

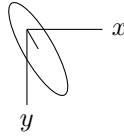
Notes April 25

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0.1 Simple Pendulum

Say we have some rigid body



Small oscillations, we have

$$\hat{r} = r(\hat{y} \cos(\theta + \varphi) + \hat{x} \sin(\theta + \varphi))$$

and

$$\vec{v} = \dot{\theta} r [-\hat{y} \sin(\theta + \varphi) + \hat{x} \cos(\theta + \varphi)]$$

then for U we write

$$U = mgd(1 - \cos \theta)$$

with

$$E = \frac{1}{2} I \dot{\theta}^2 + mgd(1 - \cos \theta) \dot{\theta} = \pm \sqrt{\frac{2}{I} (E - mgd(1 - \cos \theta))}$$

So we think about $1 - \cos \theta = 2 \sin^2(\frac{\theta}{2})$. Let θ_0 be the max value of θ , i.e. where our approximation starts to fall apart.

To solve, we can evaluate

$$T = \frac{2r_0}{\sqrt{gd}} \int_0^{\theta_0} \frac{d\theta}{\sqrt{\sin^2 \frac{\theta_0}{2} - \sin^2 \frac{\theta}{2}}}$$

Let $k = \sin \frac{\theta_0}{2}$ and $k_z = \sin \frac{\theta}{2}$, so we write

$$dz = \frac{1}{2} \frac{\cos \theta/2}{k} d\theta$$

which makes

$$T = \frac{2r}{\sqrt{gd}} \int_{z=0}^1 \frac{z k dz}{\sqrt{1 - k^2 z^2}} \frac{1}{k \sqrt{1 - z^2}}$$

in the small angle approximation, we have

$$\sqrt{1 - k^2 z^2} \approx 1 + \frac{1}{2} k^2 z^2$$

which allows us to break up T into some shit that requires a lot of trig substitution, which ultimately yields

$$T = \frac{4r_0}{\sqrt{gd}} \left(\frac{\pi}{2} + \frac{1}{2} k^2 \frac{\pi}{4} \right) = \frac{2\pi r_0}{\sqrt{gd}} \left(1 + \frac{1}{4} \sin^2 \frac{\theta_0}{2} + \dots \right)$$

which gives us a term T_0 that tells us how much variation we have from the simple harmonic oscillator

$$T_0 = \frac{2\pi r_0}{\sqrt{gd}}$$

For instance, if we plug in $\theta_0 = 23^\circ$, we get about a 1% change.

Now, we do some

0.1.1 Pertubation Theory

and write down the ODE

$$\ddot{x} + \omega_0^2 x - \lambda x^2 = 0$$

which, looking for $x(\lambda, t)$

$$x(\lambda, t) = x_0(t) + \lambda x_1(t) + \lambda^2 x_2(t) + \dots$$

$$\begin{aligned}\dot{x} &= \dot{x}_0 + \lambda \dot{x}_1 \\ \ddot{x} &= \ddot{x}_0 + \lambda \ddot{x}_1\end{aligned}$$

which gives

$$\begin{aligned}\ddot{x} + \omega_0^2 x - \lambda x^2 &= 0 \\ \ddot{x}_0 + \lambda \ddot{x}_1 + \omega_0^2 (x_0 + \lambda x_1) - \lambda (x_0 + \lambda x_1)^2 &= 0 \\ \ddot{x}_0 + \lambda \ddot{x}_1 + \omega_0^2 \lambda x_1 - \lambda x_0^2 - 2\lambda^2 x_0 x_1 - \lambda^3 x_1^2 &= 0\end{aligned}$$

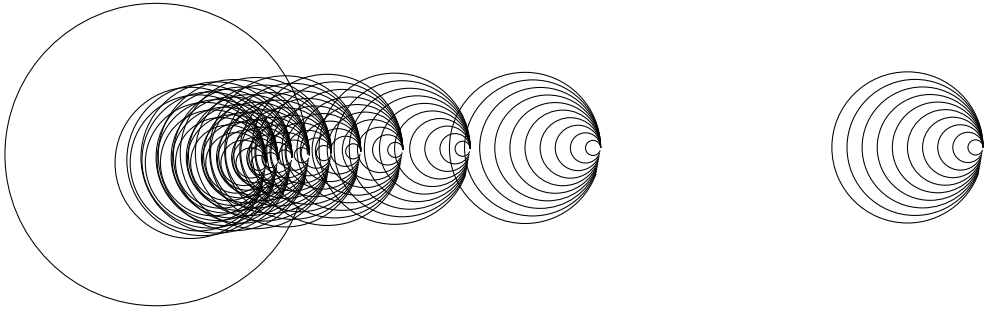
first order perturbation of the thing, so if we're forcing it

$$\ddot{x}_1 + \omega_0^2 x_1 = A^2 \cos^2 \omega_0 t$$

So, what I think he did

$$x(t) = x_0 + \lambda x_1 = A \cos \omega_0 t - \lambda \frac{A^2}{6\omega_0^2} (\cos(2\omega_0 t) - 3)$$

3-wave coupling.
Kolomogorov (1941). No natural scale/boundary condition in hte problem
Imagine some eddy, that you've just broken up



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