$$V_{LJ}(x) = \epsilon \left[\left(\frac{r_{min}}{x} \right)^{12} - 2 \left(\frac{r_{min}}{x} \right)^6 \right] \quad (3)$$

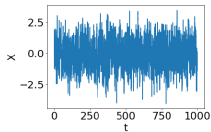
$$V_{dw-HO}(x) = k_4 \frac{x^4}{4} + k_3 \frac{x^3}{3} + k_2 \frac{x^2}{2} + k_1 x + k_0 \quad (4)$$

- Fitted displacement and position distributions

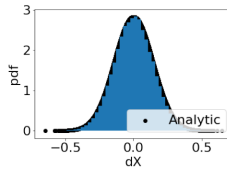
$$p(x) \propto \exp[-\beta V(x)] \quad (5)$$

$$p(x, t \mid x_0, t_0) \propto \exp[-\beta V(x - x_0)] \quad (6)$$

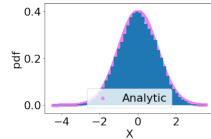
Harmonic Oscillator



(a) Trajectory



(b) Distribution of displacement



(c) Distribution of position

Figure 1.1: Stochastic trajectory of a particle in a harmonic oscillator potential

Position distribution

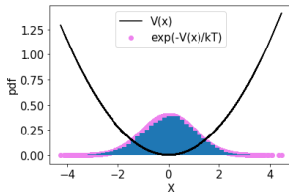
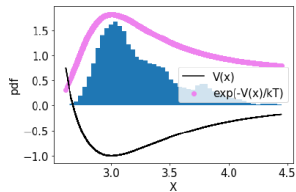
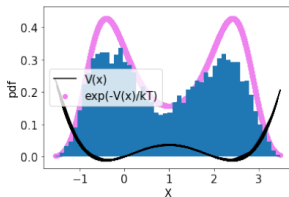
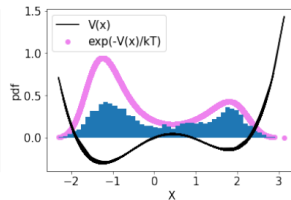
(a) Harmonic oscillator (V_{HO} , $k=1$)(b) Lennard-Jones (V_{LJ} , $\epsilon = 3$, $r_{min} = 9$)(c) Symmetric double-well HO
(V_{dw-HO} , $k_1 = 1$, $k_2 = 1$, $k_3 = -3$, $k_4 = 1$)(d) Asymmetric double-well HO
(V_{dw-HO} , $k_1 = 1$, $k_2 = -2$, $k_3 = -1$, $k_4 = 1$)

Figure 1.2: Distribution of positions from stochastic trajectories

Theoretical aspects I

- Re-written eq 6, $d\omega(t) = \xi(t)dt$

$$dx = D\beta F(x) dt + \sigma d\omega(t)$$

- Average of eq. with the help of Itô lemma [2]

$$\begin{aligned} df[x(t)] &= f'[x(t)]a[x(t), t]dt + f'[x(t)]b[x(t), t]d\omega(t) \\ &\quad + \frac{1}{2}f''[x(t)](b[x(t), t])^2dt \end{aligned}$$

- Fokker-Planck equation [2]

$$\partial_t p(x, t | x_0, t_0) = \left(\partial_x a + \frac{1}{2} \partial_x^2 b \right) p(x, t | x_0, t_0) \quad (7)$$

Theoretical aspects II

- Postulates at thd. equilibrium for conservative force fields
 1. Boltzmann probability distribution
 2. No net flux of particles (probability)
- Fluctuation dissipation theorem [2]

$$\nabla D = \mathbf{F}(\mathbf{r}) (\gamma^{-1} - D\beta) \quad (8)$$

- Smoluchowski equation

$$\partial_t p(x, t | x_0, t_0) = D (\partial_x^2 - \beta F(x)) p(x, t | x_0, t_0) \quad (9)$$

- Linear Potential $V(x) = cx$ [2]

$$p(x, t | x_0, t_0) = \frac{1}{\sqrt{4\pi D(t-t_0)}} \exp \left[-\frac{(x - x_0 + D\beta c(t-t_0))^2}{4D(t-t_0)} \right] \quad (10)$$

