

# MA-108 Ordinary Differential Equations

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D1 - Lecture 19

- ①  $L(e^{-at}f(t)) = F(s+a), \quad s > s_0 + a.$
- ②  $L(u(t-a)f(t-a)) = e^{-as}F(s), \quad s > s_0, a > 0.$
- ③  $L\left(\int_0^t f(\tau) d\tau\right) = \frac{F(s)}{s}, \quad s > \max\{0, s_0\}.$
- ④  $L(tf(t)) = -F^{(1)}(s), \quad s > s_0.$
- ⑤  $L\left(\frac{f(t)}{t}\right) = \int_s^\infty F(s')ds', \quad s > s_0.$
- ⑥ Assume  $f$  is piecewise continuous and of exponential order. Then
  - (i)  $\lim_{s \rightarrow \infty} F(s) = 0,$
  - (ii)  $\lim_{s \rightarrow \infty} sF(s)$  is bounded.
- ⑦ Assume  $f$  and  $f'$  both are piecewise continuous and of exponential order. Then  $\lim_{s \rightarrow \infty} sF(s) = f(0).$
- ⑧ If  $f$  is piecewise continuous and periodic of period  $T$ , then  $L(f(t)) = \frac{1}{1 - e^{-sT}} \int_0^T f(t)e^{-st}dt, \quad s > 0$

## Theorem

*Assume  $f$  and  $f'$  both are piecewise continuous and of exponential order. Then*

$$\lim_{s \rightarrow \infty} sF(s) = f(0).$$

**Proof.** Since

$$L(f'(t)) = sL(f(t)) - f(0)$$

Since  $f$  and  $f'$  both are piecewise continuous and of exponential order, we get

$$\lim_{s \rightarrow \infty} L(f'(t)) = 0, \text{ and } \lim_{s \rightarrow \infty} sF(s) < \infty$$

Therefore,

$$\lim_{s \rightarrow \infty} sF(s) = f(0)$$

**Ex.** Let  $f(t) = L^{-1} \left( \frac{1 - s(5 + 3s)}{s((s + 1)^2 + 1)} \right)$ . Find  $f(0)$ .

We can find  $f(t)$  by partial fraction. Hence we know that  $f$  and  $f'$  are continuous and of exponential order. Therefore,

$$\begin{aligned} f(0) &= \lim_{s \rightarrow \infty} sF(s) \\ &= \lim_{s \rightarrow \infty} \frac{1 - s(5 + 3s)}{((s + 1)^2 + 1)} \\ &= \lim_{s \rightarrow \infty} \frac{1 - 5s - 3s^2}{s^2 + 2s + 2} \\ &= -3 \end{aligned}$$

## Theorem

If  $f$  is piecewise continuous and periodic of period  $T$ , then

$$L(f(t)) = \frac{1}{1 - e^{-sT}} \int_0^T f(t) e^{-st} dt, \quad s > 0$$

**Proof.**

$$\begin{aligned} L(f(t)) &= \int_0^T f(t) e^{-st} dt + \int_T^{2T} f(t) e^{-st} dt + \dots \\ &= \int_0^T f(t) e^{-st} dt + \int_0^T f(t+T) e^{-s(t+T)} dt + \dots \\ &= \int_0^T f(t) e^{-st} dt (1 + e^{-sT} + e^{-2sT} + \dots) \\ &= \frac{1}{(1 - e^{-sT})} \int_0^T f(t) e^{-st} dt, \quad s > 0 \end{aligned}$$

**Ex.** Find the Laplace transform of periodic function

$$f(t) = \begin{cases} t, & 0 \leq t < 1 \\ 0, & 1 \leq t < 2 \end{cases}, \quad f(t+2) = f(t)$$

$$\begin{aligned} L(f(t)) &= \frac{1}{(1 - e^{-2s})} \int_0^2 f(t) e^{-st} dt \\ &= \frac{1}{(1 - e^{-2s})} \int_0^1 t e^{-st} dt \\ &= \frac{1}{(1 - e^{-2s})} \left[ t \frac{e^{-st}}{-s} \Big|_0^1 - \int_0^1 \frac{e^{-st}}{-s} dt \right] \\ &= \frac{1}{(1 - e^{-2s})} \left[ \frac{e^{-s}}{-s} - \frac{1}{s^2} (e^{-s} - 1) \right] \end{aligned}$$

**Exercise.** Find Inverse Laplace transform and varify, whether  $\lim_{s \rightarrow \infty} sF(s) = f(0)$ . If not, then state why it is not.

