## Notes on the function gsw\_enthalpy\_diff\_CT\_exact(SA,CT,p\_shallow,p\_deep)

This function,  $gsw_enthalpy_diff_CT_exact(SA,CT,p_shallow,p_deep)$ , returns the difference between the specific enthalpy of two seawater parcels, both having the same Absolute Salinity and Conservative Temperature, but having different pressures. This function uses the full TEOS-10 Gibbs function  $g(S_A,t,p)$  of IOC *et al.* (2010), being the sum of the IAPWS-09 and IAPWS-08 Gibbs functions.

This function is simply two calls to each of two GSW functions as follows,

## References

IAPWS, 2008: Release on the IAPWS Formulation 2008 for the Thermodynamic Properties of Seawater. The International Association for the Properties of Water and Steam. Berlin, Germany, September 2008, available from <a href="https://www.iapws.org">www.iapws.org</a>. This Release is referred to in the text as IAPWS-08.

IAPWS, 2009: Supplementary Release on a Computationally Efficient Thermodynamic Formulation for Liquid Water for Oceanographic Use. The International Association for the Properties of Water and Steam. Doorwerth, The Netherlands, September 2009, available from <a href="http://www.iapws.org">http://www.iapws.org</a>. This Release is referred to in the text as IAPWS-09.

IOC, SCOR and IAPSO, 2010: The international thermodynamic equation of seawater – 2010: Calculation and use of thermodynamic properties. Intergovernmental Oceanographic Commission, Manuals and Guides No. 56, UNESCO (English), 196 pp. Available from <a href="http://www.TEOS-10.org">http://www.TEOS-10.org</a>

Below, for motivation and for reference, is section 3.32 of the TEOS-10 Manual (IOC *et al.* (2010))

## 3.32 Pressure to height conversion

The vertical integral of the hydrostatic equation  $(g = -vP_z)$  can be written as

$$\int_{0}^{z} g(z') dz' = \Phi^{0} - \int_{P_{0}}^{P} v(p') dP' = -\int_{P_{0}}^{P} \hat{v}(S_{SO}, 0^{\circ}C, p') dP' + \Psi + \Phi^{0}$$

$$= -\hat{h}(S_{SO}, 0^{\circ}C, p) + \Psi + \Phi^{0},$$
(3.32.1)

where the dynamic height anomaly  $\Psi$  is expressed in terms of the specific volume anomaly  $\hat{\delta} = \hat{v}(S_A, \Theta, p) - \hat{v}(S_{SO}, 0^{\circ}C, p)$  by

$$\Psi = -\int_{P} \hat{\delta}(p') dP', \qquad (3.32.2)$$

where  $P_0 = 101\,325\,\text{Pa}$  is the standard atmosphere pressure. Writing the gravitational acceleration of Eqn. (D.3) as  $g = g(\phi, z) = g(\phi, 0)\,(1 - \gamma z)$ , the left-hand side of Eqn.

(3.32.1) becomes  $g(\phi,0)(z-\frac{1}{2}\gamma z^2)$ , and using the 48-term expression for the specific enthalpy of Standard Seawater, Eqn. (3.32.1) becomes

$$\hat{h}^{48}\left(S_{SO}, 0^{\circ}C, p\right) - \Psi - \Phi^{0} + g(\phi, 0)\left(z - \frac{1}{2}\gamma z^{2}\right) = 0.$$
 (3.32.3)

This is the equation that is solved for height z in the GSW function  $\mathbf{gsw\_z\_from\_p}$ . It is traditional to ignore  $\Psi + \Phi^0$  when converting between pressure and height, and this can be done by simply calling this function with only two arguments, as in  $\mathbf{gsw\_z\_from\_p}(p,lat)$ . Ignoring  $\Psi + \Phi^0$  makes a difference to z of up to 4m at 5000 dbar. Note that height z is negative in the ocean. When the code is called with three arguments, the third argument is taken to be  $\Psi + \Phi^0$  and this is used in the solution of Eqn. (3.32.3). Dynamic height anomaly  $\Psi$  can be evaluated using the GSW function  $\mathbf{gsw\_geo\_strf\_dyn\_height}$ . The GSW function  $\mathbf{gsw\_p\_from\_z}$  is the exact inverse function of  $\mathbf{gsw\_z\_from\_p}$ ; these functions yield outputs that are consistent with each other to machine precision.

When vertically integrating the hydrostatic equation  $P_z = -g\rho$  in the context of an ocean model where Absolute Salinity  $S_A$  and Conservative Temperature  $\Theta$  are piecewise constant in the vertical, the geopotential (Eqn. (3.24.2))

$$\Phi = \int_{0}^{z} g(z') dz' = \Phi^{0} - \int_{P_{0}}^{P} v(p') dP', \qquad (3.32.4)$$

can be evaluated as a series of exact differences. If there are a series of layers of index i separated by pressures  $p^i$  and  $p^{i+1}$  (with  $p^{i+1} > p^i$ ) then the integral can be expressed (making use of (3.7.5), namely  $h_p|_{S_{A},\Theta} = \hat{h}_p = v$ ) as a sum over n layers of the differences in specific enthalpy so that

$$\Phi = \Phi^{0} - \int_{P_{0}}^{P} v(p') dP' = \Phi^{0} - \sum_{i=1}^{n} \left[ \hat{h}(S_{A}^{i}, \Theta^{i}, p^{i+1}) - \hat{h}(S_{A}^{i}, \Theta^{i}, p^{i}) \right].$$
(3.32.5)

The difference in enthalpy at two different pressures for given values of  $S_A$  and  $\Theta$  is available in the GSW Oceanographic Toolbox via the function  $\mathbf{gsw\_enthalpy\_diff}$ . The summation of a series of such differences in enthalpy occurs in the GSW functions to evaluate two geostrophic streamfunctions from piecewise-constant vertical property profiles,  $\mathbf{gsw\_geo\_strf\_dyn\_height\_pc}$  and  $\mathbf{gsw\_geo\_strf\_isopycnal\_pc}$ .