Spatiotemporal Tiling of the Kuramoto-Sivashinsky equation

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Physics Forum, Oct 18 2018



Spatiotemporal Chaos

- What do I mean by spatiotemporal chaos?
- Why study spatiotemporal chaos?



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Doubly Periodic Kuramoto-Sivashinsky equation

$$u_t(x,t) = -\underline{u_{xx}(x,t)} - \underline{u_{xxxx}(x,t)} - \underline{u(x,t)u_x(x,t)}$$

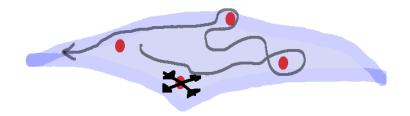
Diffusion (wrong sign) "Hyper"-diffusion Advection

Boundary conditions

$$u(x,t) = u(x,t+T_p) = u(x+L,t) = u(x+L,t+T_p)$$

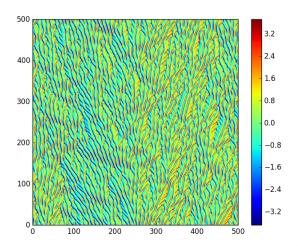


Dynamical Systems Theory: A Cartoon





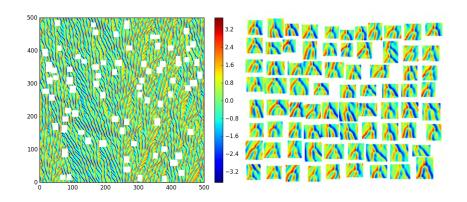
Large spatiotemporal simulation





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Shadowing





(Computational) Motivating claim

"Time domain parallel methods, such as Parareal, scale poorly in simulation of chaotic, multiscale dynamical systems, e.g., turbulent flows. This poor scalability is because the initial value problem of a chaotic dynamical system is **ill-conditioned**. A perturbation of magnitude ϵ at the beginning of the simulation can cause a difference of $\epsilon e^{\lambda t}$ at time t later, where lambda is the maximal Lyapunov exponent."

[Wang et al., 2013]



Fourier-Fourier Basis

Substituting

$$u(x,t) = \sum_{m,n} \tilde{u}_{m,n} e^{i\omega_n t + iq_m x}, \qquad (1)$$

into the Kuramoto-Sivashinsky equation gives,

$$\mathbf{F}(\mathbf{x}) = \left(i\frac{2\pi n}{T} + \left(\frac{2\pi m}{L}\right)^2 - \left(\frac{2\pi m}{L}\right)^4\right)\tilde{u}_{m,n} + \frac{1}{2}\frac{2\pi m}{L}\mathcal{F}((\mathcal{F}^{-1}(\tilde{u}_{m,n}))^2)$$
(2)

Solve nonlinear algebraic equations for $\mathbf{x}=(\tilde{\mathbf{u}},T,L)$. Similar to formulation in [López, 2015].



Disadvantages of Spatiotemporal Methods

- Strong coupling
- Many degrees of freedom
- Linearized system is ill-conditioned



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Advantages of Spatiotemporal Methods

- Can use variational methods
- Pourier-Fourier basis makes computation easier and faster
- Oan use (spatiotemporal) symmetries to reduce degrees of freedom

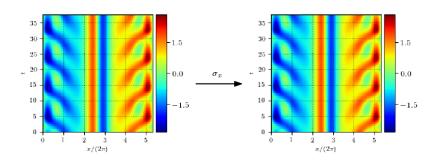


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Reflection Invariance

Subgroup $\mathbb{Z}_2\times\mathbb{I}$

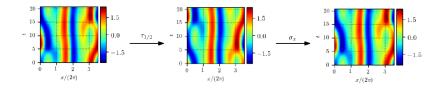
$$u(x,t) = \sigma_x u(x,t)$$



Shift-Reflection Invariance

Subgroup $\{e, \sigma_x \tau_{1/2}\}$ (itself a subgroup of Subgroup $Z_2 \times C_2$)

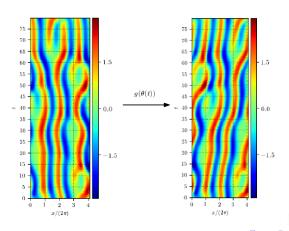
$$u(x,t) = \sigma_x \tau_{1/2} u(x,t)$$



Spatial Translation Symmetry

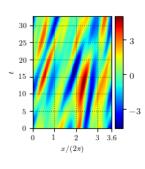
Subgroup $C_M \times \mathbb{I}$

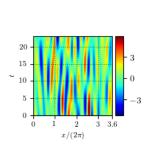
$$u(x,t) = g(\theta(t)) \circ u(x,t+T_p)$$

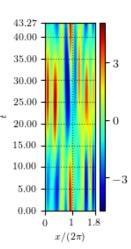




Initial conditions











Adjoint Descent Method

$$\mathbf{F}(\mathbf{x}) = \left(i\frac{2\pi n}{T} + \left(\frac{2\pi m}{L}\right)^2 - \left(\frac{2\pi m}{L}\right)^4\right)\tilde{u}_{m,n} + \frac{1}{2}\frac{2\pi m}{L}\mathcal{F}((\mathcal{F}^{-1}(\tilde{u}_{m,n}))^2)$$
(3)

$$I = \frac{1}{2} \mathbf{F}^{\top} \mathbf{F} \,. \tag{4}$$



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Adjoint Descent Method

Develop fictitious time (τ) flow by differentiation of cost function.

