

Homework 7

Quantum Mechanics

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Problem 1. 5.1

Solution.

We are concerned here with the new ground state ket $|0\rangle$ and the new ground state energy shift Δ_0 in the presence of perturbation $V = bx$.

$$|0\rangle = |0^0\rangle + \sum_{j \neq 0} |j^0\rangle \frac{V_{j0}}{E_0^0 - E_j^0} + \dots$$

$$\Delta_0 = V_{00} + \sum_{j \neq 0} \frac{|V_{j0}|^2}{E_0^0 - E_j^0} + \dots$$

$$V_{nk} = b \langle i^0 | x | j^0 \rangle = b \sqrt{\frac{\hbar}{2m\omega}} \left(\sqrt{j} \delta_{i,j-1} + \sqrt{j+1} \delta_{i,j+1} \right)$$

The lowest nonvanishing order is then V_{01} . Therefore

$$\Delta_0 = -\frac{b^2 \hbar}{2m\omega} \frac{1}{\hbar\omega} = -\frac{b^2}{2m\omega^2}$$

To solve it exactly, notice that the potential is of the form

$$V_1(x) = ax^2 + bx$$

The new potential shifts to the left by $b/2$, has a new minimum at $-b/2a$. So it is really just the original problem, we just have to make a change of coordinates and shift the equilibrium point down by $-b/2m\omega^2$. Therefore,

$$\Delta = -\frac{b}{2a} = -\frac{b}{2m\omega^2}$$

which is exactly what we got with perturbation theory. ■

Problem 2. 5.2

Solution.

In general, the first order shift in the energy levels i is

$$\Delta_i = V_{ii} = \langle i^0 | V | i^0 \rangle$$

Furthermore, the perturbation Hamiltonian is

$$V = \frac{V_0 x}{L}$$

So we just need to identify the matrix elements along the diagonal of this matrix:

$$\begin{aligned} V_{ii} &= \langle i^0 | V | i^0 \rangle = \frac{V_0}{L} \langle i^0 | x | i^0 \rangle \\ &= \frac{V_0}{L} \frac{2}{L} \int_0^L x \sin^2 \left(\frac{n\pi x}{L} \right) dx \\ &= \frac{V_0}{L} \end{aligned}$$
■

Problem 3. 5.5

Solution.

Up to order λ^2 , we have

$$\begin{aligned} |i\rangle &= |i^0\rangle + \lambda \sum_{j \neq i} \frac{V_{ij}}{E_i^0 - E_j^0} |j^0\rangle \\ &+ \lambda^2 \left(\sum_{j \neq i} \sum_{l \neq i} \frac{V_{jl} V_{li} |j^0\rangle}{(E_i^0 - E_j^0)(E_i^0 - E_l^0)} - \sum_{j \neq i} \frac{V_{ii} V_{ji} |j^0\rangle}{(E_i^0 - E_j^0)^2} \right) \end{aligned}$$

Now recall that we chose the normalization $\langle i^0|i \rangle = 1$ (which can be seen from the above equation), so what we have to calculate is $|\langle i^0|i \rangle|^2/|\langle i|i \rangle|^2$

$$\begin{aligned}\langle i|i \rangle &= 1 + \lambda^2 \left(\sum_{j \neq i} \frac{V_{ij}^*}{E_i^0 - E_j^0} \langle j^0| \right) \left(\sum_{j \neq i} \frac{V_{ij}}{E_i^0 - E_j^0} |j^0 \rangle \right) \\ &= 1 + \lambda^2 \sum_{j \neq i} \frac{V_{ij}^* V_{ij}}{(E_i^0 - E_j^0)^2}\end{aligned}$$

So the probability is just

$$|\langle i^0|i \rangle|^2/|\langle i|i \rangle|^2 = \left(1 + 2\lambda^2 \sum_{j \neq i} \frac{V_{ij}^* V_{ij}}{(E_i^0 - E_j^0)^2} \right)^{-1} + \mathcal{O}(\lambda^3)$$

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Problem 4. 5.7

Solution.

We can write this Hamiltonian as

$$H_0 = H_x + H_y$$

Given some state $|n, m\rangle$, where H_x acts on $|n\rangle$ and H_y acts on $|m\rangle$.

$$\begin{aligned}H_0 |n, m\rangle &= (H_x + H_y) |n, m\rangle \\ &= (E_x^n + E_y^m) |n, m\rangle \\ &= \hbar\omega \left(n + \frac{1}{2} + m + \frac{1}{2} \right) \\ &= \hbar\omega(n + m + 1)\end{aligned}$$

