

Generalized Statistical Complexity Measure: a new tool for Dynamical Systems.

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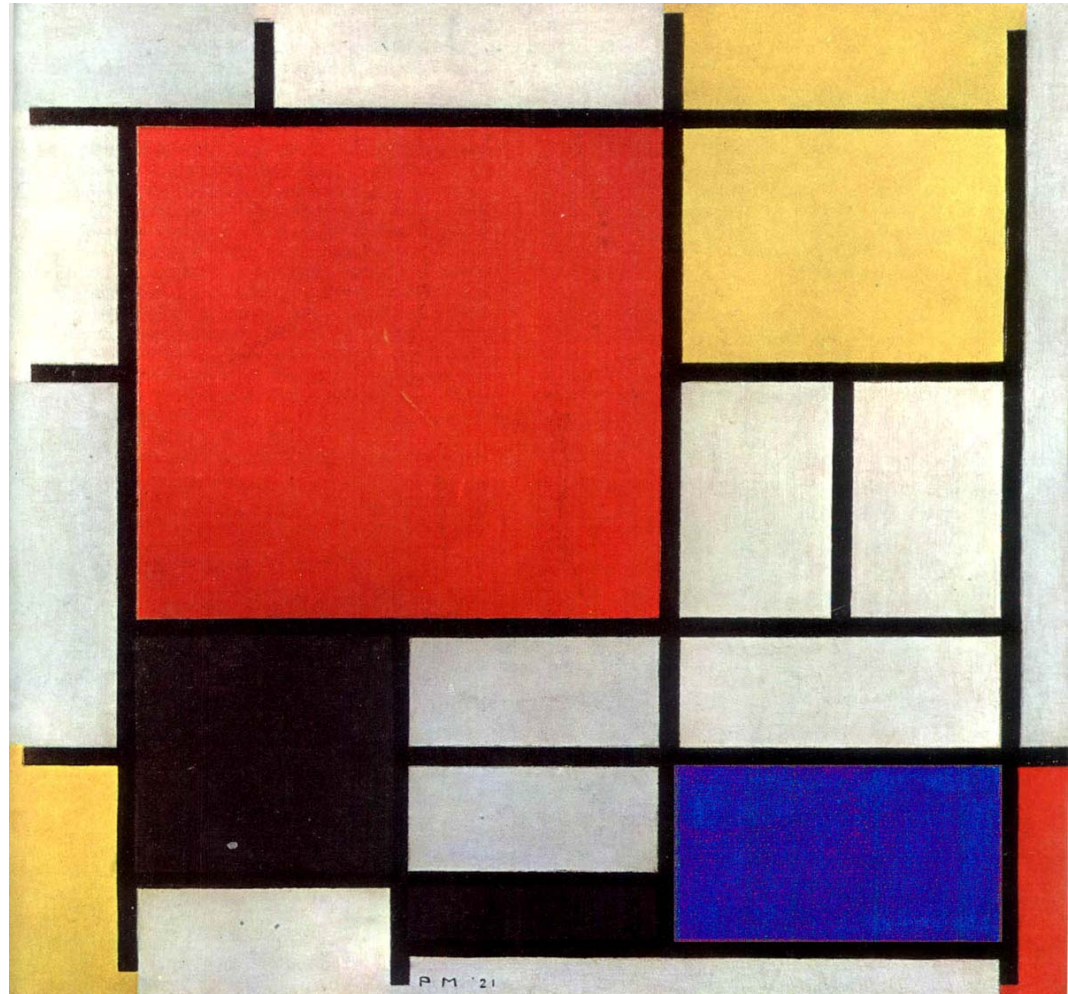
The Simple & The Complex



- **Complexity ?**

The Simple & The Complex

Mondrian

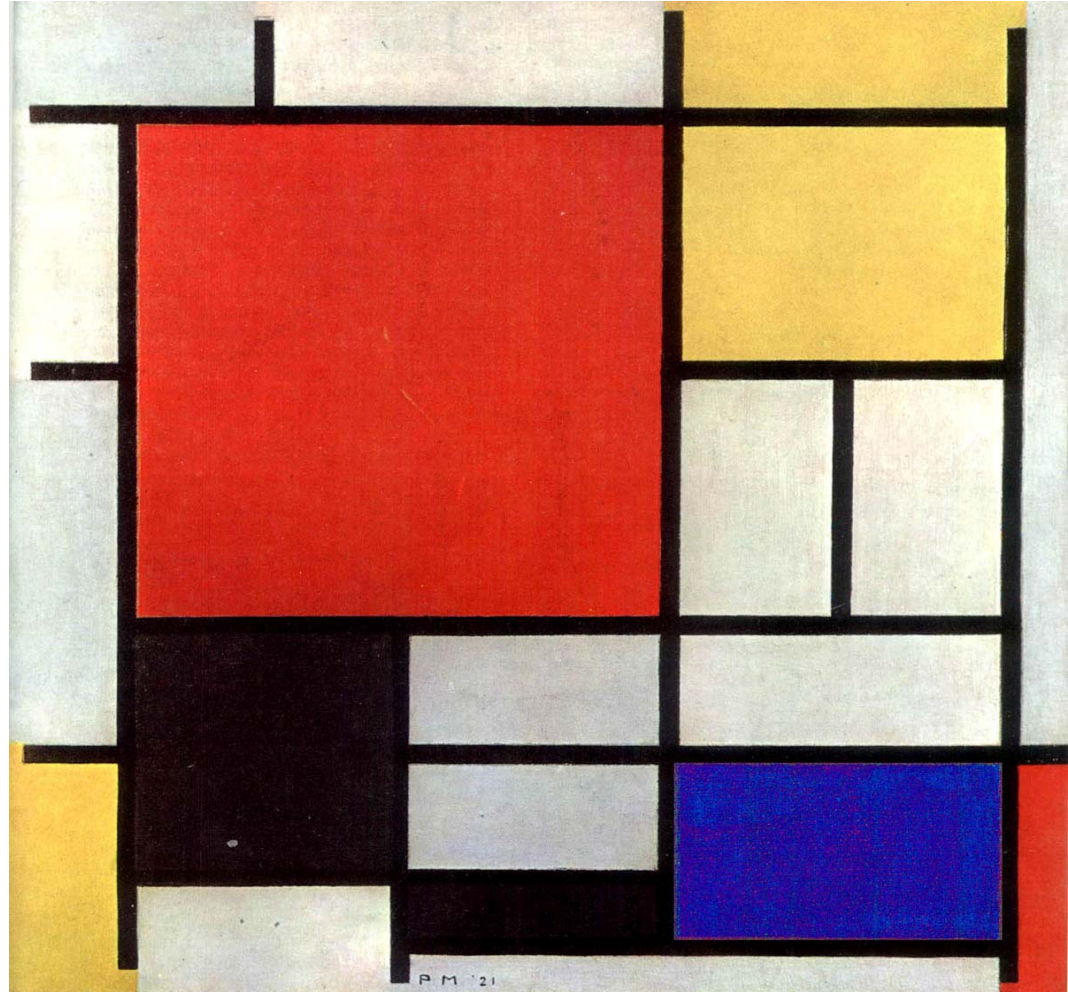


The Simple & The Complex

$$H = 0$$

$$C = 0$$

Mondrian



The Simple & The Complex

Pollock



The Simple & The Complex

H = 1

C = 0

Pollock



The Simple & The Complex

Bosch



The Simple & The Complex

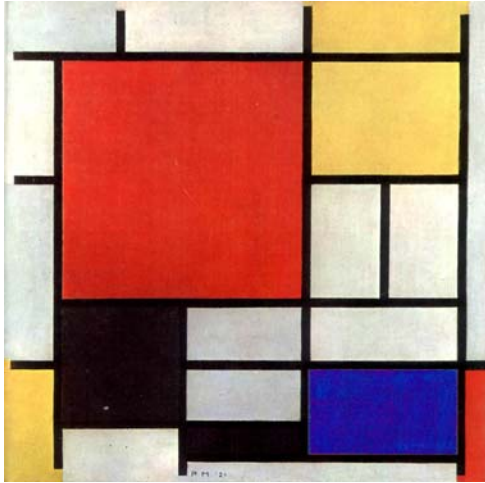
H \neq 0

C \neq 0

Bosch



The Simple & The Complex



$$H = 0$$

$$C = 0$$



$$H \neq 0$$

$$C \neq 0$$



$$H = 1$$

$$C = 0$$

Complexity

The COMPLEXITY has to do with intricate structures hidden in the dynamics, emerging from a system which itself is much simpler than its dynamics. Complexity is characterized by the paradoxical situation of complicated dynamics of simple systems.

