

Spatiotemporal Tiling of the Kuramoto-Sivashinsky equation

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Motivation

- High dimensional fluid flows are hard. Most results in “minimal flow units”¹
- Conventional computational tools are providing diminishing returns
- Spatiotemporal methods are scalable²

¹J. Jiménez and P. Moin, J. Fluid Mech. **225**, 213–240 (1991).

²Q. Wang et al., Phys. Fluids **25**, 110818 (2013).

Doubly Periodic Kuramoto-Sivashinsky equation

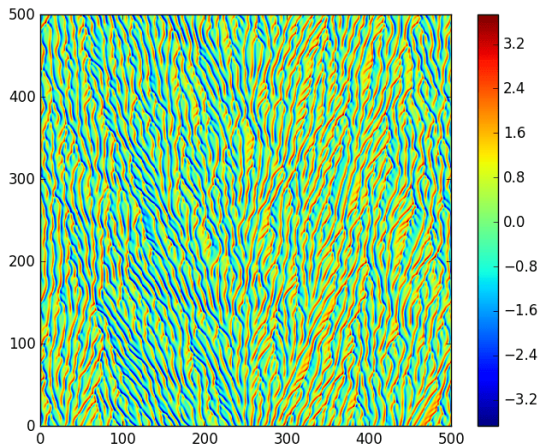
$$u_t(x, t) = -\underline{u_{xx}(x, t)} - \underline{u_{xxxx}(x, t)} - \underline{u(x, t)u_x(x, t)}$$

Diffusion (wrong sign) “Hyper”-diffusion Advection

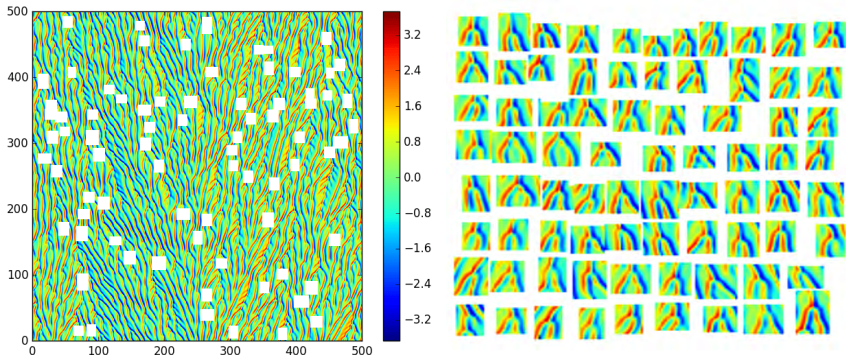
Boundary conditions

$$u(x, t) = u(x, t + T_p) = u(x + L, t) = u(x + L, t + T_p)$$

Large spatiotemporal simulation



Shadowing



Substituting

$$u(x, t) = \sum_{m,n} \tilde{u}_{m,n} e^{i\omega_n t + i q_m x}, \quad (1)$$

into the Kuramoto-Sivashinsky equation gives,

$$\mathbf{F}(\mathbf{x}) = \left(i \frac{2\pi n}{T} + \left(\frac{2\pi m}{L} \right)^2 - \left(\frac{2\pi m}{L} \right)^4 \right) \tilde{u}_{m,n} + \frac{1}{2} \frac{2\pi m}{L} \mathcal{F}((\mathcal{F}^{-1}(\tilde{u}_{m,n}))^2) \quad (2)$$

Solve nonlinear algebraic equations for $\mathbf{x} = (\tilde{\mathbf{u}}, T, L)$. Similar to formulation of Lopez³. We use real-valued Fourier basis in computations.

³V. López, *Numerical continuation of invariant solutions of the complex Ginzburg–Landau equation*, 2015.

