

Homework 1

Quantum Mechanics

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Problem 1. For the spin $1/2$ state $|+\rangle_x$, evaluate both sides of the inequality

$$\langle(\Delta A)^2\rangle\langle(\Delta B)^2\rangle \geq \frac{1}{4}|\langle[A, B]\rangle|^2$$

for the operators $A = S_x$ and $B = S_y$, and show that the inequality is satisfied. Repeat for the operators $A = S_z$ and $B = S_y$

Solution.

Let $A = S_x$ and $B = S_y$. The variance $\langle(\Delta S_x)^2\rangle$ in state $|+\rangle_x$ must be zero since $|+\rangle_x$ is an eigenvector of S_x

$$\langle(\Delta S_x)^2\rangle = \langle S_x^2\rangle - \langle S_x\rangle^2 = 0$$

Therefore, the LHS of the above inequality is zero. The commutator $[S_x, S_y] = i\hbar S_z$ and

$$\langle S_z\rangle = \langle + |_x S_z | + \rangle_x = 0$$

Clearly the inequality is satisfied since both sides are zero. Now let $A = S_z$ and $B = S_y$. Since the state is prepared in $|+\rangle_x$, the variances $\langle(\Delta S_x)^2\rangle$ and $\langle(\Delta S_y)^2\rangle$ must be $1/4$ (this is just a fair coin toss).

The commutator $[S_z, S_y] = -i\hbar S_x$ and $\langle S_x\rangle = \frac{\hbar}{2}$. The inequality then reads

$$\frac{1}{16} \geq \frac{\hbar^2}{16}$$

which is satisfied given that $\hbar \approx 10^{-34} \text{ J} \cdot \text{s}$



Problem 2. Suppose a 2×2 matrix X (not necessarily Hermitian, nor unitary) is written as

Solution.

$$\begin{aligned}\mathrm{Tr}(X) &= \mathrm{Tr}(a_0) + \mathrm{Tr}\left(\sum_k a_k \sigma_k\right) \\ &= 2a_0\end{aligned}$$

$$\begin{aligned}\mathrm{Tr}(\sigma_k X) &= \mathrm{Tr}\left(\sigma_k a_0 + \sigma_k \sum_j a_j \sigma_j\right) \\ &= \mathrm{Tr}\left(\sigma_k a_0 + \sum_j a_j \sigma_k \sigma_j\right) \\ &= \mathrm{Tr}\left(\sum_j a_j \sigma_k \sigma_j\right)\end{aligned}$$

We can write out the equation $X = a_0 + \sigma \cdot a$ explicitly

$$X = \begin{pmatrix} a_0 + a_3 & a_1 - ia_3 \\ a_1 + ia_2 & a_0 - a_3 \end{pmatrix}$$

Thus we have four equations involving X_{ij} 's and a_k for $k = (1, 2, 3)$. We can manipulate those four equations to show that

$$\begin{aligned}a_0 &= \frac{X_{11} + X_{22}}{2} \\ a_1 &= \frac{X_{12} + X_{21}}{2} \\ a_2 &= \frac{X_{21} - X_{12}}{2} \\ a_3 &= \frac{X_{11} - X_{22}}{2}\end{aligned}$$

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Problem 3.

Solution.

$$\sigma \cdot a' = \exp\left(\frac{i\sigma \cdot \hat{n}\phi}{2}\right) \sigma \cdot a \exp\left(-\frac{i\sigma \cdot \hat{n}\phi}{2}\right)$$

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Problem 4.

Solution.

$$A(|i\rangle + |j\rangle) = i|i\rangle + j|j\rangle$$

If we have degenerate eigenvalues i.e., $i = j$ then

$$A(|i\rangle + |j\rangle) = i(|i\rangle + |j\rangle)$$

and $|i\rangle + |j\rangle$ is also an eigenvector of A

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Problem 5.