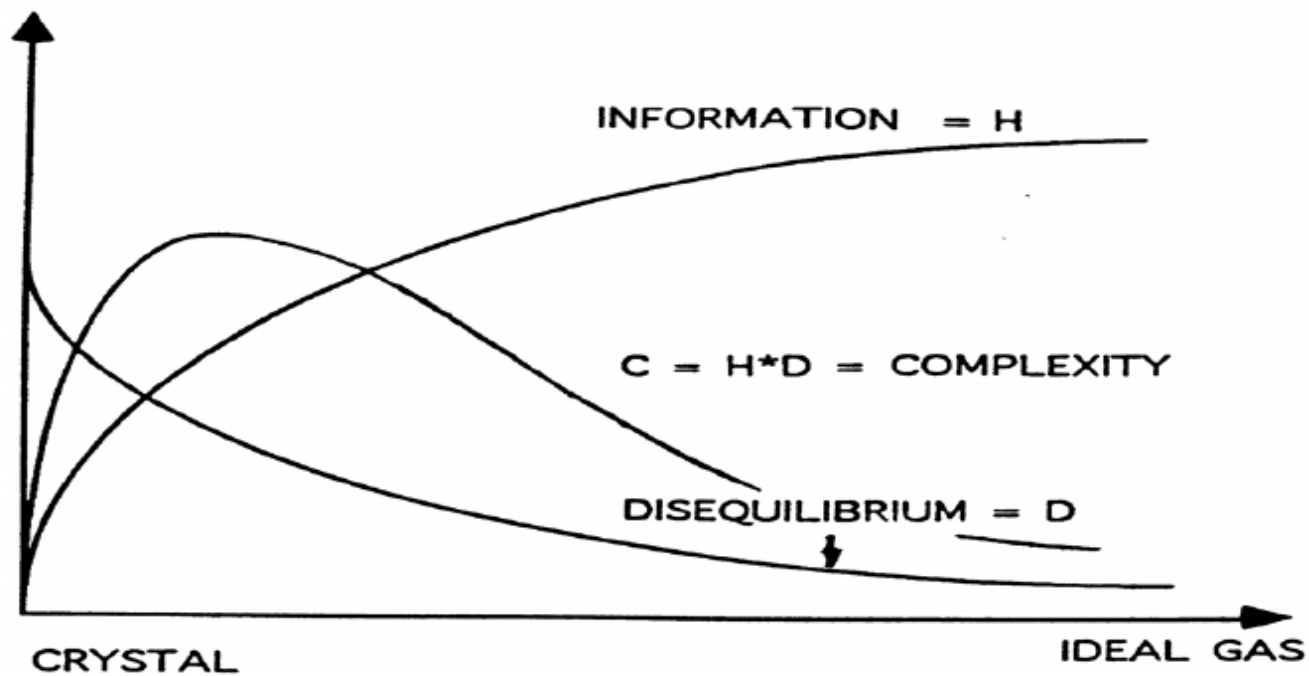
- Periodic motion it is not complex.
- White noise it is not complex.

Crystal & Ideal Gas

Crystal	Ideal Gas
<ul style="list-style-type: none">• High ordered system• Minimal information stored in the system• Probability distribution in phase space: $P_j = 1$ for $j = k$ $P_j = 0$ for $j \neq k$• Maximum Disequilibrium	<ul style="list-style-type: none">• Completely disordered system• Maximal information stored in the system.• Probability distribution in phase space: $P_j = 1 / N$ for $j = 1, \dots, N$• Minimum Disequilibrium

Statistical Complexity

$$\text{COMPLEXITY} \\ C = H \cdot Q$$



Disorder H

- We define for a given probability distribution

$$P = \{p_j, j = 1, \dots, N\} \in \Omega \subset \mathbb{R}^N$$

and its associate information measure $\mathcal{I}[P]$, an amount of “disorder” H in the fashion

$$H[P] = \mathcal{I}[P] / \mathcal{I}_{max} ,$$

where $\mathcal{I}_{max} = \mathcal{I}[P_e]$ and P_e is the probability distribution which maximize the information measure, where P_e is the equilibrium probability distribution. Then $0 \leq H \leq 1$.

Disequilibrium Q

- We define the “disequilibrium” adopting some kind of distance from the equilibrium distribution P_e of the accessible states of the system.

$$Q[P] = Q_0 \mathcal{D}[P, P_e] ,$$

where Q_0 is a normalization constant and $0 \leq Q \leq 1$.

The disequilibrium Q would reflect on the systems’s “architecture”, being different from zero if there are “privileged”, or more likely states among the accessible ones.

Selection of the information measure I

- **Shannon Entropy:**

$$S_S [P] = - \sum_i p_i \cdot \ln [p_i] .$$

- **Tsallis Entropy:**

$$S_T^{(q)} [P] = \frac{1}{q-1} \left[1 - \sum_i (p_i)^q \right] .$$

- **Escort-Tsallis Entropy:**

$$S_G^{(q)} [P] = \frac{1}{q-1} \left[1 - \left\{ \sum_i (p_i)^{1/q} \right\}^{-q} \right] .$$

- **Rényi entropy:**

$$S_R^{(q)} [P] = \frac{1}{(1-q)} \ln \left\{ \sum_{j=1}^N (p_j)^q \right\} .$$

Selection of Distance D

- **Euclidean distance:**

$$\mathcal{D}_E[P, P_e] = \|P - P_e\|_E = \sum_{j=1}^N \left\{ p_j - \frac{1}{N} \right\}^2 .$$

- **Wootters distance:**

$$\mathcal{D}_W[P_1, P_2] = \cos^{-1} \left\{ \sum_{j=1}^N \left(p_j^{(1)} \right)^{1/2} \cdot \left(p_j^{(2)} \right)^{1/2} \right\} .$$

- **Relative entropy (Kullback relative entropy):**

$$\mathcal{D}_{K_q^\kappa}[P, P_e] = K_q^{(\kappa)}[P|P_e] = S_q^{(\kappa)}[P_e] - S_q^{(\kappa)}[P] .$$

Where $\kappa = S, T, G, R$ – Shannon, Tsallis, Generalized Escort-Tsallis, Rényi, entropic functional forms.

- **Jensen divergence:**

$$\begin{aligned} \mathcal{D}_{J_q^\kappa}[P, P_e] &= \mathcal{J}_{S_q^\kappa}^{1/2}[P, P_e] = \\ &= \frac{1}{2} K_q^{(\kappa)} \left[P \mid \frac{P + P_e}{2} \right] + \frac{1}{2} K_q^{(\kappa)} \left[P_e \mid \frac{P + P_e}{2} \right] . \end{aligned}$$

Where $\kappa = S, T, G, R$ – Shannon, Tsallis, Generalized Escort-Tsallis, Rényi, entropic functional forms.

Generalized Statistical Complexity Measures

The family of *Statistical Complexity Measures*, $C_{\nu,q}^{(\kappa)}$, is defined by