Disadvantages of spatiotemporal methods

- ① High dimensionality
- ② Numerically challenging
- ③ Possible to find isolated solutions that are unimportant
- ④ Find different members of same continuous families (redundancy)

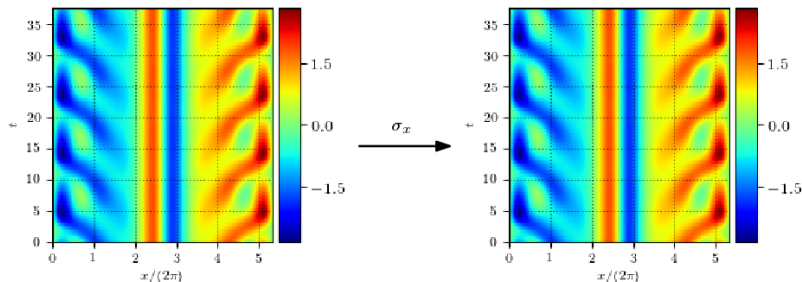
Advantages of spatiotemporal methods

- ① Can use variational methods
- ② Fourier-Fourier basis
- ③ Topological constraint (boundary conditions) provides access to potent methods
- ④ Stability does not play a role in finding solutions.
- ⑤ Can exploit spatiotemporal symmetries

Spatiotemporal Symmetries

Reflection invariance

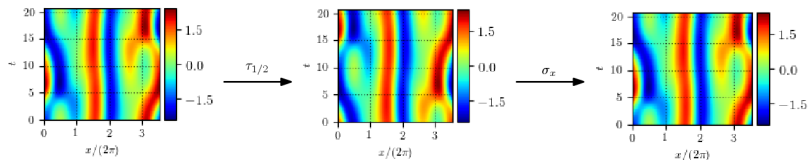
$$u(x, t) = \sigma_x u(x, t)$$



Spatiotemporal Symmetries

Shift-Reflection Invariance

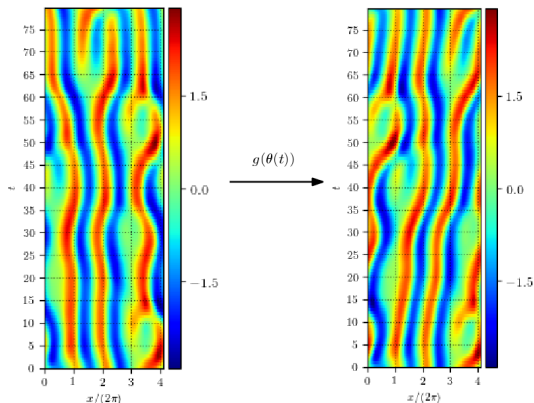
$$u(x, t) = \sigma_x \tau_{1/2} u(x, t)$$



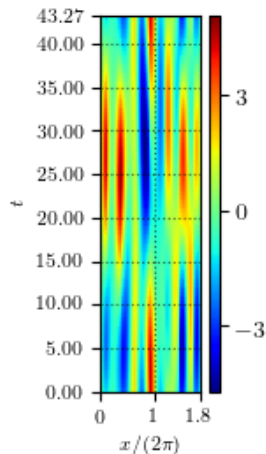
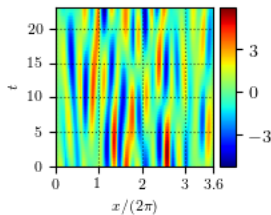
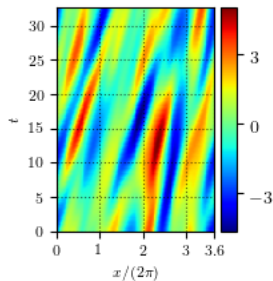
Spatiotemporal Symmetries

Spatial Translation Symmetry

$$u(x, t) = g(\theta(t)) \circ u(x, t + T_p)$$



Initial conditions



Adjoint Descent Method

$$\mathbf{F}(\mathbf{x}) = (i\frac{2\pi n}{T} + (\frac{2\pi m}{L})^2 - (\frac{2\pi m}{L})^4)\tilde{u}_{m,n} + \frac{1}{2}\frac{2\pi m}{L}\mathcal{F}((\mathcal{F}^{-1}(\tilde{u}_{m,n}))^2) \quad (3)$$

$$I = \frac{1}{2}\mathbf{F}^\top \mathbf{F}. \quad (4)$$

introduce fictitious time (τ) flow by differentiation of cost function,

$$\partial_{\tau} I = (J^{\top} \mathbf{F})^{\top} (\partial_{\tau} \mathbf{x}) \quad (5)$$

“adjoint descent” method is defined by choice⁴

$$\partial_{\tau} \mathbf{x} = -(J^{\top} \mathbf{F}) \quad (6)$$

⁴M. Farazmand, J. Fluid M. **795**, 278–312 (2016).

