EG3029 Chemical Thermodynamics

Solution Thermodynamics

Fundamental Property Relation

Total Gibbs energy for closed system

$$d(nG) = (nV)dP - (nS)dT$$

General case of a single-phase, open system

$$d(nG) = \left[\frac{\partial(nG)}{\partial P}\right]_{T,n} dP + \left[\frac{\partial(nG)}{\partial T}\right]_{P,n} dT + \sum_{i} \left[\frac{\partial(nG)}{\partial n_{i}}\right]_{P,T,n_{j}} dn_{i}$$

Chemical potential of species i

$$\mu_i = \left[\frac{\partial (nG)}{\partial n_i}\right]_{P,T,n_j}$$

Fundamental Property Relation

• Fundamental property relation for single-phase fluid systems of variable mass and composition: $\frac{d(nG) = (nV)dP - (nS)dT + \sum_{i} \mu_{i} dn_{i}}{d(nG) = (nV)dP - (nS)dT + \sum_{i} \mu_{i} dn_{i}}$

• For one mole of solution:

$$V = \left(\frac{\partial G}{\partial P}\right)_{T,x}$$

$$S = -\left(\frac{\partial G}{\partial T}\right)_{P,x}$$

Enthalpy

$$H = G + TS = G - T \left(\frac{\partial G}{\partial T} \right)_{P,x}$$

Chemical Potential and Phase Equilibria

• Closed system of two phases (α and β) in equilibrium:

$$d(nG)^{\beta} = (nV)^{\beta} dP - (nS)^{\beta} dT + \sum_{i} \mu_{i}^{\beta} dn_{i}^{\beta}$$

 $d(nG)^{\alpha} = (nV)^{\alpha} dP - (nS)^{\alpha} dT + \sum \mu_i^{\alpha} dn_i^{\alpha}$

• For system exhibiting π phases and N species:

$$\mu_i^{\alpha} = \mu_i^{\beta} = \dots = \mu_i^{\pi}$$

$$i = 1, 2, \dots, N$$

Thus, multiple phases at the same T and P are in equilibrium when the chemical potential of each species is the same in all phases.

Partial Properties

General partial molar property:

$$\overline{M}_{i} = \left[\frac{\partial(nM)}{\partial n_{i}}\right]_{P,T,n_{j}}$$

It is a response function i.e. a measure of the response of the total property *nM* to the addition at constant *T* and *P* of a differential amount of species *i* to a finite amount of solution.

Partial Properties

Equations Relating Molar and Partial Molar Properties

For the total property M:

$$M = \sum_{i} x_{i} \overline{M}_{i}$$

$$nM = \sum_{i} n_{i} \overline{M}_{i}$$

General expression for dM:

$$dM = \sum_{i} x_{i} d\overline{M}_{i} + \sum_{i} \overline{M}_{i} dx_{i}$$

Gibbs/Duhem equation

$$\left(\frac{\partial M}{\partial P}\right)_{T,x} dP + \left(\frac{\partial M}{\partial T}\right)_{P,x} dT - \sum_{i} x_{i} d\overline{M}_{i} = 0$$

and at constant *T* and *P*:

$$\sum_{i} x_{i} d\overline{M}_{i} = 0$$

Partial Properties Partial Properties in Binary Solutions

 Partial properties are readily calculated directly from an expression for the solution property as a function of composition at constant T and P:

$$\overline{M}_1 = M + x_2 \frac{dM}{dx_1}$$

$$\overline{M}_2 = M - x_1 \frac{dM}{dx_1}$$

• In derivative forms:

$$x_1 \frac{d\overline{M}_1}{dx_1} + x_2 \frac{d\overline{M}_2}{dx_1} = 0$$

$$\frac{d\overline{M}_1}{dx_1} = \frac{x_2}{x_1} \frac{d\overline{M}_2}{dx_1}$$

Ideal-Gas Mixture Model

- Useful model because it
 - Has a molecular basis
 - Approximates reality in well-defined limit of zero pressure
 - Is analytically simple
- Partial pressure:

$$p_i = \frac{y_i RT}{V^{ig}} = y_i P \qquad (i = 1, 2, \dots, N)$$

Gibbs's theorem:

A partial molar property (other than volume) of a constituent species in an ideal-gas mixture is equal to the corresponding molar property of the species as a pure ideal gas at the mixture temperature but at a pressure equal to its partial pressure in the mixture

