

Solutions - Christian Thierfelder

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1 Advanced Topics in Gravity – Exercise sheet 5 - 2025-06-24

1.1 Exercise 1 - Energy Conditions

Analyze the different energy conditions: whether they hold, are violated, and which conditions must be satisfied if needed.

1. Consider a cosmological constant such that:

$$T_{\alpha\beta}^{(\Lambda)} = -\rho_{\Lambda} g_{\alpha\beta} \quad (1.1)$$

where

$$\rho_{\Lambda} = -p_{\Lambda} = \frac{\Lambda}{8\pi G}, \quad \text{and} \quad \left(T_{\alpha\beta} - \frac{1}{2} g_{\alpha\beta} T \right) = \rho_{\Lambda} g_{\alpha\beta}.$$

2. Consider a perfect fluid

$$T_{\alpha\beta} = (\rho + p)u_{\alpha}u_{\beta} + pg_{\alpha\beta}, \quad \text{with,} \quad (1.2)$$

with $g^{\alpha\beta}u_{\alpha}u_{\beta} = -1$ and

$$T^{\alpha}_{\alpha} = -\rho + 3p, \quad \text{so} \quad \left(T_{\alpha\beta} - \frac{1}{2} g_{\alpha\beta} T \right) = p g_{\alpha\beta}$$

Note:

These two exercises may be substituted for one of the analysis topics mentioned during the lecture today, such as:

- Misner spacetime geodesics and lightcones
- Exotic matter and safety requirements for travellers in Morris-Thorne wormholes

Taking the trace of the Einstein field equations (signature $[-, +, +, +]$)

$$R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} = \kappa T_{\mu\nu} \quad (1)$$

$$\rightarrow R - \frac{1}{2} 4R = \kappa T \quad (2)$$

$$\rightarrow -R = \kappa T \quad (3)$$

$$R_{\mu\nu} = \kappa \left(T_{\mu\nu} - \frac{1}{2} T g_{\mu\nu} \right) \quad (4)$$

where we used $g_{\mu\nu}g^{\nu\rho} = \delta_{\mu}^{\rho}$ and $g_{\mu\nu}g^{\mu\nu} = \delta_{\mu}^{\mu} = 4$.

1. With EFE:

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = \kappa \left(\underbrace{T_{\mu\nu}}_{\equiv 0} - \frac{\Lambda}{\kappa}g_{\mu\nu} \right) \quad (5)$$

$$= \kappa \left(\underbrace{-\frac{\Lambda}{\kappa}g_{\mu\nu}}_{=T^\Lambda} \right) \quad (6)$$

So we see that the cosmological constant can be understood as a special fluid with $p = -\Lambda/\kappa$ and $\rho = -p = \Lambda/\kappa$

- *WEC - weak energy condition* $T_{\alpha\beta}t^\alpha t^\beta \geq 0$ for all timelike t^α ($g_{\alpha\beta}t^\alpha t^\beta < 0$)

$$T_{\alpha\beta}t^\alpha t^\beta = -\frac{\Lambda}{\kappa} \underbrace{g_{\alpha\beta}t^\alpha t^\beta}_{<0} \quad (7)$$

$$= \frac{\Lambda}{\kappa} \quad (8)$$

Holds if $\Lambda \geq 0$.

- *NEC - null energy condition* $T_{\alpha\beta}l^\alpha l^\beta \geq 0$ for all null l^α

$$T_{\alpha\beta}l^\alpha l^\beta = -\frac{\Lambda}{\kappa} \underbrace{g_{\alpha\beta}l^\alpha l^\beta}_{=0} \quad (9)$$

$$= 0 \quad (10)$$

Holds.

- *SEC - strong energy condition* $(T_{\alpha\beta} - \frac{1}{2}Tg_{\alpha\beta})u^\alpha u^\beta \geq 0$ for all unit timelike u^α

$$(T_{\alpha\beta} - \frac{1}{2}Tg_{\alpha\beta})u^\alpha u^\beta = \left[-\frac{\Lambda}{\kappa} - \frac{1}{2}4\left(-\frac{\Lambda}{\kappa}\right) \right] g_{\alpha\beta}u^\alpha u^\beta \quad (11)$$

$$= -\frac{\Lambda}{\kappa}(-1)g_{\alpha\beta}u^\alpha u^\beta \quad (12)$$

$$= \frac{\Lambda}{\kappa}g_{\alpha\beta}u^\alpha u^\beta \quad (13)$$

$$= 0 \quad (14)$$

Holds if $\Lambda \geq 0$

- *DEC - Dominant energy condition* $-T^\alpha_\beta u^\beta$ is a future directed timelike or null vector for all future directed timelike vectors u^α

For a perfect fluid the DEC holds when $\rho \geq |p|$. This means we require $\rho_\Lambda \geq |p_\Lambda| = \rho_\Lambda$ - so it holds.

