Bell's Inequality

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CHSH Inequality

Alice: Q, R Bob: S, T

Classical observables described by the joint distribution P(Q, R, S, T). Combination of correlations between Alice and Bobs measurements are bounded according to the CHSH inequality

$$|E(QS) + E(RS) + E(RT) - E(QT)| \le 2$$

Tsirelson's Inequality

Quantum mechanics breaks the CHSH inequality

For the quantum version, define 4 spin operators along arbitrary directions $Q = \vec{a} \cdot \sigma, R = \vec{r} \cdot \sigma, S = \vec{s} \cdot \sigma, T = \vec{t} \cdot \sigma.$ The book uses $\vec{q} = (0,0,1), \vec{r} = (1,0,0), \vec{s} = (-\frac{1}{\sqrt{2}},0,-\frac{1}{\sqrt{2}}), \vec{t} = (-\frac{1}{\sqrt{2}},0,\frac{1}{\sqrt{2}})$

One can prove that (and I do this partially on the next slide)

$$|\langle Q \otimes S \rangle + \langle R \otimes S \rangle + \langle R \otimes T \rangle - \langle Q \otimes T \rangle| \le 2\sqrt{2}$$

Clearly this exceeds the CHSH bound by a factor of $\sqrt{2}$.

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The Tsirelson bound

I will start from the solution to Problem 2.3 in the book:

$$(Q \otimes S + R \otimes S + R \otimes T - Q \otimes T)^2 = 4I + [Q, R] \otimes [S, T]$$

Jensen's inequality:

$$\langle (Q \otimes S + R \otimes S + R \otimes T - Q \otimes T) \rangle^{2} \leq \langle (Q \otimes S + R \otimes S + R \otimes T - Q \otimes T)^{2} \rangle$$

$$= \langle 4I + [Q, R] \otimes [S, T] \rangle$$

$$= 4 + \langle [Q, R] \otimes [S, T] \rangle$$

This gives us a very general bound on the LHS for possible $Q,R,S,T,|\psi\rangle$; however, we may still want to know how $\langle (Q\otimes S+R\otimes S+R\otimes T-Q\otimes T)\rangle$ depends on the $|\psi\rangle$ we pick.

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Assumptions and objectives

As in the book, Assume \vec{q} , \vec{r} and \vec{s} , \vec{t} are orthogonal and all \vec{q} , \vec{r} , \vec{s} , \vec{t} live in a plane (say x-z)

Now draw a set of pure states $|\psi\rangle$ possibly with different degrees of entanglement

Question: How does the correlation function vary across states with different degrees of entanglement?

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The Tsirelson bound

Using that
$$(\sigma \cdot \vec{a})(\sigma \cdot \vec{b}) = \vec{a} \cdot \vec{b} + i\sigma \cdot (\vec{a} \times \vec{b})$$

$$[A, B] = i\sigma \cdot (\vec{a} \times \vec{b} - \vec{b} \times \vec{a}) = 2i\sigma \cdot (\vec{a} \times \vec{b})$$

Let $\vec{n} = \vec{q} \times \vec{r}$ and $\vec{m} = \vec{s} \times \vec{t}$. Per our constraints, we could have $\vec{n} = \vec{m} = \hat{y}$.

$$\begin{split} \langle [Q,R] \otimes [S,T] \rangle &= -4 \, \langle \psi | \, \sigma \cdot \vec{n} \otimes \sigma \cdot \vec{m} \, | \psi \rangle \\ &= -4 \, \langle \psi | \, \sigma \cdot \hat{y} \otimes \sigma \cdot \hat{y} \, | \psi \rangle \\ &= -4 \, \langle \psi | \, \sigma_{1y} \otimes \sigma_{2y} \, | \psi \rangle \end{split}$$

$$\langle (Q \otimes S + R \otimes S + R \otimes T - Q \otimes T) \rangle \leq 2\sqrt{1 - \langle \psi | \sigma_{1y} \otimes \sigma_{2y} | \psi \rangle}$$

