

Hodge Numbers from Picard–Fuchs Equations

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Abstract. Given a variation of Hodge structure over \mathbb{P}^1 with Hodge numbers $(1, 1, \dots, 1)$, we show how to compute the degrees of the Deligne extension of its Hodge bundles, following Eskin–Kontsevich–Möller–Zorich, by using the local exponents of the corresponding Picard–Fuchs equation. This allows us to compute the Hodge numbers of Zucker’s Hodge structure on the corresponding parabolic cohomology groups. We also apply this to families of elliptic curves, K3 surfaces and Calabi–Yau threefolds.

Key words: variation of Hodge structures; Calabi–Yau manifolds

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1 Introduction

The goal of this paper is to compute the Hodge numbers of the parabolic cohomology groups of a variation of Hodge structure (VHS) in the case where we know the Picard–Fuchs equation.

In more detail, let us assume that C is a smooth quasiprojective curve which bears an \mathbb{R} -VHS of weight k . If \mathbb{V} is the underlying local system and $j: C \hookrightarrow \overline{C}$ is the embedding of C into its smooth completion, then Zucker [31] showed that $H^1(\overline{C}, j_* \mathbb{V})$ (which we will call the *parabolic cohomology* of \mathbb{V}) bears a pure Hodge structure of weight $k + 1$. Moreover, Zucker also showed that this Hodge structure is closely related to the filtrands of the Leray spectral sequence when our VHS comes from a family of manifolds, so one can use the Hodge structure on the parabolic cohomology groups to understand the Hodge structure of fibrations and vice versa.

More recently, following work of Morrison and Walcher [26, 30], interest has arisen in determining when the Hodge structure on the parabolic cohomology of specific families of Calabi–Yau threefolds admits rational $(2, 2)$ classes, since such classes correspond to possible normal functions. Morrison and Walcher [26] gave an exciting interpretation of such normal functions of families of Calabi–Yau threefolds in terms of D -branes.

Following this, del Angel, Müller-Stach, van Straten, and Zuo [3] sought further examples of such normal functions, by studying Hodge structures on the parabolic cohomology of families of

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Calabi–Yau threefolds whose underlying VHS’s are pull-backs of VHS’s on the thrice-punctured sphere with $b_3 = 4$; such VHS’s with $b_3 = 4$ were classified by Doran and Morgan [12], who found 14 cases. They reduced these computations to the computation of the degree of the quasi-canonical extension of the Hodge bundles of these VHS’s. However, it appears as if the authors of [3] were not able to compute these degrees in great generality, thus many entries in their tables are left blank. Further progress was made by Holborn and Müller-Stach [23], however it appears that they were still not able to complete the table from [3]. Some related computations of degrees of quasi-canonical extensions of Hodge bundles have also been carried out by Green, Griffiths, and Kerr [17].

Recently, a preprint of Eskin, Kontsevich, Möller, and Zorich [13, Section 6] suggested a way to compute the degrees of Hodge bundles of \mathbb{R} -VHS’s of $(1, 1, \dots, 1)$ -type, when the Picard–Fuchs operators controlling these \mathbb{R} -VHS’s are known. The goal of this paper is to present the technique of [13] in greater generality¹ and apply it to an array of situations. In particular, we offer a completion of the tables of [3].

The application most germane to the previous work of the authors is to the study of threefolds fibred by lattice polarized K3 surfaces. In [7, 8, 9, 10], we constructed Calabi–Yau threefolds by first choosing convenient families of lattice polarized K3 surfaces $\mathcal{X} \rightarrow \mathbb{P}^1$, then performing base change along suitable maps $g: \mathbb{P}^1 \rightarrow \mathbb{P}^1$. One of the main results of this work [10] is a classification of all Calabi–Yau threefolds fibred by M_n -polarized K3 surfaces, where M_n denotes the rank 19 lattice $M_n := E_8 \oplus E_8 \oplus H \oplus \langle -2n \rangle$. In order to restrict the number of cases that we needed to check geometrically, the results and techniques of this paper helped immensely.

