

The Linear Advection Equation and The Fourier Transform

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The Simplest Model For a Wave is:

$$\eta_t + c\eta_x = 0$$

$$\eta(x, 0) = \eta_0(x)$$

To solve it, we introduce the coordinate transformation

$$\xi = x - ct, \quad \chi = x + ct$$

$$\hookrightarrow \partial_t = -c\partial_\xi + c\partial_\chi$$

$$\partial_x = \partial_\xi + \partial_\chi$$

$$\hookrightarrow \eta_t + c\eta_x = 0$$

$$\hookrightarrow -c\eta_\xi + c\eta_\chi + c(\eta_\xi + \eta_\chi) = 0$$

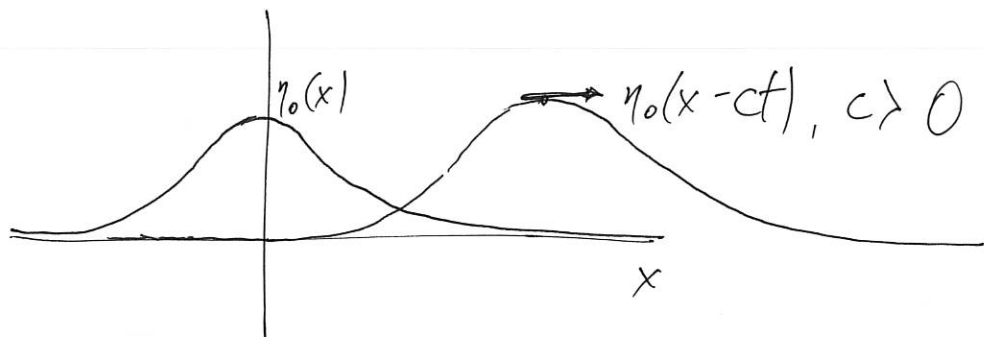
$$\hookrightarrow \mathcal{L}c\eta_\chi = 0$$

$$\hookrightarrow \eta = \eta(\xi) = \eta(x - ct)$$

if we match this at $t=0$ for $\eta_0(x) \rightarrow$

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$$\eta(x, t) = \eta_0(x - ct)$$



\hookrightarrow Initial condition propagates without change of shape.

A slightly more complicated model is

$$\eta_{tt} = c^2 \eta_{xx}$$

$$\eta(x, 0) = \eta_0(x)$$

$$\eta_t(x, 0) = v_0(x)$$

Again, let $\xi = x - ct$; $\chi = x + ct$

$$\hookrightarrow \partial_t^2 = (-c\partial_\xi + c\partial_\chi)^2 = c^2\partial_\xi^2 - 2c^2\partial_\xi\partial_\chi + c^2\partial_\chi^2$$

$$\partial_x^2 = (\partial_\xi + \partial_X)^2 = \partial_\xi^2 + 2\partial_\xi \partial_X + \partial_X^2$$

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$$\hookrightarrow 4c^2 \partial_{\xi X}^2 \eta = 0$$

