### Thermodynamics and Kinetics

#### C. Paolucci

University of Notre Dame

Department of Chemical & Biomolecular Engineering

• You used GAMESS to calculate the internal (ground state) electronic energy for various molecules under the Born-Oppenheimer approximation.

$$\bullet \ 2\,\mathrm{H^+} + \mathrm{O^{8+}} + 10\,\mathrm{e^-} \xrightarrow{E^{elec}} \mathrm{H_2O}$$

• Even at 0 K molecules have vibrational energy, called the zero-point vibrational energy (ZPVE). You calculated this from a frequency calculation.

$$ZPVE = \frac{1}{2}h\sum_{i=1}^{3n-6}\nu_i$$

• You used GAMESS to calculate the internal (ground state) electronic energy for various molecules under the Born-Oppenheimer approximation.

$$\bullet \ 2\,\mathrm{H^+} + \mathrm{O^{8+}} + 10\,\mathrm{e^-} \xrightarrow{E^{elec}} \mathrm{H_2O}$$

• Even at 0 K molecules have vibrational energy, called the zero-point vibrational energy (ZPVE). You calculated this from a frequency calculation.

$$ZPVE = \frac{1}{2}h\sum_{i=1}^{3n-6}\nu_i$$

• You used GAMESS to calculate the internal (ground state) electronic energy for various molecules under the Born-Oppenheimer approximation.

$$\bullet \ 2\,\mathrm{H^+} + \mathrm{O^{8+}} + 10\,\mathrm{e^-} \xrightarrow{E^{elec}} \mathrm{H_2O}$$

• Even at 0 K molecules have vibrational energy, called the zero-point vibrational energy (ZPVE). You calculated this from a frequency calculation.

$$ZPVE = \frac{1}{2}h\sum_{i=1}^{3n-6}\nu_i$$

• You used GAMESS to calculate the internal (ground state) electronic energy for various molecules under the Born-Oppenheimer approximation.

• 
$$2 H^+ + O^{8+} + 10 e^- \xrightarrow{E^{elec}} H_2O$$

• Even at 0 K molecules have vibrational energy, called the zero-point vibrational energy (ZPVE). You calculated this from a frequency calculation.

$$ZPVE = \frac{1}{2}h\sum_{i=1}^{3n-6}\nu_i$$

.

# What is the energy of a molecule?

- $E^0$  is the minimum energy of a molecule.
- Higher energies are possible through electronic or vibrational excitations as well as translational and rotational motion. If we assume these contributions are separable

$$E = E^0 + E^{elec'} + E^{vib} + E^{rot} + E^{tran}$$

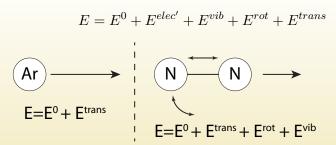
# What is the energy of a molecule?

- $E^0$  is the minimum energy of a molecule.
- Higher energies are possible through electronic or vibrational excitations as well as translational and rotational motion. If we assume these contributions are separable

$$E = E^0 + E^{elec'} + E^{vib} + E^{rot} + E^{trans}$$

# What is $\overline{the\ energy\ of\ a\ molecule?}$

- $E^0$  is the minimum energy of a molecule.
- Higher energies are possible through electronic or vibrational excitations as well as translational and rotational motion. If we assume these contributions are separable



• We need to specify these quantities to describe the microscopic state of the molecule.

- Macroscopically we want the average value of these quantities over all possible states of an ensemble of molecules at some specified externally imposed set of thermodynamic conditions. (ensemble average)
- This is the domain of statistical mechanics.
- The relative probability of a molecule being in a state with energy  $E_i$  above  $E^0$  is given by the Boltzmann factor

$$P_i \propto e^{-E_i \beta}, \qquad \beta = \frac{1}{k_B T}$$

$$Q(N, V, T) = \sum_{i} e^{-E_{i}\beta}$$

- Macroscopically we want the average value of these quantities over all possible states of an ensemble of molecules at some specified externally imposed set of thermodynamic conditions. (ensemble average)
- This is the domain of statistical mechanics.
- The relative probability of a molecule being in a state with energy  $E_i$  above  $E^0$  is given by the Boltzmann factor

$$P_i \propto e^{-E_i \beta}, \qquad \beta = \frac{1}{k_B T}$$

$$Q(N, V, T) = \sum_{i} e^{-E_{i}\beta}$$

- Macroscopically we want the average value of these quantities over all possible states of an ensemble of molecules at some specified externally imposed set of thermodynamic conditions. (ensemble average)
- This is the domain of statistical mechanics.
- The relative probability of a molecule being in a state with energy  $E_i$  above  $E^0$  is given by the Boltzmann factor

