PH108

Lecture 23:

Maxwell Equations: Electromagnnetic Waves

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Supplementary reading: Griffiths Sec 9.1, 9.2 (not 9.3)

Maxwell Equations : time dependent \vec{E} , \vec{B}

$$\vec{\nabla} \cdot \vec{E} = \frac{\rho}{\epsilon_0}$$

$$\vec{\nabla} \cdot \vec{B} = 0$$

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

$$\vec{\nabla} \times \vec{B} = \mu_0 \vec{J} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

$$\vec{\nabla} \cdot \vec{J} + \frac{\partial \rho}{\partial t} = 0$$

<u>Physics motivation</u>: measurable quantities (force) should not depend on observer frame of reference

time dependent \vec{E} , \vec{B} with NO ρ , \vec{J}

$$\vec{\nabla} \cdot \vec{E} = 0$$

$$\vec{\nabla} \cdot \vec{B} = 0$$

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

$$\vec{\nabla} \times \vec{B} = \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

These are FOUR coupled first order differential equations

Use
$$\vec{\nabla} \times (\vec{\nabla} \times \vec{X}) = \vec{\nabla}(\vec{\nabla} \cdot \vec{X}) - \nabla^2 \vec{X}$$
 to uncouple them

$2^{\rm nd}$ order differential equation for \vec{E}

$$\vec{\nabla} \cdot \vec{E} = 0$$

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} \qquad \vec{\nabla} \times \vec{B} = \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

$$\vec{\nabla} \times (\vec{\nabla} \times \vec{E}) = -\frac{\partial (\vec{\nabla} \times \vec{B})}{\partial t} = -\mu_0 \epsilon_0 \frac{\partial^2 \vec{E}}{\partial t^2}$$

$$\vec{\nabla} (\vec{\nabla} \cdot \vec{E}) - \nabla^2 \vec{E}; \qquad \nabla^2 \vec{E} = \mu_0 \epsilon_0 \frac{\partial^2 \vec{E}}{\partial t^2}$$

2^{nd} order differential equation for \vec{E} and \vec{B}

These are called wave equations

LHS has space [L]⁻²
$$[\mu_0 \epsilon_0] = [L^{-2}T^2] \rightarrow \frac{1}{velocity^2}$$
 RHS has time [T]⁻²

 \vec{E} , \vec{B} "waves" "travel" in free space with velocity $v = \frac{1}{\sqrt{\mu_0 \epsilon_0}}$

$$\nabla^{2}\vec{B} = \mu_{0}\epsilon_{0}\frac{\partial^{2}\vec{B}}{\partial t^{2}}$$

$$\nabla^{2}\vec{E} = \mu_{0}\epsilon_{0}\frac{\partial^{2}\vec{E}}{\partial t^{2}}$$

What does it mean: \vec{E} , \vec{B} "travel"?

$\vec{E}(\vec{r})$ is vector field

For every point \vec{r} in space, there is a vector \vec{E} at time t

at time $t + \Delta t$ the vectors \vec{E} are translated by $\Delta \vec{r}$

Velocity of travel is
$$v = \frac{\Delta \vec{r}}{\Delta t}$$

... Similarly $\vec{B}(\vec{r})$ is vector field ...

Example: $\vec{E} = \hat{z} E_0 \sin(y - vt)$

Electric field is in \hat{z} direction

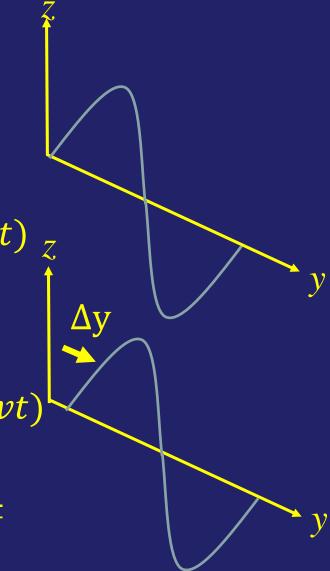
At
$$t + \Delta t$$
:

$$\sin[(y + \Delta y) - v(t + \Delta t)] = \sin(y - vt)$$

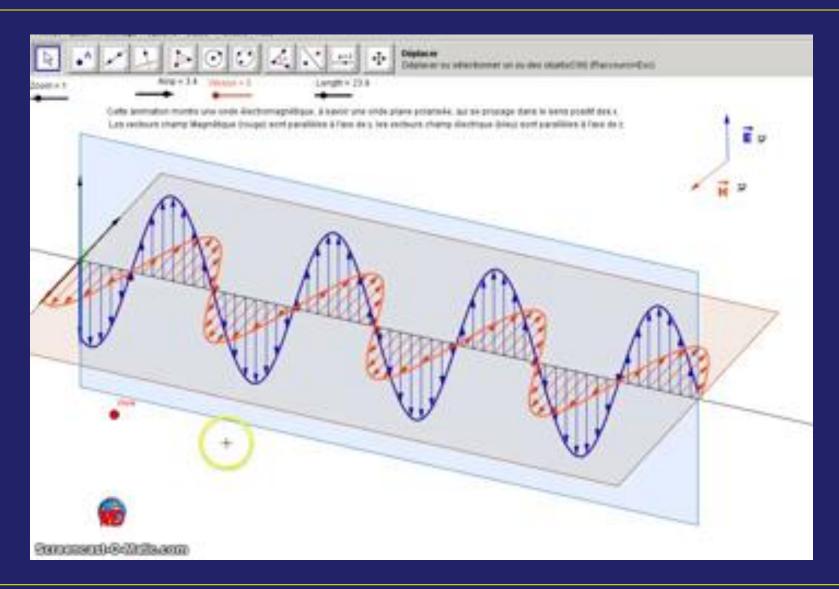
$$IF \Delta y = v\Delta t$$

It is easy to show that $\vec{E} = \hat{z}E_0\sin(y - vt)$ satisfies Maxwell's equations