$$C_{\nu,q}^{(\kappa)}[P] = H_q^{(\kappa)}[P] \cdot Q_q^{(\nu)}[P]$$

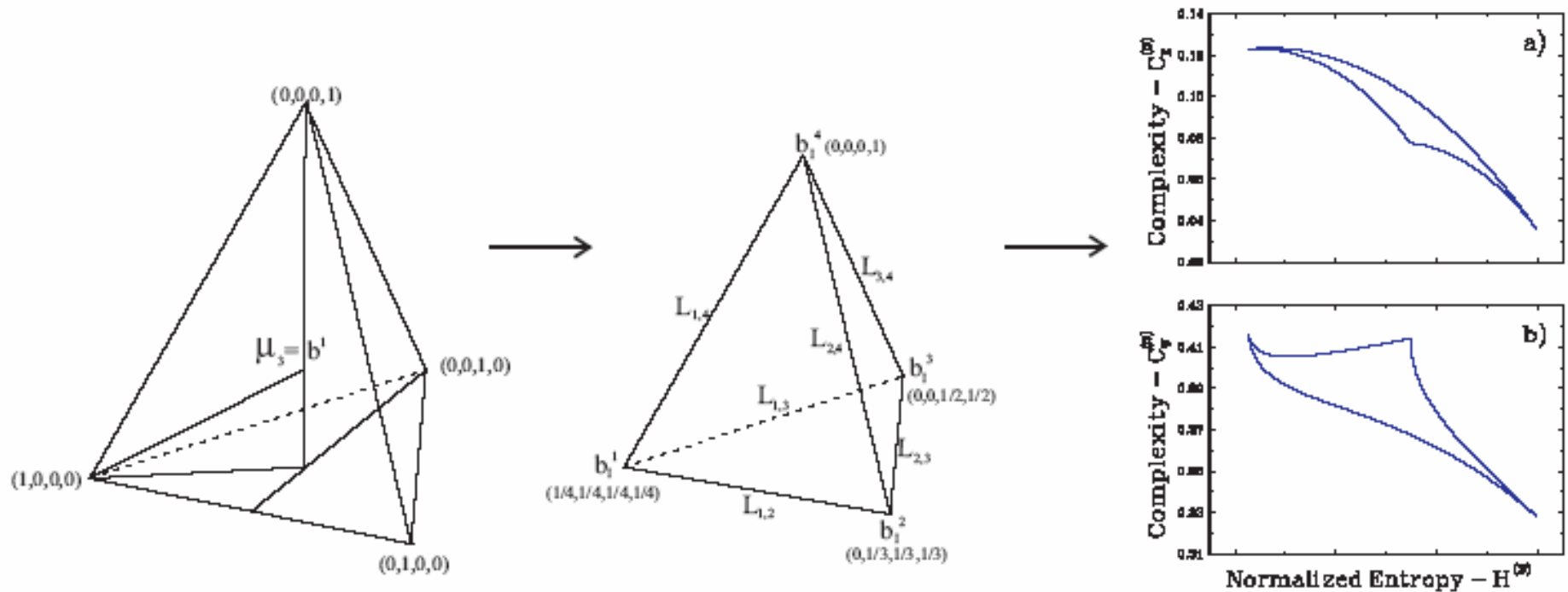
This quantity reflects on the interplay between the amount of information stored in the system and its disequilibrium.

- $\kappa = S, T, G, R$: Shannon, Tsallis, Generalized Escort-Tsallis, Rényi, for a fixed q .

In Shannons instance ($\kappa = S$) we have, of course, $q = 1$.

- $\nu = E, W, K, J$: Euclidean, Wootters, Kullback, Jensen.

Maximum and Minimum of Generalized Statistical Complexity Measures



- (i) Probability subspace Ω for $N=4$: $\Omega \equiv \Delta^3$ (3-simplex) in an hyperplane of dimension 3. Dotted lines effect the barycentric subdivision with μ_3 the Ω -barycenter.
- (ii) Sub-simplex Δ_I^3 .
- (iii) Maximum and minimum of complexity as function of H obtained by consecutive borders of the sub-simplex Δ_I^3 .

Application to Logistic Map

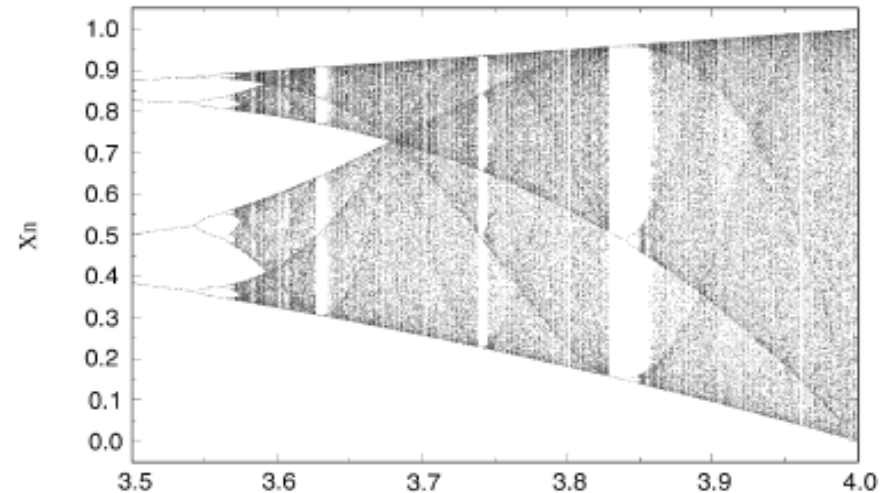
The Logistic Map, $F : x_n \rightarrow x_{n+1}$ is described by the ecologically motivated, dissipative system described by the first order difference equation

$$x_{n+1} = r \cdot x_n \cdot (1 - x_n)$$

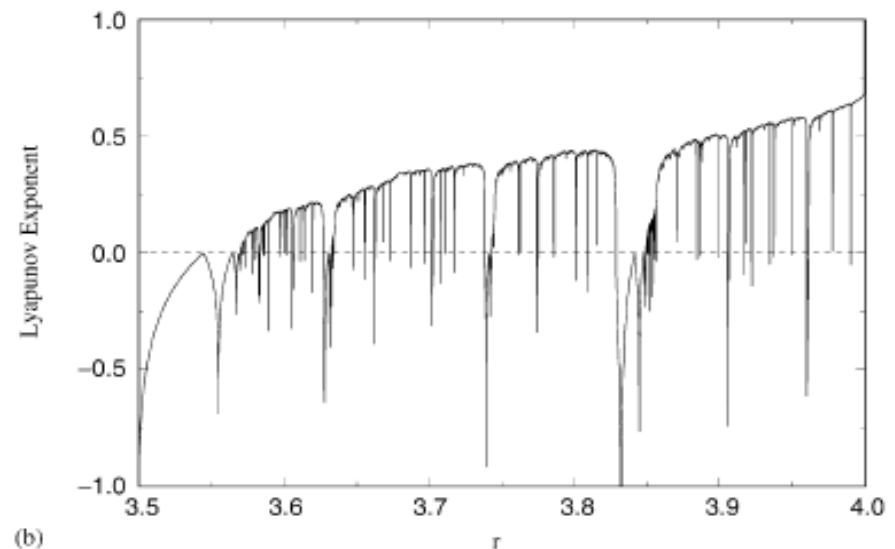
with $0 \leq x_n \leq 1$ and $0 < r \leq 4$.

Binary treatment (symbolic dynamics) of the logistic map:

For each parameter value, r , the dynamics of the logistic map was reduced to a binary sequence (0 if $x \leq \frac{1}{2}$; 1 if $x > \frac{1}{2}$) and binary strings of length 12 were considered as states of the system. The concomitant probabilities are assigned according to the frequency of occurrence after running over at least 2^{22} iterations.

