

# Criticality & information transfer in neural networks

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# What is Criticality?

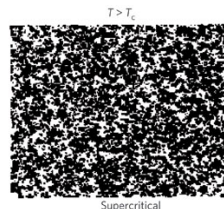
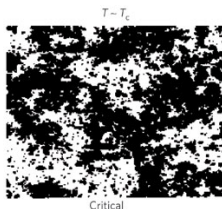
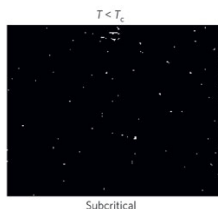
- Anything with a power law.

# What is Criticality?

- But really...

# What is Criticality?

- Near a critical phase transition, as we vary a parameter of the system the correlation lengths diverge in space and the dynamics slow. The spatial correlations take the form of power law distributions.
- The canonical example



# The case for Criticality in Neuroscience

From Chialvo, Emergent Complex Dynamics, Nature Physics 6 2010:

- Brain has necessary hardware to show criticality
- Healthy brains lack a preferred time scale
- fMRI degree distributions are power-law
- Neuronal avalanches have power law size distributions
- All known models that display complex emergent behavior display criticality

These arguments all make the case for a resemblance but lack any reason why an organism would want a critical nervous system.

# Hypothesis

The dynamics of systems near a critical phase transition enhance information propagation through the system and naturally enact certain encoding strategies.

To this end we (are) examining simple systems that display a critical transition:

Does the Ising model naturally perform compressive sensing but only near the phase transition?

Is dynamical synchronization to a measured signal optimized near criticality?

# Neural network: dynamical model

A model of  $N$  interconnected “neurons” that obey the dynamics:

$$\frac{dV_i}{dt} = -V_i + \sum_j J_{ij} S_j$$

where  $S_i$  is the “activity” of a neuron, defined as

$$S_i \equiv \tanh(g \tilde{J} V_i).$$

Actually it might as well be any sigmoidal function of  $V_i$ . The sigmoidal behavior limits the coupling to avoid runaway unstable solutions (also biophysically motivated, like a saturation of the interneuronal coupling).

## Random neural network: dynamical model

The randomness lies in  $J_{ij}$ . Each entry is a Gaussian-distributed random variable with the statistics

$$\begin{aligned}\langle J_{ij} \rangle &= 0 \\ \langle J_{ij} J_{kl} \rangle &= \frac{\tilde{J}^2}{N} \delta_{ik} \delta_{jl}.\end{aligned}$$

In general, the dynamics of this kind of system are non-relaxational (i.e. the system does not converge to a global minimum of an energy function) since the couplings are asymmetric.

In the large- $N$  limit, the dynamics can be reduced to the form

$$\frac{dV_i}{dt} = -V_i + \eta_i$$

and there exist a broad range of chaotic solutions.



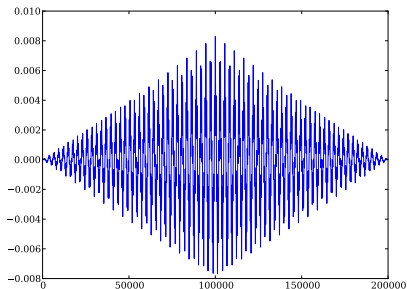
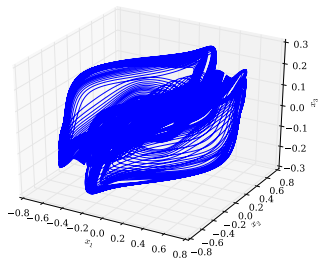
# Random neural network: dynamical model

With these definitions, the quantity  $g\tilde{J}$  acts like a control parameter for the dynamics of the quantity