$$P_i \propto e^{-E_i\beta}, \qquad \beta = \frac{1}{k_B T}$$

$$Q(N,V,T) = \sum_i e^{-E_i\beta}$$

- Macroscopically we want the average value of these quantities over all possible states of an ensemble of molecules at some specified externally imposed set of thermodynamic conditions. (ensemble average)
- This is the domain of statistical mechanics.
- The relative probability of a molecule being in a state with energy  $E_i$  above  $E^0$  is given by the Boltzmann factor

$$P_i \propto e^{-E_i\beta}, \qquad \beta = \frac{1}{k_B T}$$

$$Q(N, V, T) = \sum_{i} e^{-E_{i}\beta}$$

- Macroscopically we want the average value of these quantities over all possible states of an ensemble of molecules at some specified externally imposed set of thermodynamic conditions. (ensemble average)
- This is the domain of statistical mechanics.
- The relative probability of a molecule being in a state with energy  $E_i$  above  $E^0$  is given by the Boltzmann factor

$$P_i \propto e^{-E_i \beta}, \qquad \beta = \frac{1}{k_B T}$$

 $\bullet$  Sum over all states is the (in this canonical, N,V,T free variables) partition function

$$Q(N, V, T) = \sum_{i} e^{-E_{i}\beta}$$

 $\bullet$  All thermodynamic quantities can be derived from Q

- Macroscopically we want the average value of these quantities over all possible states of an ensemble of molecules at some specified externally imposed set of thermodynamic conditions. (ensemble average)
- This is the domain of statistical mechanics.
- The relative probability of a molecule being in a state with energy  $E_i$  above  $E^0$  is given by the Boltzmann factor

$$P_i \propto e^{-E_i \beta}, \qquad \beta = \frac{1}{k_B T}$$

• Sum over all states is the (in this canonical, N, V, T free variables) partition function

$$Q(N, V, T) = \sum_{i} e^{-E_{i}\beta}$$

 $U = -\left. \frac{\partial \ln Q}{\partial \beta} \right|_{NV}$ 

.

- Macroscopically we want the average value of these quantities over all possible states of an ensemble of molecules at some specified externally imposed set of thermodynamic conditions. (ensemble average)
- This is the domain of statistical mechanics.
- The relative probability of a molecule being in a state with energy  $E_i$  above  $E^0$  is given by the Boltzmann factor

$$P_i \propto e^{-E_i \beta}, \qquad \beta = \frac{1}{k_B T}$$

• Sum over all states is the (in this canonical, N, V, T free variables) partition function

$$Q(N, V, T) = \sum_{i} e^{-E_{i}\beta}$$

$$G = -\frac{\ln Q}{\beta} + \frac{V}{\beta} \left. \frac{\partial \ln Q}{\partial V} \right|_{NT}$$

.

- The good news: if we know Q we can find any thermodynamic quantity.
- The bad news: can be difficult to calculate for some systems

- The good news: if we know Q we can find any thermodynamic quantity.
- The bad news: can be difficult to calculate for some systems

- The good news: if we know Q we can find any thermodynamic quantity.
- The bad news: can be difficult to calculate for some systems

No intermolecular interactions Only sum over states of individual molecule Need all inter/intramolecular states, MC/MD to get thermodynamic quantities

No rotation or translation. Sum over vibrational modes

- The good news: if we know Q we can find any thermodynamic quantity.
- The bad news: can be difficult to calculate for some systems

No intermolecular interactions Only sum over states of individual molecule Need all inter/intramolecular states, MC/MD to get thermodynamic quantities

No rotation or translation. Sum over vibrational modes

• Here we will focus on an ideal gas.

• For an ideal gas of N indistinguishable molecules

$$Q = \frac{q^N(V, T)}{N!}$$

 Once again assuming decoupling the molecular partition function is

$$q = q_{trans} q_{vib} q_{rot} q_{elec}$$

• We can treat each of these pieces with models you have seen before. GAMESS chooses 298.15 K by default.

• For an ideal gas of N indistinguishable molecules

$$Q = \frac{q^N(V, T)}{N!}$$

• Once again assuming decoupling the molecular partition function is

$$q = q_{trans} q_{vib} q_{rot} q_{elec}$$

• We can treat each of these pieces with models you have seen before. GAMESS chooses 298.15 K by default.