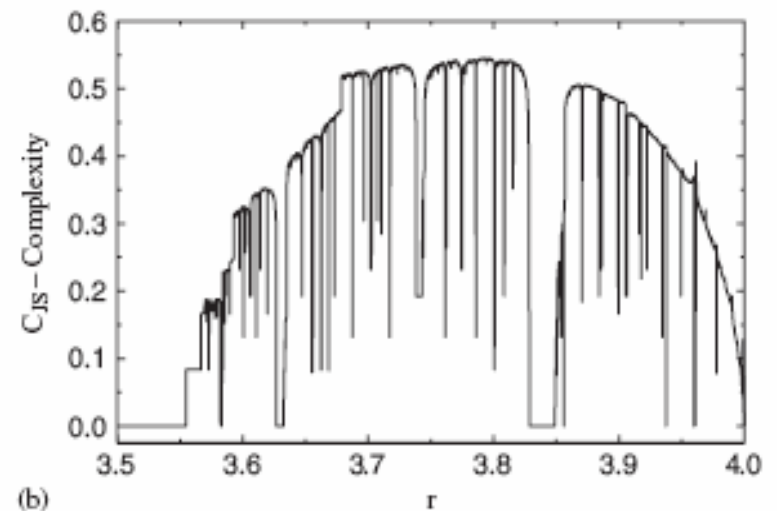
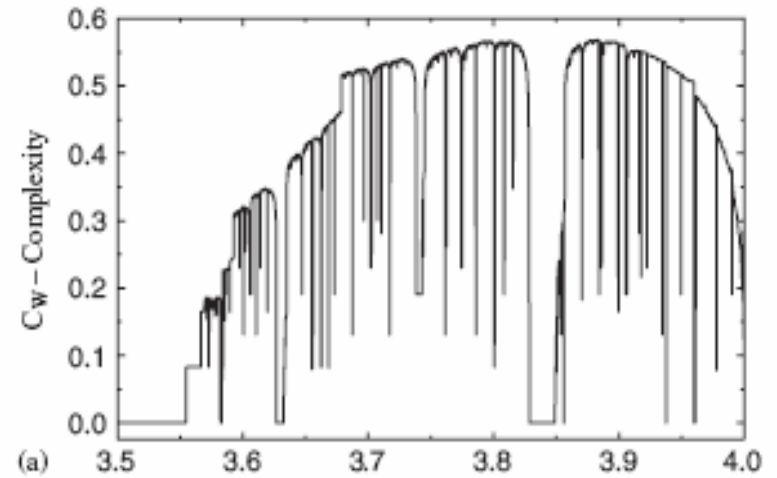
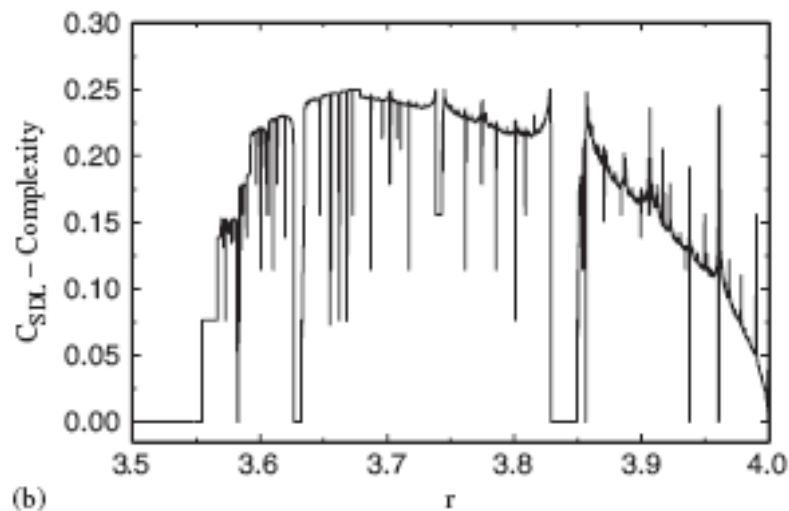
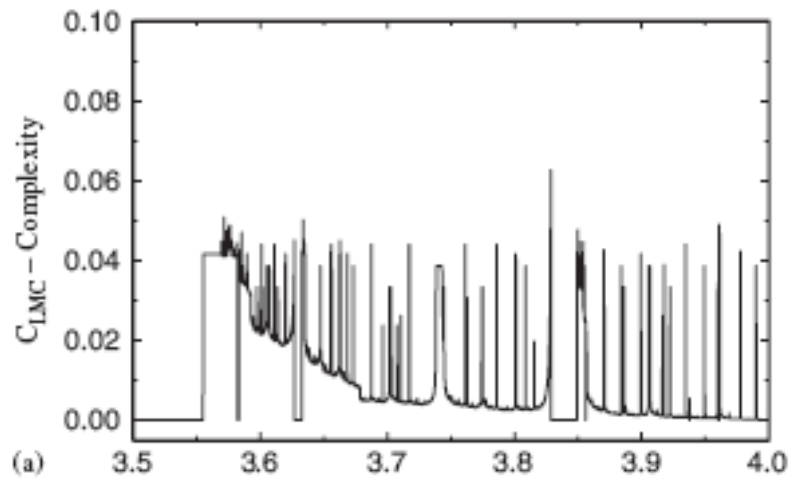


(a)



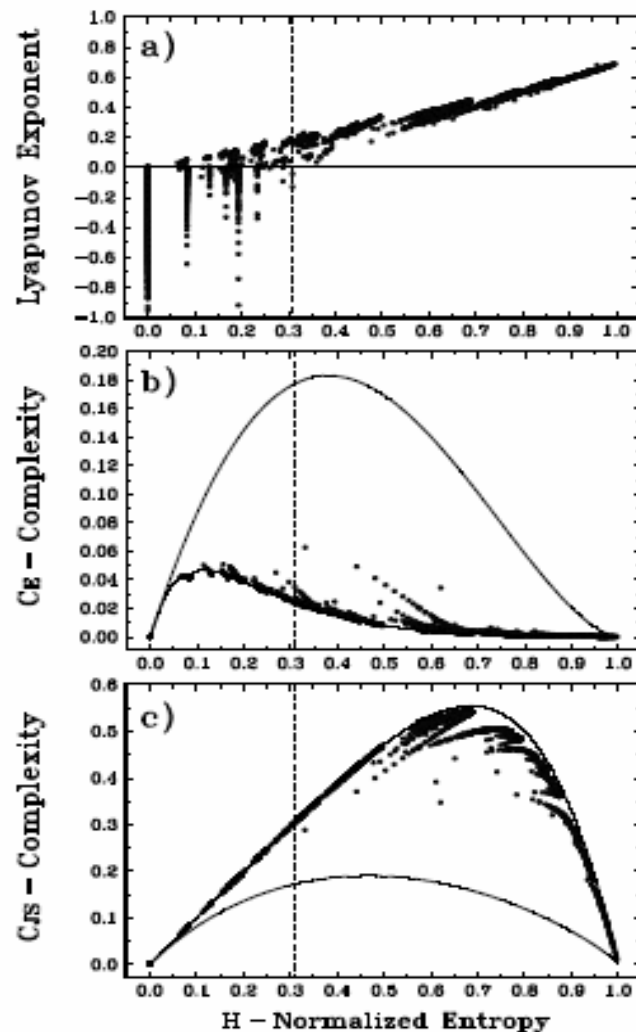
(b)

Application to Logistic Map



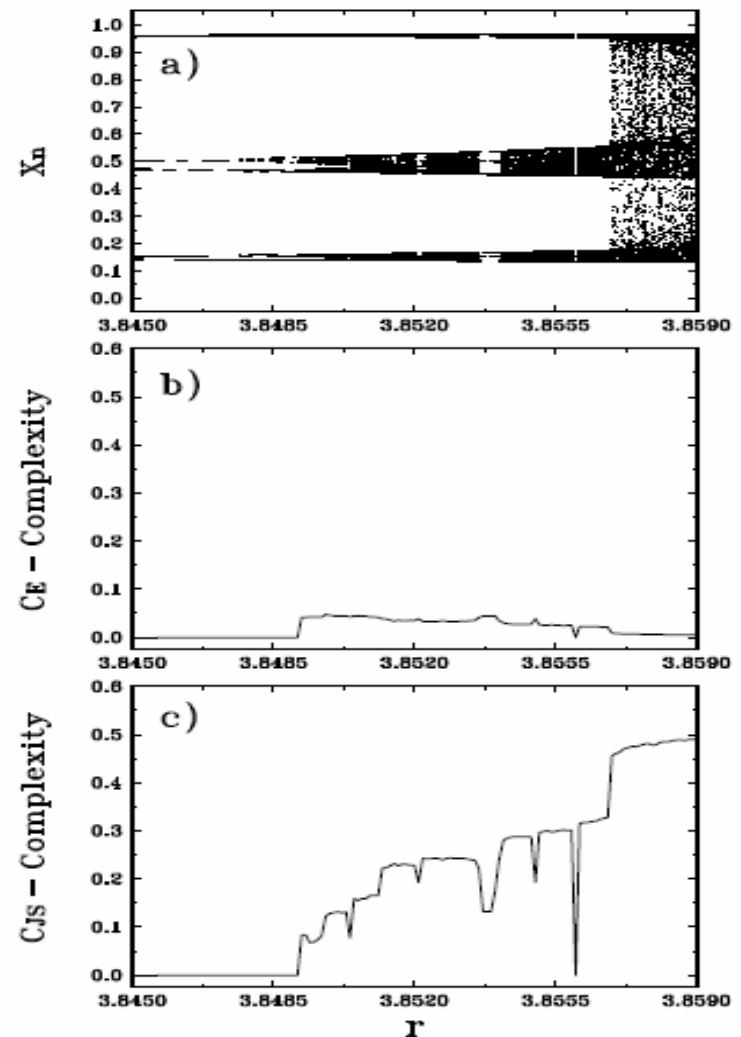
Application to Logistic Map

Notice that, for the case of periodic windows, if $H < \mathcal{H} \approx 0.3$, we can ascertain that $\Lambda < 0$, while for $H > \mathcal{H}$ we see that $\Lambda > 0$, which entails chaotic behavior. The LMC statistical complexity is larger for periodic than for chaotic motion, which is wrong!. The Jensen-Shannon statistical complexity measure, C_{JS} , on the other hand, behaves in opposite manner, and is also different for distinct degrees of periodicity.