In more generality, our results allow us to compute the Hodge numbers of the parabolic cohomology $H^1(\mathbb{P}^1, j_* \mathbb{V})$, where $\mathcal{X} \rightarrow \mathbb{P}^1$ is any smooth projective threefold fibred by K3 surfaces over \mathbb{P}^1 with generically Picard rank 19 fibres, and \mathbb{V} is the corresponding local system of cohomology over the complement of the critical values of the fibration. By the Leray spectral sequence [31, Section 15], this is a direct summand of the Hodge structure on $H^3(\mathcal{X}, \mathbb{Q})$, and its complement has no $(3, 0)$ or $(0, 3)$ part, so in fact this computes $h^{3,0}(\mathcal{X})$ and places a lower bound on $h^{2,1}(\mathcal{X})$.

Finally, we note that in the case of an elliptic fibration, the degree of the Hodge bundle can be computed geometrically by putting the fibration into Weierstrass normal form

$$Y^2 = X^3 + g_2 X + g_3$$

and computing the degrees of g_2 and g_3 . In the case of K3 surface fibrations (and to an even greater extent for Calabi–Yau threefold fibrations) we do not have such a normal form available, because there are many families of K3 surfaces polarized by rank 19 lattices over \mathbb{P}^1 . One can construct normal forms corresponding to each polarizing lattice, as we have done in [9, 10] for some specifically chosen lattices, but this approach quickly becomes intractable: given an arbitrary rank 19 lattice L , it is very difficult to find explicit representatives for all L -polarized K3 surfaces (see [5] for some discussion of this problem). The work in this paper allows us to bypass such considerations in the presence of a known Picard–Fuchs equation. In a sense this is antipodal to the work of Fujino [16], who computes the Hodge bundles of K3 surface fibrations with only unipotent monodromy.

1.1 Structure of this paper

Section 2 begins with a brief discussion of the necessary background in Hodge theory and the theory of ordinary differential equations, following [13], which we require in order to state the

¹As stated, the results of [13] apply only to VHS’s of $(1, 1, 1, 1)$ -type, though it is clear, even in their exposition, that their approach is more general.

main results in Section 2.4. Section 2 concludes with a discussion of how these results are affected by base change.

The remainder of the paper presents some applications of this theory. Firstly, in Section 3, we apply the results of Section 2 to the case of elliptic fibrations. Here we show how to compute the Hodge numbers of an elliptic surface over \mathbb{P}^1 directly from its Picard–Fuchs equation.

In Section 4, we move on to analyze $(1, 1, 1)$ -type \mathbb{R} -VHS's over \mathbb{P}^1 , which arise in the context of K3 fibrations. We again compute their Hodge numbers using the corresponding Picard–Fuchs equations, proving a Hodge bundle formula for such fibrations. In particular, we can use this to place constraints on the possible K3 fibrations that can arise on a Calabi–Yau threefold, which allows us to recover an approximate form of a result from [7].

Finally, in Section 5, we will consider the case of Calabi–Yau threefold fibrations over \mathbb{P}^1 . Here we achieve our main goal and complete the computations of [3] and [23]. Finally, we conclude with a result that allows us to constrain the possible Calabi–Yau fourfolds which may admit fibrations by the quintic mirror Calabi–Yau threefold. The results we use to do this apply in greater generality; we expect similar results to hold for any Calabi–Yau threefolds X with $(1, 1, 1, 1)$ -type Hodge structure on $H^3(X, \mathbb{Q})$.

2 Background and important results

In this section we will develop the necessary background material and state a number of results that are necessary for the subsequent computations in Sections 3, 4 and 5.

2.1 Variations of Hodge structure

Let us begin with a real variation of polarizable Hodge structure over a quasi-projective curve C . We denote this \mathbb{R} -VHS by $(\mathbb{V}, \mathcal{F}^\bullet, \nabla)$, where \mathbb{V} is the underlying real local system, \mathcal{F}^\bullet is the Hodge filtration on $\mathcal{V} := \mathbb{V} \otimes \mathcal{O}_C$ and ∇ is the Gauss–Manin connection. We require that ∇ satisfy Griffiths transversality [18, 19, 22]; in other words, $\nabla(\mathcal{F}^i)$ should be contained in $\mathcal{F}^{i+1} \otimes \Omega_C$.