• For an ideal gas of N indistinguishable molecules

$$Q = \frac{q^N(V, T)}{N!}$$

• Once again assuming decoupling the molecular partition function is

$$q = q_{trans} q_{vib} q_{rot} q_{elec}$$

• We can treat each of these pieces with models you have seen before. GAMESS chooses 298.15 K by default.

- Translation: Particle in a 3-D box (have to choose a standard state pressure, GAMESS chooses 1 atm)
- Rotation: Rigid Rotor
- Vibration: Harmonic Oscillator
- $\bullet$  Electronic: Sum over excited electronic states, usually only relevant at high T

- Translation: Particle in a 3-D box (have to choose a standard state pressure, GAMESS chooses 1 atm)
- Rotation: Rigid Rotor
- Vibration: Harmonic Oscillator
- $\bullet$  Electronic: Sum over excited electronic states, usually only relevant at high T
- Typical values

$$q_{trans} \approx 10^{30}$$
  $q_{rot} \approx 100$   $q_{vib} \approx 1$   $q_{elec} \approx g_0$ 

### Ideal Gas Example

- GAMESS will give you these values making the key assumptions on the previous slide after a frequency calculation.
- Lets look at Argon

### Ideal Gas Example

#### • Ar gas

```
THERMOCHEMISTRY AT T= 298.15 K
USING IDEAL GAS, RIGID ROTOR, HARMONIC NORMAL MODE APPROXIMATIONS.
P= 1.01325E+05 PASCAL.
ALL FREQUENCIES ARE SCALED BY 1.00000
                    LN O
ELEC. 1.00000E+00 0.000000
TRANS. 9.92422E+06 16.110489
ROT. 1.00000E+00 0.000000
VIB. 1.00000E+00 0.000000
TOT. 9.92422E+06 16.110489
         E H G CV CP S
      KJ/MOL KJ/MOL KJ/MOL-K J/MOL-K J/MOL-K
ELEC. 0.000 0.000 0.000 0.000 0.000 0.000
TRANS. 3.718 6.197 -39.937 12.472 20.786 154.735
ROT. 0.000 0.000 0.000 0.000 0.000 0.000
VIB. 0.000 0.000 0.000 0.000 0.000 0.000
TOTAL 3.718 6.197 -39.937 12.472 20.786
                                           154.735
VIB. THERMAL CORRECTION E(T) - E(0) = H(T) - H(0) = 0.000 \text{ J/MOL}
```

• Note: GAMESS scales  $q_{trans}$  by  $2.46 \times 10^{25}$ 

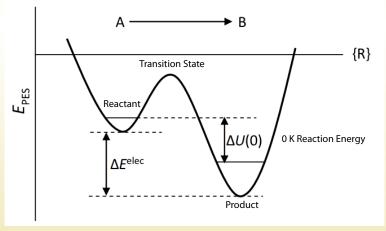
### Ideal Gas Example

#### • Now N<sub>2</sub>

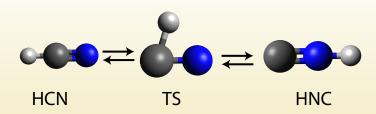
```
THERMOCHEMISTRY AT T= 298.15 K
USING IDEAL GAS, RIGID ROTOR, HARMONIC NORMAL MODE APPROXIMATIONS.
P= 1.01325E+05 PASCAL.
ALL FREOUENCIES ARE SCALED BY 1.00000
THE MOMENTS OF INERTIA ARE (IN AMU*BOHR**2)
    0.00000
            30.68954
                       30.68954
THE ROTATIONAL SYMMETRY NUMBER IS 2.0
THE ROTATIONAL CONSTANTS ARE (IN GHZ)
             58.75254
                       58.75254
THE HARMONIC ZERO POINT ENERGY IS (SCALED BY 1.000)
      0.006598 HARTREE/MOLECULE 1448.038767 CM**-1/MOLECULE
      4.140155 KCAL/MOL 17.322407 KJ/MOL
                     LN O
ELEC. 1.00000E+00 0.000000
                    15.577756
TRANS. 5.82551E+06
ROT. 5.28208E+01
                      3.966904
      1.00000E+00 0.000001
VTB.
       3.07708E+08 19.544662
TOT.
          E H G CV
                                            CP S
       KJ/MOL KJ/MOL KJ/MOL J/MOL-K
                                         J/MOL-K J/MOL-K
     0.000 0.000 0.000 0.000 0.000 0.000
ELEC.
TRANS.
      3.718 6.197 -38.616 12.472 20.786 150.306
ROT. 2.479 2.479 -9.834 8.314 8.314
                                                41.297
VIB. 17.322 17.322 17.322 0.001 0.001
       23.520
                25.999
                      -31.128
                               20.787 29.102
                                                 191.603
TOTAL
VIB. THERMAL CORRECTION E(T) - E(0) = H(T) - H(0) =
                                          0.030 J/MOL
```