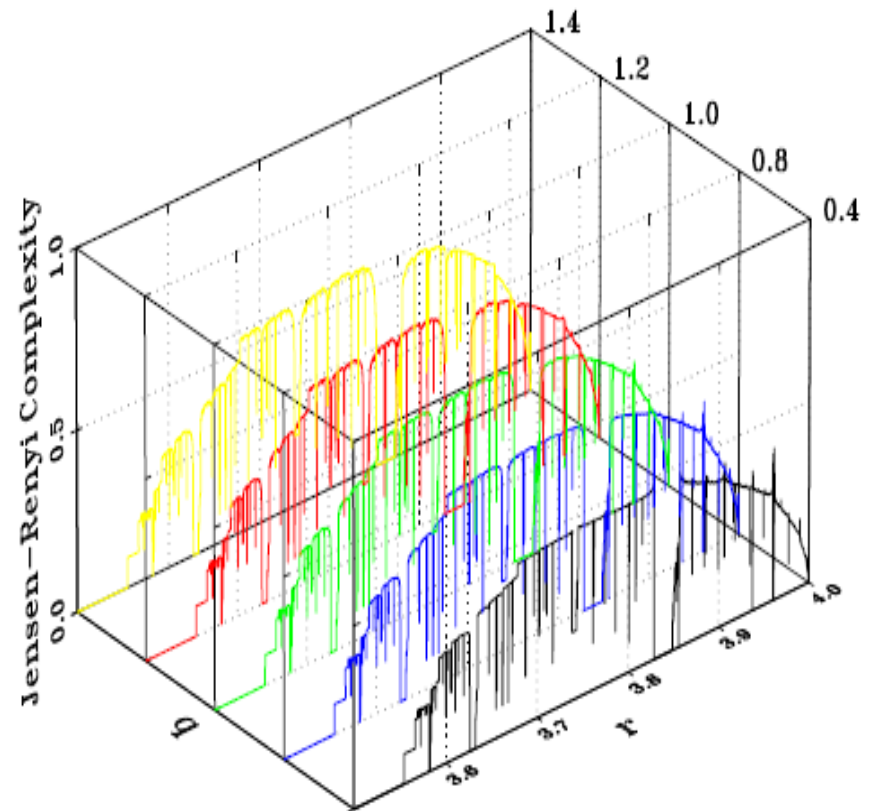
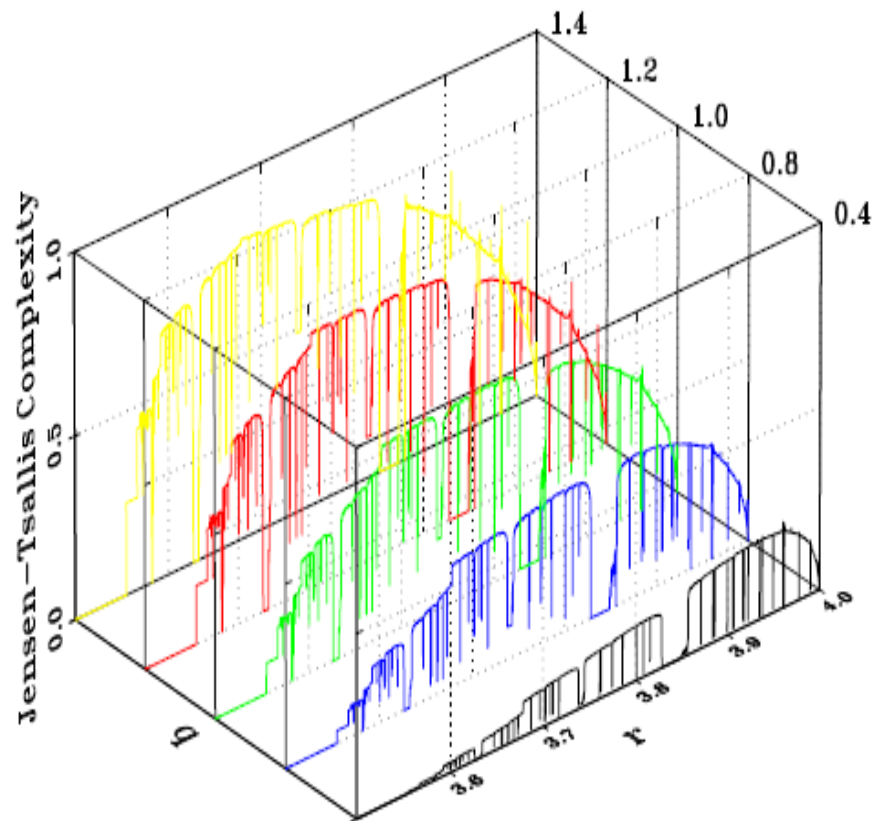


Application to Logistic Map

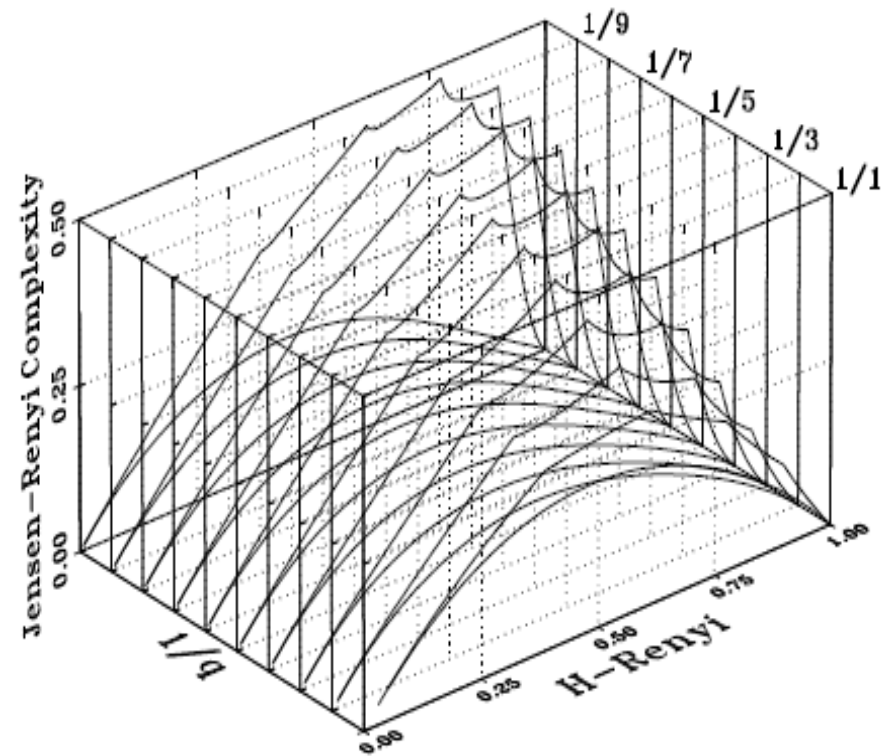
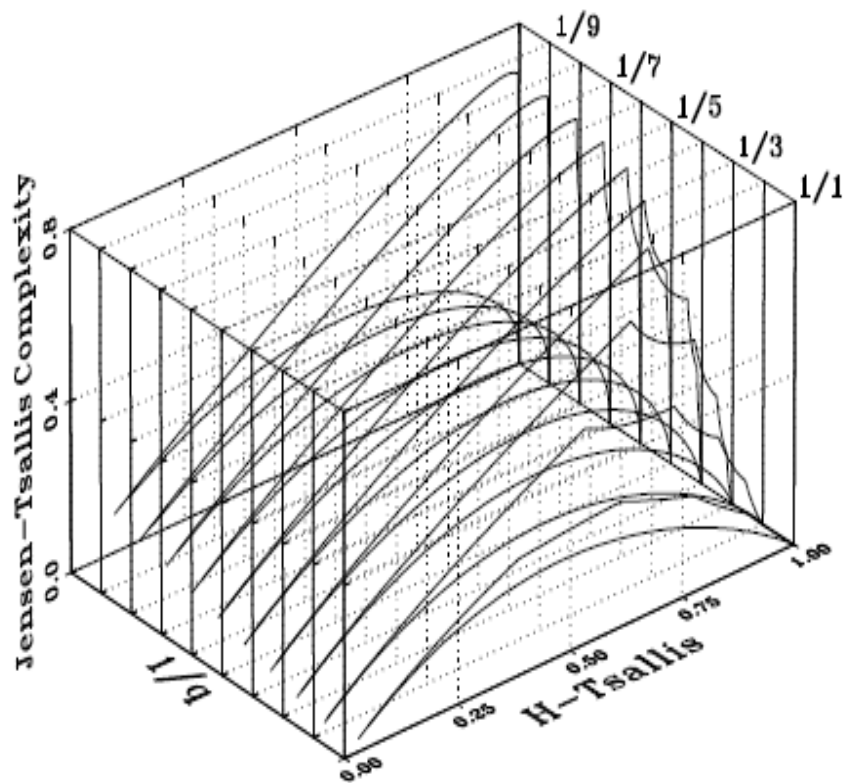
Summing up: the Jensen-Shannon statistical complexity measure *i)* becomes intensive, *ii)* is able to distinguish among distinct degrees of periodicity, and *iii)* yields a better description of dynamical features (a better grasp of dynamical details).