Following Deligne [4], we can canonically extend the bundle \mathcal{V} to a bundle on the smooth completion \overline{C} of C , as follows. Near a point $p \in \Delta = \overline{C} \setminus C$, we may choose a chart so that p is the center of a disc D . Let V_0 be a fibre of \mathbb{V} near p and suppose that monodromy around p acts on V_0 by a quasi-unipotent transformation T . Define subspaces W_α of V_0 by

$$W_\alpha := \{v \in V_0 : (T - \zeta_\alpha)^k v = 0\},$$

where $\zeta_\alpha = e^{2\pi i \alpha}$ and α is chosen to be in the interval $[0, 1)$. These vector spaces are zero for all but finitely many values of $\alpha \in [0, 1)$, and we have a direct sum decomposition $V_0 = \bigoplus_\alpha W_\alpha$. The vector spaces W_α over the different points of the punctured disc $D \setminus \{p\}$ define a sub-bundle \mathbb{W}_α of \mathbb{V} .

Let T_α denote the action of monodromy on \mathbb{W}_α ; note that $\zeta_\alpha^{-1} T_\alpha$ is unipotent by construction. Define

$$N_\alpha := \frac{-1}{2\pi i} \log(T_\alpha),$$

where the branch of the logarithm is chosen so that the unique eigenvalue of N_α is in the interval $[0, 1)$. Then for any section v of $\mathbf{e}^* \mathbb{W}_\alpha$, where $\mathbf{e}(z) := \exp(2\pi i z)$ denotes the complex exponential, we may define

$$\tilde{v}(z) = \exp(2\pi i z(N_\alpha - \alpha))v(z),$$

where $z \in \mathbb{H}$ is an element of the upper half-plane. Since $v(z+1) = T(v(z))$, we see that $\tilde{v}(z)$ is invariant under the translation $z \mapsto z+1$, so \tilde{v} defines a holomorphic section of \mathcal{V} on the

punctured disc D^* . Deligne's canonical extension of \mathcal{V} is the bundle on D which is the \mathcal{O}_D -module spanned by all such vectors over all $\alpha \in [0, 1)$. This is denoted $\bar{\mathcal{V}}$.

The Hodge filtration \mathcal{F}^\bullet extends to a Hodge filtration on $\bar{\mathcal{V}}$, which we denote by $\bar{\mathcal{F}}^\bullet$. The graded pieces of $\bar{\mathcal{F}}^\bullet$ will be denoted $\mathcal{E}^{p, \ell-p} = \bar{\mathcal{F}}^p / \bar{\mathcal{F}}^{p-1}$, where ℓ is the length of the Hodge filtration.

Following [13, Section 2], there is a parabolic filtration of the fibre V_p of $\bar{\mathcal{V}}$ over each point $p \in \Delta$, given by $V^{\geq \beta} := \bigoplus_{\alpha \geq \beta} V_\alpha$, where V_α is the subspace of V_p spanned by local sections of $\bar{\mathcal{V}}$ coming from sections of $e^* \mathbb{W}_\alpha$. The *parabolic degree* of $\bar{\mathcal{V}}$ is given by

$$\deg_{\text{par}} \bar{\mathcal{V}} := \deg \bar{\mathcal{V}} + \sum_{\substack{p \in \Delta \\ \alpha \in [0, 1)}} \alpha \dim V_\alpha. \quad (2.1)$$

This parabolic filtration on the fibres of $\bar{\mathcal{V}}$ over points in Δ extends to a parabolic filtration on the fibres of $\mathcal{E}^{p, \ell-p}$; we denote its graded pieces by $\mathcal{E}_\alpha^{p, \ell-p}$. The connection ∇ extends to a connection on $\bar{\mathcal{V}}$ with logarithmic poles along Δ or, in other words,

$$\bar{\nabla}: \bar{\mathcal{V}} \longrightarrow \bar{\mathcal{V}} \otimes \Omega_C(\Delta).$$

