

A Pair of Orthogonal Polynomial Sequences on $[0, \infty]$ That Uniformly Converge to The Bessel functions of the First Kind of Orders 0 and 1

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Abstract

The even-indexed orthonormalized Fourier transforms of the Chebyshev polynomials of the first kind form a basis in a reproducing-kernel Hilbert space for the Bessel function of the first kind J_0 and likewise for the odd-indexed functions which form a basis that reproduces $J_0 = -J_1$. Suprisingly, such a basis for these functions was not known to exist before this.

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1 The Type-I Chebyshev Polynomials $T_n(x)$

Let T_n be the Chebyshev polynomials of the first kind, also said to be of Type-I, defined by

$$\begin{aligned}
 T_n(x) &= {}_2F_1\left(n, -n \middle| \frac{1}{2} - \frac{x}{2}\right) \\
 &= \int_{-\infty}^{\infty} e^{ixy} \hat{T}_n(y) dy \\
 &= \int_{-\infty}^{\infty} e^{ixy} \frac{i}{y} (e^{-iy} F_n^+(y) - e^{iy} (-1)^n F_n^-(y)) dy \\
 &= \int_{-\infty}^{\infty} e^{ixy} \int_{-\infty}^{\infty} e^{-iyz} T_n(z) dz dy
 \end{aligned} \tag{1}$$

where ${}_2F_1$ is the (Gauss) hypergeometric function. [?, (13.140)]

1.1 The Fourier Transforms $\hat{T}_n(y)$ of $T_n(x)$

The functions $\hat{T}_n(y)$ are Fourier transforms of $T_n(x)$ defined by

$$\begin{aligned}\hat{T}_n(y) &= \int_{-\infty}^{\infty} e^{-ixy} T_n(x) dx = \int_{-1}^1 e^{-ixy} T_n(x) dx \\ &= \int_{-\infty}^{\infty} e^{-ixy} {}_2F_1\left(n, -n \middle| \frac{1}{2} - \frac{x}{2}\right) dx \\ &= \frac{i}{y} (e^{-iy} F_n^+(y) - e^{i(\pi n + y)} F_n^-(y))\end{aligned}\tag{2}$$

where

$$F_n^{\pm}(y) = {}_3F_1\left(1, n, -n \middle| \frac{\pm iy}{2}\right)\tag{3}$$

The L^2 norm of $\hat{T}_n(y)$ is

$$\|\hat{T}_n\| = \sqrt{\int_{-\infty}^{\infty} \hat{T}_n(y)^2 dy} = \sqrt{\frac{4(-1)^n \pi - (2n^2 - 1)}{4n^2 - 1}}\tag{4}$$

then define the normalized Fourier transforms $Y_n(y)$ of $T_n(x)$ by

$$\begin{aligned}Y_n(y) &= \frac{\hat{T}_n(y)}{\|\hat{T}_n\|} \\ &= \frac{i}{y} \left(\frac{e^{-iy} F_n^+(y) - e^{i(\pi n + y)} F_n^-(y)}{\sqrt{\frac{4(-1)^n \pi - (2n^2 - 1)}{4n^2 - 1}}} \right)\end{aligned}\tag{5}$$

1.2 Orthogonalizing $Y_n(y)$ Via The Gram-Schmidt Process

Apply the Gram-Schmidt process to the normalized Fourier transforms of the Type I Chebyshev polynomials $Y_n(y)$ to get $\psi_n(y)$

$$Y_n^{\perp}(y) = Y_n(y) - \sum_{m=1}^{n-1} \frac{\langle Y_n(y), Y_m^{\perp}(y) \rangle}{\langle Y_m^{\perp}(y), Y_m^{\perp}(y) \rangle} Y_m^{\perp}(y)\tag{6}$$

Let

$$A_{k,n} = -(-1)^{n+\binom{k}{2}} (k-2n+1)! 2^{2n-1-k} \binom{k+1}{k-2n+1} \binom{2k+2-2n}{k+1}\tag{7}$$

and

$$B_{k,n} = \frac{(-1)^{n+\binom{k}{2}} 2^{k-2n} (k-n)! \binom{\frac{1}{2}-n+k}{k-2n}}{n!}\tag{8}$$

then defined the associated functions

$$\Psi_n^{\sin}(y) = \frac{\sin(y)\sqrt{2n-1}}{x^n\sqrt{\pi}} \sum_{k=0}^{n-2} x^{2k} A_{k,n-2} \quad (9)$$

and

$$\Psi_n^{\cos}(y) = \frac{\cos(y)\sqrt{2n-1}}{x^n\sqrt{\pi}} \sum_{k=0}^{n-2} x^{2k+1} B_{k,n-2} \quad (10)$$

then $Y_n^\perp(y)$ can be expressed as

$$Y_n^\perp(y) = \Psi_n^{\sin}(y) + \Psi_n^{\cos}(y) \quad (11)$$

so that

$$\psi_n(y) = \lambda_h Y_n^\perp(y) \quad (12)$$

-3	1	0	0	0	0	0	...
15	-6	0	0	0	0	0	...
105	-45	1	0	0	0	0	...
-945	420	-15	0	0	0	0	...
-10395	4725	-210	1	0	0	0	...
135135	-62370	3150	-28	0	0	0	...
2027025	-945945	51975	-630	1	0	0	...
-34459425	16216200	-945945	13860	-45	0	0	...
-654729075	310134825	-18918900	315315	-1485	1	0	...
13749310575	-6547290750	413513100	-7567560	45045	-66	0	...
...

Table 1. The first 10 row-vectors of $A_{k,n}$ matrix

2 An Integral Covariance Operator

The eigenvalues λ_k of the integral covariance operator

$$Tf(x) = \int_0^\infty J_0(x-y)f(x)dx \quad (13)$$

where J_0 is the Bessel function of the first kind of order 0 are given by

$$\begin{aligned} \lambda_k &= \sqrt{\frac{4k+1}{\pi}} (k+1)^{2-\frac{1}{2}} \\ &= \sqrt{\frac{4k+1}{\pi}} \frac{\Gamma\left(k+\frac{1}{2}\right)^2}{\Gamma(k+1)^2} \end{aligned} \quad (14)$$

which, together with the eigenfunctions

$$\psi_k(x) = e^{i\pi k} \sqrt{\frac{4k+1}{\pi}} j_{2k}(y) \quad (15)$$

solve the characteristic(eigenfunction) equation

$$\psi_k(y) = \lambda_k \int_0^\infty J_0(x-y) \psi_k(x) dx \quad (16)$$

so that

$$J_0(x) = \sum_{k=0}^{\infty} \lambda_k \psi_k(y)$$