$$\partial_{\tau} I = (J^{\top} \mathbf{F})^{\top} (\partial_{\tau} \mathbf{x}) \tag{5}$$

"adjoint descent" method defined by choice [Farazmand, 2016],

$$\partial_{\tau} \mathbf{x} = -(J^{\top} \mathbf{F}) \tag{6}$$



Backtracking Gauss-Newton Method

Solve $\mathbf{F}(\mathbf{x}^*) = 0$ via linearized system

$$\mathbf{F}(\mathbf{x}^*) = \mathbf{F}(\mathbf{x}) + J\Delta\mathbf{x} \tag{7}$$

$$J = \begin{bmatrix} \frac{\partial \mathbf{F}}{\partial \mathbf{x}} & \frac{\partial \mathbf{F}}{\partial T} & \frac{\partial \mathbf{F}}{\partial L} \end{bmatrix}$$
 (8)

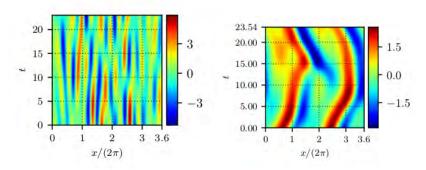
$$J\Delta \mathbf{x} = -\mathbf{F}(\mathbf{x}) \tag{9}$$

- ullet Least-squares solver, minimizes norm of $\Delta {f x}$
- Line-search and or backtracking improves convergence.

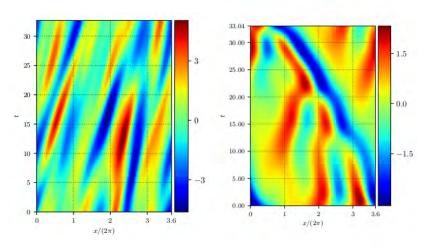


Shift-reflect symmetry

A couple thousand new invariant 2-torus solutions, examples:

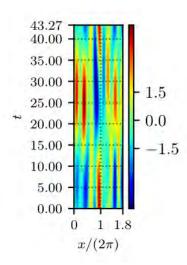


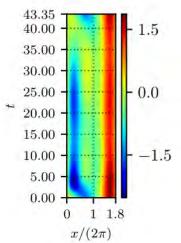
Spatial translation symmetry





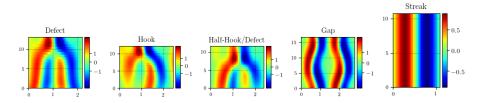
Reflection symmetry





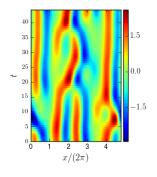


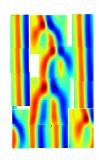
Qualitative guesses for tiles





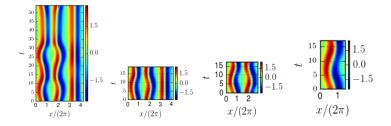
Qualitative tiling of actual solution





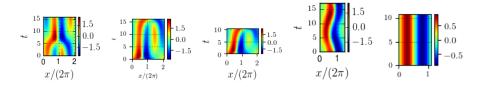


Finding tiles by converging subdomains





Tiles

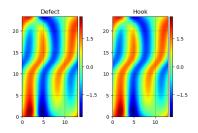




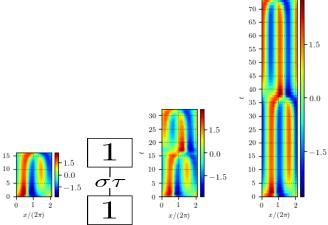


Continuous families and numerical continuation

Evidence that "hook" equals "defect". Numerically continue each to the same domain size, quotient continuous symmetries.

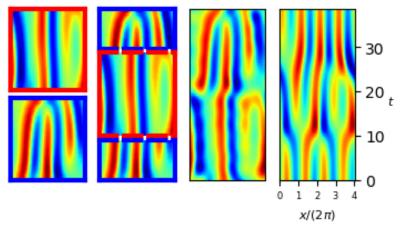


Homoclinic connection from tiles



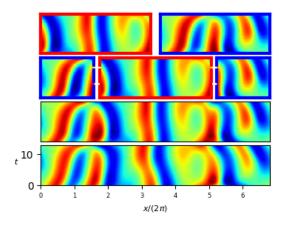


Gluing (larger) solutions, time example





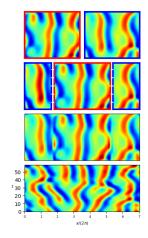
Gluing (larger) solutions, space example





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Gluing (larger) solutions, space example





Summary and important contributions

- New nonlinear optimization routines
- Larger (domain size) periodic solutions
- Connection between small domains and large domains (shadowing)
 via tiles
- Spatiotemporal gluing of solutions



Future Work I

- Need to find all "tiles", hope for small finite number.
- Need to deal with continuous families.
- Need to figure out spatiotemporal symbolic dynamics



References I

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An adjoint-based approach for finding invariant solutions of Navier-Stokes equations. *J. Fluid M.*, 795:278–312.

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[Wang et al., 2013] Wang, Q., Gomez, S. A., Blonigan, P. J., Gregory, A. L., and Qian, E. Y. (2013).

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Thank you for listening!

Questions?