This connection is horizontal with respect to the Hodge filtration $\bar{\mathcal{F}}^\bullet$, so we obtain morphisms of bundles

$$\theta_{i-1}: \mathcal{E}^{\ell-i, i} \longrightarrow \mathcal{E}^{\ell-i-1, i+1} \otimes \Omega_C(\Delta).$$

This map also respects the local filtration at each point $p \in \Delta$, i.e.,

$$\bar{\nabla}: V_p^{\geq \beta} \longrightarrow V_p^{\geq \beta} \otimes \Omega_C(\Delta).$$

The following useful result appears as part of the proof of [13, Theorem 6.1].

Lemma 2.1. *Assume that $\bar{C} = \mathbb{P}^1$. Then*

$$\deg_{\text{par}} \mathcal{E}^{p, q} = -\deg_{\text{par}} \mathcal{E}^{q, p}.$$

The following two examples will be useful in later portions of this paper.

Example 2.2 (elliptic curves). Let us take a degeneration of elliptic curves over the unit disc in \mathbb{C} , with central fibre of Kodaira type IV*. In this case, one can check (see [25, Table VI.2.1]) that the local monodromy action on the first integral cohomology of a general fibre has matrix

$$\begin{pmatrix} -1 & -1 \\ 1 & 0 \end{pmatrix}.$$

This matrix has eigenvalues equal to $e^{2\pi i/3}$ and $e^{4\pi i/3}$. Therefore, the spaces $V_{1/3}$ and $V_{2/3}$ are both 1-dimensional. We can ask how these spaces interact with the canonical extension of the Hodge filtration. We know that the map $\nabla: \bar{\mathcal{F}}^0 \rightarrow \bar{\mathcal{F}}^1 \otimes \Omega_D(0)$ respects the filtration $\mathbb{C}^2 = V^{\geq 1/3} \supset V^{\geq 2/3} = \mathbb{C}$ and thus the image of ∇ in $\bar{\mathcal{F}}^1 \otimes \Omega_D(0)$ must be $V^{\geq 2/3} = V_{2/3}$, since ∇ is not zero. Therefore $\mathcal{E}_{2/3}^{0,1} = \mathcal{E}^{0,1}$ and $\mathcal{E}_{1/3}^{0,1} = 0$. Similar arguments show that $\mathcal{E}_{1/3}^{1,0} = \mathcal{E}^{1,0}$ and $\mathcal{E}_{2/3}^{1,0} = 0$.

Using Kodaira's classification of germs of one parameter degenerations of elliptic curves, one can compute parabolic filtrations on the Hodge bundles corresponding to all non-unipotent degenerations of elliptic curves in the same way. The result is given in Table 1, which lists the values of α such that $\mathcal{E}_\alpha^{1,0} = \mathcal{E}^{1,0}$ and $\mathcal{E}_\alpha^{0,1} = \mathcal{E}^{0,1}$ for each type of singular fibre.

Using this, the parabolic degrees of $\mathcal{E}^{1,0}$ and $\mathcal{E}^{0,1}$ can be determined easily once one knows their degrees. Furthermore, for an elliptic fibration over \mathbb{P}^1 , we may combine this with equation (2.1) and Lemma 2.1 to show that $-\deg \mathcal{E}^{1,0} - \deg \mathcal{E}^{0,1}$ is equal to the number of fibres around which monodromy is not unipotent. This is precisely [23, Remark 2.4].

Fibre type	$\mathcal{E}^{1,0}$	$\mathcal{E}^{0,1}$
I_n	0	0
I_n^*	1/2	1/2
II or II*	1/6	5/6
III or III*	1/4	3/4
IV or IV*	1/3	2/3

Table 1. Values of α such that $\mathcal{E}_\alpha^{p,q} = \mathcal{E}^{p,q}$ for degenerations of elliptic curves.