Ideal-Gas Mixture Model

Properties of ideal-gas mixtures

$$H^{ig} = \sum_{i} y_i H_i^{ig}$$

$$S^{ig} = \sum_{i} y_i S_i^{ig} - R \sum_{i} y_i \ln y_i$$

$$G^{ig} = \sum_{i} y_{i} G_{i}^{ig} + RT \sum_{i} y_{i} \ln y_{i}$$

Ideal gas at constant T:

$$dG_i^{ig} = V_i^{ig}dP = \frac{RT}{P}dP = RT d \ln P$$

and after integration:

$$\mu_i^{ig} = \overline{G}_i^{ig} = \Gamma_i(T) + RT d \ln(y_i P)$$

$$G^{ig} = \sum_{i} y_{i} \Gamma_{i}(T) + RT \sum_{i} y_{i} \ln(y_{i}P)$$

Fugacity and Fugacity Coefficient Pure Species

- Chemical potential μ provides fundamental criterion for phase equilibria (also for reaction equilibria), but the determination of its absolute value is not possible
- Fugacity f

$$G_i = \Gamma_i(T) + RT \ln f_i$$

$$f_i^{ig} = P$$

$$G_i - G_i^{ig} = G_i^R = RT \ln \frac{f_i}{P} = RT \ln \phi_i$$

 $\phi_i = \frac{f_i}{P}$

Residual Gibbs energy

Fugacity coefficient

Fugacity and Fugacity Coefficient VLE of Pure Species

Saturated vapour and saturated liquid in equilibrium:

For a pure species coexisting liquid and vapour phases are in equilibrium when they have the same temperature, pressure and fugacity.

Fugacity of pure liquid

$$f_i^l(P) = \frac{f_i^v(P_i^{sat})}{P_i^{sat}} \frac{f_i^l(P_i^{sat})}{f_i^v(P_i^{sat})} \frac{f_i^l(P)}{f_i^l(P_i^{sat})} P_i^{sat}$$

$$f_i = \phi_i^{sat} P_i^{sat} \exp \frac{V_i^l (P - P_i^{sat})}{RT}$$

 V_i^I is assumed constant

Fugacity and Fugacity Coefficient **Species in Solution**

 For species i in a mixture of real gases or in a solution of liquids:

$$\left(\mu_i = \Gamma_i(T) + RT \ln \hat{f}_i\right)$$

Phases in equilibrium:

$$\left(\hat{f}_i^{\alpha} = \hat{f}_i^{\beta} = \dots = \hat{f}_i^{\pi}\right) \qquad i = 1, 2, \dots, N$$

Thus, multiple phases at the same T and P are in equilibrium when the fugacity of each constituent species is the same in all phases.

• Fugacity coefficient for species in solution: $\hat{\phi}_i = \frac{\hat{f}_i}{v_i P}$

$$\left[\hat{\phi_i} = \frac{\hat{f_i}}{y_i P}\right]$$

Ideal-Solution Model

A solution is ideal when:

$$\mu_i^{id} = \overline{G}_i^{id} = G_i(T, P) + RT \ln x_i$$

Total values of properties:

$$G^{id} = \sum_{i} x_i G_i + RT \sum_{i} x_i \ln x_i$$

$$V^{id} = \sum_{i} x_i V_i$$

$$S^{id} = \sum_{i} x_i S_i - R \sum_{i} x_i \ln x_i$$

$$H^{id} = \sum_{i} x_i H_i$$

Lewis/Randall rule

$$\mu_i = \Gamma_i(T) + RT \ln \hat{f}_i$$

$$G_i = \Gamma_i(T) + RT \ln f_i$$



$$\hat{f}_i^{id} = x_i f_i$$

Applies to each species in an ideal solution at all conditions of *P*, *T*, and composition

• If M represents the molar (or unit-mass) value of any extensive thermodynamic property (e.g. V, U, H, S, G, ...), then an excess property M^E is defined: $M^E = M - M^{id}$ Mactual value Mid value in ideal solution

Fundamental excess property relation

$$d\left(\frac{nG^{E}}{RT}\right) = \frac{nV^{E}}{RT}dP - \frac{nH^{E}}{RT^{2}}dT + \sum_{i} \frac{\overline{G}_{i}^{E}}{RT}dn_{i}$$

