Sinusoidal steady state analysis (Phasors)

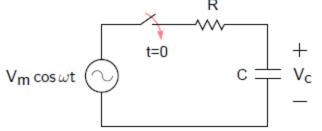
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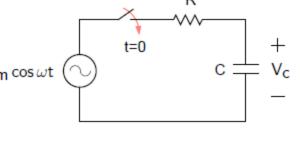
Reference: L. Bobrow

Sinusoidal signals (AC signals)

- Response to repetitive, sinusoidal (AC) signal
 - Ubiquitous
 - Ordinary household voltage
 - Radio (AM/FM)
 - Television etc.
- In a lot of circuit applications
 - Power consumed/supplied is of importance
 - Instantaneous power (ratings in circuits)

Sinusoidal analysis (time-domain, steady state)





$$RCV_c' + V_c = V_m \cos \omega t$$

$$V_c(t) = V_f(t) + V_n(t)$$

$$V_n(t) = Ae^{-\frac{t}{RC}}$$

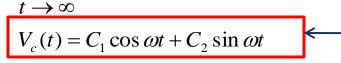
$$V_f(t) = C_1 \cos \omega t + C_2 \sin \omega t$$

Substituting

$$\omega RC(-C_1 \sin \omega t + C_2 \cos \omega t) + C_1 \cos \omega t + C_2 \sin \omega t = V_m \cos \omega t$$

C₁ an C₂ can be found by equating coefficients on both sides

$$V_c(t) = Ae^{-\frac{t}{RC}} + C_1 \cos \omega t + C_2 \sin \omega t$$



Sinusoidal steady-state response

- We assume the forced response of RLC ckts to a sinusoidal input to be a sinusoid
 - Each of the elements R, L, C will respond with a sinusoid to a sinusoidal input; no change in frequency ω
 - Sum of sinusoidal outputs is also a sinusoid
- Any circuit containing resistors, capacitors, inductors, sinusoidal voltage and current sources (of the same frequency), dependent (linear) sources behaves in a similar manner, viz., each current and voltage in the circuit becomes purely sinusoidal as t \rightarrow infinity

Real sinusoid -> Complex Sinusoid

RLC ckt Output Input $A\cos(\omega t + \theta)$ $B\cos(\omega t + \phi)$ $A\sin(\omega t + \theta)$ $B\sin(\omega t + \phi)$ Shift by 90deg $KB\sin(\omega t + \phi)$ $KA\sin(\omega t + \theta)$ Scale by K, linearity Superposition $B\cos(\omega t + \phi) + KB\sin(\omega t + \phi)$ $A\cos(\omega t + \theta) + KA\sin(\omega t + \theta)$ K=j $A\cos(\omega t + \theta) + jA\sin(\omega t + \theta)$ $B\cos(\omega t + \phi) + jB\sin(\omega t + \phi)$ $Ae^{j(\omega t+\theta)}$ Euler's formula $Be^{j(\omega t + \varphi)}$

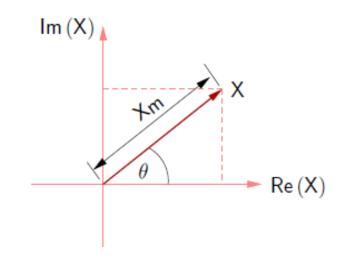
- Finding soln (B, φ) to real sinusoid is equivalent to finding the soln (B, φ) to the complex sinusoid (B, φ)
- Using complex numbers makes differential eqn soln easier

Phasors: Definition

$$\mathbf{X} = X_m e^{j\theta} = X_m \angle \theta$$

$$x(t) = \text{Re}[\mathbf{X}e^{j\omega t}] = \text{Re}[X_m e^{j\theta}e^{j\omega t}] = X_m \cos(\omega t + \theta)$$

$$X_m \angle \theta = X_m \cos \theta + X_m j \sin \theta$$



- Phasor is a complex number
- The term ωt is implicit
- Can also be written in the polar or rectangular form
- Also called frequency domain representation
- Significant simplification of sinusoidal steady state analysis

Phasor: Examples

Time domain	Frequency domain
$v_1(t)=3.2\cos(\omega t+30^\circ) V$	$V_1 = 3.2 \angle 30^\circ = 3.2 \exp{(j\pi/6)} V$
$i(t) = -1.5 \cos (\omega t + 60^{\circ}) A$ = $1.5 \cos (\omega t + \pi/3 - \pi) A$ = $1.5 \cos (\omega t - 2\pi/3) A$	$I = 1.5 \angle (-2\pi/3) A$
$v_2(t) = -0.1\cos(\omega t) V$ $= 0.1\cos(\omega t + \pi) V$	$V_2 = 0.1 \angle \pi V$
$i_2(t) = 0.18 \sin(\omega t) A$ = $0.18 \cos(\omega t - \pi/2) A$	$I_2 = 0.18 \angle (-\pi/2) A$
$i_3(t) = \sqrt{2}\cos(\omega t + 45^\circ) A$	$I_3 = 1 + j 1 A$ = $\sqrt{2} \angle 45^{\circ} A$