IF
$$\vec{B} = \hat{x}B_0\sin(y - vt)$$
 and $v = \frac{1}{\sqrt{\mu_0\epsilon_0}}$



\vec{E} , \vec{B} wave propagation in free space



General Solutions to the wave equation

$$\vec{E}(r,t) = \vec{E}_0 e^{i(\vec{k}\cdot\vec{r}-\omega t)}$$

 \vec{k} points in the direction of travel

$$\vec{B}(r,t) = \vec{B}_0 e^{i(\vec{k}\cdot\vec{r}-\omega t)}$$

 $v = \frac{\omega}{|\vec{k}|}$ is the velocity

$$v = \frac{1}{\sqrt{\mu_0 \epsilon_0}} \equiv c$$
 in free space : velocity of light in free space

$$v = \frac{1}{\sqrt{\mu\epsilon}}$$
 in a real material \equiv velocity of light in material

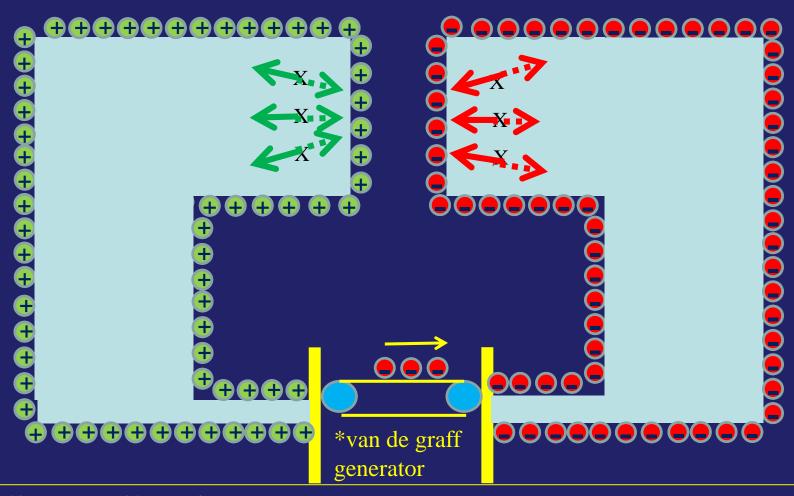
... more in Griffiths Sec 9.1, 9.2

Why are we doing wave in the last lecture?



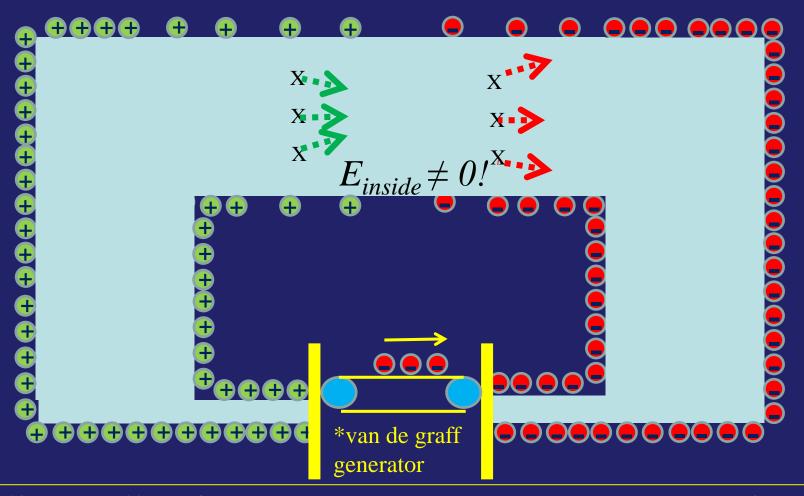
Walk into a room: flick the switch→ Light bulb goes ON

No free charge: EM wave propagation in a conductor



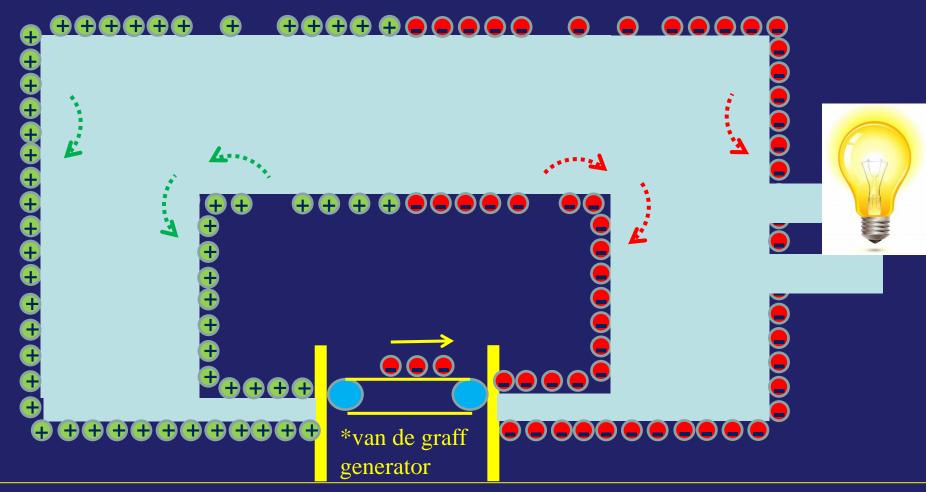
Switch closed

For a short time, $E_{inside} \neq 0$ $E = E(r, t) \sim \delta(r, t)$



Electric field travels in the conductor

 $E(r,t) \sim \delta(r,t)$ travels at velocity $v = \frac{1}{\sqrt{\mu\epsilon}} \rightarrow$ bulb turns ON



THANK YOU