M in Relation to G	M^R in Relation to G^R	M^E in Relation to G^E
$V = (\partial G/\partial P)_{T,x} (11.4)$	$V^R = (\partial G^R / \partial P)_{T,x}$	$V^E = (\partial G^E / \partial P)_{T,x}$
$S = -(\partial G/\partial T)_{P,x} $ (11.5)	$S^R = -(\partial G^R/\partial T)_{P,x}$	$S^E = -(\partial G^E/\partial T)_{P,x}$
$H = G + TS$ $= G - T(\partial G/\partial T)_{P,x}$ $= -RT^{2} \left[\frac{\partial (G/RT)}{\partial T} \right]_{P,x}$	$H^{R} = G^{R} + TS^{R}$ $= G^{R} - T(\partial G^{R}/\partial T)_{P,x}$ $= -RT^{2} \left[\frac{\partial (G^{R}/RT)}{\partial T} \right]_{P,x}$	$H^{E} = G^{E} + TS^{E}$ $= G^{E} - T(\partial G^{E}/\partial T)_{P,x}$ $= -RT^{2} \left[\frac{\partial (G^{E}/RT)}{\partial T} \right]_{P,x}$
$C_P = (\partial H/\partial T)_{P,x}$ $= -T(\partial^2 G/\partial T^2)_{P,x}$	$C_P^R = (\partial H^R / \partial T)_{P,x}$ $= -T(\partial^2 G^R / \partial T^2)_{P,x}$	$C_P^E = (\partial H^E / \partial T)_{P,x}$ $= -T(\partial^2 G^E / \partial T^2)_{P,x}$

Relations with activity coefficient:

$$\left(\gamma_i = \frac{\hat{f}_i}{x_i f_i} \right) \qquad \left(\overline{G}_i^E = RT \ln \gamma_i \right)$$

$$\overline{G}_i^E = RT \ln \gamma_i$$

Effects of pressure and temperature

$$\left(\frac{\partial \ln \gamma_i}{\partial P}\right)_{T,x} = \frac{\overline{V_i}^E}{RT}$$

$$\left(\frac{\partial \ln \gamma_i}{\partial T}\right)_{P,x} = -\frac{\overline{H}_i^E}{RT^2}$$

Summability and Gibbs/Duhem equations

$$\frac{G^E}{RT} = \sum_{i} x_i \ln \gamma_i$$

$$\sum_{i} x_{i} d \ln \gamma_{i} = 0$$

const T, P

- Excess properties can be determined
 - G^E from VLE data
 - $-H^{E}$ from mixing experiments

$$-S^{E} \text{ from } S^{E} = \frac{H^{E} - G^{E}}{T}$$

All excess properties become zero as either species approaches purity

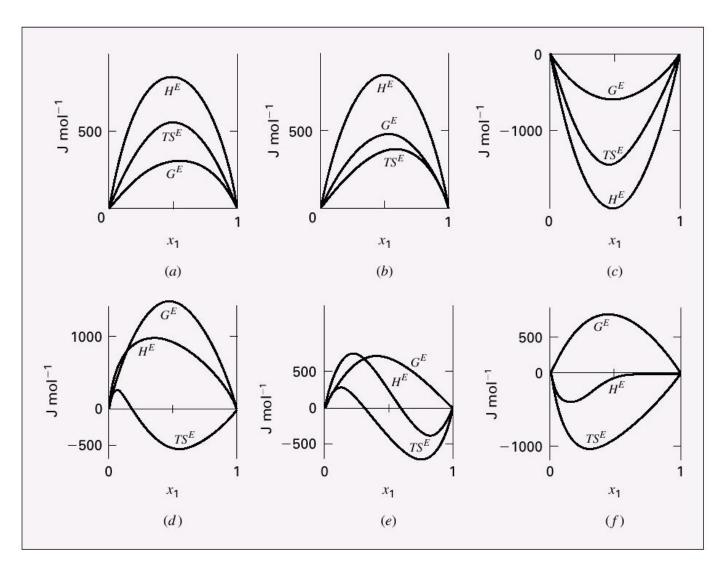
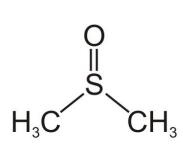
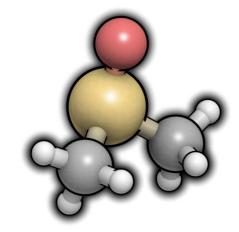