Monodromy matrix	$\mathcal{E}^{2,0}$	$\mathcal{E}^{1,1}$	$\mathcal{E}^{0,2}$
T_{un}	1/2	1/2	1/2
T_{un}^2	0	0	0
T_i	1/4	1/2	3/4
$-T_i$	1/4	0	3/4
T_i^2	1/2	0	1/2
T_ω	1/6	1/2	5/6
T_ω^2	1/3	0	2/3
T_{nod}	0	1/2	0

Table 2. Values of α such that $\mathcal{E}_\alpha^{p,q} = \mathcal{E}^{p,q}$ for $(1, 1, 1)$ -VHS's.

Example 2.3 ($(1, 1, 1)$ variations of Hodge structure). Let us take a degeneration \mathcal{X} of K3 surfaces over the complex unit disc D , so that the restriction of \mathcal{X} to $D^* := D \setminus \{0\}$ is an M -polarized family of K3 surfaces, in the sense of [7, Definition 2.1], for some rank 19 lattice M . The transcendental lattices of the general fibres give rise to a VHS of type $(1, 1, 1)$ over D^* .

One can generally describe all possible monodromy matrices that can underlie such a $(1, 1, 1)$ -VHS. Up to sign, they are just symmetric squares of parabolic and elliptic elements of $\text{SL}_2(\mathbb{R})$; this comes from the fact that $\text{O}(2, 1)$ is, up to sign, a symmetric square of the group $\text{SL}_2(\mathbb{R})$.

By [23, Lemma 3.1], the monodromy of \mathcal{X} around $0 \in D$ can be written in Jordan canonical form as either

$$\begin{pmatrix} \lambda & 1 & 0 \\ 0 & \lambda & 1 \\ 0 & 0 & \lambda \end{pmatrix} \quad \text{or} \quad \begin{pmatrix} \lambda_1 & 0 & 0 \\ 0 & \lambda_2 & 0 \\ 0 & 0 & \lambda_3 \end{pmatrix}$$

for some $\lambda, \lambda_i \in \mathbb{C}$ roots of unity with $\sum_{i=1}^3 \lambda_i \in \mathbb{Z}$. Cases that are relevant to our work [10] are

$$\begin{aligned} T_{\text{un}} &= \begin{pmatrix} -1 & 1 & 0 \\ 0 & -1 & 1 \\ 0 & 0 & -1 \end{pmatrix}, & T_i &= \begin{pmatrix} \sqrt{-1} & 0 & 0 \\ 0 & -\sqrt{-1} & 0 \\ 0 & 0 & -1 \end{pmatrix}, \\ T_\omega &= \begin{pmatrix} -e^{2\pi i/3} & 0 & 0 \\ 0 & -e^{4\pi i/3} & 0 \\ 0 & 0 & -1 \end{pmatrix}, & T_{\text{nod}} &= \begin{pmatrix} -1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}, \end{aligned}$$

and their powers.

As in the previous example, we can compute the filtration on the degenerate Hodge structure from these matrices. The result is given in Table 2, which lists the values of α such that $\mathcal{E}_\alpha^{2,0} = \mathcal{E}^{2,0}$, $\mathcal{E}_\alpha^{1,1} = \mathcal{E}^{1,1}$, and $\mathcal{E}_\alpha^{0,2} = \mathcal{E}^{0,2}$ for each type of VHS.

For a fibration of K3 surfaces over \mathbb{P}^1 , we may combine this with equation (2.1) and Lemma 2.1 to show that $-\deg \mathcal{E}^{2,0} - \deg \mathcal{E}^{0,2}$ computes the number of fibres around which we have monodromy matrix with Jordan normal form T_{un} , T_i , $-T_i$, T_i^2 , T_ω or T_ω^2 . If we adjust the definition

of Π in [23, Section 3] to be the set of points in Δ whose monodromy matrices have at least two eigenvalues not equal to 1, then [23, Remark 3.4] holds.

