

what is 'chaos'?

a field theorist stroll through Bernoullistan

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ChaosBook.org/overheads/spatiotemporal
→ Chaotic field theory slides

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Mephistopheles knocks at Faust's door and says, "Du mußt es dreimal sagen!"

. "You have to say it three times"

— Johann Wolfgang von Goethe

. *Faust I - Studierzimmer 2. Teil*

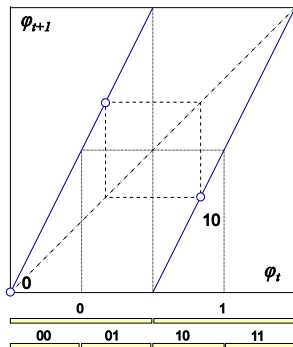
- 1 what is this about
- 2 coin toss
- 3 temporal cat
- 4 spatiotemporal cat
- 5 bye bye, dynamics

(1) coin toss, if you are stuck in XVIII century

time-evolution formulation

fair coin toss

Bernoulli map



$$\phi_{t+1} = \begin{cases} 2\phi_t \\ 2\phi_t \pmod{1} \end{cases}$$

\Rightarrow fixed point $\bar{0}$, 2-cycle $\bar{01}$, \dots

a coin toss

the essence of deterministic chaos

what is (mod 1) ?

map with integer-valued 'stretching' parameter $s \geq 2$:

$$x_{t+1} = s x_t$$

(mod 1) : subtract the integer part $m_t = \lfloor s x_t \rfloor$
so fractional part ϕ_{t+1} stays in the unit interval $[0, 1)$

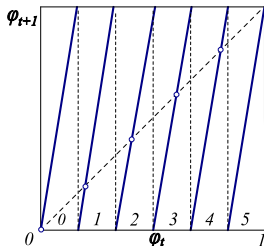
$$\phi_{t+1} = s \phi_t - m_t, \quad \phi_t \in \mathcal{M}_{m_t}$$

m_t takes values in the s -letter alphabet

$$m \in \mathcal{A} = \{0, 1, 2, \dots, s-1\}$$

a fair dice throw

slope 6 Bernoulli map



$$\phi_{t+1} = 6\phi_t - m_t, \quad \phi_t \in \mathcal{M}_{m_t}$$

6-letter alphabet

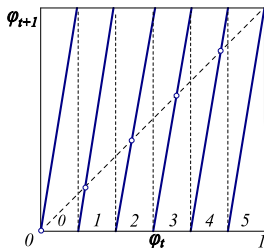
$$m_t \in \mathcal{A} = \{0, 1, 2, \dots, 5\}$$

6 subintervals $\{\mathcal{M}_0, \mathcal{M}_1, \dots, \mathcal{M}_5\}$

what is chaos ?

a fair dice throw

6 subintervals $\{\mathcal{M}_{m_t}\}$, 6^2 subintervals $\{\mathcal{M}_{m_1 m_2}\}, \dots$



each subinterval contains a periodic point, labeled by $M = m_1 m_2 \dots m_n$

$N_n = 6^n - 1$ **unstable** orbits

definition : chaos is

positive Lyapunov ($\ln s$) - positive entropy ($\frac{1}{n} \ln N_n$)

definition : chaos is

positive **Lyapunov** ($\ln s$) - positive **entropy** ($\frac{1}{n} \ln N_n$)

- **Lyapunov** : how fast is local escape?
- **entropy** : how many ways of getting back?

