Homework 2

Quantum Mechanics
August 29th, 2022 — Clayton Seitz

Problem 1. *Problem 1.12 from Sakurai*

Solution.

If we choose the representation such that $|1\rangle = \frac{1}{\sqrt{2}}\begin{pmatrix}1\\1\end{pmatrix}$ and $|2\rangle = \frac{1}{\sqrt{2}}\begin{pmatrix}1\\-1\end{pmatrix}$ then we can use the definition of the outer product to show that

$$H = a\begin{pmatrix}1 & 1\\1 & -1\end{pmatrix}$$

The energy eigenvalues are then found by

$$\begin{aligned}\det(H-\lambda I) &= \det\begin{pmatrix}a-\lambda & a\\ a & -a-\lambda\end{pmatrix}\\ &= (a-\lambda)(-a-\lambda)-a^2\\ &= \lambda^2-2a^2=0\end{aligned}$$

therefore $E_\pm = \pm a\sqrt{2}$. The + eigenvector $|\psi_+\rangle$ is given by the system

$$\begin{pmatrix}a-E_+ & a\\ a & -a-E_+\end{pmatrix}\begin{pmatrix}\psi_1^+\\ \psi_2^+\end{pmatrix} = \begin{pmatrix}a-a\sqrt{2} & a\\ a & -a-a\sqrt{2}\end{pmatrix}\begin{pmatrix}\psi_1^+\\ \psi_2^+\end{pmatrix}=0$$

$$\begin{aligned}(1-\sqrt{2})\psi_1^+ + \psi_2^+ &= 0\\ \psi_1^+ - (1+\sqrt{2})\psi_2^+ &= 0\end{aligned}$$

The second equation is just the first multiplied by $(1-\sqrt{2})$ so we can choose $\psi_1^+ = 1$ giving $\psi_2^+ = \sqrt{2}-1$

The eigenvector $|\psi_-\rangle$ is found similarly

$$\begin{pmatrix} a - E_- & a \\ a & -a - E_- \end{pmatrix} \begin{pmatrix} \psi_1^- \\ \psi_2^- \end{pmatrix} = \begin{pmatrix} a + a\sqrt{2} & a \\ a & -a + a\sqrt{2} \end{pmatrix} \begin{pmatrix} \psi_1^- \\ \psi_2^- \end{pmatrix} = 0$$

$$\begin{aligned} (1+\sqrt{2})\psi_1^+ + \psi_2^+ &= 0 \\ \psi_1^+ + (-1+\sqrt{2})\psi_2^+ &= 0 \end{aligned}$$

Similar to before, the second equation is $(-1+\sqrt{2})$ multiplied by the first, allowing us to set $\psi_1^- = 1$ and $\psi_2^- = -(1+\sqrt{2})$, giving a $|\psi_-\rangle$ that is orthogonal to $|\psi_+\rangle$

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Problem 2. *Problem 1.13 from Sakurai*

Solution.

Writing H out in matrix form gives

$$\begin{aligned} H &= \frac{H_{11}}{2}\begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix} + \frac{H_{22}}{2}\begin{pmatrix} 1 & -1 \\ -1 & 1 \end{pmatrix} + H_{12}\begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \\ &= \frac{H_{11}+H_{22}}{2} I + \frac{H_{11}-H_{12}}{2}\sigma_x + H_{12}\sigma_z \\ &= aI + b\sigma_x + c\sigma_z \end{aligned}$$

where have made appropriate substitutions of constants for brevity. Now this implies,

$$\begin{aligned} H|\psi\rangle &= (aI + b\sigma_x + c\sigma_z)|\psi\rangle \\ &= a|\psi\rangle + (b\sigma_x + 0\sigma_y + c\sigma_z)|\psi\rangle \end{aligned}$$

Any $|\psi\rangle$ is an eigenvector under the identity operation, so what we are really after is an eigenvector of the operator $\boldsymbol{\sigma} \cdot \boldsymbol{a}$ for $\boldsymbol{a} = (b, 0, c)$. In other words, if $|\psi\rangle$ is an eigenvector of $\boldsymbol{\sigma} \cdot \boldsymbol{a}$ then it is also an eigenvector of H. It is useful to work with the unit vector in the direction of $\boldsymbol{a}$ which is $\hat{\boldsymbol{n}} = (b/\sqrt{b^2+c^2}, 0, c/\sqrt{b^2+c^2})$. We already know the eigenvectors of $\boldsymbol{\sigma} \cdot \hat{\boldsymbol{n}}$

$$
\begin{aligned}
|\psi_+\rangle &= \cos\frac{\beta}{2}\,|+\rangle + \exp(i\alpha)\sin\frac{\beta}{2}\,|-\rangle \\
|\psi_-\rangle &= -\sin\frac{\beta}{2}\,|+\rangle + \exp(i\alpha)\cos\frac{\beta}{2}\,|-\rangle
\end{aligned}
$$

where we take the definition that α is the polar angle and β the azimuthal angle. Therefore

$$
\begin{aligned}
&\alpha = 0 \\
&\beta = \arctan\left(\frac{n_z}{n_x}\right) = \arctan\left(\frac{c}{b}\right) = \arctan\left(\frac{2H_{12}}{H_{11} - H_{12}}\right)
\end{aligned}
$$

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Problem 3. *Problem 1.15 from Sakurai*