$$(i) \quad \frac{s}{(s^2 + a^2)^2}, \quad (ii) \quad \frac{s}{(s^2 - a^2)^2}$$

$$(iii) \quad \frac{s^2}{(s + 1)^3}, \quad (iv) \quad \frac{1}{\sqrt{s + 1}}$$

$$(v) \quad \frac{s}{(s - a)^{3/2}}, \quad (vi) \quad \frac{s}{(s + 4)^6}$$

$$(vii) \quad \frac{e^{-s}}{s^5}, \quad (viii) \quad \frac{e^{-2s}}{(s + 1)^2}$$

$$(ix) \quad \ln \left( 1 + \frac{a^2}{s^2} \right), \quad (x) \quad \ln \left( 1 - \frac{a^2}{s^2} \right),$$

$$(xi) \quad \frac{1}{s(1 - e^{-s})}, \quad (xii) \quad \frac{1}{s(1 + e^{-s})}$$

$$(xiii) \quad \frac{s}{(s^2 + 1)^{3/2}}, \quad (xiv) \quad \frac{1}{(s + 1)(1 - e^{-2s})}$$

# Constant coefficient equations with Impulses

In this section, we will consider ODE  $ay'' + by' + cy = f(t)$ ,  $a, b, c \in \mathbb{R}$ , and  $f(t)$  represents a force that is very large for a short time and zero otherwise. Such forces are **impulsive**.

Impulsive forces occur when two objects collide, e.g. the force due to hammer blow. Since it will be difficult to know  $f(t)$ , how do we solve our ODE.

If  $f$  is an integrable function and  $f(t) = 0$  outside  $[t_0, t_0 + h]$ , then  $I = \int_{t_0}^{t_0+h} f(t) dt$  is called the **total impulse** of  $f$ . This is total area under the graph of  $f$ .

The idea is that if  $h$  is very small, then  $y(t)$  will be sensitive to  $I$ , but rather insensitive to shape of  $f$ . Therefore, if we vary the shape of  $f$ , but keep the area same, then  $y(t)$  will change very little.



# Constant coefficient equations with Impulses

Thus we replace  $f$  by a simple rectangular pulse  $f(t) = I/h$  on  $[t_0, t_0 + h]$  and zero outside  $[t_0, t_0 + h]$ .

Now the ODE gets simplified, but the solution still depends on  $h$ , which may be difficult to find for  $f$ . So we take the ideal situation, and make  $h \rightarrow 0$  to eliminate  $h$ .

Define a rectangular pulse having unit impulse as

$$D(t - t_0; h) = \begin{cases} 1/h, & t \in [t_0, t_0 + h] \\ 0, & \text{otherwise} \end{cases}$$

Here  $D$  is for physicist P.A. Dirac, who developed the idea of impulsive forces in 1929.

$$\delta(t - t_0) := \lim_{h \rightarrow 0} D(t - t_0; h) = \begin{cases} \infty, & t = t_0 \\ 0, & \text{otherwise} \end{cases}$$

$\delta(t - t_0)$  is not an ordinary function, but we can think of it as idealised *point unit impulse* focused at  $t = t_0$ .

# Constant coefficient equations with Impulses

Let us prove that if  $g$  is continuous at  $t = t_0$ , then

$$\lim_{h \rightarrow 0} \int_0^{\infty} g(t) D(t - t_0; h) dt = g(t_0)$$

**Proof.** we may assume that  $g(t)$  is continuous on  $[t_0, t_0 + h]$ , by taking  $h$  small. Then

$$\begin{aligned} \lim_{h \rightarrow 0} \int_0^{\infty} g(t) D(t - t_0; h) dt &= \lim_{h \rightarrow 0} \int_{t_0}^{t_0+h} g(t) \frac{1}{h} dt \\ &= \lim_{h \rightarrow 0} \frac{1}{h} g(t_0 + h_1) h, \text{ for some } h_1 \in [0, h], \text{ by MVT} \\ &= g(t_0) \end{aligned}$$

In particular,  $\int_0^{\infty} \delta(t - t_0) dt = 1$ , take  $g = 1$ .

The “generalised function”  $\delta(t - t_0)$  is called **Dirac delta function** or **unit impulse function**.

Let's try to define the meaning of solution of IVP

$$ay'' + by' + cy = \delta(t - t_0), \quad y(0) = 0, \quad y'(0) = 0, \quad t_0 > 0 \quad (*)$$

### Theorem

Fix  $t_0 \geq 0$ . For each  $h > 0$ , let  $y_h$  be the solution of IVP

$$ay_h'' + by_h' + cy_h = D(t - t_0; h), \quad y_h(0) = 0, \quad y_h'(0) = 0, \quad \text{Then}$$

$$\lim_{h \rightarrow 0+} y_h(t) = u(t - t_0)w(t - t_0),$$

where

$$w(t) = L^{-1} \left( \frac{1}{as^2 + bs + c} \right)$$

Therefore the solution of  $(*)$  is defined as

$$y(t) = u(t - t_0)w(t - t_0), \quad \text{where } w(t) = L^{-1} \left( \frac{1}{as^2 + bs + c} \right)$$

Note  $w(t) = L^{-1} \left( \frac{1}{as^2 + bs + c} \right)$  is the solution of the IVP

$$aw'' + bw' + cw = 0, \quad w(0) = 0, \quad w'(0) = \frac{1}{a}$$

In fact  $w(t)$  is defined on  $(-\infty, \infty)$ , and is given by

$$w = \frac{e^{m_2 t} - e^{m_1 t}}{a(m_2 - m_1)}, \quad w = \frac{1}{a} t e^{m_1 t}, \quad \frac{1}{a\omega} e^{\lambda t} \sin \omega t$$

depending on whether the characteristic polynomial  $p(m) = am^2 + bm + c$  has distinct real roots  $m_1, m_2$ , or repeated real roots  $m_1, m_1$ , or complex conjugate roots  $\lambda \pm i\omega$ .