Damped Gauss-Newton Method

Solve $\mathbf{F}(\mathbf{x}^*) = 0$ via linearized system

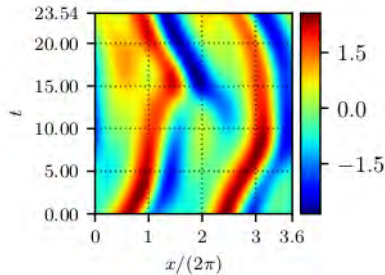
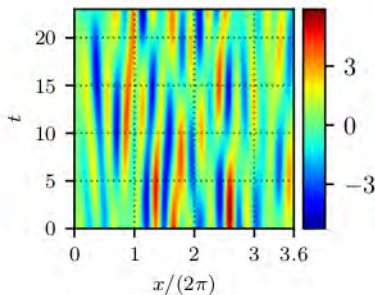
$$\mathbf{F}(\mathbf{x}^*) = \mathbf{F}(\mathbf{x}) + J\delta\mathbf{x} \quad (7)$$

$$J = \begin{bmatrix} \frac{\partial \mathbf{F}}{\partial \mathbf{x}} & \frac{\partial \mathbf{F}}{\partial T} & \frac{\partial \mathbf{F}}{\partial L} \end{bmatrix} \quad (8)$$

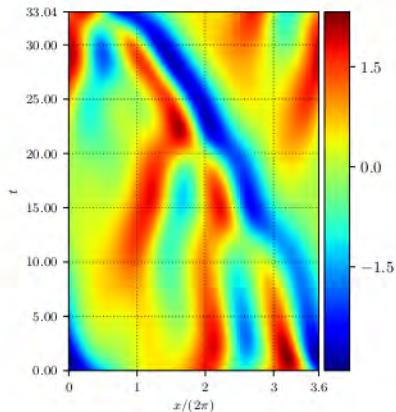
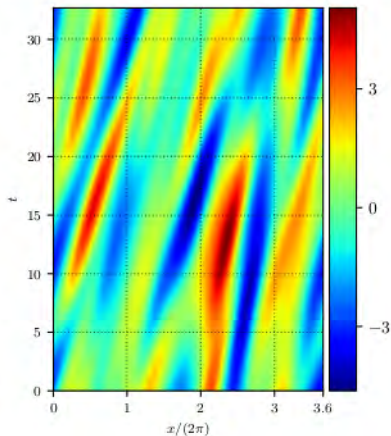
$$J\delta\mathbf{x} = -\mathbf{F}(\mathbf{x}) \quad (9)$$

Shift-reflect symmetry

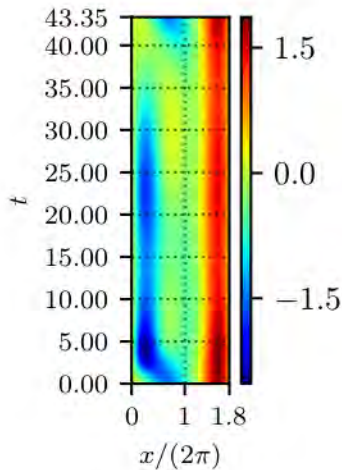
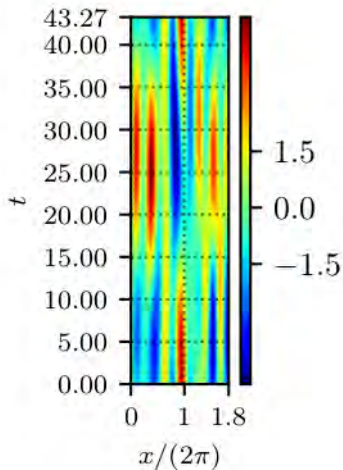
A couple thousand new invariant 2-torus solutions, examples:



