

# Homework 2

Quantum Mechanics

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## Problem 1. 2.2

**Solution.**

The matrix representation of  $A$  is

$$A = \begin{pmatrix} \langle 0| A |0\rangle & \langle 0| A |1\rangle \\ \langle 1| A |0\rangle & \langle 1| A |1\rangle \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

In the output basis

$$A = \begin{pmatrix} \langle 0| A |0\rangle & \langle 0| A |1\rangle \\ \langle 1| A |0\rangle & \langle 1| A |1\rangle \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

We can choose a different basis, say  $|+\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$ ,  $|-\rangle = \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)$ .

$$U = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$$

In this basis  $A$  takes the form:

$$A' = UA = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & -1 \\ 1 & 1 \end{pmatrix}$$

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## Problem 2. 2.9

**Solution.**

$$\sigma_z = |1\rangle \langle 1| - |0\rangle \langle 0|$$

$$\sigma_x = |1\rangle \langle 0| + |0\rangle \langle 1|$$

$$\sigma_y = i|0\rangle \langle 1| - i|1\rangle \langle 0|$$

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**Problem 3. 2.12**

**Solution.** A matrix is diagonalizable if and only if the algebraic multiplicity equals the geometric multiplicity of each eigenvalue. It is easy to show that the characteristic equation here is  $(1 - \lambda)^2 = 0$  which only has one solution. ■

**Problem 4. 2.17**

**Solution.**

If  $H$  is normal, it must be diagonalizable and has the eigendecomposition

$$H = U\Lambda U^\dagger$$

where  $U$  is some unitary matrix. The conjugate transpose is

$$H^\dagger = U^\dagger \Lambda^\dagger U$$

If  $H = H^\dagger$ , and  $\Lambda$  is diagonal, then

$$U^\dagger \Lambda^\dagger U = U\Lambda U^\dagger$$

which means  $\Lambda = \Lambda^\dagger$  i.e. the eigenvalues are real. Furthermore, if  $\Lambda$  is diagonal and purely real, then clearly  $H = H^\dagger$ . ■

**Problem 5. 2.18**

**Solution.** For a unitary matrix  $U^\dagger U = I$ , so for an eigenvector  $|\alpha\rangle$ ,

$$\langle\alpha| U^\dagger U |\alpha\rangle = \langle\alpha| I |\alpha\rangle = 1$$

and  $\langle\alpha| U^\dagger U |\alpha\rangle = \lambda^* \lambda$ , so  $\lambda^* \lambda = 1$ . ■

**Problem 6. 2.24**

**Solution.** ■

**Problem 7. Gram-Schmidt**

**Solution.** It suffices to show that the following matrix has nonzero determinant:

$$A = \begin{pmatrix} 1 & 1 & 1 \\ 0 & 1 & 1 \\ -1 & 0 & 1 \end{pmatrix}$$

And it is straightforward to show that

$$\det(A) = 1$$

Therefore these vectors are indeed linearly independent, but not orthogonal. We can make them orthogonal using the Gram-Schmidt procedure. Let  $|0\rangle, |1\rangle, |2\rangle$  be our non-orthogonal basis vectors.

$$|0'\rangle = |0\rangle$$

For the second basis vector we have

$$\begin{aligned} |1'\rangle &\propto |1\rangle - \langle 0'|1\rangle |0\rangle \\ &= |1\rangle - |0\rangle \end{aligned}$$

and with the appropriate normalization we get

$$|1'\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ 1 \\ -1 \end{pmatrix}$$

For the third basis vector we have

$$\begin{aligned} |2'\rangle &\propto |2\rangle - \langle 0|2\rangle |0\rangle - \langle 1'|2\rangle |1'\rangle \\ &= |2\rangle \end{aligned}$$

and with the appropriate normalization we get

$$|2'\rangle = \frac{1}{\sqrt{3}} \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix}$$

The physical interpretation of this relative to the standard basis for  $\mathbb{R}^3$  is that we rotated the standard basis  $\phi = -\pi/4$  and  $\theta = \pi/4$ . ■

**Problem 8.** *Normal matrix parameterization*

**Solution.**

$$\begin{aligned} (a \cdot \sigma)(a^* \cdot \sigma) &= (a^* \cdot \sigma)(a \cdot \sigma) \\ \Rightarrow a \cdot a^* + i(a \times a^*) \cdot \sigma &= a^* \cdot a + i(a^* \times a) \cdot \sigma \\ \Rightarrow (a \times a^*) \cdot \sigma &= (a^* \times a) \cdot \sigma \\ \Rightarrow (a \times a^*) &= (a^* \times a) \end{aligned}$$

which occurs when  $a = a^*$  and the vector is strictly real. If this is satisfied, the operator is normal, and has a spectral decomposition:

$$\begin{aligned} A &= \lambda_1 |\lambda_1\rangle \langle \lambda_1| + \lambda_2 |\lambda_2\rangle \langle \lambda_2| \\ &= \lambda_1 P_1 + \lambda_2 P_2 \end{aligned}$$
■

**Problem 9.** *2.26*

**Solution.** Writing out  $|\psi\rangle^{\otimes 2}$  explicitly, we have

$$|\psi\rangle^{\otimes 2} = \frac{1}{2} (|00\rangle + |01\rangle + |10\rangle + |11\rangle)$$

or in terms of tensor products we have

$$|\psi\rangle^{\otimes 2} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix} \otimes \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix} = \frac{1}{2} \begin{pmatrix} 1 \\ 1 \\ 1 \\ 1 \end{pmatrix}$$