Figure 11.4: Excess properties at 50°C for six binary liquid systems: (a) chloroform(1)/n-heptane(2); $(b)\ acetone(1)/methanol(2); (c)\ acetone(1)/chloroform(2); (d)\ ethanol(1)/n-heptane(2); (e)\ ethanol(1)/chloroform(2); (f)\ ethanol(1)/water(2).$

DMSO:

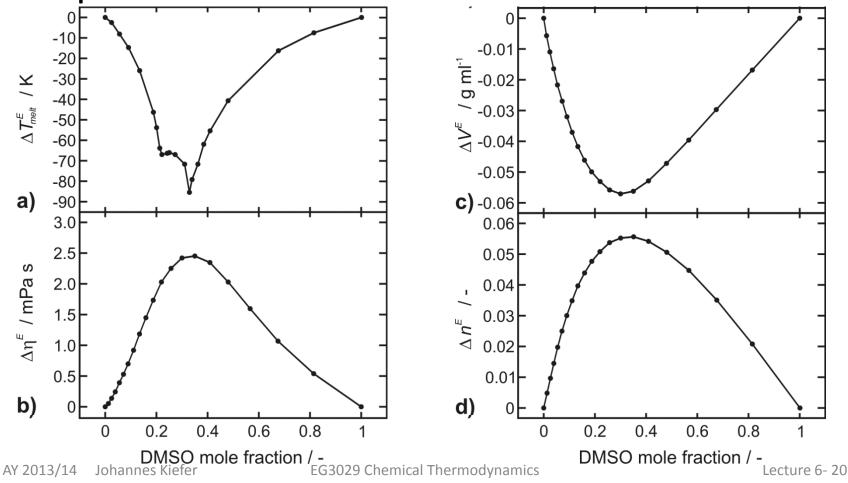
- Polyfunctional molecule
- Properties:
 - Melting point: 18 degC
 - Density: 1.1 g/cm³
 - Boiling point: 189 degC
 - Refractive index: 1.479
 - Viscosity: 1.996 cP
- Applications:
 - Solvent
 - Antifreeze
 - ...



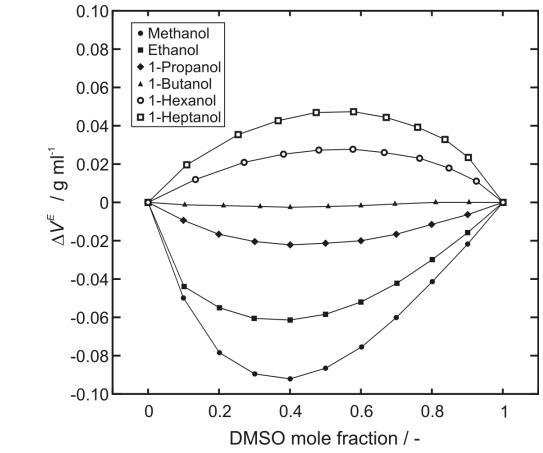


J. Kiefer, K. Noack, B. Kirchner Hydrogen Bonding in Mixtures of Dimethyl Sulfoxide and Cosolvents *Current Physical Chemistry* **1** (2011) 340-351.

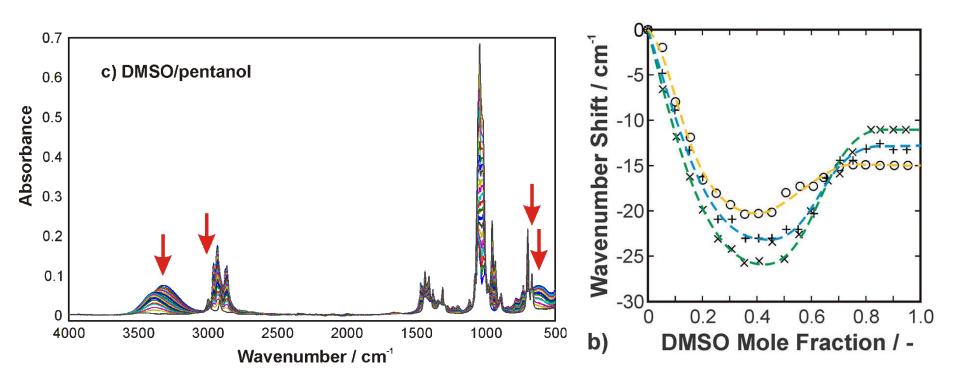
Aqueous DMSO solutions



Alcoholic DMSO solutions



Alcoholic DMSO solutions



Property Changes of Mixing

 Excess properties approach zero for ideal solutions, but thermodynamic properties might still change upon mixing

