Recovering the Metric - An Overview of Lie Algebras

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Conversely, if $[\ ,\]$ is skew-symmetric, then [x,x]+[x,x]=0 implies that 2[x,x]=0. Now we see that this implies [x,x]=0 so long as our field is not of characteristic 2, for in those spaces 2=0 and we can deduce nothing about [x,x].

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A vector, z, of \mathfrak{g} is said to be in the center of \mathfrak{g} if

$$[x,z] = \mathbf{0} \ \forall x \in \mathfrak{g}.$$

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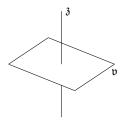


Figure: The Heisenberg Algebra, h

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j-maps

The equation that we used to define j_z ,

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This makes the computation of many complicated geometric objects—such as curvatures—into "simple" calculations in linear algebra!

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It's nice to see that we can *encode* the bracket into a matrix representation, but how can we use that encoding to as the Lie Bracket "map"?

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So it will be enough find a way to write all Lie Brackets as a linear combination of basis brackets.

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 with $x = w_1 e_1 + w_2 e_2$ $y = v_1 e_1 + v_2 e_2$ So $[x, y] = [w_1 e_1 + w_2 e_2, v_1 e_1 + v_2 e_2]$ $= [w_1 e_1, v_1 e_1 + v_2 e_2] + [w_2 e_2, v_1 e_1 + v_2 e_2$

$$= [w_1e_1, v_1e_1 + v_2e_2] + [w_2e_2, v_1e_1 + v_2e_2]$$

$$= [w_1e_1, v_1e_1] + [w_1e_1, v_2e_2] + [w_2e_2, v_1e_1] + [w_2e_2, v_2e_2]$$

$$= w_1v_1[e_1, e_1] + w_1v_2[e_1, e_2] + w_2v_1[e_2, e_1] + w_2v_2[e_2, e_2]$$

$$= w_1v_1L_{11} + w_1v_2L_{12} + w_2v_1L_{21} + w_2v_2L_{22}$$

This means that the Lie Bracket is fully described by the matrix, L

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$$x^{T}L = (w_1w_2)\begin{pmatrix} L_{11} & L_{12} \\ L_{21} & L_{22} \end{pmatrix} = (w_1L_{11} + w_2L_{12}, w_1L_{12} + w_2L_{22})$$

Generalizing the Lie Bracket

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What is the center of \mathbb{R}^3 with the cross product?

β₃'s Lie Bracket

How could the three dimensional Heisenberg Algebra be represented?

\$\mathbf{h}_3\end{array}\text{s Lie Bracket}

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What is the center of \mathfrak{h}_3 ? The center is the span of $\{e_3\}$.

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$$E = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

Recall that for any $z \in \mathfrak{z}$, the map $j_z : \mathfrak{v} \to \mathfrak{v}$ is defined by the identity

$$\langle y, j_z(x) \rangle_{\mathfrak{v}} = \langle z, [x, y] \rangle_{\mathfrak{z}}.$$
 (7)

We wish to find a matrix J_z that represents the map j_z . Since j_z maps v to itself, J_z should be a matrix in $\mathbb{R}^{n \times n}$. Rewriting (7) in matrix form, we get

$$\langle y, j_z(x) \rangle_v = \langle z, [x, y] \rangle_{\delta}$$

$$y^T E(J_z x) = z^{t}(x^T L y)$$

$$y^T (EJ_z) x = z^{t}(-y^T L x)$$

$$= -z^{t}(y^T L x)$$

$$= -y^T(z^{t} L) x$$

$$= y^T(z^{t}(-L)) x$$

$$= y^T(z^{t} L^T) x.$$

The equation $y^T(EJ_z)x = y^T(z^tL^T)x$ must hold for arbitrary vectors x and y in v, implying that

$$EJ_z = z^{\mathrm{t}}L^{\mathrm{T}}.$$

Since $det(E) \neq 0$, it is invertible and

$$J_z = E^{-1}(z^{\mathrm{t}}L^{\mathrm{T}}). \tag{8}$$

This gives us a method to compute the j-map for any z using only matrix computations. Furthermore, if $z=z_k$ is a basis vector of \mathfrak{z} , then $z_k^{\mathfrak{t}}L^T$ is just $(L^k)^T$. Thus, for basis vectors z_k , the map j_{z_k} is represented by the matrix

$$J_{z_k} = E^{-1} (L^k)^{\mathrm{T}}.$$

Since the j-maps are linear, knowing how they act on a basis of \mathfrak{z} is good enough. Indeed, we should be able to define a stack of j-maps by

$$J = E^{-1}L^{\mathrm{T}}. (9)$$

The *j*-map j_z should then be $J_z = z^t J$.

What we did

Using the programming language, sage, we created a program that would compute the j-maps of an arbitrary Lie algebra.

TikZ

Acknowledgements

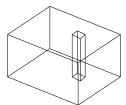


Figure : The L-Stack

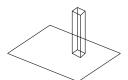


Figure : The L-Stack