So the energies of the three lowest states are $\hbar\omega, 2\hbar\omega, 3\hbar\omega$. There is a degeneracy - two unique $|n, m\rangle$ have the same eigenvalue for the first excited state. For example $|0, 1\rangle$ and $|1, 0\rangle$ have the same energy. Now, we are given the perturbation

$$V = \delta m \omega^2 xy$$

and we need to find a general matrix representation in the unperturbed $|n, m\rangle$ basis. I will use a shorthand:

$$\begin{aligned} V_{ij} &= \delta m \omega^2 \langle i | xy | j \rangle \\ &= \delta m \omega^2 \langle i_x | x | j_x \rangle \langle i_y | y | j_y \rangle \\ &= \delta \frac{\hbar \omega}{2} \left(\sqrt{j_x} \delta_{i_x, j_x-1} + \sqrt{j_x+1} \delta_{i_x, j_x+1} \right) \left(\sqrt{j_y} \delta_{i_y, j_y-1} + \sqrt{j_y+1} \delta_{i_y, j_y+1} \right) \end{aligned}$$

First, we consider the ground state. To first order,

$$\Delta_i = V_{ii} = 0$$

This can be seen from the above equation, where the matrix elements of V_{ij} are always zero when $i = j$. For the first excited state, there is a degeneracy, so we need to diagonalize V in the $|1, 0\rangle, |0, 1\rangle$ subspace. The matrix is

$$\begin{aligned} V' &= \begin{pmatrix} \langle 0, 1 | V | 0, 1 \rangle & \langle 0, 1 | V | 1, 0 \rangle \\ \langle 1, 0 | V | 0, 1 \rangle & \langle 1, 0 | V | 1, 0 \rangle \end{pmatrix} \\ &= \delta \frac{\hbar \omega}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \end{aligned}$$

which is a familiar matrix, so we can immediately write

$$\Delta_1^1 = \pm \delta \frac{\hbar \omega}{2}$$

and the eigenvectors are $\frac{1}{\sqrt{2}} (|0, 1\rangle \pm |1, 0\rangle)$

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Problem 5. 5.12a

Solution.

We need to find the eigenvectors of V in the degenerate subspace. If they are all independent, then the three-fold degeneracy is lifted. I will suppress the quantum numbers n and m_s for brevity

$$V = \begin{pmatrix} \langle 1, -1 | V | 1, -1 \rangle & \langle 1, -1 | V | 1, 0 \rangle & \langle 1, -1 | V | 1, 1 \rangle \\ \langle 1, 0 | V | 1, -1 \rangle & \langle 1, 0 | V | 1, 0 \rangle & \langle 1, 0 | V | 1, 1 \rangle \\ \langle 1, 1 | V | 1, -1 \rangle & \langle 1, 1 | V | 1, 0 \rangle & \langle 1, 1 | V | 1, 1 \rangle \end{pmatrix}$$

It is immediately obvious that the elements along the diagonal must be zero, by symmetry. Also, the matrix is Hermitian (which can be verified by inspecting the spherical harmonics $Y_l^m(\theta, \phi)$). Therefore, we just need to find $\langle 1, -1 | V | 1, 0 \rangle$, $\langle 1, -1 | V | 1, 1 \rangle$ and $\langle 1, 0 | V | 1, 1 \rangle$. Write,

$$\begin{aligned}\langle 1, -1 | x^2 | 1, 0 \rangle &= \int r^2 \sin^2 \theta \cos^2 \phi Y_1^{-1}(\theta, \phi) Y_1^0(\theta, \phi) r dr d\theta d\phi \\ &\propto \int r^2 \sin^2 \theta \cos^2 \phi \sin \theta e^{i\phi} \cos \theta r dr d\theta d\phi = 0\end{aligned}$$

due to the θ integral. It is similar for $\langle y^2 \rangle$ - the integral over θ is again zero and $\langle 1, -1 | V | 1, 0 \rangle = 0$. Moving on,

$$\begin{aligned}\langle 1, 0 | x^2 | 1, 1 \rangle &= \int r^2 \sin^2 \theta \cos^2 \phi Y_1^0(\theta, \phi) Y_1^1(\theta, \phi) r dr d\theta d\phi \\ &\propto \int r^2 \sin^2 \theta \cos^2 \phi \sin \theta e^{-i\phi} \cos \theta r dr d\theta d\phi = 0\end{aligned}$$

So it is basically the same as before and $\langle 1, 0 | V | 1, 1 \rangle = 0$. Finally,

$$\begin{aligned}\langle 1, -1 | x^2 | 1, 1 \rangle &= \int r^2 \sin^2 \theta \cos^2 \phi Y_1^{-1}(\theta, \phi) Y_1^1(\theta, \phi) r dr d\theta d\phi \\ &\propto \int r^3 \sin^4 \theta \cos^2 \phi dr d\theta d\phi\end{aligned}$$

$$\begin{aligned}\langle 1, -1 | y^2 | 1, 1 \rangle &= \int r^2 \sin^2 \theta \cos^2 \phi Y_1^{-1}(\theta, \phi) Y_1^1(\theta, \phi) r dr d\theta d\phi \\ &\propto \int r^3 \sin^2 \theta \cos^4 \phi dr d\theta d\phi\end{aligned}$$

We are not asked to work this integral out in detail, but it is nonzero and we will ultimately get a matrix like

$$V = \begin{pmatrix} 0 & 0 & \alpha \\ 0 & 0 & 0 \\ \alpha & 0 & 0 \end{pmatrix}$$

since α is real. Finally,

$$\det(V - \lambda I) = \det \begin{pmatrix} -\lambda & 0 & \alpha \\ 0 & -\lambda & 0 \\ \alpha & 0 & -\lambda \end{pmatrix} = -\lambda^3 + \lambda\alpha^2 = 0$$

which has three solutions $\lambda = 0, \pm\alpha$. So the degeneracy is broken. ■

Problem 6. 5.24

Solution.

