# Bell's Inequality

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### CHSH and Tsirelson's Inequalities

Alice: Q, R Bob: S, T

Classical observables distributed according to P(Q, R, S, T). Combination of correlations between Alice and Bobs measurements are bounded according to the CHSH inequality

$$|E(QS) + E(RS) + E(RT) - E(QT)| \le 2$$

For the quantum version, define 4 spin operators along arbitrary directions  $Q = \vec{q} \cdot \sigma, R = \vec{r} \cdot \sigma, S = \vec{s} \cdot \sigma, T = \vec{t} \cdot \sigma.$ 

The book uses 
$$\vec{q} = (0,0,1), \vec{r} = (1,0,0), \vec{s} = (-\frac{1}{\sqrt{2}},0,-\frac{1}{\sqrt{2}}), \vec{t} = (-\frac{1}{\sqrt{2}},0,\frac{1}{\sqrt{2}})$$

$$|\langle Q \otimes S \rangle + \langle R \otimes S \rangle + \langle R \otimes T \rangle - \langle Q \otimes T \rangle| \le 2\sqrt{2}$$

## Assumptions and objectives

As in the book, Assume  $\vec{q}, \vec{r}$  and  $\vec{s}, \vec{t}$  are orthogonal and all  $\vec{q}, \vec{r}, \vec{s}, \vec{t}$  live in a plane (say x-z)

Now draw a set of pure states  $|\psi\rangle$  possibly with different degrees of entanglement

Question: How does the Tsirelson bound vary across states with different degrees of entanglement?

#### The Tsirelson bound

Solution to Problem 2.3 in the book:

$$(Q \otimes S + R \otimes S + R \otimes T - Q \otimes T)^2 = 4I + [Q, R] \otimes [S, T]$$

$$\langle (Q \otimes S + R \otimes S + R \otimes T - Q \otimes T) \rangle^{2} \leq \langle (Q \otimes S + R \otimes S + R \otimes T - Q \otimes T)^{2} \rangle$$

$$= \langle 4I + [Q, R] \otimes [S, T] \rangle$$

$$= 4 + \langle [Q, R] \otimes [S, T] \rangle$$

for fixed Q, R, S, T.

#### The Tsirelson bound

Using that 
$$(\sigma \cdot a)(\sigma \cdot b) = a \cdot b + i\sigma \cdot (a \times b)$$

$$[A, B] = i\sigma \cdot (\vec{a} \times \vec{b} - \vec{b} \times \vec{a}) = 2i\sigma \cdot (\vec{a} \times \vec{b})$$

Let  $\vec{n} = \vec{q} \times \vec{r}$  and  $\vec{m} = \vec{s} \times \vec{t}$ . Per our constraints, we could have  $\vec{n} = \vec{m} = \hat{y}$ .

$$\langle [Q, R] \otimes [S, T] \rangle = -4 \langle \psi | \sigma \cdot \vec{n} \otimes \sigma \cdot \vec{m} | \psi \rangle$$

$$= -4 \langle \psi | \sigma \cdot \hat{z} \otimes \sigma \cdot \hat{z} | \psi \rangle$$

$$= -4 \langle \psi | \sigma_{1v} \otimes \sigma_{2v} | \psi \rangle$$

$$\langle (Q \otimes S + R \otimes S + R \otimes T - Q \otimes T) \rangle \leq 2\sqrt{1 - \langle \psi | \sigma_{1y} \otimes \sigma_{2y} | \psi \rangle}$$