Solution. We will make use of the following representations of the spin operators

$$\begin{aligned} S_x &= \frac{\hbar}{2} (|+\rangle \langle -| + |- \rangle \langle +|) \\ S_y &= \frac{i\hbar}{2} (-|+\rangle \langle -| + |- \rangle \langle +|) \\ S_z &= \frac{\hbar}{2} (|+\rangle \langle +| - |- \rangle \langle -|) \end{aligned}$$

$$\begin{aligned} [S_x, S_y] &= \frac{i\hbar^2}{2} (|+\rangle \langle -| + |- \rangle \langle +|) (-|+\rangle \langle -| + |- \rangle \langle +|) \\ &\quad - (-|+\rangle \langle -| + |- \rangle \langle +|) (|+\rangle \langle -| + |- \rangle \langle +|) \\ &= \frac{i\hbar^2}{2} (|+\rangle \langle +| - |- \rangle \langle -|) \\ &= i\hbar S_z \end{aligned}$$

Flipping the order of the commutator always flips the sign of the result i.e. $[S_i, S_j] = -[S_j, S_i]$. Thus for $[S_y, S_x]$ we would get $-i\hbar S_z$.

$$\begin{aligned}
[S_y, S_z] &= \frac{i\hbar^2}{4} (-|+\rangle \langle -| + |- \rangle \langle +|) (|+\rangle \langle +| - |- \rangle \langle -|) \\
&\quad - (|+\rangle \langle +| - |- \rangle \langle -|) (|+\rangle \langle -| + |- \rangle \langle +|) \\
&= \frac{i\hbar^2}{4} (|+\rangle \langle -| + |- \rangle \langle +|) \\
&= i\hbar S_x
\end{aligned}$$

$$\begin{aligned}
[S_z, S_x] &= \frac{\hbar^2}{4} (|+\rangle \langle +| - |- \rangle \langle -|) (|+\rangle \langle -| + |- \rangle \langle +|) \\
&\quad - (|+\rangle \langle -| + |- \rangle \langle +|) (|+\rangle \langle +| - |- \rangle \langle -|) \\
&= -\frac{\hbar^2}{4} (-|+\rangle \langle -| + |- \rangle \langle +|) \\
&= i\hbar S_y
\end{aligned}$$

For the anticommutator relations, all we need to prove is that $S_i S_j = -S_j S_i$ when $i \neq j$. Of course, when $i = j$ we will always have $\{S_i, S_j\} = 2S_i^2 = \frac{\hbar^2}{2}$ since $S_i^2 = I \quad \forall i$ ■

Problem 6.

Solution.

We would like to find a representation for the state $|\mathbf{S} \cdot \hat{n}; +\rangle$ in the S_z basis. We first write the operator $\mathbf{S} \cdot \hat{n}$ explicitly in this basis

$$\begin{aligned}
\mathbf{S} \cdot \hat{n} &= \sin \beta \cos \alpha S_x + \sin \beta \sin \alpha S_y + \cos \beta S_z \\
&= \frac{\hbar}{2} \begin{pmatrix} \cos \beta & \sin \beta \exp(-i\alpha) \\ \sin \beta \exp(i\alpha) & -\cos \beta \end{pmatrix}
\end{aligned}$$

As usual, we find the eigenvalues of this operator by solving the characteristic equation:

$$\begin{aligned}
\det(\mathbf{S} \cdot \hat{n} - \lambda I) &= \left(\frac{\hbar}{2} \cos \beta - \lambda \right) \left(-\frac{\hbar}{2} \cos \beta - \lambda \right) - \frac{\hbar^2}{4} \sin^2 \beta \\
&= \lambda^2 - \frac{\hbar^2}{4} = 0
\end{aligned}$$

Therefore $\lambda = \pm \frac{\hbar}{2}$ as expected. Let ψ_1 and ψ_2 represent the components of the eigenket $|\mathbf{S} \cdot \hat{n}; +\rangle$ of this operator. We then need to solve the following system for the components ψ_1 and ψ_2

$$\begin{aligned}\psi_1 \cos \beta + \psi_2 \sin \beta \exp(-i\alpha) &= \psi_1 \\ \psi_1 \sin \beta \exp(i\alpha) - \psi_2 \cos \beta &= \psi_2\end{aligned}$$

The system does not have a real solution. But we can make a lucky guess that $\psi_1 = \cos \frac{\beta}{2}$ and $\psi_2 = \sin \frac{\beta}{2} \exp(i\alpha)$

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