2.2 Picard–Fuchs equation of a variation of Hodge structure

Now let us take a look at how the Picard–Fuchs equation of a variation of Hodge structure arises; more details for everything in this section and the next may be found in [13, Section 6.4]. We start with the same basic setup as before: a variation of Hodge structure $(\mathbb{V}, \mathcal{F}^\bullet, \nabla)$ over a quasi-projective curve C with its canonical extension to \overline{C} . For the sake of simplicity, we will assume that $\overline{C} = \mathbb{P}^1$.

If we now choose a global section² σ of the Serre twist $\overline{\mathcal{F}}^\ell(i)$, where ℓ is the length of the Hodge filtration and i is some integer, then we can build a differential equation associated to σ . Taking a polarization $\langle -, - \rangle$ on \mathbb{V} , we can assign to σ a set of multivalued functions

$$s_v(t) = \langle \sigma, v \rangle,$$

where v is some flat local section on \mathbb{V} (or any flat local section v of $\mathcal{V} := \mathbb{V} \otimes \mathcal{O}_C$). The functions $s_v(t)$ give multivalued meromorphic functions on \mathbb{P}^1 , with monodromy occurring only around points in $\Delta = \overline{C} \setminus C$. One can then produce a global differential equation L_σ , of rank at most the rank of \mathbb{V} on C , whose solution set is the set of functions $s_v(t)$. The differential operator L_σ is a Fuchsian ODE (which will be discussed in the following section), called the *Picard–Fuchs equation*.

The equation L_σ is defined as follows. For an appropriate choice of affine coordinate t on $\overline{C} = \mathbb{P}^1$, we may assume that $\infty \in \Delta$, and therefore that the vector field d/dt is a global vector field on $C = \mathbb{P}^1 \setminus \Delta$. If α and β are sections of \mathcal{V} , then we have

$$\frac{d}{dt} \langle \alpha, \beta \rangle = \langle \nabla_{d/dt} \alpha, \beta \rangle \pm \langle \alpha, \nabla_{d/dt} \beta \rangle,$$

where $\nabla_{d/dt} \alpha$ is defined to be $\nabla(\alpha) \in \mathcal{V} \otimes \Omega_C$ paired with d/dt . Therefore,

$$\frac{ds_v(t)}{dt} = \langle \nabla_{d/dt} \sigma, v \rangle,$$

since v is chosen to be a flat local section of \mathcal{V} .

Since $\nabla_{d/dt}^i(\sigma)$ all live in \mathcal{V} , if $r = \text{rank } \mathcal{V}$, then $\sigma, \nabla_{d/dt} \sigma, \dots, \nabla_{d/dt}^r \sigma$ must be linearly dependent over the field of meromorphic functions on C . Therefore, we can write an equation

$$\nabla_{d/dt}^n \sigma + f_1(t) \nabla_{d/dt}^{n-1} \sigma + \dots + f_n(t) \sigma = 0$$

for some minimal $n \leq r$. Therefore, for any flat local section v of \mathcal{V} , we obtain the equation

$$\frac{d^n s_v(t)}{dt^n} + f_1(t) \frac{d^{n-1} s_v(t)}{dt^{n-1}} + \dots + f_n(t) s_v(t) = 0,$$

by linearity of the operator $\langle -, v \rangle$. In other words, $s_v(t)$ is annihilated by the differential operator

$$L_\sigma = \frac{d^n}{dt^n} + f_1(t) \frac{d^{n-1}}{dt^{n-1}} + \dots + f_n(t).$$

On the other hand, to a rank r ordinary differential equation L with regular singularities, we can assign to L its solution sheaf $\text{Sol}(L)$, which is a subsheaf of \mathcal{O}_C spanned over \mathbb{C} by the local solutions of L . The solution sheaf of L_σ is precisely the subsheaf of \mathcal{O}_C spanned locally by the functions $s_v(t)$.

²More generally, if such a section does not exist, one may instead construct a Picard–Fuchs D -module; however this will not be relevant to our work.

