

A Uniformly Convergent Orthonormal Expansion for the Bessel Function of the First Kind of Order 0

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Theorem 1

Let $\psi_n(y)$ be defined as

$$\begin{aligned}
 \psi_n(y) &= (-1)^n \sqrt{\frac{4n+1}{\pi}} j_{2n}(y) \\
 &= (-1)^n \sqrt{\frac{4n+1}{\pi}} \sqrt{\frac{\pi}{2y}} J_{2n+\frac{1}{2}}(y) \\
 &= (-1)^n \sqrt{\frac{(4n+1)\pi}{\pi 2y}} J_{2n+\frac{1}{2}}(y) \\
 &= (-1)^n \sqrt{\frac{(4n+1)}{2y}} J_{2n+\frac{1}{2}}(y) \\
 &= (-1)^n \sqrt{\frac{2n+\frac{1}{2}}{y}} J_{2n+\frac{1}{2}}(y)
 \end{aligned} \tag{1}$$

where J_ν denotes the Bessel function of the first kind and j_n the spherical Bessel function. Then

$$\begin{aligned}
 J_0(x) &= \sum_{n=0}^{\infty} \psi_n(x) \int_0^{\infty} J_0(y) \psi_n(y) dy \\
 &= \sum_{n=0}^{\infty} (-1)^n \sqrt{\frac{4n+1}{\pi}} j_{2n}(x) \sqrt{\frac{2n+\frac{1}{2}}{\pi}} \frac{\Gamma(n+\frac{1}{2})^2}{\Gamma(n+1)^2} \\
 &= \frac{1}{2} \frac{1}{\sqrt{\pi x}} \sum_{n=0}^{\infty} (-1)^n (4n+1) \frac{\Gamma(n+\frac{1}{2})^2}{\Gamma(n+1)^2} J_{2n+\frac{1}{2}}(x) \\
 &= \frac{1}{\sqrt{4\pi x}} \sum_{n=0}^{\infty} \frac{(-1)^n (4n+1) \Gamma(n+\frac{1}{2})^2}{\Gamma(n+1)^2} J_{2n+\frac{1}{2}}(x)
 \end{aligned} \tag{2}$$

with uniform convergence $\forall x \in \mathbb{C}$. Moreover, $\{\psi_n\}$ forms an orthonormal system in $L^2([0, \infty))$ satisfying

$$\int_0^{\infty} \psi_m(t) \psi_n(t) dt = \delta_{mn} = \begin{cases} 1 & m = n \\ 0 & m \neq n \end{cases} \tag{3}$$

Proof.

Step 1: Orthonormality of $\psi_n(y)$.

Substituting the definition:

$$\begin{aligned}
\langle \psi_m, \psi_n \rangle &= (-1)^{m+n} \sqrt{\frac{(4m+1)(4n+1)}{\pi^2}} \frac{\pi}{2} \int_0^\infty \frac{J_{2m+\frac{1}{2}}(y) J_{2n+\frac{1}{2}}(y)}{y} dy \\
&= (-1)^{m+n} \sqrt{\frac{(4m+1)(4n+1)}{4\pi}} \cdot \frac{\delta_{mn}}{2n+\frac{1}{2}} \\
&= \delta_{mn}
\end{aligned} \tag{4}$$

The crucial integral follows from Bessel orthogonality:

$$\int_0^\infty \frac{J_\mu(y) J_\nu(y)}{y} dy = \frac{2}{\pi} \frac{\delta_{\mu\nu}}{\mu + \nu} \tag{5}$$

Step 2: Expansion Coefficients. Using the orthonormal basis:

$$c_n = \int_0^\infty J_0(y) \psi_n(y) dy = (-1)^n \sqrt{\frac{4n+1}{2}} \frac{\Gamma(n+\frac{1}{2})^2}{\Gamma(n+1)^2} \tag{6}$$

This emerges from evaluating:

$$\int_0^\infty J_0(y) \frac{J_{2n+\frac{1}{2}}(y)}{\sqrt{y}} dy = \frac{\Gamma(n+\frac{1}{2})^2}{\sqrt{2} \Gamma(n+1)^2} \tag{7}$$

Step 3: Series Simplification. Substitute c_n into the expansion:

$$J_0(x) = \sum_{n=0}^\infty \frac{(-1)^n (4n+1)}{2\sqrt{\pi x}} \frac{\Gamma(n+\frac{1}{2})^2}{\Gamma(n+1)^2} J_{2n+\frac{1}{2}}(x) \tag{8}$$

TODO: write correct proof, its not due to Wierstrauss, hint, prove that we can always choose an N for a given epsilon such that including that many terms in the expansion results in an error less than epsilon and recognizing that the contribution from the n-th term can be no greater than $\int_0^\infty J_0(y) \psi_n(y) dy$ \square