Application to Logistic Map



Application to Logistic Map



$$N = 6$$

Probability distribution P

Given a time series

$$X = \{x_j, j = 1, \dots, N\} \in \mathbb{R}^N$$

we can define the associate probability distribution function based on

- Histogram of amplitudes.
- Binary representation.
- Frequency (Fourier Transform).
- Frequency bands (Wavelet Transform).
- Ordinal Patterns (Attractor representation).

Probability distribution

Band-Pompe Methodology:

Given the time-series $\{x_t, t = 1, \dots, T\}$ and an embedding dimension $d > 1$, we are interested in *ordinal patterns* of order d generated by

$$(s) \mapsto (x_{s-(d-1)}, x_{s-(d-2)}, \dots, x_{s-1}, x_s)$$

which assign to each time s the d -dimensional vector of values at times $s, s-1, \dots, s-(d-1)$.

Clearly, the greater the d -value, the more information on the past our vectors are able to yield.

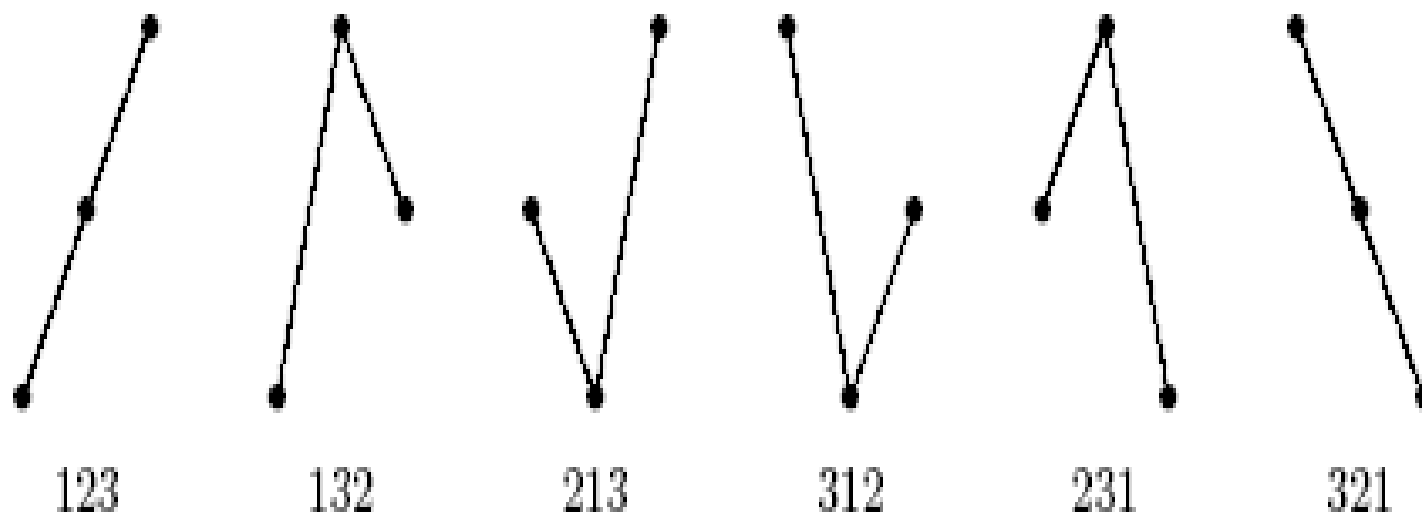
By the *ordinal pattern* related to the time (s) we mean the permutation $\pi = (r_0, r_1, \dots, r_{d-1})$ of $(0, 1, \dots, d-1)$ defined by

$$x_{s-r_{d-1}} \leq x_{s-r_{d-2}} \leq \dots \leq x_{s-r_1} \leq x_{s-r_0}$$

Note that the underlying probability distribution is “extracted” by appropriate consideration regarding causal effects in the system’s dynamics.

Probability distribution

Case $D_e = 3$, the number of patterns will be $D_e! = 3! = 6$.
Graphically one have:

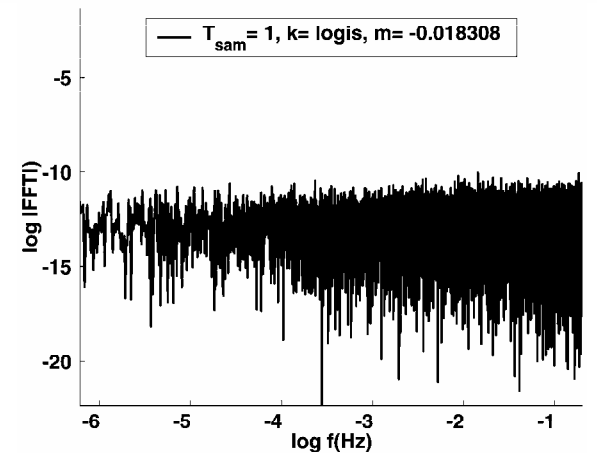
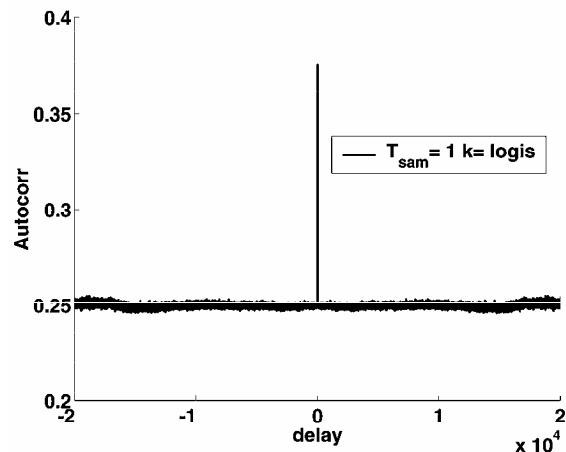
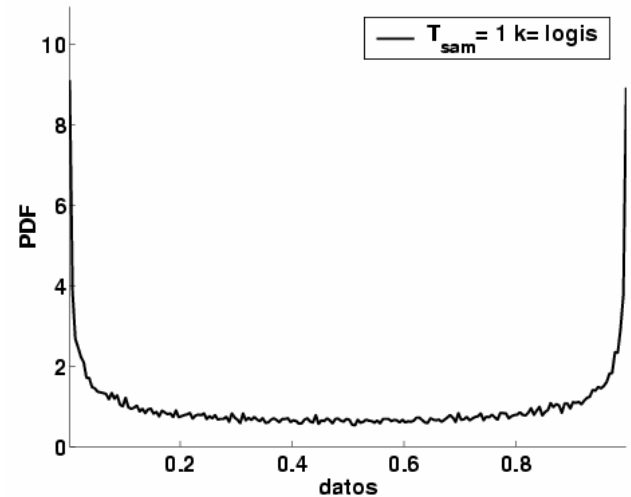
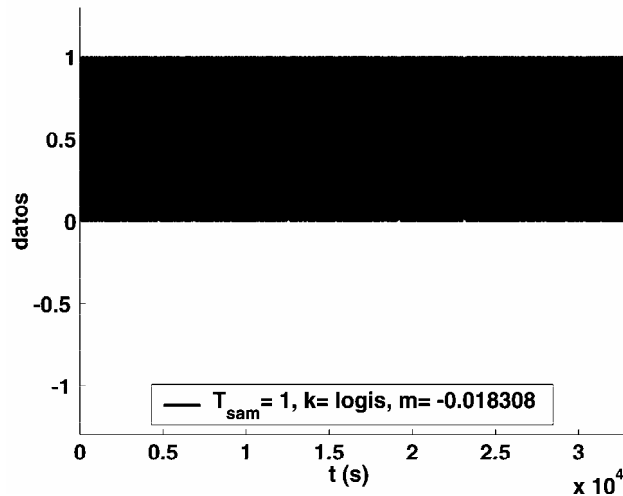


Chaos, Noise & Noise $1/f^k$

Although being of a quite different physical origin, time series arising from *chaotic systems* share with those generated by *stochastic processes* several properties that make them almost undistinguishable:

- a wide-band power spectrum,
- power spectrum of type $1/f^k$,
- a delta-like autocorrelation function,
- an irregular behavior of the measured signals, etc.

