

# Stone's Theorem, Shift Group, and Fourier Transform

## Definitions

**Shift Group:** For  $f \in L^2(\mathbb{R})$ , define the family of unitary operators  $(S_t)_{t \in \mathbb{R}}$  by

$$(S_t f)(x) = f(x + t).$$

**Generator of Shift Group:** Define  $A = \frac{d}{dx}$  on the domain

$$D(A) = \{f \in L^2(\mathbb{R}) : f' \in L^2(\mathbb{R})\},$$

where  $f'$  is in the distributional sense.

**Momentum Operator:** Define  $P = -iA = -i \frac{d}{dx}$  on the same domain  $D(P) = D(A)$ .

**Fourier Transform:**

$$\mathcal{F}[f](\omega) = \hat{f}(\omega) := \int_{-\infty}^{\infty} f(x) e^{-i\omega x} dx.$$

## Theorems and Proofs

**Theorem 1.** [Stone's Theorem Applied to Shift Group] The strongly continuous unitary group  $(S_t)$  on  $L^2(\mathbb{R})$  has a densely defined skew-adjoint generator  $A = \frac{d}{dx}$  such that  $S_t = e^{tA}$ . The generator satisfies

$$A f = \lim_{h \rightarrow 0} \frac{S_h f - f}{h}$$

in the  $L^2$  topology on the domain  $D(A)$ .

**Proof.** Let  $f \in D(A)$ . Then

$$\frac{S_h f(x) - f(x)}{h} = \frac{f(x+h) - f(x)}{h} \rightarrow f'(x)$$

as  $h \rightarrow 0$  in  $L^2$  norm. Thus, the infinitesimal generator of  $S_t$  is  $A = \frac{d}{dx}$ .

To verify  $A$  is skew-adjoint, for  $f, g \in D(A)$ :

$$\langle A f, g \rangle = \int_{-\infty}^{\infty} f'(x) \overline{g(x)} dx \quad (1)$$

$$= f(x) \overline{g(x)} \Big|_{-\infty}^{\infty} - \int_{-\infty}^{\infty} f(x) \overline{g'(x)} dx \quad (2)$$

$$= 0 - \int_{-\infty}^{\infty} f(x) \overline{g'(x)} dx \quad (3)$$

$$= -\langle f, A g \rangle \quad (4)$$

Therefore  $A^* = -A$ , confirming  $A$  is skew-adjoint.  $\square$

**Theorem 2.** *[Relation Between Generators] The shift group is generated by both the skew-adjoint operator  $A = \frac{d}{dx}$  and the self-adjoint momentum operator  $P = -iA$ :*

$$S_t = e^{tA} = e^{-itP}$$

**Proof.** Since  $P = -iA$ , we have  $-itP = -it(-iA) = -i^2tA = tA$ . Therefore:

$$e^{-itP} = e^{tA}$$

For  $f \in D(A)$ , using the Taylor expansion:

$$e^{tA} f(x) = \sum_{n=0}^{\infty} \frac{t^n}{n!} A^n f(x) = \sum_{n=0}^{\infty} \frac{t^n}{n!} f^{(n)}(x) = f(x+t) = S_t f(x) \quad \square$$

**Theorem 3.** *[Complex Exponentials Are Eigenfunctions] For any  $\omega \in \mathbb{R}$ :*

$$1. A e^{i\omega x} = i\omega e^{i\omega x} \text{ (eigenvalue } i\omega \text{ for skew-adjoint } A)$$

$$2. P e^{i\omega x} = \omega e^{i\omega x} \text{ (eigenvalue } \omega \text{ for self-adjoint } P)$$

**Proof.** Direct calculations:

$$1. A e^{i\omega x} = \frac{d}{dx} e^{i\omega x} = i\omega e^{i\omega x}$$

$$2. P e^{i\omega x} = -i \frac{d}{dx} e^{i\omega x} = -i(i\omega) e^{i\omega x} = \omega e^{i\omega x} \quad \square$$

**Theorem 4.** *[Spectral Decomposition via Fourier Transform] Under the Fourier transform  $\mathcal{F}$ :*

1. *The self-adjoint momentum operator becomes multiplication by  $\omega$ :  
 $\mathcal{F}[Pf](\omega) = \omega \hat{f}(\omega)$*
2. *The shift group becomes multiplication by a phase:  $\mathcal{F}[S_t f](\omega) = e^{i\omega t} \hat{f}(\omega)$*

**Proof.** For part 1, if  $f \in D(P)$ :

$$\mathcal{F}[Pf](\omega) = \int_{-\infty}^{\infty} (-i f'(x)) e^{-i\omega x} dx \quad (5)$$

$$= -i \int_{-\infty}^{\infty} f'(x) e^{-i\omega x} dx \quad (6)$$

Integration by parts (boundary terms vanish):

$$= -i [0 + i\omega \hat{f}(\omega)] = \omega \hat{f}(\omega) \quad (7)$$

For part 2:

$$\mathcal{F}[S_t f](\omega) = \int_{-\infty}^{\infty} f(x+t) e^{-i\omega x} dx \quad (8)$$

Let  $u = x + t$ , so  $x = u - t$ ,  $dx = du$ :

$$= \int_{-\infty}^{\infty} f(u) e^{-i\omega(u-t)} du \quad (9)$$

$$= e^{i\omega t} \int_{-\infty}^{\infty} f(u) e^{-i\omega u} du = e^{i\omega t} \hat{f}(\omega) \quad (10)$$

□

**Theorem 5.** *[Eigenfunction Property of Shift Group] Complex exponentials are eigenfunctions of the shift group:*

$$S_t e^{i\omega x} = e^{i\omega t} e^{i\omega x}$$

*with eigenvalue  $e^{i\omega t}$ .*

**Proof.**

$$S_t e^{i\omega x} = e^{i\omega(x+t)} = e^{i\omega x} e^{i\omega t} = e^{i\omega t} e^{i\omega x}$$

□

**Corollary 6.** *[Consistency Check] The eigenvalue relationships are consistent:*

$$S_t e^{i\omega x} = e^{tA} e^{i\omega x} = e^{t(i\omega)} e^{i\omega x} = e^{i\omega t} e^{i\omega x}$$

*since  $A$  has eigenvalue  $i\omega$  on  $e^{i\omega x}$ .*

## Conclusion

Stone's theorem ensures that the shift group  $(S_t)$  has a **\*\*skew-adjoint generator\*\***  $A = \frac{d}{dx}$ , whose eigenfunctions are the complex exponentials  $e^{i\omega x}$  with purely imaginary eigenvalues  $i\omega$ . The related **\*\*self-adjoint momentum operator\*\***  $P = -iA$  has the same eigenfunctions but with real eigenvalues  $\omega$ .

The Fourier transform provides the spectral decomposition that diagonalizes both operators:

- $P$  becomes multiplication by  $\omega$  (real eigenvalues)
- $S_t$  becomes multiplication by  $e^{i\omega t}$  (unitary eigenvalues on the unit circle)

This mathematical structure underlies all of Fourier analysis: complex exponentials are the fundamental building blocks because they are precisely the functions that transform simply under shifts, making them the natural basis for analyzing translation-invariant systems. The distinction between the skew-adjoint generator  $A$  (with imaginary eigenvalues) and the self-adjoint momentum operator  $P$  (with real eigenvalues) is crucial for understanding why unitary groups arise from self-adjoint operators via Stone's theorem.