• Now we can use this information to calculate reaction energies. First lets just start with 0 K

• Now we can use this information to calculate reaction energies. First lets just start with 0 K



- Now we can use this information to calculate reaction energies. First lets just start with 0 K
- $HCN \Longrightarrow TS^{\ddagger} \Longrightarrow HNC$



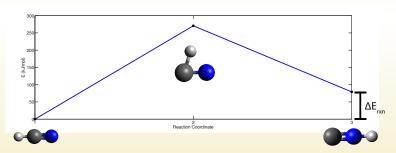
$$\bullet$$
 HCN  $\Longrightarrow$  TS<sup>‡</sup>  $\Longrightarrow$  HNC

$$\Delta E_{rxn} = \sum E_{products} - \sum E_{reactants}$$

•

 $\bullet$  HCN  $\Longrightarrow$  TS<sup>‡</sup>  $\Longrightarrow$  HNC

$$\Delta E_{rxn} = \sum E_{products} - \sum E_{reactants}$$



• We can do thermodynamic quantities

$$\Delta H_{rxn}^{\circ}(T) = \Delta E^{elec} + \Delta Z PV E + \Delta H^{\circ}(T)_{trans} + \Delta H^{\circ}(T)_{rot} + \Delta H^{\circ}(T)_{vib}$$

Here the  $\Delta$  is product(s) minus reactant(s).

•  $\Delta H_{rxn}^{\circ} = +77.65 \ kJ/mol$ , endothermic reaction

0

 $\bullet$  HCN  $\rightleftharpoons$  TS<sup>‡</sup>  $\rightleftharpoons$  HNC

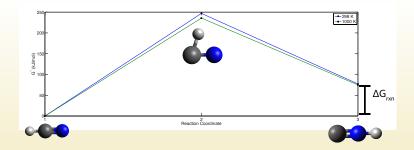
	E (kJ/mol)	ZPVE	$G_{tot}^{\circ}(298K)$	$G_{tot}^{\circ}(1000 \text{ K})$
HCN	-240693.2	50.0165	-0.76	-160.462
HNC	-240612.4	48.2648	-2.593	-164.758
TS	-240403.5	30.9667	-24.35	-195.324

 $\bullet$  HCN  $\Longrightarrow$  TS<sup>‡</sup>  $\Longrightarrow$  HNC

	E (kJ/mol)	ZPVE	$G_{tot}^{\circ}(298K)$	$G_{tot}^{\circ}(1000 \text{ K})$
HCN	-240693.2	50.0165	-0.76	-160.462
HNC	-240612.4	48.2648	-2.593	-164.758
TS	-240403.5	30.9667	-24.35	-195.324

• Here  $G_{tot}^{\circ}(T)$  is the Gibbs free energy contribution from translation, rotation, and vibration at the specified temperature and standard state.

 $\bullet$  HCN  $\Longrightarrow$  TS<sup>‡</sup>  $\Longrightarrow$  HNC



- $HCN \Longrightarrow TS^{\ddagger} \Longrightarrow HNC$
- Can calculate Equilibrium constant

$$K = \frac{q_{HNC}}{q_{HCN}} e^{-\Delta E_{rxn}/k_B T} = e^{-\Delta G_{rxn}^{\circ}(T)/k_B T}$$

## Gibbs Free Energy

- $\bullet$  HCN  $\rightleftharpoons$  TS<sup>‡</sup>  $\rightleftharpoons$  HNC
- Can calculate Equilibrium constant

$$K = \frac{q_{HNC}}{q_{HCN}} e^{-\Delta E_{rxn}/k_B T} = e^{-\Delta G_{rxn}^{\circ}(T)/k_B T}$$

• We did calculations at two Temperatures we could have also used the Van't Hoff Equation if the temperature range is small enough that  $\Delta H_{rxn}^{\circ} \approx \text{constant}$ 

•

$$\ln \frac{K(T_2)}{K(T_1)} = -\frac{\Delta H_{rxn}^{\circ}}{R} \left( \frac{1}{T_2} - \frac{1}{T_1} \right)$$