Spatial translation symmetry

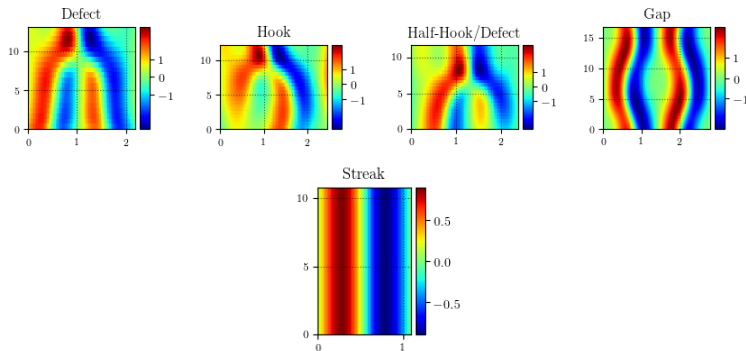


Reflection symmetry

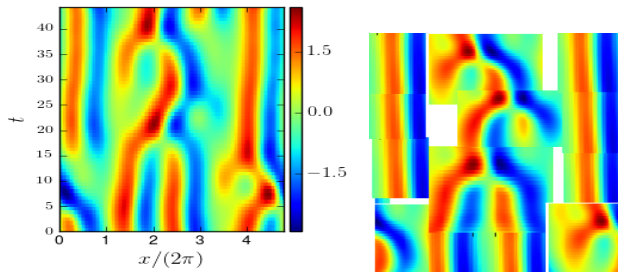


Qualitative guesses for tiles

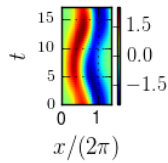
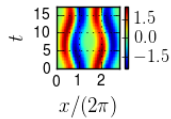
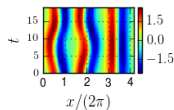
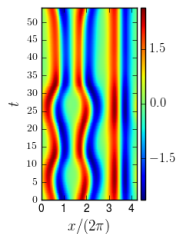
Sift through collection to find frequently occurring spatiotemporal patterns



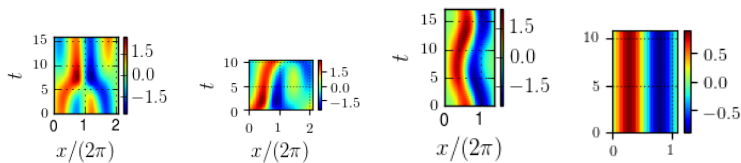
Qualitative tiling of actual solution



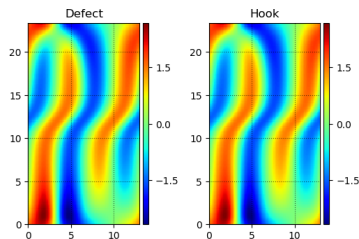
Finding tiles by converging subdomains



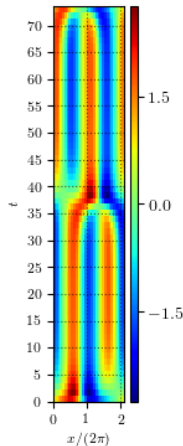
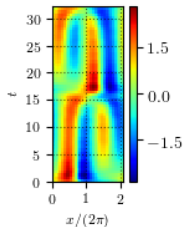
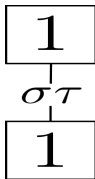
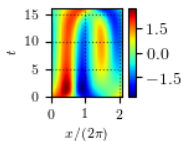
Converged tiles



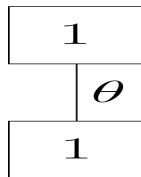
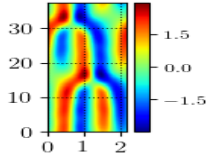
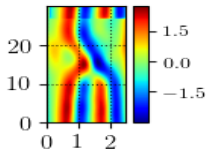
Continuous families and numerical continuation



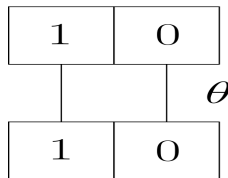
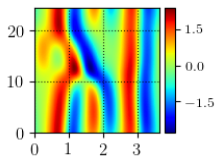
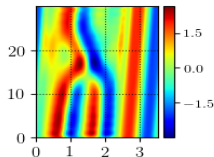
How do we use tiles?