# Quantifying entanglement: partial traces

$$\begin{aligned} \operatorname{Tr}_{A}(\rho_{AB}) &= \sum_{ijkl} \rho_{ij}^{kl} \operatorname{Tr}_{A}(|i\rangle \langle k|) \otimes |j\rangle \langle l| \\ &= \sum_{i} \left( \sum_{jl} \rho_{ij}^{il} |j\rangle \langle l| \right) \\ &= (\rho_{00}^{00} + \rho_{10}^{10}) |0\rangle \langle 0| + (\rho_{00}^{01} + \rho_{10}^{11}) |0\rangle \langle 1| + (\rho_{01}^{00} + \rho_{11}^{10}) |1\rangle \langle 0| + (\rho_{01}^{01} + \rho_{11}^{11}) |1\rangle \langle 1| \end{aligned}$$

$$\begin{aligned} \operatorname{Tr}_{B}(\rho_{AB}) &= \sum_{ijkl} \rho_{ij}^{kl} |i\rangle \langle k| \otimes \operatorname{Tr}_{B}(|j\rangle \langle l|) \\ &= \sum_{j} \left( \sum_{ik} \rho_{ij}^{kj} |i\rangle \langle k| \right) \\ &= (\rho_{00}^{00} + \rho_{01}^{01}) |0\rangle \langle 0| + (\rho_{00}^{10} + \rho_{01}^{11}) |0\rangle \langle 1| + (\rho_{10}^{00} + \rho_{11}^{01}) |1\rangle \langle 0| + (\rho_{10}^{10} + \rho_{11}^{11}) |1\rangle \langle 1| \end{aligned}$$

## Reduced density matrices for an arbitrary state

$$\operatorname{Tr}_{A}(\rho_{AB}) = \begin{pmatrix} \rho_{00}^{00} + \rho_{10}^{10} & \rho_{00}^{01} + \rho_{10}^{11} \\ \rho_{00}^{00} + \rho_{11}^{10} & \rho_{01}^{01} + \rho_{11}^{11} \end{pmatrix} = \begin{pmatrix} |\alpha|^{2} + |\gamma|^{2} & \alpha\beta^{*} + \gamma\delta^{*} \\ \beta\alpha^{*} + \delta\gamma^{*} & |\beta|^{2} + |\delta|^{2} \end{pmatrix}$$

$$\operatorname{Tr}_{\mathcal{B}}(\rho_{AB}) = \begin{pmatrix} \rho_{00}^{00} + \rho_{01}^{01} & \rho_{00}^{10} + \rho_{01}^{11} \\ \rho_{10}^{00} + \rho_{11}^{01} & \rho_{10}^{10} + \rho_{11}^{11} \end{pmatrix} = \begin{pmatrix} |\alpha|^2 + |\beta|^2 & \alpha\gamma^* + \beta\delta^* \\ \gamma\alpha^* + \delta\beta^* & |\gamma|^2 + |\delta|^2 \end{pmatrix}$$

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### Definition of entanglement entropy

The entanglement entropy of a bipartite system is the Von Neumann entropy of either reduced density matrix (it doesn't matter which one we choose)

$$S(\rho_A) = -\text{Tr}(\rho_A \log \rho_A) = -\sum_{x} \lambda_x \log \lambda_x$$

for eigenvalues  $\lambda_x$  of  $\rho_A$ . This tells us: do the reduced states  $\rho_A$  and  $\rho_B$  contain all the information in  $\rho_{AB}$ ?

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## Entanglement entropy and Tsirelson's bound

Sanity check:

Draw random reals  $a, b, c, d, e, f, g, h \sim U([0, 1]^8)$ 

$$\mathsf{Construct}\,\left|\psi\right\rangle = \left(\mathsf{a}+\mathsf{i}\mathsf{b}\right)\left|\mathsf{00}\right\rangle + \left(\mathsf{c}+\mathsf{i}\mathsf{d}\right)\left|\mathsf{01}\right\rangle + \left(\mathsf{e}+\mathsf{i}\mathsf{f}\right)\left|\mathsf{10}\right\rangle + \left(\mathsf{g}+\mathsf{i}\mathsf{h}\right)\left|\mathsf{11}\right\rangle$$

Normalize 
$$|\psi\rangle o \frac{|\psi\rangle}{\sum_n |c_n|^2}$$

Compute  $S(\rho_A)$  and  $2\sqrt{1-\langle\psi|\,\sigma_{1y}\otimes\sigma_{2y}\,|\psi\rangle}$ 

## Entanglement entropy and Tsirelson's bound