2.3 Fuchsian ODE's

Let L be a linear ordinary differential equation in one variable, written as

$$\frac{d^n}{dt^n} + \sum_{i=1}^n f_i(t) \left(\frac{d}{dt} \right)^{n-i},$$

where the f_i are meromorphic functions in t . This equation can be rewritten in terms of the operator $\delta_t = t \frac{d}{dt}$ as

$$\delta_t^n + \sum_i^n g_i(t) \delta_t^{n-i} \quad (2.2)$$

for some meromorphic functions g_1, \dots, g_n . Indeed, one can check that

$$t^n \left(\frac{d}{dt} \right)^n = \delta_t(\delta_t - 1) \cdots (\delta_t - n + 1),$$

so it is easy to translate between the two expressions for L .

A point $a \in \mathbb{P}^1$ is called *nonsingular* if $f_i(t)$ is holomorphic at a , and is a *regular singular point* if $(t-a)^i f_i(t)$ is holomorphic at a , for all $i \in \{1, \dots, n\}$. Note that 0 is a regular singular point if and only if the coefficients $g_i(t)$ in equation (2.2) are all holomorphic at 0. A linear differential equation of one variable with only regular singularities at all points in \mathbb{P}^1 is called a *Fuchsian ODE*. Regular singular points can be divided into *apparent singularities*, around which monodromy acts as the identity on the solution sheaf, and *actual singularities*, where it does not.

Let L be a Fuchsian ODE with a regular singular point at $0 \in \mathbb{P}^1$. The polynomial

$$T^n + \sum_{i=1}^n g_i(0) T^{n-i}$$

is called the *indicial equation* of L at 0, and its roots are called the *characteristic exponents* of L at 0. For an arbitrary point $a \in \mathbb{P}^1$, one may compute the indicial equation and characteristic exponents of L at a by making the variable change $t \mapsto t - a$. We will denote the characteristic exponents of L at a by μ_1^a, \dots, μ_n^a ; throughout the paper we will assume that the μ_i^a are ordered by magnitude, so that $\mu_1^a \leq \mu_2^a \leq \dots \leq \mu_n^a$. Often, the data of the regular singular points a_i and their characteristic exponents $\mu_i^{a_j}$ are arranged into a matrix known as the *Riemann scheme* of L , which looks like

$$\left(\begin{array}{cccc} a_1 & a_2 & \cdots & a_k \\ \mu_1^{a_1} & \mu_1^{a_2} & \cdots & \mu_1^{a_k} \\ \vdots & \vdots & \ddots & \vdots \\ \mu_n^{a_1} & \mu_n^{a_2} & \cdots & \mu_n^{a_k} \end{array} \right).$$

Information about the monodromy of the solution sheaf of L can be read off from the Riemann scheme. For instance, if a is a regular singular point whose characteristic exponents are all integers, then monodromy around a is unipotent, i.e., all of its eigenvalues are 1. As a partial converse, if a is an apparent singularity, then the characteristic exponents at a are all integers.

Finally, note that it also makes sense to compute the characteristic exponents at a nonsingular point a of L . If one does so, one obtains $(\mu_1^a, \mu_2^a, \dots, \mu_n^a) = (j, j+1, \dots, j+n-1)$, for some integer j .

Remark 2.4. It is often natural to scale the solutions to a differential equation by a meromorphic function $h(t)$ on \mathbb{P}^1 to obtain a *twisted* differential equation, denoted hL . This operation acts on the characteristic exponents of L by an overall scaling ${}^h\mu_i^a = \mu_i^a + \text{ord}_a(h)$, where $\text{ord}_a(h)$ is the order of h at a . Importantly, the differences between the characteristic exponents are not affected by the process of twisting. In the main results of this paper, we will only need to know the differences between the characteristic exponents, so knowing L_σ up to twist is enough.

This is important in applications, since one often only knows the Picard–Fuchs equation of a variation of Hodge structure up to twist: for instance, this is all that the Griffiths–Dwork method [20, 21] for computing Picard–Fuchs equations of hypersurfaces guarantees. In Theorem 2.7, only the difference in characteristic exponents is relevant; therefore, our Hodge bundle computations can be done with hL for any meromorphic function h .