\Rightarrow **ergodicity**

the precise sense in which **dice throw**
is an example of deterministic chaos

(2) field theorist's chaos

lattice formulation

lattice Bernoulli

recast the time-evolution Bernoulli map

$$\phi_{t+1} = s\phi_t - m_t$$

as 1-step difference equation on the temporal lattice

$$-\phi_{t+1} + s\phi_t = m_t, \quad \phi_t \in [0, 1)$$

field ϕ_t , source m_t

on each site t of a 1-dimensional lattice $t \in \mathbb{Z}$

write an n -sites lattice segment as

the field configuration and the symbol block

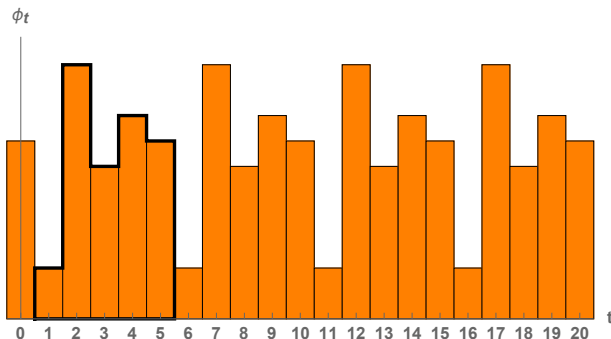
$$\Phi = (\phi_{t+1}, \dots, \phi_{t+n}), \quad \mathbf{M} = (m_{t+1}, \dots, m_{t+n})$$

‘M’ for ‘marching orders’ : come here, then go there, ...

scalar field theory on 1-dimensional lattice

write a periodic field over n -sites Bravais cell as
the [field configuration](#) and the [symbol block](#) (sources)

$$\Phi = (\phi_{t+1}, \dots, \phi_{t+n}), \quad \mathbf{M} = (m_{t+1}, \dots, m_{t+n})$$



‘M’ for ‘marching orders’ : come here, then go there, ...

think globally, act locally

Bernoulli condition at every lattice site t , local in time

$$-\phi_{t+1} + s\phi_t = m_t$$

is enforced by the global equation

$$(-r + s1) \Phi = M,$$

$[n \times n]$ shift matrix

$$r_{jk} = \delta_{j+1,k}, \quad r = \begin{pmatrix} 0 & 1 & & & \\ & 0 & 1 & & \\ & & & \ddots & \\ & & & 0 & 1 \\ 1 & & & & 0 \end{pmatrix}$$

compares the neighbors

think globally, act locally

solving the lattice Bernoulli system

$$\mathcal{J}\Phi = M,$$

$$[n \times n] \text{ Hill matrix} \quad \mathcal{J} = -r + s \mathbf{1},$$

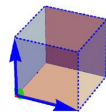
is a search for zeros of the function

$$F[\Phi] = \mathcal{J}\Phi - M = 0$$

the entire **global lattice state** Φ_M is now

a single **fixed point** $(\phi_1, \phi_2, \dots, \phi_n)$

in the n -dimensional unit hyper-cube



$$\Phi \in [0, 1]^n$$

orbit stability

Hill matrix

solving a nonlinear

$$F[\Phi] = 0 \quad \text{fixed point condition}$$

with Newton method requires evaluation of the $[n \times n]$

Hill matrix

$$\mathcal{J}_{ij} = \frac{\delta F[\Phi]_i}{\delta \phi_j}$$

what does this global Hill matrix do?

- 1 fundamental fact !
- 2 global stability of lattice state Φ , perturbed everywhere

(1)

fundamental fact

(1) fundamental fact

to satisfy the fixed point condition

$$\mathcal{J}\Phi - \mathbf{M} = 0$$

the Hill matrix \mathcal{J}

- 1 stretches the unit hyper-cube $\Phi \in [0, 1)^n$ into the n -dimensional **fundamental parallelepiped**
- 2 maps each periodic point $\Phi_{\mathbf{M}} \Rightarrow$ integer lattice \mathbb{Z}^n point
- 3 then translate by integers $\mathbf{M} \Rightarrow$ into the origin

hence $N_n = \text{total } \# \text{ solutions} = \# \text{ integer lattice points within the fundamental parallelepiped}$

the **fundamental fact**¹ : **Hill determinant** counts solutions

$$N_n = \text{Det } \mathcal{J}$$

$\#$ integer points in fundamental parallelepiped = its volume

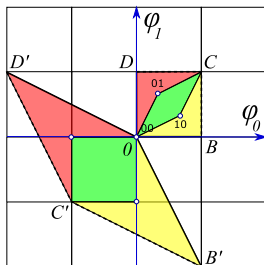
¹M. Baake et al., J. Phys. A **30**, 3029–3056 (1997).