# Gibbs Free Energy

•  $HCN \Longrightarrow TS^{\ddagger} \Longrightarrow HNC$ 

• Can calculate Equilibrium constant

$$K = \frac{q_{HNC}}{q_{HCN}} e^{-\Delta E_{rxn}/k_B T} = e^{-\Delta G_{rxn}^{\circ}(T)/k_B T}$$

• We did calculations at two Temperatures we could have also used the Van't Hoff Equation if the temperature range is small enough that  $\Delta H_{rxn}^{\circ} \approx \text{constant}$ 

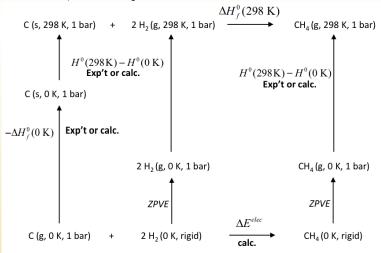
$$\ln \frac{K(T_2)}{K(T_1)} = -\frac{\Delta H_{rxn}^{\circ}}{R} \left( \frac{1}{T_2} - \frac{1}{T_1} \right)$$

• For arbitrary activities (or pressures assuming Raoult's Law)

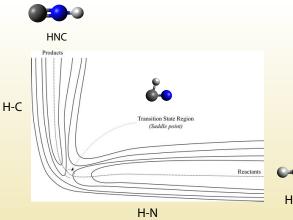
$$\Delta G_{rxn} = \Delta G_{rxn}^{\circ} + RT \ln \frac{a_{HNC}}{a_{HCN}}$$

## Thermodynamic Cycle

• We can use this approach to calculate all sorts of thermodynamic properties under all different kinds of conditions, for example the heat of formation of methane:

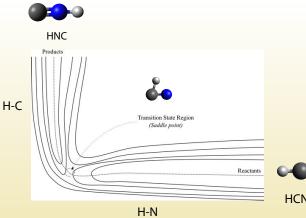


- Existence of a potential energy surface.
- 2 Existence of a dividing surface, point of no return.
- Existence of a critical TS point on that surface
- Quasi-equilibrium between reactants and
- (a) Harmonic PES near the TS

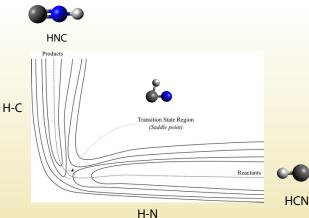




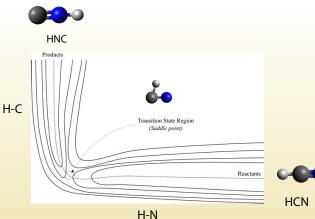
- Existence of a potential energy surface.
- Existence of a dividing surface, point of no return.



- Existence of a potential energy surface.
- Existence of a dividing surface, point of no return.
- Existence of a critical TS point on that surface.

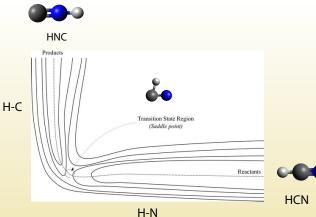


- Existence of a potential energy surface.
- Existence of a dividing surface, point of no return.
- Existence of a critical TS point on that surface.
- Quasi-equilibrium between reactants and TS



#### Transition state theory assumes

- Existence of a potential energy surface.
- 2 Existence of a dividing surface, point of no return.
- Sexistence of a critical TS point on that surface.
- Quasi-equilibrium between reactants and TS
- 6 Harmonic PES near the TS



PAOLUCCI

• Under these assumptions one can show ( $\Delta n$  is change in moles)

$$k_f(T) = \frac{k_B T}{h} \frac{q_{\ddagger}}{q_{reactants}} (P/RT)^{\Delta n} e^{-\Delta E_a/k_B T}$$

• Could group all the leading terms into A(T)

$$k_f(T) = A(T)e^{-\Delta E_a/k_B T}$$

- Looks familiar right?
- equivalently

$$k_f(T) = \frac{k_B T}{h} e^{\frac{-\Delta^{\ddagger} G^{\circ}(T)}{k_B T}}$$

$$E_{app} = RT^2 \frac{\partial \ln k}{\partial T}$$

• Under these assumptions one can show ( $\Delta n$  is change in moles)

$$k_f(T) = \frac{k_B T}{h} \frac{q_{\ddagger}}{q_{reactants}} (P/RT)^{\Delta n} e^{-\Delta E_a/k_B T}$$