Phasor representation/Impedance of resistor

$$\begin{split} i(t) &= I_m \cos(\omega t + \theta) \\ v(t) &= Ri(t) = RI_m \cos(\omega t + \theta) = V_m \cos(\omega t + \theta) \\ \text{Re}[V_m e^{j(\omega t + \theta)}] &= \text{Re}[RI_m e^{j(\omega t + \theta)}] \end{split}$$

$$\mathbf{V} = R\mathbf{I}$$

Phasor representation

$$\mathbf{Z} = \mathbf{V/I} = R + j0$$

Z is the impedance of the resistor

Phasor representation/Impedance of capacitor

Time domain

Phasor



$$v(t) = V_m \cos(\omega t + \theta)$$

$$i(t) = C\frac{dv(t)}{dt} = -C\omega V_m \sin(\omega t + \theta) = C\omega V_m \cos(\omega t + \theta + \frac{\pi}{2})$$

$$\mathbf{V} = V_m \angle \theta$$

$$\mathbf{I} = C\omega V_{m} \angle (\theta + \frac{\pi}{2}) = \omega C V_{m} e^{j\theta} e^{j\frac{\pi}{2}} = j\omega C \mathbf{V}$$

Phasor representation

$$\mathbf{Z} = \mathbf{V/I} = 0 + \frac{1}{j\omega C}$$

$$\mathbf{Y} = \mathbf{I/V} = 0 + j\omega C$$

Z is the impedance of the capacitor

Y is the admittance of the capacitor

Phasor representation/Impedance of inductor

Time domain

Phasor

$$i(t) = I_m \cos(\omega t + \theta)$$

$$v(t) = L \frac{di(t)}{dt} = -L\omega I_m \sin(\omega t + \theta) = L\omega I_m \cos(\omega t + \theta + \frac{\pi}{2})$$

$$\mathbf{V} = L\omega I_m \angle (\theta + \frac{\pi}{2}) = \omega L I_m e^{j\theta} e^{j\frac{\pi}{2}} = j\omega L \mathbf{I}$$

Phasor representation

$$\mathbf{Z} = \mathbf{V}/\mathbf{I} = 0 + i\omega L$$

 $\mathbf{I} = I_m \angle \theta$

$$\mathbf{Y} = \mathbf{I/V} = 0 + \frac{1}{j\omega L}$$

Z is the impedance of the capacitor

Y is the admittance of the capacitor

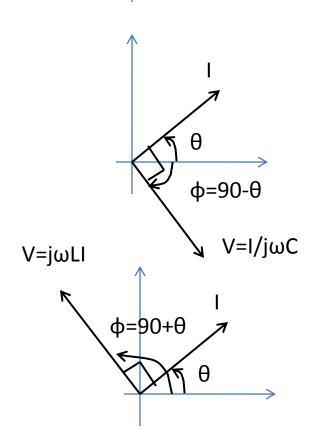
Lead and Lag

$$V = RI$$

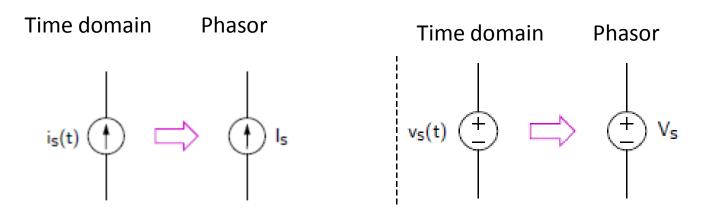
$$\mathbf{I} = C\omega V_m \angle (\theta + \frac{\pi}{2}) = \omega C V_m e^{j\theta} e^{j\frac{\pi}{2}} = j\omega C \mathbf{V}$$

$$\mathbf{V} = L\omega I_{m} \angle (\theta + \frac{\pi}{2}) = \omega L I_{m} e^{j\theta} e^{j\frac{\pi}{2}} = j\omega L \mathbf{I}$$

- Resistor → in phase
- Capacitor current *leads* V by 90deg
- Inductor → voltage *leads* I by 90deg



Sources



- An independent sinusoidal current source, $i_s(t) = I_m \cos(\omega t + \theta)$, can be represented by the phasor $I_m < \theta$ (i.e., a constant complex number)
- An independent sinusoidal voltage source, $v_s(t) = V_m \cos(\omega t + \theta)$, can be represented by the phasor $V_m < \theta$ (i.e., a constant complex number)
- Dependent (linear) sources can be treated in the sinusoidal steady state in the same manner as a resistor, i.e., by the corresponding phasor relationship. For example, for a CCVS, we have,