Bayesian Inference

[1]

$$p(\theta | X) = \frac{p(X | \theta) \pi(\theta)}{p(X)} = \frac{\mathcal{L}_n(\theta) \pi(\theta)}{c_n} \quad (11)$$

$$c_n = p(X_1, \dots, X_n) = \int \mathcal{L}_n(\theta) \pi(\theta) d\theta$$

- Parameters: $\theta = \{k_i\}$, $i = 1 \dots n$ from $F(\{k_i\})$, $X = dx_{i+1,i}$
- The likelihood: $\mathcal{L}_n(\theta) = \mathcal{N}(dx_{i+1,i} - F_i, 2D\Delta t)$
- The prior: $\pi(\theta) = \frac{1}{n}$

Results I

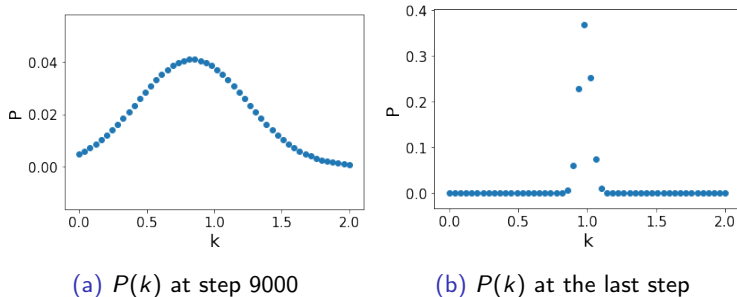
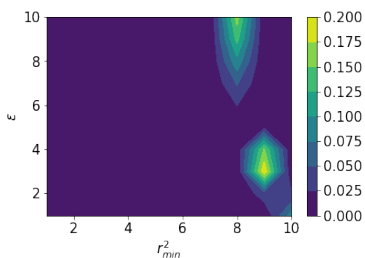
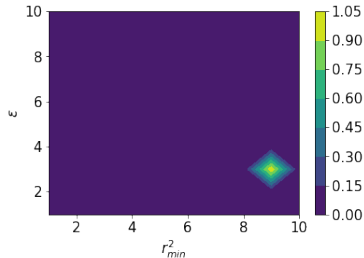


Figure 1: Evolution of $P(k)$ for the HO case (Correct value $k = 1$)

Results II



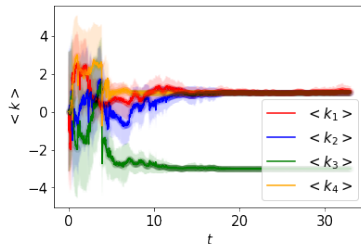
(a) $P(\epsilon; r_{min}^2)$ at step 40



(b) $P(\epsilon; r_{min}^2)$ at the last step

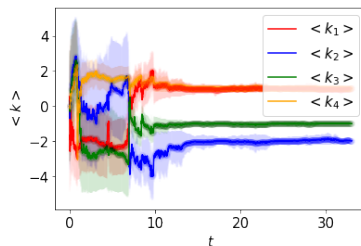
Figure 2: Evolution of $P(\epsilon; r_{min}^2)$ for the LJ case (Correct values $\epsilon = 3; r_{min}^2 = 9$)

Results III



(a)

$k_1 = 1, k_2 = 1, k_3 = -3, k_4 = 1$



(b)

$k_1 = 1, k_2 = -2, k_3 = -1, k_4 = 1$

Figure 3: Evolution of the inferred parameters and the diminishing standard deviation in time from the trajectories in the double-well HO

Conclusions

Advantages

- It works!
- No complicated Smoluchowski solution for $p(dX|F)$ needed
- Not all the data necessary

Disadvantages

- Slow with more parameters
- Analytic form of the forces needed
- A good guess of parameter intervals

References



H. Liu and L. Wasserman.

"Statistical Machine Learning", Chapter 12 .
2014.



K. Schulten.

"Lecture Notes PHYS498", Chapters 2,4,5.