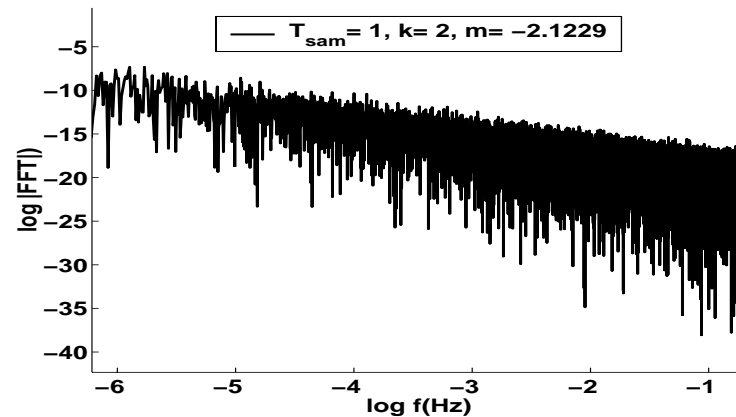
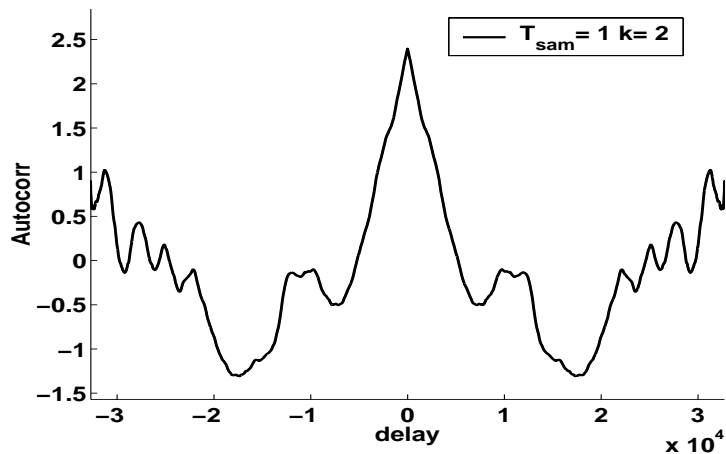
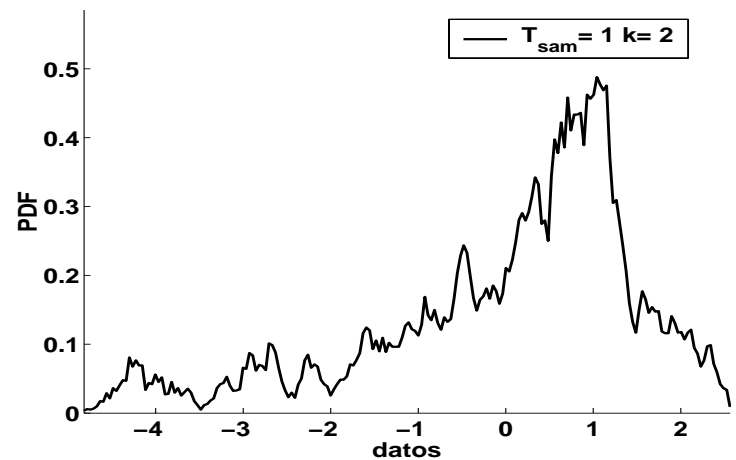
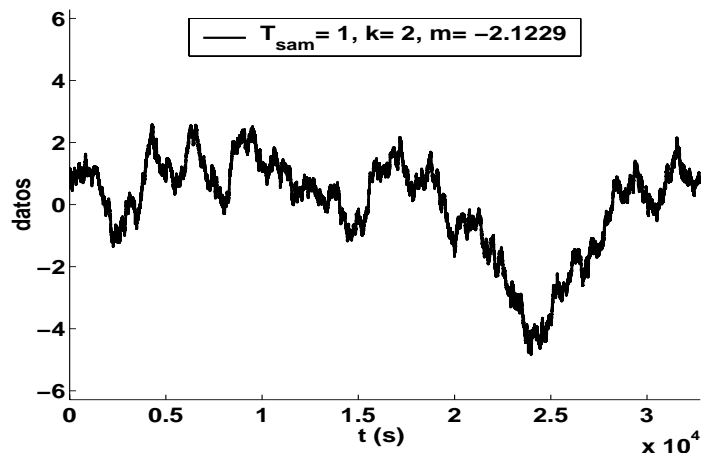
Chaos, Noise & Noise 1 / f κ



Logistic Map: $X(n+1) = R X(n) (1 - X(n))$; $R = 4$

$$H_S = 0.638781 \quad C_{JS}^{(MPR)} = 0.477845$$

Chaos, Noise & Noise $1/f^k$



Noise $1/f^k : k = 2$

$$H_S = 0.858 \quad C_{JS}^{(MPR)} = 0.232$$

Chaos, Noise & Noise 1 / f^k

Chaotic systems:

- The Logistic Map defined by:

$$x_{n+1} = r x_n (1 - x_n) .$$

Note that for $r = 4$ this map has a non uniform natural invariant probability density function (PDF).

- The Skew Tent Map: one has

$$\begin{cases} x/\omega & \text{for } x \in [0, \omega] \\ (1-x)/(1-\omega) & \text{for } x \in [\omega, 1] \end{cases} .$$

For any ω -value this map has a uniform natural invariant PDF ($\omega = 0.1847$ is here considered).

- Henon's Map: it is a 2D extension of the Logistic Map given by:

$$\begin{cases} x_{n+1} = 1 - a x_n^2 + y_n \\ y_{n+1} = b x_n \end{cases} .$$

The values used here, $a = 1.4$ and $b = 0.3$, correspond to a chaotic attractor with a non-smooth PDF.

Chaos, Noise & Noise 1 / f^k

- **The Lorenz Map of Rossler's oscillator:** for the 3D continuous Rossler oscillator one has

$$\begin{cases} \dot{x} = -y - z \\ \dot{y} = x + a y \\ \dot{z} = b + z(x - c) \end{cases}.$$

where $a = 0.2$, $b = 0.2$, and $c = 5.7$ correspond to a chaotic attractor. The Lorenz map is obtained by storing only x -minimal values.

- **Schuster Maps:** Schuster and coworkers introduced a class of maps which generate intermittent signals with chaotic bursts that also display $1/f^z$ noise

$$x_{n+1} = x_n + x_n^z, \quad \text{Mod } 1.$$

In particular, results for $z = 5/2$, 2 and $3/2$ are reported.