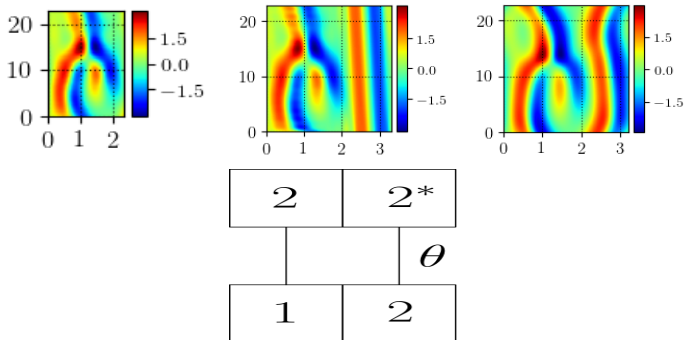
Spatiotemporal symbolic dynamics



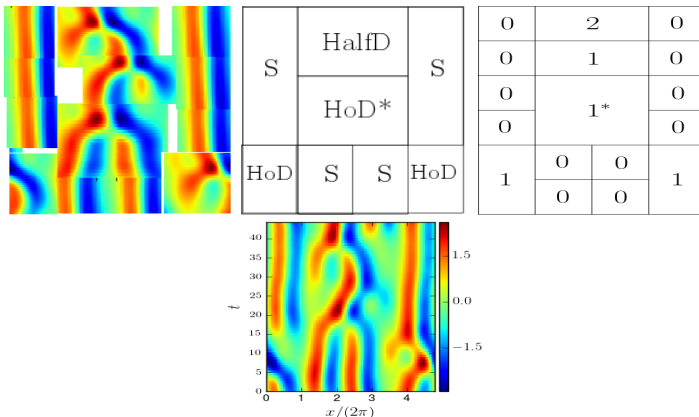
Spatiotemporal symbolic dynamics



Spatiotemporal symbolic dynamics

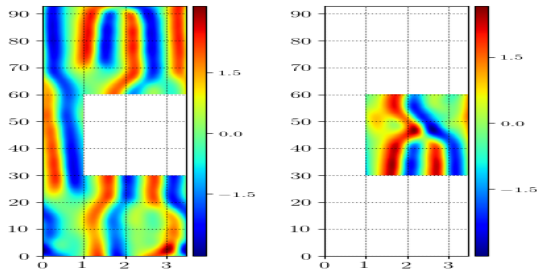


Quantitative tiling of actual solution

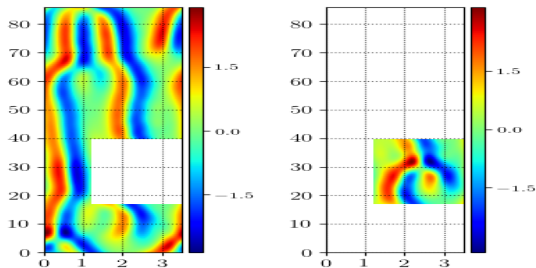


⁴Each symbol is technically paired with spatiotemporal parameters (T, L)

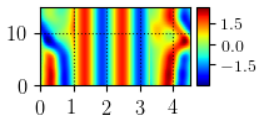
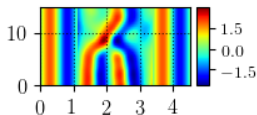
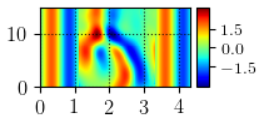
New solutions from gluing together old