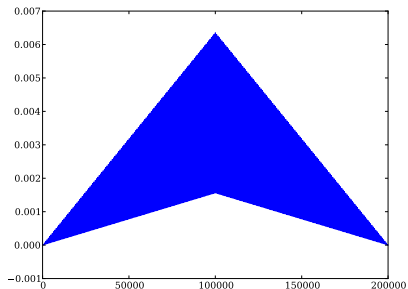
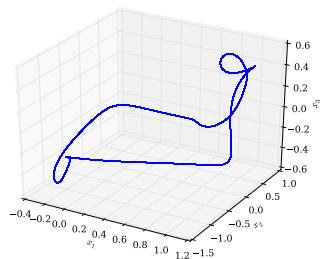
$$\Delta(t) \equiv \langle V_i(t_0 + t)V_i(t_0) \rangle$$

Roughly, when  $g\tilde{J}$  crosses above 1, the dynamics undergo a qualitative change in behavior that looks much like the phase transition in the Ising model;  $\Delta$  acts as a qualifier for the behavior.

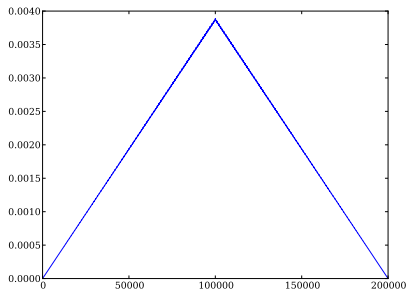
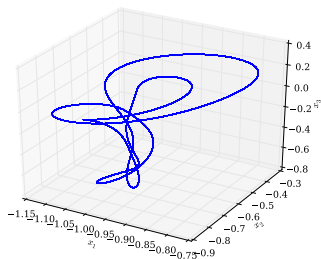
$$g\tilde{J} = 1.1$$



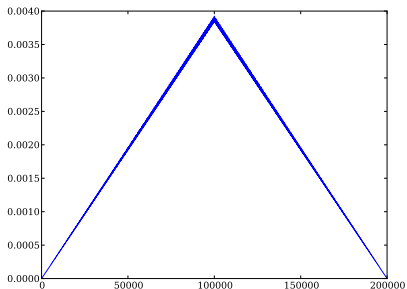
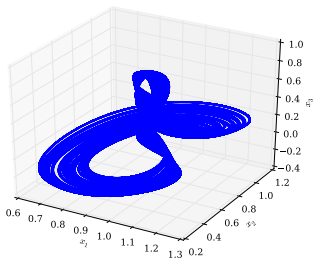
$$g\tilde{J} = 1.2$$



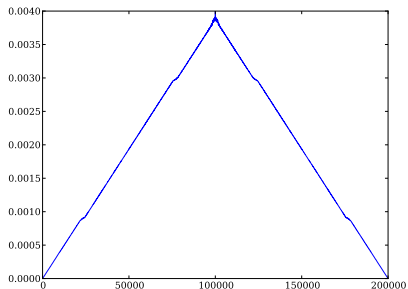
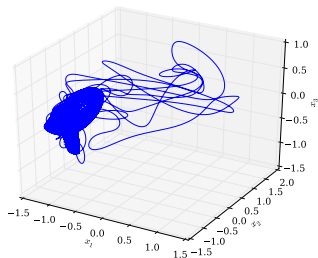
$$g\tilde{J} = 1.26$$



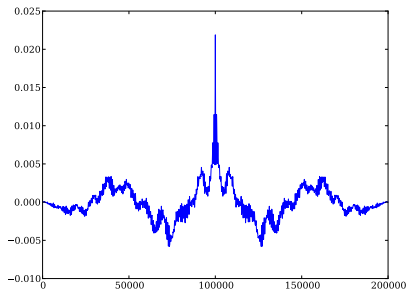
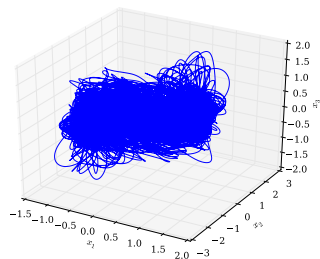
$$g\tilde{J} = 1.266$$



$$g\tilde{J} = 1.2668$$



$$g\tilde{J} = 1.27$$



$$g\tilde{J} = 1.3$$