2. **The problem is a bit unclear - do we consider just a ideal fluid or fluid plus cosmological constant (there is a p_Λ in the last equation).** We therefore consider the more complicated case and can set $\Lambda = 0$ if needed

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = \kappa \left(T_{\mu\nu}^{\text{PF}} - \frac{\Lambda}{\kappa}g_{\mu\nu} \right) \quad (15)$$

$$= \kappa \left((\rho + p)u_\mu u_\nu + pg_{\mu\nu} - \frac{\Lambda}{\kappa}g_{\mu\nu} \right) \quad (16)$$

$$= \kappa \left((\rho + p)u_\mu u_\nu + \left[p - \frac{\Lambda}{\kappa} \right] g_{\mu\nu} \right) \quad (17)$$

$$(18)$$

then

$$T_{\mu\nu} = (\rho + p)u_\mu u_\nu + \left[p - \frac{\Lambda}{\kappa}\right] g_{\mu\nu} \quad (19)$$

$$\rightarrow T^\mu_\mu = (\rho + p) \underbrace{u^\mu u_\mu}_{=-1} + 4 \left[p - \frac{\Lambda}{\kappa}\right] \quad (20)$$

$$= 3p - \rho - 4 \frac{\Lambda}{\kappa} \quad (21)$$

$$\rightarrow T_{\mu\nu} - \frac{1}{2} g_{\mu\nu} T = (\rho + p)u_\mu u_\nu + \left[p - \frac{\Lambda}{\kappa}\right] g_{\mu\nu} - \frac{1}{2} (3p - \rho - 4 \frac{\Lambda}{\kappa}) g_{\mu\nu} \quad (22)$$

$$= (\rho + p)u_\mu u_\nu + \left[-\frac{1}{2}p + \frac{1}{2}\rho + \frac{\Lambda}{\kappa}\right] g_{\mu\nu} \quad (23)$$

- *WEC - weak energy condition* $T_{\alpha\beta} t^\alpha t^\beta \geq 0$ for all timelike t^α ($g_{\alpha\beta} t^\alpha t^\beta < 0$)

$$T_{\alpha\beta} t^\alpha t^\beta = (\rho + p) \underbrace{u_\alpha u_\beta t^\alpha t^\beta}_{=(u_\alpha t^\alpha)^2 \geq 0} + \left[p - \frac{\Lambda}{\kappa}\right] \underbrace{g_{\alpha\beta} t^\alpha t^\beta}_{< 0} \quad (24)$$

Holds if $p + \rho \geq 0$ and $p - \frac{\Lambda}{\kappa} \geq 0$.

- *NEC - null energy condition* $T_{\alpha\beta} l^\alpha l^\beta \geq 0$ for all null l^α
See above - holds if $p + \rho \geq 0$
- *SEC - strong energy condition* $(T_{\alpha\beta} - \frac{1}{2} T g_{\alpha\beta}) u^\alpha u^\beta \geq 0$ for all unit timelike u^α

$$(T_{\alpha\beta} - \frac{1}{2} T g_{\alpha\beta}) u^\alpha u^\beta = (\rho + p) \underbrace{u_\alpha u_\beta u^\alpha u^\beta}_{=(u^2)^2 = 1} + \left[-\frac{1}{2}p + \frac{1}{2}\rho + \frac{\Lambda}{\kappa}\right] \underbrace{g_{\alpha\beta} u^\alpha u^\beta}_{=-1} \quad (25)$$

$$= (\rho + p) + \left[\frac{1}{2}p - \frac{1}{2}\rho - \frac{\Lambda}{\kappa}\right] \quad (26)$$

$$= \frac{1}{2} \left(3p + \rho + 2 \frac{\Lambda}{\kappa}\right) \quad (27)$$

Holds if $3p + \rho + 2 \frac{\Lambda}{\kappa} \geq 0$

- *DEC - Dominant energy condition* $-T^\alpha_\beta u^\beta$ is a future directed timelike or null vector for all future directed timelike vectors u^α
For $\Lambda = 0$ holds if $\rho \geq |p|$.

3. Safety requirements for travelers in Morris-Thorne wormholes

The general safety requirements (derived from what we know about Schwarzschild/Kerr BHs) are (independent of specific solution/shape of the wormhole)

- Small tidal forces - to avoid squeezing/stretching
- Small acceleration - to avoid trauma
- No even horizon I - to avoid signal loss between brain and rest of the body - or weird quantum effects on atoms (nucleus inside EH - electrons outside)
- No even horizon II - to be able to go back and forth (no sure if this counts as safety requirement)
- Solution must be stable with respect to external perturbations - wormhole should not collapse when traveler passes through
- Short transition time to other side - to avoid starvation/suffocation
- Large enough passage - so whole body/ship fits through
- Low radiation levels (Doppler shift of surrounding particles)

I just found the paper - but haven't read it yet.