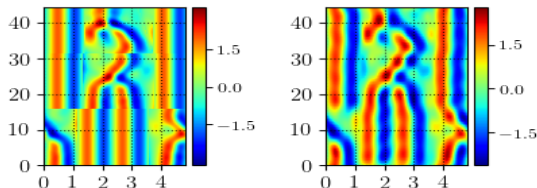
New solutions from gluing together old



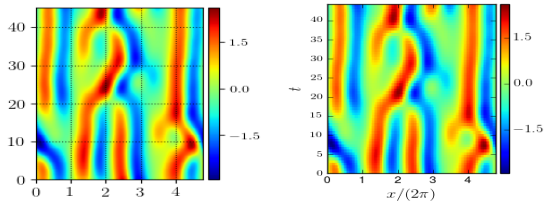
New solutions from gluing together old



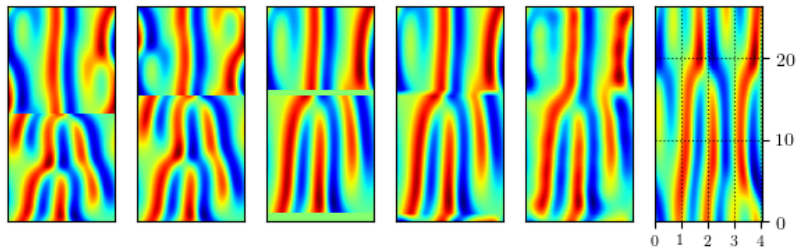
New solutions from gluing together old



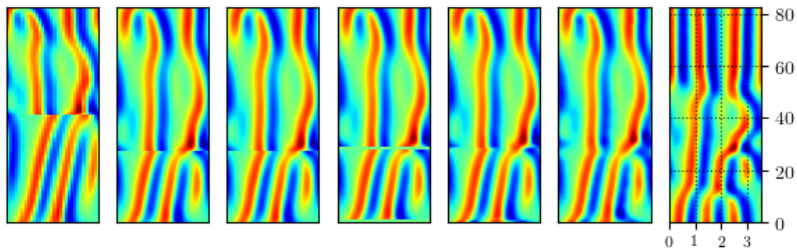
New solutions from gluing together old



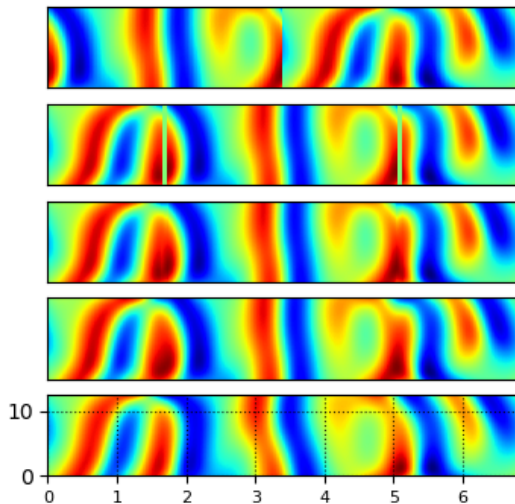
Gluing (larger) solutions, pre-periodic orbit time example



Gluing (larger) solutions, relative periodic orbit time example



Gluing (larger) solutions, pre-periodic orbit space example



Future Work

- Investigate Lagrangian formulation and application of Hill's formula^{5,6}
- Complete spatiotemporal symbolic dynamics
- Continuing improvement of numerical methods
- Application of theory to PDEs with more continuous spatial dimensions such as quasi 2-D Kolmogorov flow^{7,8}.

⁵N. H. Ibragimov and T. Kolsrud, Nonlin. Dyn. **36**, 29–40 (2004).

⁶S. V. Bolotin and D. V. Treschev, Russ. Math. Surv. **65**, 191 (2010).

⁷B. Suri et al., Phys. Rev. E **98**, 023105 (2018).

⁸G. J. Chandler and R. R. Kerswell, J. Fluid Mech. **722**, 554–595 (2013).

Conclusion

- Developed spatiotemporal methods for finding invariant 2-tori of the Kuramoto-Sivashinsky equation
- Beginning to develop symbolic dynamics for the spatiotemporal Kuramoto-Sivashinsky equation
- Developed new method for finding larger (domain) solutions by combining smaller (domain) solutions