example : fundamental parallelepiped for $n = 2$

Hill matrix for $s = 2$; unit square basis vectors ; their images :

$$\mathcal{J} = \begin{pmatrix} 2 & -1 \\ -1 & 2 \end{pmatrix}; \quad \Phi_B = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \rightarrow \Phi_{B'} = \mathcal{J} \Phi_B = \begin{pmatrix} 2 \\ -1 \end{pmatrix} \dots,$$

Bernoulli periodic points of period 2



$$N_2 = 3$$

fixed point Φ_{00}

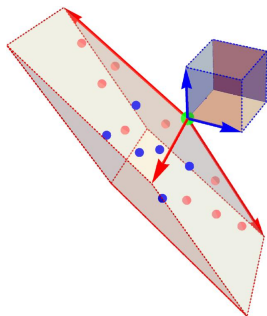
2-cycle Φ_{01}, Φ_{10}

square $[0BCD] \Rightarrow \mathcal{J} \Rightarrow$ fundamental parallelepiped $[0B'C'D']$

fundamental fact for any n

an $n = 3$ example

\mathcal{T} [unit hyper-cube] = [fundamental parallelepiped]



unit hyper-cube $\Phi \in [0, 1)^3$

$n > 3$ cannot visualize

a periodic point \Rightarrow integer lattice point : ● on a face, ● in the interior

(2)

orbit stability

(2) orbit stability vs. temporal stability

Hill matrix

$\mathcal{J}_{ij} = \frac{\delta F[\Phi]_i}{\delta \phi_j}$ stability under **global** perturbation of the whole orbit
for n large, a huge $[dn \times dn]$ matrix

temporal Jacobian matrix

J propagates **initial** perturbation n time steps
small $[d \times d]$ matrix

J and \mathcal{J} are related by²

Hill's 1886 remarkable formula

$$|\text{Det } \mathcal{J}_M| = |\det(\mathbf{1} - J_M)|$$

\mathcal{J} is **huge**, even ∞ -dimensional matrix

J is **tiny**, few degrees of freedom matrix

²G. W. Hill, Acta Math. 8, 1–36 (1886).

field theorist's chaos

definition : chaos is

expanding	Hill determinants	$\text{Det } \mathcal{J}$
exponential \sharp	field configurations	N_n

the precise sense in which
a (discretized) field theory is deterministically chaotic

note : there is no 'time' in this definition

periodic orbit theory

volume of a periodic orbit

Ozorio de Almeida and Hannay³ 1984 :

of periodic points is related to a Jacobian matrix by

principle of uniformity

“periodic points of an ergodic system, counted with their natural weighting, are uniformly dense in phase space”

where

‘natural weight’ of periodic orbit M

$$\frac{1}{|\det(1 - J_M)|}$$

³A. M. Ozorio de Almeida and J. H. Hannay, J. Phys. A **17**, 3429 (1984).

periodic orbits partition lattice states into neighborhoods

how come **Hill determinant** $\text{Det } \mathcal{J}$ counts periodic points ?