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Quantifying entanglement: partial traces

$$\begin{aligned} \operatorname{Tr}_{A}(\rho_{AB}) &= \sum_{ijkl} \rho_{ij}^{kl} \operatorname{Tr}_{A}(|i\rangle \langle k|) \otimes |j\rangle \langle I| \\ &= \sum_{i} \left(\sum_{jl} \rho_{ij}^{il} |j\rangle \langle I| \right) \\ &= (\rho_{00}^{00} + \rho_{10}^{10}) |0\rangle \langle 0| + (\rho_{00}^{01} + \rho_{10}^{11}) |0\rangle \langle 1| + (\rho_{01}^{00} + \rho_{11}^{10}) |1\rangle \langle 0| + (\rho_{01}^{01} + \rho_{11}^{11}) |1\rangle \langle 1| \end{aligned}$$

$$\begin{aligned} \operatorname{Tr}_{B}(\rho_{AB}) &= \sum_{ijkl} \rho_{ij}^{kl} |i\rangle \langle k| \otimes \operatorname{Tr}_{B}(|j\rangle \langle l|) \\ &= \sum_{j} \left(\sum_{ik} \rho_{ij}^{kj} |i\rangle \langle k| \right) \\ &= (\rho_{00}^{00} + \rho_{01}^{01}) |0\rangle \langle 0| + (\rho_{00}^{10} + \rho_{01}^{11}) |0\rangle \langle 1| + (\rho_{10}^{00} + \rho_{11}^{01}) |1\rangle \langle 0| + (\rho_{10}^{10} + \rho_{11}^{11}) |1\rangle \langle 1| \end{aligned}$$

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Reduced density matrices for an arbitrary state

$$\operatorname{Tr}_{A}(\rho_{AB}) = \begin{pmatrix} \rho_{00}^{00} + \rho_{10}^{10} & \rho_{00}^{01} + \rho_{11}^{11} \\ \rho_{01}^{00} + \rho_{11}^{10} & \rho_{01}^{01} + \rho_{11}^{11} \end{pmatrix} = \begin{pmatrix} |\alpha|^{2} + |\gamma|^{2} & \alpha\beta^{*} + \gamma\delta^{*} \\ \beta\alpha^{*} + \delta\gamma^{*} & |\beta|^{2} + |\delta|^{2} \end{pmatrix}$$

$$\operatorname{Tr}_{B}(\rho_{AB}) = \begin{pmatrix} \rho_{00}^{00} + \rho_{01}^{01} & \rho_{00}^{10} + \rho_{01}^{11} \\ \rho_{10}^{00} + \rho_{11}^{01} & \rho_{10}^{10} + \rho_{11}^{11} \end{pmatrix} = \begin{pmatrix} |\alpha|^{2} + |\beta|^{2} & \alpha\gamma^{*} + \beta\delta^{*} \\ \gamma\alpha^{*} + \delta\beta^{*} & |\gamma|^{2} + |\delta|^{2} \end{pmatrix}$$

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Definition of entanglement entropy

The entanglement entropy of a bipartite system is the Von Neumann entropy of either reduced density matrix (it doesn't matter which one we choose)

$$S(\rho_A) = -\text{Tr}(\rho_A \log \rho_A) = -\sum_x \lambda_x \log \lambda_x$$

for eigenvalues λ_x of ρ_A . This tells us: do the reduced states ρ_A and ρ_B contain all the information in ρ_{AB} ?

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Entanglement entropy and Tsirelson's bound

Sanity check:

Draw random reals $a, b, c, d, e, f, g, h \sim U([0, 1]^8)$

Construct
$$|\psi\rangle=\left(a+ib\right)|00\rangle+\left(c+id\right)|01\rangle+\left(e+if\right)|10\rangle+\left(g+ih\right)|11\rangle$$

Normalize
$$|\psi\rangle o \frac{|\psi\rangle}{\sum_n |c_n|^2}$$

Compute
$$S(\rho_A)$$
 and $2\sqrt{1-ra{\psi}\sigma_{1y}\otimes\sigma_{2y}\ket{\psi}}$

Entanglement entropy and Tsirelson's bound