Solution. After the first measurement along $+\hat{z}$, all of our atoms are prepared in the $|+\rangle$ state in the S_z basis. At the next apparatus oriented along $\hat{n}$, more atoms will be filtered out since $|+\rangle$ is not an eigenket of the $\mathbf{S}\cdot\hat{n}$ operator. Recall that $|+\rangle_n$ is

$$
|+\rangle_n = \cos\frac{\beta}{2}\,|+\rangle + \sin\frac{\beta}{2}\,|-\rangle
$$

The probability the state $|+\rangle$ survives is given by the inner product

$$
\begin{aligned}
|\langle+|+\rangle_n|^2 &= |\langle+|\cos\frac{\beta}{2}\,|+\rangle + \langle+|\sin\frac{\beta}{2}\,|-\rangle\,|^2 \\
&= \cos^2\frac{\beta}{2}
\end{aligned}
$$

After this, all atoms are in the $|+\rangle_n$ state. We then filter the atoms one more time with an apparatus along $-\hat{z}$. The fraction that survive this one is given by

$$
\begin{aligned}
|\langle-|+\rangle_n|^2 &= |\langle-|\cos\frac{\beta}{2}\,|+\rangle + \langle-|\sin\frac{\beta}{2}\,|-\rangle\,|^2 \\
&= \sin^2\frac{\beta}{2}
\end{aligned}
$$

Therefore the fraction output is $\cos^2\frac{\beta}{2}\sin^2\frac{\beta}{2}$. We can maximize this function by setting $\beta = \pi/2$

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Problem 4. *Problem 1.16 from Sakurai*

Solution.
We have the observable

$$O = \frac{1}{\sqrt{2}}\begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}$$

$$\begin{aligned}\det(O - \lambda I) &= \det\begin{pmatrix} -\lambda & \frac{1}{\sqrt{2}} & 0 \\ \frac{1}{\sqrt{2}} & -\lambda & \frac{1}{\sqrt{2}} \\ 0 & \frac{1}{\sqrt{2}} & -\lambda \end{pmatrix} \\ &= -\lambda\left(\lambda^2 - \frac{1}{2}\right) - \frac{1}{\sqrt{2}}\left(-\frac{\lambda}{\sqrt{2}}\right) \\ &= -\lambda^3 + \lambda = 0\end{aligned}$$

Clearly our eigenvalues are $\lambda = \pm 1$

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Problem 5. *Problem 1.23 from Sakurai*

Solution. For the ground state, the position space wavefunction $|\psi\rangle$ is a solution to the eigenvalue equation

$$\begin{aligned}H|\psi\rangle &= \left[\frac{\boldsymbol{p}^2}{2m} + \boldsymbol{V}(x)\right]|\psi\rangle \\ &= -\frac{\hbar^2}{2m}\frac{\partial^2|\psi\rangle}{\partial x^2} + V(x)|\psi\rangle \\ &= E|\psi\rangle\end{aligned}$$

We set the boundary conditions $\psi(0) = 0$ and $\psi(a) = 0$ since the wavefunction must vanish at the two walls. Since $V(x) = 0$ inside the well, Schrodinger's equation reduces to

$$H|\psi\rangle = -\frac{\hbar^2}{2m}\frac{\partial^2|\psi\rangle}{\partial x^2} = E|\psi\rangle$$

This equation has the general solution

$$|\psi\rangle = A\exp(ikx) + B\exp(-ikx)$$

Given our boundary condition $\psi(a) = 0$, the wavelength must satisfy $ka = n\pi$ which means that $k = \frac{n\pi}{a}$ for integer $n > 0$, which gives us the solution

$$|\psi\rangle = A\sin\left(\frac{n\pi x}{a}\right)$$

It is straightforward to show that

$$\langle\psi|\psi\rangle = \frac{2}{a}\int_0^a \sin^2\left(\frac{n\pi x}{a}\right)dx = 1$$

Giving the eigenkets

$$|\psi\rangle = \sqrt{\frac{2}{a}}\sin\left(\frac{n\pi x}{a}\right)$$

the variance in position when $n = 1$ is

$$\begin{aligned}
\langle(\Delta x)^2\rangle &= \langle x^2\rangle - \langle x\rangle^2 \\
&= \langle\psi| x^2 |\psi\rangle - (\langle\psi| x |\psi\rangle)^2 \\
&= x^2\langle\psi|\psi\rangle - (x\langle\psi|\psi\rangle)^2 \\
&= \frac{2}{a}\int_0^a x^2\sin^2\left(\frac{\pi x}{a}\right) - \left(\frac{2}{a}\int_0^a x\sin^2\left(\frac{\pi x}{a}\right)\right)^2
\end{aligned}$$

We can immediately write the value of $(\langle x\rangle)^2$ based on the symmetry of the wavefunction

$$\left(\int_0^a x\sin^2\left(\frac{\pi x}{a}\right)dx\right)^2 = \frac{a^2}{4}$$

The term $\langle x^2 \rangle$ is given by the integral

$$\int_0^a x^2 \sin^2(\alpha x)\, dx$$

$$\begin{aligned}
\langle (\Delta p)^2 \rangle &= \langle p^2 \rangle - \langle p \rangle^2 \\
&= \langle \psi | \, p^2 \, | \psi \rangle - (\langle \psi | \, p \, | \psi \rangle)^2 \\
&= \langle \psi | \, \hbar^2 \frac{\partial^2}{\partial x^2} \, | \psi \rangle - \left(\langle \psi | - i\hbar \frac{\partial}{\partial x} \, | \psi \rangle \right)^2 \\
&= -c\alpha^2 \hbar^2 \int_0^a \sin(\alpha x) dx
\end{aligned}$$

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Problem 6. *Problem 1.24 from Sakurai*

Solution.

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