

Last time we took the schrodinger equation and transformed it to the radial equation in center of mass frame.

$$\left( = \frac{\hbar^2}{m} \frac{1}{r^2} \partial_r r^2 \partial_r + \frac{1}{4} m \omega^2 r^2 + \frac{ke^2}{r} \right) + \left( -\frac{\hbar^2}{4m} \frac{1}{R^2} \partial_R R^2 \partial_R + m \omega^2 R^2 + \frac{\hbar l(l+1)}{R^2} \right) \psi(r) \theta(R) = (E_r + E_R) \psi(r) \theta(R) \quad (1)$$

we are only interested in the

$$\left( = \frac{\hbar^2}{m} \frac{1}{r^2} \partial_r r^2 \partial_r + \frac{1}{4} m \omega^2 r^2 + \frac{ke^2}{r} \right) \psi(r) = E_r \psi(r) \quad (2)$$

we can make the equation dimensionless which gives

$$\left( -\partial_\rho^2 + \rho^2 + \frac{\beta}{\rho} \right) u(\rho) = \lambda u(\rho) \quad \left( -\partial_\rho^2 + V(\rho) \right) u(\rho) = \lambda u(\rho) \quad (3)$$

Discretize

$$u(\rho) \rightarrow u(\rho_i) = u_i \quad \rho \rightarrow \rho_i = \rho_0 + i \hbar h = \frac{b-a}{n+1} a = 0 \quad (4)$$

the linear algebra equation is

$$\hat{A} = \begin{pmatrix} \frac{2}{h^2} + V_1 & -\frac{1}{h^2} & & & \\ -\frac{1}{h^2} & \frac{2}{h^2} + V_2 & -\frac{1}{h^2} & & \\ & & \ddots & & \\ & & & -\frac{1}{h^2} & \frac{2}{h^2} + V_{N-1} \end{pmatrix} \quad (5)$$

$$\hat{A}x = \lambda x \quad (6)$$

$$S^T A x = \lambda S^T x \quad (7)$$

$$S^T A S (S^T x) = \lambda (S^T x) \quad (8)$$

Eigenvalues unchanged under orthogonal transforms

Eigenvectors change, but the length is preserved

Dot product between vectors is conserved