$$\hookrightarrow \partial_X (\partial_\xi \eta) = 0$$

$$\hookrightarrow \partial_\xi \eta = f(\xi)$$

$$\hookrightarrow \eta = \int^\xi f(s) ds + g(X)$$

$$= \tilde{f}(\xi) + g(X)$$

$$\text{or: } \eta(x, t) = f(x - ct) + g(x + ct)$$

$$\hookrightarrow \eta_0 = f(x) + g(x)$$

$$v_0 = -cf'(x) + cg'(x)$$

$$\hookrightarrow \eta_0' = f' + g'$$

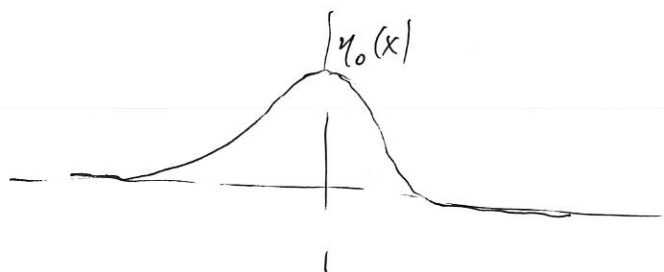
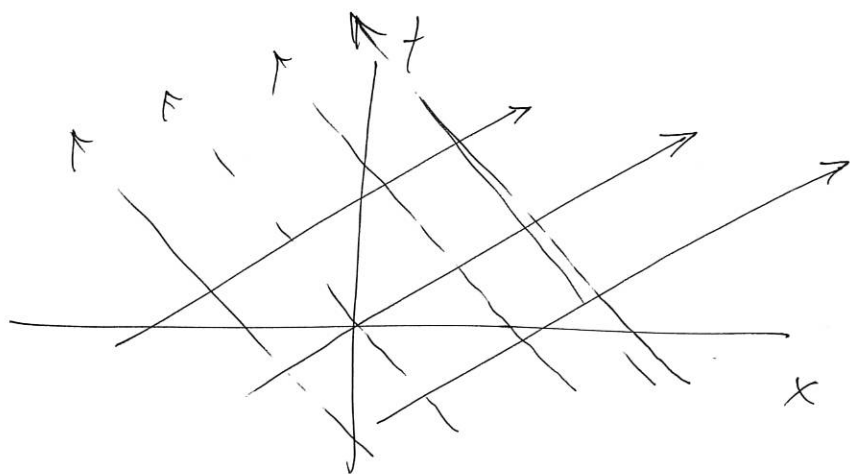
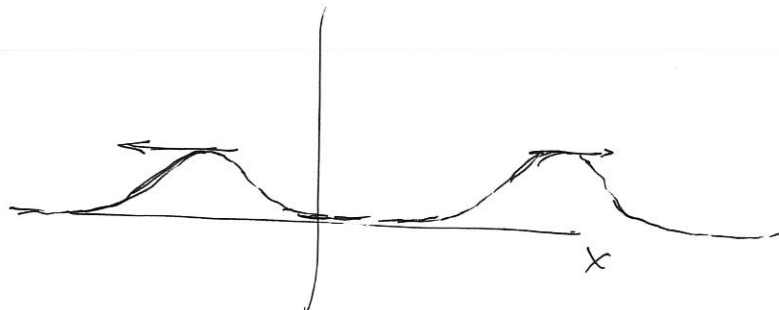
$$v_0 = -cf' + cg'$$

$$\rightarrow c\eta_0' + v_0 = 2cg'$$

$$c\eta_0' - v_0 = 2cf'$$

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$$u(x,t) = \frac{1}{2} u_0(x-ct) + \frac{1}{2} u_0(x+ct) + \frac{1}{2c} \int_{x-ct}^{x+ct} v_0(s) ds$$


 t


so now information propagates in both directions, but again, functions essentially just get transported about.

Now let's see all of this from a very different viewpoint.

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The Fourier Transform:

$$\hat{f}(k) = \int_{\mathbb{R}} e^{-ikx} f(x) dx$$

And its inverse

$$f(x) = \frac{1}{2\pi} \int_{\mathbb{R}} e^{ikx} \hat{f}(k) dk$$

The most important property of a Fourier-Transform is what it does to derivatives

$$\begin{aligned} (\partial_x f)^\wedge &= \int_{\mathbb{R}} e^{-ikx} \partial_x f dx = \cancel{e^{-ikx} f} \Big|_{-\infty}^{\infty} + ik \int_{-\infty}^{\infty} e^{-ikx} f(x) dx \\ &= ik \hat{f}(k) \end{aligned}$$

$$\text{i.e. } (\partial_x)^\wedge = ik \longrightarrow (\partial_x^n)^\wedge = (ik)^n$$

or derivatives become polynomials upon ~~differentiation~~ transformation.

So, let's go back to

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$$\eta_t + c\eta_x = 0$$

$$\eta(x, 0) = \eta_0(x)$$

We transform $\eta_t + c\eta_x = 0$

$$\hookrightarrow \hat{\eta}_t + ikc\hat{\eta} = 0$$

$$\hookrightarrow \hat{\eta}_t = -ikc\hat{\eta}$$

$$\hookrightarrow \hat{\eta}(k, t) = \hat{\eta}(k, 0) e^{-ikct}$$

$$\text{or } \hat{\eta}(k, t) = \hat{\eta}_0(k) e^{-ikct}$$

$$\hookrightarrow \eta(x, t) = \frac{1}{2\pi} \int_{\mathbb{R}} \hat{\eta}(k, t) e^{ikx} dk$$

$$= \frac{1}{2\pi} \int_{\mathbb{R}} \hat{\eta}_0(k) e^{ik(x-ct)} dk = \eta_0(x-ct).$$

So, same answer, but now we have a very different understanding of what is going on.