

## Abstract

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# 1 Stationary Perturbation Theory

## 1.1 Formalism for non-degenerate eigenstates

We can try to solve unsolvable problems by dividing the hamiltonian in two, taking

$$\hat{H} = \hat{H}_0 + \lambda \hat{H}_1$$

with  $\hat{H}_0$  being some hamiltonian we can solve exactly, with

$$H_0 \left| \varphi_n^{(0)} \right\rangle = E_n^{(0)} \left| \varphi_n^{(0)} \right\rangle$$

### 1.1.1 First order energy correction

Let's take our first order expansion of  $|n\rangle = \left| n^{(0)} \right\rangle + \lambda \left| n^{(1)} \right\rangle \dots$

$$(H_0 + \lambda H_1) \sum_{i=0}^{\infty} \lambda^i \left| n^{(i)} \right\rangle = \sum_{i=0}^{\infty} \lambda^i E_n^{(0)} \sum_{i=0}^{\infty} \lambda^i \left| n^{(i)} \right\rangle$$

The first order correction comes out to be that the first order correction is  $E_n^{(1)} = \left\langle n^{(0)} \right| H_1 \left| n^{(0)} \right\rangle$ . Stated more consicely, it is

$$E_n^{(1)} = \langle H_1 \rangle_{(0)}$$

### 1.1.2 First order state correction

If we want to find the first order correction to our state  $|n\rangle$ , we should take some  $\sum_m c_m \left| m^{(0)} \right\rangle$ , where  $|m\rangle$  is the “natural basis” given by solving the eigenvalue equation of the unperturbed case. We can take

$$\left\langle m^{(0)} \right| H_0 \left| n^{(1)} \right\rangle + \left\langle m^{(0)} \right| H_1 \left| n^{(0)} \right\rangle = \left\langle m^{(0)} \right| n^{(1)} \right\rangle E_n^{(0)}$$

Of course,  $\left\langle m^{(0)} \right| n^{(1)} \right\rangle$  is the  $c_1$  correction we wanted from the sum above, so we can left multiply and get

$$E_0 \left\langle m^{(0)} \right| n^{(1)} \right\rangle + \left\langle m^{(0)} \right| H_1 \left| n^{(0)} \right\rangle = \left\langle m^{(0)} \right| n^{(1)} \right\rangle E_n^{(0)}$$

which gives

$$\left\langle m^{(0)} \right| n^{(1)} \right\rangle E_n^{(0)} = \frac{\left\langle m^{(0)} \right| H_1 \left| n^{(0)} \right\rangle}{E_n^{(0)} - E_m^{(0)}}$$

We're still missing the  $\left| n^{(0)} \right\rangle$  term, but if we use the fact that  $|n\rangle$  is normalized, we should have

$$\langle n | n \rangle = 1 = \langle n^{(0)} | n^{(0)} \rangle + 2\lambda \operatorname{Re} \langle n^{(0)} | n^{(1)} \rangle + O(\lambda^2)$$

This yields that  $2\lambda \operatorname{Re} \langle n^{(0)} | n^{(1)} \rangle \approx 0$  to order  $\lambda^2$ , which lets us write down

$$|n\rangle = \left| n^{(0)} \right\rangle + i\alpha \left| n^{(0)} \right\rangle + \lambda \sum_{m \neq n} \frac{\left\langle m^{(0)} \right| H_1 \left| n^{(0)} \right\rangle}{E_n^{(0)} - E_m^{(0)}}$$

If we want to correct the parallel component, we make a small angle approximation to write

$$|n\rangle = (1 + i\alpha) \left| n^{(0)} \right\rangle + \dots = e^{i\alpha} \left| n^{(0)} \right\rangle + \dots$$

which lets us simplify even further by writing

$$|n\rangle = \left| n^{(0)} \right\rangle + (1 - i\alpha)\lambda \sum_{m \neq n} \frac{\left\langle m^{(0)} \right| H_1 \left| n^{(0)} \right\rangle}{E_n^{(0)} - E_m^{(0)}}$$

and since  $\alpha$  is of order  $\lambda$ , we can just drop it, since thier product is of order  $\lambda^2$  Giving the final first order correction to the state is just (dropping  $\lambda$  by setting it equal to one, and letting  $H_1$  be small compared to the energy difference)

$$|n\rangle = \left| n^{(0)} \right\rangle + \sum_{m \neq n} \left( \frac{\left\langle m^{(0)} \right| H_1 \left| n^{(0)} \right\rangle}{E_n^{(0)} - E_m^{(0)}} \right) \left| m^{(0)} \right\rangle$$

## 1.2 Example: Harmonic Oscillation + Perturbations

## 1.3 Degenerate Perturbation Theory

## 1.4 More Examples!