Spectral Analysis of the Ornstein-Uhlenbeck Process on R

1 Covariance Function and Spectral Density

The covariance function of the Ornstein-Uhlenbeck (OU) process is given by:

$$C(x) = \sigma^2 e^{-\alpha |x|} \tag{1}$$

where σ^2 is the variance and $\alpha > 0$ is the mean reversion rate.

The spectral density $S(\omega)$ is the Fourier transform of C(x):

$$S(\omega) = \int_{-\infty}^{\infty} C(x) e^{-i\omega x} dx$$

$$= \sigma^{2} \int_{-\infty}^{\infty} e^{-\alpha|x|} e^{-i\omega x} dx$$

$$= \sigma^{2} \left(\int_{0}^{\infty} e^{-(\alpha+i\omega)x} dx + \int_{0}^{\infty} e^{-(\alpha-i\omega)x} dx \right)$$

$$= \sigma^{2} \left[\frac{1}{\alpha+i\omega} + \frac{1}{\alpha-i\omega} \right]$$

$$= \frac{2\sigma^{2}\alpha}{\alpha^{2}+\omega^{2}}$$

$$(2)$$

2 Orthogonal Polynomials

The polynomials orthogonal with respect to the spectral density $S(\omega)$ are related to the Routh-Romanovski polynomials. Let $x = \omega / \alpha$, then the weight function becomes:

$$w(x) = \frac{1}{1 + x^2} \tag{3}$$

The Routh-Romanovski polynomials $R_n(x)$ are defined by the recurrence relation:

$$R_0(x) = 1$$

 $R_1(x) = x$
 $R_{n+1}(x) = x R_n(x) - k_n R_{n-1}(x)$ for $n > 1$ (4)

where k_n is a coefficient that can vary depending on the specific normalization used. In some sources, $k_n = n^2$, but this may differ.

These polynomials satisfy the orthogonality relation:

$$\int_{-\infty}^{\infty} R_m(x) R_n(x) w(x) dx = h_n \delta_{mn}$$
(5)

where h_n is a normalization constant and δ_{mn} is the Kronecker delta.

The first few polynomials are:

$$R_0(x) = 1$$

 $R_1(x) = x$
 $R_2(x) = x^2 - 1$
 $R_3(x) = x^3 - 3x$
 $R_4(x) = x^4 - 6x^2 + 3$ (6)

The polynomials orthogonal with respect to $S(\omega)$ are:

$$P_n(\omega) = R_n\left(\omega/\alpha\right) \tag{7}$$

3 Fourier Transforms of Orthogonal Polynomials

The Fourier transforms of the first few Routh-Romanovski polynomials are:

$$r_{0}(t) = \sqrt{2\pi} \,\delta(t)$$

$$r_{1}(t) = i\sqrt{2\pi/\alpha} \,\frac{d}{dt} [e^{-\alpha|t|}]$$

$$r_{2}(t) = -\sqrt{2\pi/\alpha^{2}} \,\frac{d^{2}}{dt^{2}} [e^{-\alpha|t|}] - \sqrt{2\pi} \,\delta(t)$$
(8)

where $\delta(t)$ is the Dirac delta function.

4 Entropy Integral and Non-Compactness

To show that the covariance operator of the OU process is not compact on $L^2(\mathbb{R})$, we analyze the ϵ -entropy integral.

The ϵ -covering number $N(\epsilon)$ is related to the spectral density:

$$N(\epsilon) \approx \int_{-\infty}^{\infty} \max\left(1, \sqrt{\frac{S(\omega)}{\epsilon^2}}\right) d\omega$$
 (9)

For large ω , $S(\omega) \sim 2 \sigma^2 \alpha / \omega^2$, so:

$$N(\epsilon) \approx 2 \int_0^\infty \max\left(1, \frac{\sqrt{2\sigma^2 \alpha}}{\epsilon \omega}\right) d\omega$$
 (10)

Let $\omega_{\epsilon} = \sqrt{2 \sigma^2 \alpha} / \epsilon$. Then:

$$N(\epsilon) \approx 2 \left[\omega_{\epsilon} + \int_{\omega_{\epsilon}}^{\infty} \frac{\sqrt{2\sigma^{2}\alpha}}{\epsilon \omega} d\omega \right]$$

$$= 2 \left[\frac{\sqrt{2\sigma^{2}\alpha}}{\epsilon} + \frac{\sqrt{2\sigma^{2}\alpha}}{\epsilon} \log \left(\frac{\infty}{\omega_{\epsilon}} \right) \right]$$

$$\approx \frac{C}{\epsilon} \log \left(\frac{1}{\epsilon} \right)$$
(11)

where C is a constant depending on σ and α .

The ϵ -entropy $H(\epsilon)$ is defined as $\log(N(\epsilon))$, so:

$$H(\epsilon) \approx \log\left(\frac{C}{\epsilon}\right) + \log\left(\log\left(\frac{1}{\epsilon}\right)\right)$$
 (12)

The entropy integral is:

$$\int_0^1 H(\epsilon) d\epsilon \approx \int_0^1 \left[\log \left(\frac{C}{\epsilon} \right) + \log \left(\log \left(\frac{1}{\epsilon} \right) \right) \right] d\epsilon \tag{13}$$

The second term causes the divergence. Using the change of variable $u = \log(1/\epsilon)$:

$$\int_0^1 \log\left(\log\left(\frac{1}{\epsilon}\right)\right) d\epsilon = \int_0^\infty \log\left(u\right) e^{-u} du = \infty \tag{14}$$

Therefore:

$$\int_{0}^{1} H(\epsilon) d\epsilon = \infty \tag{15}$$

This divergence of the entropy integral demonstrates that the covariance operator of the OU process is not compact on $L^2(\mathbb{R})$.