Writing out  $|\psi\rangle^{\otimes 3}$  explicitly, we have

$$|\psi\rangle^{\otimes 3} = \frac{1}{2^{3/2}} (|000\rangle + |001\rangle + |100\rangle + |010\rangle + |101\rangle + |111\rangle + |110\rangle + |011\rangle)$$

or in terms of tensor products we have

$$|\psi\rangle^{\otimes 3} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix} \otimes \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix} \otimes \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix} = \frac{1}{2^{3/2}} \begin{pmatrix} 1 \\ 1 \\ 1 \\ 1 \\ 1 \\ 1 \\ 1 \\ 1 \end{pmatrix}$$

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### Problem 10. 2.27

**Solution.**

$$X \otimes Z = \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \\ 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \end{pmatrix}$$

$$I \otimes X = \begin{pmatrix} 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix}$$

$$X \otimes I = \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{pmatrix}$$

Clearly from this last result, the tensor product does not necessarily commute.

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**Problem 11. 2.33**

**Solution.** Consider the outer product representation of the Hadamard gate

$$\begin{aligned} H &= \frac{1}{\sqrt{2}} (|0\rangle \langle 0| + |0\rangle \langle 1| + |1\rangle \langle 0| - |1\rangle \langle 1|) \\ &= \frac{1}{\sqrt{2}} \sum_{x,y} (-1)^{x \cdot y} |x\rangle \langle y| \end{aligned}$$

To see how this generalizes, consider:

$$\begin{aligned} |0\rangle \langle 0| \otimes |0\rangle \langle 0| &= (|0\rangle \otimes |0\rangle)(\langle 0| \otimes \langle 0|) \\ &= |00\rangle \otimes \langle 00| \end{aligned}$$

This is generally true of the tensor product, so multiplying sums like the one above  $n$  times will give a similar expression for vectors  $\mathbf{x}$  and  $\mathbf{y}$ :

$$H^{\otimes n} = \frac{1}{\sqrt{2}} \sum_{\mathbf{x}, \mathbf{y}} (-1)^{\mathbf{x} \cdot \mathbf{y}} |\mathbf{x}\rangle \langle \mathbf{y}|$$

For  $n = 2$ , we just need to evaluate this sum for all possible binary strings of length 2 and add them up according to this scheme, which gives

$$H^{\otimes 2} = \frac{1}{2} \begin{pmatrix} 1 & 1 & 1 & 1 \\ 1 & -1 & 1 & -1 \\ 1 & 1 & -1 & -1 \\ 1 & -1 & -1 & 1 \end{pmatrix}$$

■

**Problem 12. 2.34**

**Solution.** We can compute these functions of matrices if we have its spectral decomposition. For an arbitrary matrix  $A$ , with spectral decomposition, the square root of  $A$  is simply

$$\sqrt{A} = \sum_n \sqrt{\lambda_n} |\lambda_n\rangle \langle \lambda_n|$$

and its logarithm is

$$\log A = \sum_n \log \lambda_n |\lambda_n\rangle \langle \lambda_n|$$

Therefore we start by finding the eigendecomposition of the matrix. The characteristic equation is

$$\lambda^2 - 8\lambda + 7 = 0$$

so the two eigenvalues are  $\lambda_1 = 1$ ,  $\lambda_2 = 7$ . The eigenvectors are:

$$|\lambda_1\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \end{pmatrix}, |\lambda_2\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix}$$

Thus for its square root we get

$$\begin{aligned} \sqrt{A} &= \sum_n \sqrt{\lambda_n} |\lambda_n\rangle \langle \lambda_n| \\ &= |\lambda_1\rangle \langle \lambda_1| + \sqrt{7} |\lambda_2\rangle \langle \lambda_2| \\ &= \frac{1}{2} \begin{pmatrix} 1 + \sqrt{7} & -1 + \sqrt{7} \\ -1 + \sqrt{7} & 1 + \sqrt{7} \end{pmatrix} \end{aligned}$$

and the logarithm is

$$\begin{aligned} \log A &= \log 7 |\lambda_2\rangle \langle \lambda_2| \\ &= \frac{\log 7}{2} \begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix} \end{aligned}$$

■

### Problem 13. 2.35

#### Solution.

First recall that  $(\vec{a} \cdot \vec{\sigma})^k$  is identity for even  $k$  and is just  $(\vec{a} \cdot \vec{\sigma})$  for odd  $k$ .

$$\begin{aligned}
\exp(i\theta \vec{a} \cdot \vec{\sigma}) &= \sum_{k=0}^{\infty} \frac{1}{k!} (i\theta \vec{a} \cdot \vec{\sigma})^k \\
&= \sum_{k=0}^{\infty} \frac{(i\theta)^k}{k!} (\vec{a} \cdot \vec{\sigma})^k \\
&= \sum_{k \text{ odd}}^{\infty} \frac{(-1)^k \theta^k}{k!} (\vec{a} \cdot \vec{\sigma})^k + i \sum_{k \text{ even}}^{\infty} \frac{(-1)^k (\theta)^k}{k!} (\vec{a} \cdot \vec{\sigma})^k \\
&= \cos \theta I + i \sin \theta (\vec{a} \cdot \vec{\sigma})
\end{aligned}$$

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**Problem 14.** 2.39

**Solution.**

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