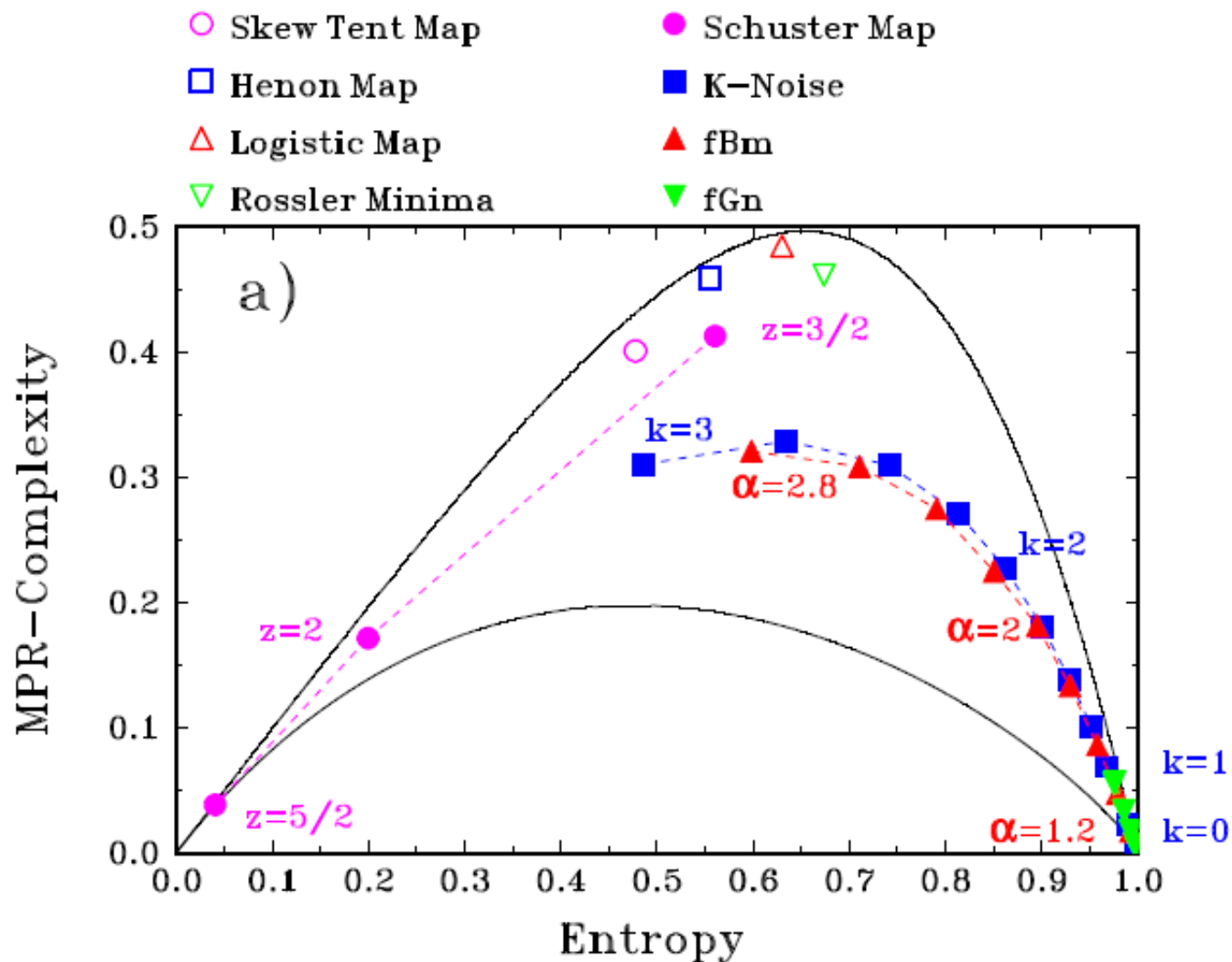
Chaos, Noise & Noise 1 / f^k

Stochastic processes :

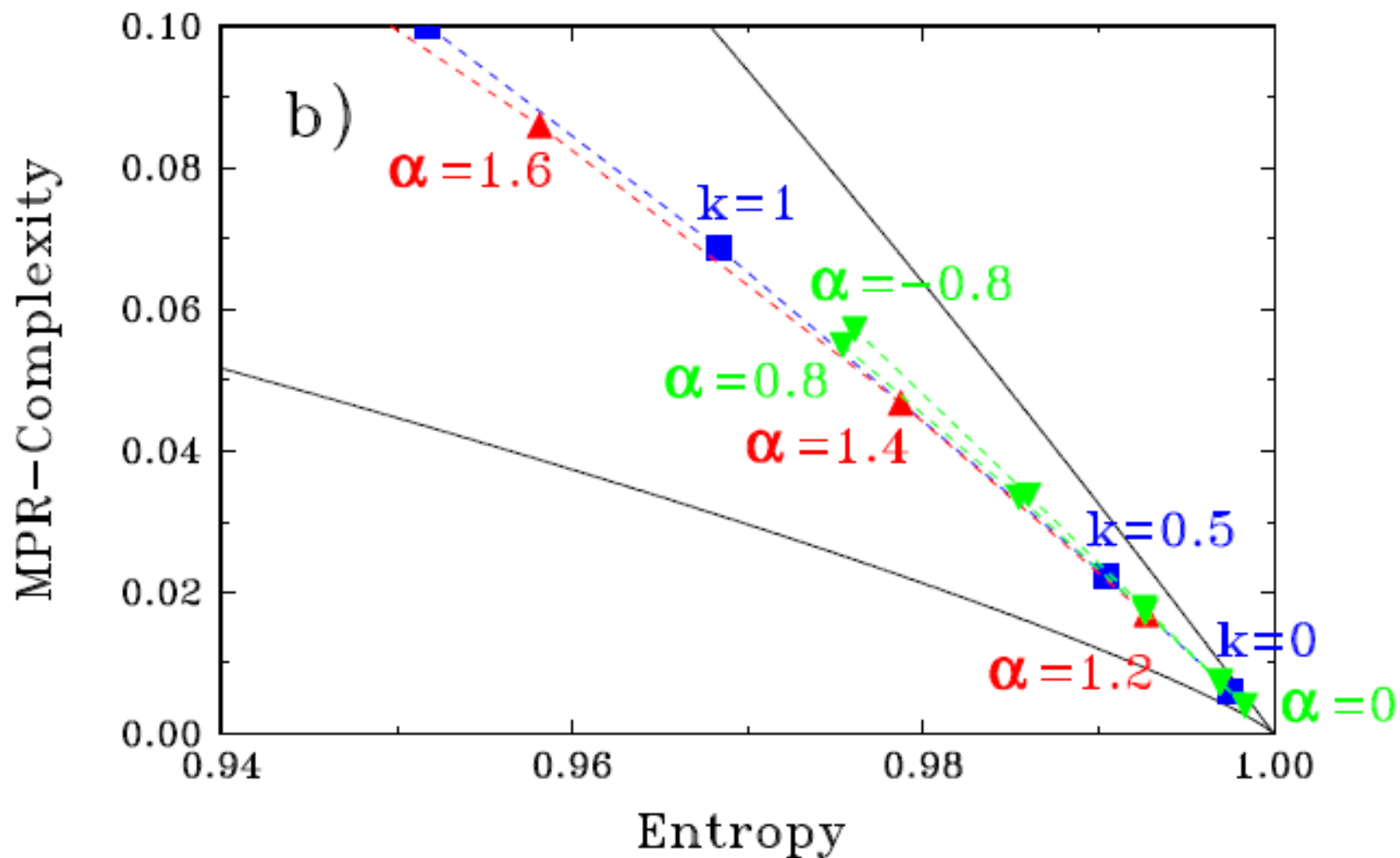
- Noises with f^{-k} Power Spectrum. The ensuing time series x_i has the desired PS and, by construction, is representative of non-Gaussian noises.
- Fractional Brownian motion (fBm) and fractional Gaussian noise (fGn): fBm is the only family of processes which is (a) Gaussian, (b) self-similar, and (c) endowed with stationary increments. The fBm and fGn are continuous but non-differentiable processes (in the classical sense). As a non-stationary process, they do not possess a spectrum defined in the usual sense; however, it is possible to define a *generalized power spectrum* of the form: $\Phi \propto |f|^{-\alpha}$, with $\alpha = 2\mathcal{H} + 1$, $1 < \alpha < 3$ for fBm and, $\alpha = 2\mathcal{H} - 1$, $-1 < \alpha < 1$, for fGn.

Hurst's \mathcal{H} parameter defines two distinct regions in the interval $(0, 1)$. When $\mathcal{H} > 1/2$, consecutive increments tend to have the same sign so that these processes are *persistent*. For $\mathcal{H} < 1/2$, on the other hand, consecutive increments are more likely to have opposite signs, and we say that they are *anti-persistent*.

Chaos, Noise & Noise $1 / f^k$



Chaos, Noise & Noise $1 / f^k$



Other Applications

- Characterization of laser propagation through turbulent media.
- Encryption test of pseudo-aleatory messages embedded on chaotic laser signals.
- fBm and fGn dynamics.
- Characterization of Pseudo Random Number Generators and randomization of chaotic series.
- Stochastic resonance and coherent resonance.
- Analysis of Classical-Quantum transition problem.
- Study of mode and genre of literary texts. Internal chronological development of authorial styles.
- Alterations in the erythrocyte due to different illness.
- Brain electrical activity: EEG background distinction, epileptic activity, sleep, ERP, EP.
- Quantitative brain maturation.
- Identification of biomarkers that correlate with cancer progression.

Thanks !!!!



... and see you in Newcastle, Australia

ENTROPY AND STATISTICAL COMPLEXITY PUBLICATIONS:

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