• Could group all the leading terms into A(T)

$$k_f(T) = A(T)e^{-\Delta E_a/k_B T}$$

- Looks familiar right?
- equivalently

$$k_f(T) = \frac{k_B T}{h} e^{\frac{-\Delta^{\ddagger} G^{\circ}(T)}{k_B T}}$$

• Under these assumptions one can show ( $\Delta n$  is change in moles)

$$k_f(T) = \frac{k_B T}{h} \frac{q_{\ddagger}}{q_{reactants}} (P/RT)^{\Delta n} e^{-\Delta E_a/k_B T}$$

• Could group all the leading terms into A(T)

$$k_f(T) = A(T)e^{-\Delta E_a/k_B T}$$

- Looks familiar right?
- equivalently

$$k_f(T) = \frac{k_B T}{h} e^{\frac{-\Delta^{\ddagger} G^{\circ}(T)}{k_B T}}$$



• Under these assumptions one can show ( $\Delta n$  is change in moles)

$$k_f(T) = \frac{k_B T}{h} \frac{q_{\ddagger}}{q_{reactants}} (P/RT)^{\Delta n} e^{-\Delta E_a/k_B T}$$

• Could group all the leading terms into A(T)

$$k_f(T) = A(T)e^{-\Delta E_a/k_B T}$$

- Looks familiar right?
- equivalently

$$k_f(T) = \frac{k_B T}{h} e^{\frac{-\Delta^{\dagger} G^{\circ}(T)}{k_B T}}$$



• Under these assumptions one can show ( $\Delta n$  is change in moles)

$$k_f(T) = \frac{k_B T}{h} \frac{q_{\ddagger}}{q_{reactants}} (P/RT)^{\Delta n} e^{-\Delta E_a/k_B T}$$

• Could group all the leading terms into A(T)

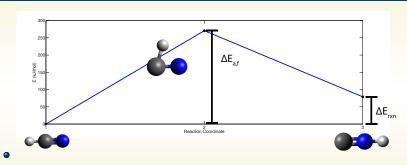
$$k_f(T) = A(T)e^{-\Delta E_a/k_BT}$$

- Looks familiar right?
- equivalently

$$k_f(T) = \frac{k_B T}{h} e^{\frac{-\Delta^{\ddagger} G^{\circ}(T)}{k_B T}}$$

$$E_{app} = RT^2 \frac{\partial \ln k}{\partial T}$$

## $\overline{Kinetics}$

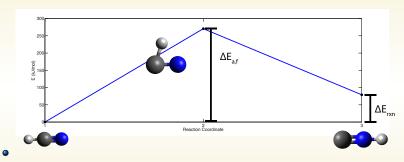


$$\bullet \ k_r = k_f/K$$

$$r = k_f C_{HCN} - \frac{k_f}{K} C_{HNC}$$

• Could stick in a mass balance, e.g unsteady CSTR

$$\frac{dC_{HCN}}{dt} = \frac{C_{HCN,0} - C_{HCN}}{\tau} - \tau$$

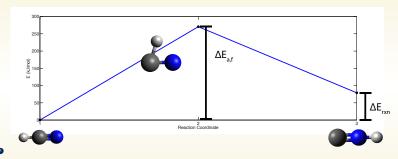


$$\bullet \ k_r = k_f/K$$

$$r = k_f C_{HCN} - \frac{k_f}{K} C_{HNC}$$

• Could stick in a mass balance, e.g unsteady CSTR

$$\frac{dC_{HCN}}{dt} = \frac{C_{HCN,0} - C_{HCN}}{\tau} - r$$

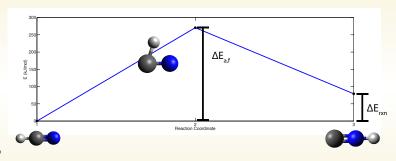


$$k_r = k_f/K$$

$$r = k_f C_{HCN} - \frac{k_f}{K} C_{HNC}$$

• Could stick in a mass balance, e.g unsteady CSTR

$$\frac{dC_{HCN}}{dt} = \frac{C_{HCN,0} - C_{HCN}}{\tau} -$$



• 
$$k_r = k_f/K$$

$$r = k_f C_{HCN} - \frac{k_f}{K} C_{HNC}$$

ullet Could stick in a mass balance, e.g unsteady CSTR

$$\frac{dC_{HCN}}{dt} = \frac{C_{HCN,0} - C_{HCN}}{\tau} - r$$