$$\Delta G^{id} = RT \sum_{i} x_{i} \ln x_{i}$$

$$\Delta S^{id} = -R \sum_{i} x_{i} \ln x_{i}$$

$$\Delta V^{id} = 0$$

$$\Delta H^{id} = 0$$

Worked example: mixing process for a binary solution

- Alternative to equations of state
- Typically more accurate when strong intermolecular interactions are present
- Commonly used:
 - Margules equations
 - Van Laar equations
 - Wilson equations
 - Non-Random-Two-Liquids (NRTL)
 - UNIversal QUAsi Chemical (UNIQUAC)

$$\gamma_i = \frac{\hat{f}_i}{x_i f_i} = \frac{\hat{f}_i}{\hat{f}_i^{id}}$$

- 2 parameter models:
 - Margules equations for binary systems

$$\ln \gamma_1 = x_2^2 \cdot \left[A_{12} + 2(A_{21} - A_{12}) x_1 \right]$$

$$\ln \gamma_2 = x_1^2 \cdot \left[A_{21} + 2(A_{12} - A_{21}) x_2 \right]$$

Van Laar equations

$$\ln \gamma_1 = A'_{12} \cdot \left(1 + \frac{A'_{12}x_1}{A'_{21}x_2}\right)^{-2} \qquad ; \qquad \ln \gamma_2 = A'_{21} \cdot \left(1 + \frac{A'_{21}x_1}{A'_{12}x_2}\right)^{-2}$$

- 2 parameter models:
 - Wilson equations

$$\frac{G^{E}}{RT} = -x_{1} \ln(x_{1} + x_{2}\Lambda_{12}) - x_{2} \ln(x_{2} + x_{1}\Lambda_{21})$$

$$\ln \gamma_{1} = -\ln(x_{1} + x_{2}\Lambda_{12}) + x_{2} \left(\frac{\Lambda_{12}}{x_{1} + x_{2}\Lambda_{12}} - \frac{\Lambda_{21}}{x_{2} + x_{1}\Lambda_{21}}\right)$$

$$\ln \gamma_{2} = -\ln(x_{2} + x_{1}\Lambda_{21}) + x_{1} \left(\frac{\Lambda_{12}}{x_{1} + x_{2}\Lambda_{12}} - \frac{\Lambda_{21}}{x_{2} + x_{1}\Lambda_{21}}\right)$$

- 3 parameter models:
 - Non-Random-Two-Liquids (NRTL)

$$\frac{G^{E}}{x_{1}x_{2}RT} = \frac{G_{21}\tau_{21}}{x_{1} + x_{2}G_{21}} + \frac{G_{12}\tau_{12}}{x_{2} + x_{1}G_{12}}$$

$$\ln \gamma_{1} = x_{2}^{2} \left[\tau_{21} \left(\frac{G_{21}}{x_{1} + x_{2}G_{21}} \right)^{2} + \frac{G_{12}\tau_{12}}{\left(x_{2} + x_{1}G_{12} \right)^{2}} \right]$$

$$\ln \gamma_{2} = x_{1}^{2} \left[\tau_{12} \left(\frac{G_{12}}{x_{2} + x_{1}G_{12}} \right)^{2} + \frac{G_{21}\tau_{21}}{\left(x_{1} + x_{2}G_{21} \right)^{2}} \right]$$

UNIversal QUAsi Chemical (UNIQUAC): greater complexity

Worked example

A liquid binary mixture is in equilibrium with its vapour at 144°C. The liquid mole fraction of component 1 is $x_1 = 0.6$. In addition, the following information is given:

- $\ln \gamma_1 = Ax_2^2$ and $\ln \gamma_2 = Ax_1^2$
- at 144°C, the saturation pressures are P_1^{sat} = 75.20 kPa and P_2^{sat} = 31.66 kPa
- the system forms an azeotrope at 144°C for which $x_1^{az} = y_1^{az} = 0.294$
- the modified Raoult's law is valid
- a) Determine the equilibrium pressure of the system.

[9 marks]

b) Determine the composition of the vapour.

[2 marks]

c) Draw a schematic *P-xy* diagram showing the azeotropic behaviour. Indicate how the *P-x* relationship would look like for Raoult's law.

[6 marks]

d) Which is the main problem that is associated with azeotropic mixtures in technical processes? How can it be overcome (give one example)?

[3 marks]