This means, if  $t_0 > 0$ , then the solution

$$\boxed{y(t) = u(t - t_0)w(t - t_0)} \text{ of the IVP}$$

$$ay'' + by' + cy = \delta(t - t_0), \quad y(0) = 0, \quad y'(0) = 0, \quad t_0 > 0 \quad (*)$$

is defined on  $(-\infty, \infty)$  and has the following properties.

$$\begin{aligned} y(t) &= 0 \text{ for all } t < t_0 \\ ay'' + by' + cy &= 0 \text{ on } (-\infty, t_0) \text{ and } (t_0, \infty) \\ y'(t_0-) &= 0, \quad y'(t_0+) = \frac{1}{a} \end{aligned}$$

When  $t_0 = 0$ ,  $y'(0-)$  is not defined, so in this case

$$y(t) = u(t - t_0)w(t - t_0) = u(t)w(t)$$

is a solution of

$$ay'' + by' + cy = \delta(t), \quad y(0) = 0, \quad y'(0+) = 0$$

**Ex.** Solve  $y'' + 2y' + y = \delta(t - t_0)$ ,  $y(0) = 0$ ,  $y'(0) = 0$ . Here

$$\begin{aligned}w(t) &= L^{-1} \left( \frac{1}{s^2 + 2s + 1} \right) \\&= L^{-1} \left( \frac{1}{(s + 1)^2} \right) \\&= e^{-t}t\end{aligned}$$

Therefore, the solution is given by

$$\begin{aligned}y(t) &= u(t - t_0)w(t - t_0) \\&= u(t - t_0)e^{-(t-t_0)}(t - t_0)\end{aligned}$$

**Ex.** Solve

$$y'' + 6y' + 5y = 3e^{-2t} + 2\delta(t - 1), \quad y(0) = -3, \quad y'(0) = 2.$$

If  $y_1(t)$  is a solution of

$$y'' + 6y' + 5y = 3e^{-2t}, \quad y(0) = -3, \quad y'(0) = 2, \text{ then}$$

$$y_1(t) = c_1 e^{-t} + c_2 e^{-5t} - e^{-2t}, \text{ where } c_1 = -\frac{5}{2} \text{ and } c_2 = \frac{1}{2}.$$

The solution of IVP is  $y(t) = y_1 + y_2$ , where  $y_2$  is a solution of

$$y'' + 6y' + 5y = 2\delta(t - 1), \quad y(0) = 0, \quad y'(0) = 0. \text{ Hence}$$

$$\begin{aligned} w(t) &= 2L^{-1} \left( \frac{1}{s^2 + 6s + 5} \right) \\ &= \frac{1}{2} \left( \frac{1}{s+1} - \frac{1}{s+5} \right) \\ &= \frac{1}{2} (e^{-t} - e^{-5t}) \end{aligned}$$

$$y(t) = -\frac{5}{2}e^{-t} + \frac{1}{2}e^{-5t} - e^{-2t} + \frac{1}{2}u(t-1) (e^{-(t-1)} - e^{-5(t-1)})$$

**Exercise.** Solve the following problems.

- ①  $y'' + 3y' + 2y = 6e^{2t} + 2\delta(t - 1), \quad y(0) = 2, \quad y'(0) = -6.$
- ②  $y'' + y = \sin 3t + 2\delta(t - \pi/2), \quad y(0) = 1, \quad y'(0) = -1.$
- ③  $y'' + 2y' + 2y = \delta(t - \pi) - 3\delta(t - 2\pi), \quad y(0) = -1, \quad y'(0) = 2.$
- ④  $y'' + 4y = f(t) + \delta(t - 2\pi), \quad y(0) = 0, \quad y'(0) = 1$  and
$$f(t) = \begin{cases} 1, & 0 \leq t < \pi/2 \\ 2, & t \geq \pi/2 \end{cases}$$
- ⑤  $y'' + 4y' + 4y = -\delta(t), \quad y(0) = 1, \quad y'(0+) = 5.$
- ⑥ Find a solution not involving unit step function which represents  $y$  on each subinterval of  $[0, \infty)$  on which the forcing function is zero.
  - (a)  $y'' - y = \sum_{k=1}^{\infty} \delta(t - k), \quad y(0) = 0, \quad y'(0) = 1$
  - (b)  $y'' - 3y' + 2y = \sum_{k=1}^{\infty} \delta(t - k), \quad y(0) = 0, \quad y'(0) = 1$