‘principle of uniformity’ is in⁴

periodic orbit theory

known as the **flow conservation** sum rule :

$$\sum_{\mathbf{M}} \frac{1}{|\det(1 - \mathcal{J}_{\mathbf{M}})|} = \sum_{\mathbf{M}} \frac{1}{|\text{Det } \mathcal{J}_{\mathbf{M}}|} = 1$$

sum over periodic points $\Phi_{\mathbf{M}}$ of period n

state space is divided into

neighborhoods of periodic points of period n

⁴P. Cvitanović, “Why cycle?”, in *Chaos: Classical and Quantum*, edited by P. Cvitanović et al. (Niels Bohr Inst., Copenhagen, 2020).

periodic orbit counting

how come a $\text{Det } \mathcal{J}$ counts periodic points ?

flow conservation sum rule :

$$\sum_{\Phi_M \in \text{Fix} f^n} \frac{1}{|\text{Det } \mathcal{J}_M|} = 1$$

Bernoulli system 'natural weighting' is simple :

the determinant $\text{Det } \mathcal{J}_M = \text{Det } \mathcal{J}$ the same for all periodic points, whose number thus verifies the **fundamental fact**

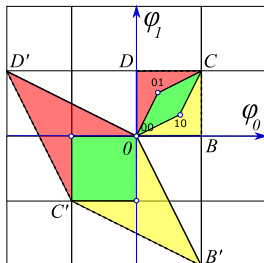
$$N_n = |\text{Det } \mathcal{J}|$$

the number of Bernoulli periodic lattice states

$$N_n = |\text{Det } \mathcal{J}| = s^n - 1 \quad \text{for any } n$$

remember the fundamental fact?

period 2 example



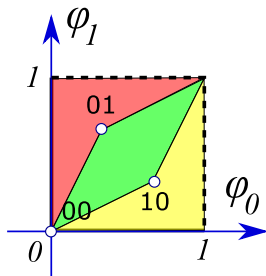
fixed point	Φ_{00}
2-cycle	Φ_{01}, Φ_{10}

$$\mathcal{T}[\text{unit hyper-cube}] = [\text{fundamental parallelepiped}]$$

look at preimages of the fundamental parallelepiped :

example : lattice states of period 2

unit hypercube, partitioned



fixed point Φ_{00}
2-cycle Φ_{01}, Φ_{10}

flow conservation sum rule

$$\frac{1}{|\text{Det } \mathcal{J}_{00}|} + \frac{1}{|\text{Det } \mathcal{J}_{01}|} + \frac{1}{|\text{Det } \mathcal{J}_{10}|} = 1$$

sum over periodic points Φ_M of period $n = 2$

state space is divided into

neighborhoods of periodic points of period n

Amazing! I did not understand a single word.
—Fritz Haake 1988

zeta function

periodic orbit theory, version (1) : counting lattice states

topological zeta function

$$1/\zeta_{\text{top}}(z) = \exp \left(- \sum_{n=1}^{\infty} \frac{z^n}{n} N_n \right)$$

- ① weight $1/n$ as by (cyclic) translation invariance, n lattice states are equivalent
- ② zeta function counts **orbits**, one per each set of equivalent lattice states

Bernoulli topological zeta function

counts **orbits**, one per each set of lattice states $N_n = s^n - 1$

$$1/\zeta_{\text{top}}(z) = \exp \left(- \sum_{n=1}^{\infty} \frac{z^n}{n} N_n \right) = \frac{1 - sz}{1 - z}$$

numerator $(1 - sz)$ says that Bernoulli orbits are built from s fundamental **primitive** lattice states,

the fixed points $\{\phi_0, \phi_1, \dots, \phi_{s-1}\}$

every other lattice state is built from their concatenations and repeats.

solved!

this is 'periodic orbit theory'

And if you don't know, now you know

summary : think globally, act locally

the problem of enumerating and determining all **lattice states** stripped to its essentials :

- 1 each solution is a zero of the global **fixed point** condition

$$F[\Phi] = 0$$

- 2 **global stability** : the Hill matrix

$$\mathcal{J}_{ij} = \frac{\delta F[\Phi]_i}{\delta \phi_j}$$

- 3 **fundamental fact** : the number of period- n orbits

$$N_n = |\text{Det } \mathcal{J}|$$

- 4 **zeta function** $1/\zeta_{\text{top}}(z)$: all predictions of the theory

next : a kicked rotor

Du mußt es dreimal sagen!

— Mephistopheles

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- 5 bye bye, dynamics