We have seen before that we can either express the state in the basis of $\mathbf{J}^2, \mathbf{L}^2, \mathbf{S}^2, J_z$ or $\mathbf{L}^2, \mathbf{S}^2, L_z, S_z$. In the former, our perturbation Hamiltonian V reads

$$V = \frac{A}{2\hbar^2} (\mathbf{J}^2 - \mathbf{L}^2 - \mathbf{S}^2) + \frac{B}{\hbar} (J_z + S_z)$$

The only hiccup of this representation is that S_z is not diagonal in the $|lsjm\rangle$ basis, so we will have to specify those matrix elements explicitly. For an arbitrary state $|lsjm\rangle$

$$\begin{aligned} \langle lsjm | V | lsjm \rangle &= \langle lsjm | \frac{A}{2\hbar^2} (\mathbf{J}^2 - \mathbf{L}^2 - \mathbf{S}^2) + \frac{B}{\hbar} (J_z + S_z) | lsjm \rangle \\ &= \frac{A}{2} (j(j+1) - l(l+1) - s(s+1)) + Bm + B \langle lsjm | S_z | lsjm \rangle \end{aligned}$$

We can write these states alternatively as

$$\begin{aligned} \left| j = l \pm \frac{1}{2}, m \right\rangle &= \pm \sqrt{\frac{l \pm m + \frac{1}{2}}{2l+1}} \left| m_l = m - \frac{1}{2}, m_s = \frac{1}{2} \right\rangle \\ &\quad + \sqrt{\frac{l \mp m + \frac{1}{2}}{2l+1}} \left| m_l = m + \frac{1}{2}, m_s = -\frac{1}{2} \right\rangle \end{aligned}$$

which means that

$$\begin{aligned}
\left\langle j = l \pm \frac{1}{2}, m \left| S_z \right| j = l \pm \frac{1}{2}, m \right\rangle &= \frac{\hbar}{2} (|c_+|^2 - |c_-|^2) \\
&= \frac{\hbar}{2} \frac{1}{2l+1} \left(\left(l \pm m + \frac{1}{2} \right) - \left(l \mp m + \frac{1}{2} \right) \right) \\
&= \pm \frac{m\hbar}{2l+1}
\end{aligned}$$

So we can write down the general matrix element as

$$\begin{aligned}
\langle l s j m | V | l s j m \rangle &= \langle l s j m | \frac{A}{2\hbar^2} (\mathbf{J}^2 - \mathbf{L}^2 - \mathbf{S}^2) + \frac{B}{\hbar} (J_z + S_z) | l s j m \rangle \\
&= \frac{A}{2} (j(j+1) - l(l+1) - s(s+1)) + Bm \pm B \frac{m}{2l+1} \\
&= \frac{A}{2} (j(j+1) - l(l+1) - s(s+1)) + Bm \left(1 \pm \frac{1}{2l+1} \right)
\end{aligned}$$

Now, let's write out the 8 states. We have two states for $l = 0$ of form $|l s j m\rangle$

$$\left| 0, \frac{1}{2}, \frac{1}{2}, \frac{1}{2} \right\rangle, \left| 0, \frac{1}{2}, \frac{1}{2}, -\frac{1}{2} \right\rangle$$

There are also six states for $l = 1$

$$\left| 1, \frac{1}{2}, \frac{3}{2}, \frac{3}{2} \right\rangle, \left| 1, \frac{1}{2}, \frac{3}{2}, -\frac{3}{2} \right\rangle, \left| 1, \frac{1}{2}, \frac{3}{2}, \frac{1}{2} \right\rangle, \left| 1, \frac{1}{2}, \frac{3}{2}, -\frac{1}{2} \right\rangle, \left| 1, \frac{1}{2}, \frac{3}{2}, \frac{1}{2} \right\rangle, \left| 1, \frac{1}{2}, \frac{3}{2}, -\frac{1}{2} \right\rangle$$

We are given the hint that two matrices will not be diagonal. We can spot Suppose we group together the $l = 0$ states into a 2x2 matrix

$$V_1 = \begin{pmatrix} \frac{3A}{4} + B & 0 \\ 0 & \frac{3A}{4} \end{pmatrix}$$

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