 $v(t) = r i_c(t)$ in the time domain

 $V = r I_c$ in the frequency domain

Phasor addition

$$\begin{split} v(t) &= v_1(t) + v_2(t) = V_{m1} \cos(\omega t + \theta_1) + V_{m2} \cos(\omega t + \theta_2) \\ \mathbf{V} &= \mathbf{V_1} + \mathbf{V_2} = V_{m1} e^{j\theta_1} + V_{m2} e^{j\theta_2} \\ \tilde{v}(t) &= \text{Re}[(V_{m1} e^{j\theta_1} + V_{m2} e^{j\theta_2}) e^{j\omega t}] = V_{m1} \cos(\omega t + \theta_1) + V_{m2} \cos(\omega t + \theta_2) = v(t) \end{split}$$

- Addition of sinusoidal quantities in the time domain can be replaced by addition of the corresponding phasors in the sinusoidal steady state.
- KCL and KVL in time domain hold for phasors also

For meshes
$$\sum v_k(t) = 0$$

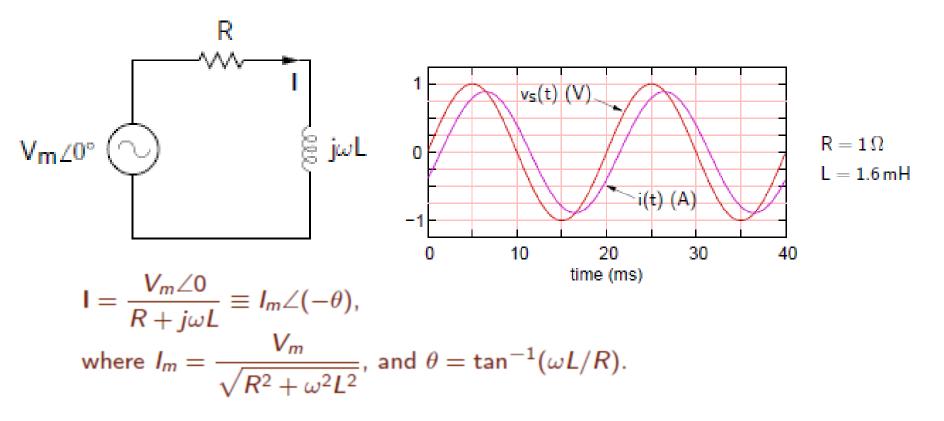
$$\sum i_k(t) = 0$$

$$\sum \mathbf{I_k} = 0$$

Sinusoidal Steady State Analysis

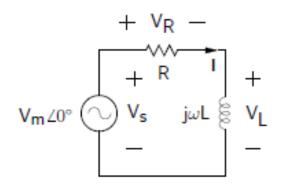
- Resistors, capacitors, and inductors can be described by V = Z I in the frequency domain, which is similar to V = R I in DC conditions (except that we are dealing with complex numbers in the frequency domain)
- An independent sinusoidal source in the frequency domain behaves like a DC source, e.g., V_s=constant (a complex number)
- For dependent sources, a time-domain relationship such as $i(t)=\beta i_c(t)$ translates to $l=\beta l_c$ in the frequency domain
- Circuit analysis in the sinusoidal steady state using phasors is therefore very similar to DC circuits with independent and dependent sources, and resistors
- Series/parallel formulas for resistors, nodal analysis, mesh analysis, Thevenin's and Norton's theorems can be directly applied
- No differential equations!
- This is called the frequency domain approach

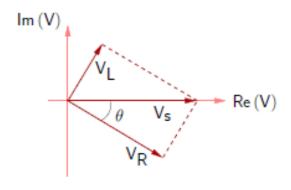
Example: RL circuit



- In the time domain, $i(t) = I_m \cos(\omega t \theta)$, which lags the source voltage since the peak (or zero) of i(t) occurs $t = \theta/\omega$ seconds after that of the source voltage.
- For R = 1Ω , L = 1.6 mH, f = 50 Hz, θ = 26.6° , t_{lag} = 1.48 ms.

Example: RL Circuit (contd)





$$\begin{split} \mathbf{I} &= \frac{V_m \angle 0}{R + j\omega L} \equiv I_m \angle (-\theta), \\ \text{where } I_m &= \frac{V_m}{\sqrt{R^2 + \omega^2 L^2}}, \text{ and } \theta = \tan^{-1}(\omega L/R). \end{split}$$

$$KVL: V_s = V_R + V_L$$

$$V_R = I \times R = R I_m \angle (-\theta)$$
,
 $V_L = I \times j\omega L = \omega I_m L \angle (-\theta + \pi/2)$,

- If R >> $j\omega L$, $\theta \rightarrow 0$, $|V_R| \approx |V_S| = V_m$.
- If R << $j\omega L$, $\theta \rightarrow \pi/2$, $|V_L| \approx |V_S| = V_m$.