# Homework 7

#### **Quantum Mechanics**

November 3, 2022

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# **Problem 1.** 5.1

### Solution.

We are concerned here with the new ground state ket  $|0\rangle$  in the presence of  $H_1$  and the new ground state energy shift  $\Delta_0$ .

$$|0\rangle = |0^{0}\rangle + \sum_{k\neq 0} |k^{0}\rangle \frac{V_{k0}}{E_{0}^{0} - E_{k}^{0}} + \dots$$

$$\Delta_0 = V_{00} + \sum_{k \neq 0} \frac{|V_{k0}|^2}{E_0^0 - E_k^0} + \dots$$

$$V_{nk} = b \left\langle n^0 \right| x \left| k^0 \right\rangle = b \sqrt{\frac{\hbar}{2m\omega}} \left( \sqrt{k} \delta_{n,k-1} + \sqrt{k+1} \delta_{n,k+1} \right)$$

The lowest nonvanishing order is then  $V_{01}$ . Therefore

$$\Delta_0 = -\frac{b^2 \hbar}{2m\omega} \frac{1}{\hbar \omega} = -\frac{b^2}{2m\omega^2}$$

To solve it exactly, notice that the potential is of the form

$$V_1(x) = ax^2 + bx$$

The new potential shifts to the left by b/2, has a new minimum at -b/2a, and the gradient has changed:

$$V'(x) = 2ax \rightarrow V'_1(x) = 2ax + b$$

This change in the gradient will not change the energy differences w.r.t. the original problem (why?) so we have really just shifted the equilibrium point down by  $-b/2m\omega^2$ .

$$\Delta = -\frac{b}{2a} = -\frac{b}{2m\omega^2}$$

which is exactly what we got with perturbation theory.

## **Problem 2.** 5.2

**Solution**. The perturbation Hamiltonian is

$$H_1 = \frac{Vx}{L}$$

$$V_{nk} = \langle n^0 | H_1 | k^0 \rangle = \frac{V}{L} \langle n^0 | x | k^0 \rangle$$
$$= \frac{V}{L} \int_0^L x \sin\left(\frac{n\pi x}{L}\right) \sin\left(\frac{k\pi x}{L}\right) dx$$

$$\Delta_0 = V_{00} + \sum_{k \neq 0} \frac{|V_{k0}|^2}{E_0^0 - E_k^0} + \dots$$

### **Problem 3.** 5.5

**Solution**. What we are asked to find is  $|\langle n^0|n\rangle|^2/|\langle n^0|n\rangle|^2$ . Up to order  $\lambda^2$ , we have

$$|n\rangle = |n^{0}\rangle + \lambda \sum_{k \neq n} |k^{0}\rangle \frac{V_{kn}}{E_{n}^{0} - E_{k}^{0}} + \lambda^{2} \left(\sum_{k \neq n} \sum_{k \neq l} \frac{V_{kl}V_{ln} |k^{0}\rangle}{(E_{n}^{0} - E_{k}^{0})(E_{n}^{0} - E_{l}^{0})} - \sum_{k \neq n} \frac{V_{nn}V_{kn} |k^{0}\rangle}{(E_{n}^{0} - E_{k}^{0})^{2}}\right)$$

Therefore,

$$\langle n^0 | n \rangle = 1 + \lambda^2 \sum_{k \neq n} \sum_{k \neq l} \frac{V_{kl} V_{ln}}{(E_n^0 - E_k^0)(E_n^0 - E_l^0)} \delta_{nk}$$

# **Problem 4.** 5.7

#### Solution.

We can write this Hamiltonian as

$$H_0 = H_x + H_y$$

Given some state  $|n, m\rangle$ , where  $H_x$  acts on  $|n\rangle$  and  $H_y$  acts on  $|m\rangle$ .

$$H_0 |n, m\rangle = (H_x + H_y) |n, m\rangle$$

$$= (E_x^n + E_y^m) |n, m\rangle$$

$$= \hbar\omega \left(n + \frac{1}{2} + m + \frac{1}{2}\right)$$

$$= \hbar\omega (n + m + 1)$$

So the energies of the three lowest states are  $\hbar\omega$ ,  $2\hbar\omega$ ,  $3\hbar\omega$ . There is a degeneracy - two unique  $|n,m\rangle$  have the same eigenvalue for the first excited state. For example  $|0,1\rangle$  and  $|1,0\rangle$  have the same energy. Now, the perturbation Hamiltonian is

$$H_0 = H_x + H_y + \delta m\omega^2 xy$$

First, to zeroth order, the energy eigenkets will not change because the diagonal of V will always be zero. This should be clear from the following calculations. For the ground state, recall

$$V_{nk} = \langle n^0 | x | k^0 \rangle = \sqrt{\frac{\hbar}{2m\omega}} \left( \sqrt{k} \delta_{n,k-1} + \sqrt{k+1} \delta_{n,k+1} \right)$$

Therefore,

$$V_{0k} = \delta m\omega^{2} \langle 0, 0 | xy | n, m \rangle$$
$$= \frac{\hbar}{2m\omega} \delta m\omega^{2} \langle n, m | 1, 1 \rangle$$

and  $\langle 0,0|1,1\rangle = 0$  so  $V_{00}$  is zero (I have just used that x can be written in terms of  $a_x$  and  $a_x^{\dagger}$  and the same for y). To first order,  $V_{(n,m),0}$  must all be zero due to orthogonality. So there is no first order shift for the ground state. For the first excited state,

$$V_{1k} = \delta m\omega^2 \langle 0, 1 | xy | n, m \rangle$$
$$= \frac{\hbar}{2m\omega} \delta m\omega^2 \sqrt{2}$$

$$\Delta_1 = \frac{V_{1k}^2}{2\hbar\omega} + \frac{V_{1k}^2}{\hbar\omega}$$

**Problem 5.** 5.12a

Solution.

**Problem 6.** 5.24

Solution.