Homework 6

Quantum Mechanics

October 28th, 2022

C Seitz

Problem 1. Problem 3.12 from Sakurai

Solution.

In general the ensemble average of an operator [A] is defined as

$$[A] = \sum_{i} w_i \langle \alpha_i | A | \alpha_i \rangle$$

where $\sum_{i} w_i = 1$

$$[\sigma_x] = a \langle +| \sigma_x | + \rangle + (1-a) \langle -; y | \sigma_x | -; y \rangle$$

$$= a \langle +| (|+\rangle \langle -| +|-\rangle \langle +|) | + \rangle + (1-a) \langle -; y | (|+\rangle \langle -| +|-\rangle \langle +|) | -; y \rangle$$

$$= 0$$

$$[\sigma_y] = a \langle + | \sigma_y | + \rangle + (1 - a) \langle -; y | \sigma_y | -; y \rangle$$

$$= ai \langle + | (|+\rangle \langle -|-|-\rangle \langle +|) | + \rangle + i(1 - a) \langle -; y | (|+\rangle \langle -|-|-\rangle \langle +|) | -; y \rangle$$

$$= i(1 - a) \langle -; y | \left(-\frac{i}{\sqrt{2}} | + \rangle - \frac{1}{\sqrt{2}} | - \rangle \right)$$

$$= -i(1 - a) \langle -; y | +; y \rangle = 0$$

$$[\sigma_z] = a \langle +| (|-\rangle \langle -|-|+\rangle \langle +|) |+\rangle + i(1-a) \langle -; y| (|-\rangle \langle -|-|+\rangle \langle +|) |-; y\rangle$$
$$= -a + i(1-a) \langle -; y| \left(-\frac{i}{\sqrt{2}} |+\rangle - \frac{1}{\sqrt{2}} |-\rangle\right)$$

Problem 2. Problem 3.13 from Sakurai

Solution.

The state vector in the S_z basis has the form

$$|\alpha\rangle = c_+ |+\rangle + c_- |-\rangle$$

First note that

$$\langle S_z \rangle = |c_+|^2 - |c_-|^2 |c_+|^2 + |c_-|^2 = 1$$

Together, these equations tell us the magnitude of each complex component.

$$|c_{+}|^{2} = \frac{\langle S_{z} \rangle + 1}{2} \quad |c_{-}|^{2} = \frac{1 - \langle S_{z} \rangle}{2}$$

$$\langle S_x \rangle = \langle \alpha | (|+\rangle \langle -|+|-\rangle \langle +|) (c_+ |+\rangle + c_- |-\rangle)$$

$$= \langle \alpha | (c_- |+\rangle + c_+ |-\rangle)$$

$$= (c_+^* \langle +|+c_-^* \langle -|) (c_- |+\rangle + c_+ |-\rangle)$$

$$= c_+^* c_- + c_-^* c_+$$

$$= |c_+||c_-| (e^{i(\theta-\phi)} + e^{i(\phi-\theta)})$$

$$= 2|c_+||c_-| \cos(\theta - \phi)$$

$$\langle S_{y} \rangle = \langle \alpha | ((i \mid +) \langle -| - i \mid -) \langle +|) (c_{+} \mid +) + c_{-} \mid -\rangle)$$

$$= i \langle \alpha | (c_{-} \mid +) - c_{+} \mid -\rangle)$$

$$= i (c_{+}^{*} \langle +| + c_{-}^{*} \langle -|) (c_{-} \mid +) - c_{+} \mid -\rangle)$$

$$= c_{+}^{*} c_{-} - c_{-}^{*} c_{+}$$

$$= |c_{+}| |c_{-}| (e^{i(\theta - \phi)} - e^{i(\phi - \theta)})$$

$$= 2i |c_{+}| |c_{-}| \sin(\theta - \phi)$$

So $\langle S_x \rangle$ gives us the phase difference of c_+ and c_- . Then the sign of $\langle S_y \rangle$ tells us whether θ or ϕ is larger, since sine is odd. This is all we can hope to extract from the expectation values, since multiplying by a global phase $e^{i\delta} |\alpha\rangle$ has no effect on the expectation values.

Problem 3. Problem 3.14 from Sakurai

Solution.

$$\rho = \sum_{i} w_{i} |\psi_{i}\rangle \langle \psi_{i}|$$

$$= \frac{1}{3} (|\alpha\rangle \langle \alpha| + |\beta\rangle \langle \beta| + |2\rangle \langle 2|)$$

We can write this out explicitly in the subspace spanned by $|0,1,2\rangle$

$$|\alpha\rangle\langle\alpha| = \frac{1}{2} \begin{pmatrix} 1 & 1 & 0 \\ 1 & 1 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad |\beta\rangle\langle\beta| = \frac{1}{2} \begin{pmatrix} 0 & 0 & 0 \\ 0 & 1 & 1 \\ 0 & 1 & 1 \end{pmatrix} \quad |2\rangle\langle2| = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$\rho = \frac{1}{6} \begin{pmatrix} 1 & 1 & 0 \\ 1 & 2 & 1 \\ 0 & 1 & 2 \end{pmatrix}$$

Now recall that $H = \hbar\omega(N + \frac{1}{2})$ which is

$$H = \left(\frac{\hbar\omega}{2}\right) \mathbb{I}_{3\times 3} + \begin{pmatrix} 0 & 0 & 0\\ 0 & 1 & 0\\ 0 & 0 & 2 \end{pmatrix}$$

$$[H] = \text{Tr}(\rho H) = \hbar \omega \text{Tr}(\rho N + \rho/2) = \hbar \omega \left(\text{Tr}(\rho N) + \text{Tr}(\rho/2) \right) = \frac{17\hbar \omega}{12}$$

Problem 4. Problem 3.15 from Sakurai

Solution.

$$\rho(t_0) = \sum_{i} w_i |\psi_i; t_0\rangle \langle \psi_i; t_0|$$

In the Schrodinger picture, the coefficients of the state vectors evolve. Therefore,

$$\rho(t) = \sum_{i} w_{i} \mathcal{U}(t, t_{0}) |\psi_{i}; t_{0}\rangle \langle \psi_{i}; t_{0}| \mathcal{U}^{\dagger}(t, t_{0})$$

$$= \mathcal{U}(t, t_{0}) \left(\sum_{i} w_{i} |\psi_{i}; t_{0}\rangle \langle \psi_{i}; t_{0}|\right) \mathcal{U}^{\dagger}(t, t_{0})$$

$$= \mathcal{U}(t, t_{0})\rho(t_{0})\mathcal{U}^{\dagger}(t, t_{0})$$

For a pure ensemble in state $|\psi_i\rangle$, the density matrix is

$$\rho(t_0) = |\psi_i; t_0\rangle \langle \psi_i; t_0|$$

At a later time, the density matrix is

$$\rho(t) = \mathcal{U}(t, t_0) \rho(t_0) \mathcal{U}^{\dagger}(t, t_0)$$

$$= \mathcal{U}(t, t_0) |\psi_i; t_0\rangle \langle \psi_i; t_0| \mathcal{U}^{\dagger}(t, t_0)$$

$$= |\psi_i; t_0\rangle \langle \psi_i; t_0|$$

since the eigenkets are stationary

Problem 5. Problem 3.16 from Sakurai

Solution.

Problem 6. Problem 3.40 from Sakurai

Solution.