Remark 2.5. One can also twist by multivalued functions h on \mathbb{P}^1 . A particular choice of (possibly) multivalued function h can be made so that hL has $f_{n-1}(t) = 0$. This is called the *projective normal form* of L . The process of twisting by a multivalued function *does* affect the differences of characteristic exponents in a significant way. For this reason, we will avoid using the projective normal form in this paper.

Remark 2.6. The characteristic exponents of L at a point a allow us to compute local solutions for L near a . We remark that in general, the solutions of $Lf = 0$ around the point a have the form $t^{\mu_i^a} P_i(t)$ for some holomorphic functions $P_i(t)$, modulo terms with logarithmic singularities at a . For instance, if all μ_i^a 's are the same, then there is an ordered basis of local solutions $\{s_i(t)\}$ around a given inductively by

$$s_i(t) = t^{\mu_i^a} P_i(t) + \sum_{j=1}^{i-1} \frac{s_j(t) \log t}{j!},$$

where the P_i are holomorphic functions which do not vanish at a .

If none of the differences between characteristic exponents are integers, then there are no terms with logarithmic singularities and the local solutions around a all have the form $t^{\mu_i^a} P_i(t)$, for some holomorphic functions $P_i(t)$. In this case, if L is the Picard–Fuchs equation of a variation of Hodge structure $(\mathbb{V}, \mathcal{F}^\bullet, \nabla)$, then Deligne's canonical extension $\bar{\mathbb{V}}$ of \mathbb{V} is spanned by the set of local sections $\{t^{\lfloor \mu_i^a \rfloor} P_i(t)\}$ and the subspaces spanned by $t^{\lfloor \mu_i^a \rfloor} P_i(t)$ are the parabolic filtrands V_α , where $\alpha := \mu_i^a - \lfloor \mu_i^a \rfloor$. A more in-depth analysis of the Frobenius method shows that this is true in general.

2.4 Important results

Here we state some important results which, in combination, will allow us to compute Hodge numbers of a variation of Hodge structure of $(1, 1, \dots, 1)$ -type if we know the Picard–Fuchs equation L_σ . The first is a result of Eskin, Kontsevich, Möller, and Zorich [13], which says that the characteristic exponents of the Picard–Fuchs equation L_σ corresponding to a variation of Hodge structure of $(1, 1, \dots, 1)$ -type can be used to compute the Kodaira–Spencer maps on the bundles $\mathcal{E}^{p,q}$.

Theorem 2.7 ([13, Lemma 6.3]). *Let $(\mathbb{V}, \mathcal{F}^\bullet, \nabla)$ be a polarized \mathbb{R} -VHS of $(1, 1, \dots, 1)$ -type over a Zariski open subset $\mathbb{P}^1 \setminus \Delta$, let $\bar{\mathcal{F}}^\bullet$ be its quasi-canonical extension to \mathbb{P}^1 , let $\mathcal{E}^{p,q}$ be the graded pieces of $\bar{\mathcal{F}}^\bullet$, and let L_σ be the Picard–Fuchs differential equation associated to a meromorphic section of σ of $\mathcal{E}^{\ell,0}$. Then for $i < \ell/2$, the Kodaira–Spencer map $\theta_i: \mathcal{E}^{\ell-i,i} \rightarrow \mathcal{E}^{\ell-1-i,1+i} \otimes \Omega_{\mathbb{P}^1}(\Delta)$ is described as follows.*

1. *If p is a nonsingular point of L_σ , then θ_i is a local isomorphism at p .*

2. If p is an apparent singularity of L_σ , then θ_i has cokernel of length $\mu_{i+2}^p - \mu_{i+1}^p - 1$.
3. If p is an actual singularity of L_σ , then θ_i has cokernel of length $\lfloor \mu_{i+2}^p \rfloor - \lfloor \mu_{i+1}^p \rfloor$.

Remark 2.8. We record the observation that Eskin, Kontsevich, Möller, and Zorich use Theorem 2.7 to compute the parabolic degrees of the Hodge bundles of hypergeometric variations of Hodge structure. Specifically, they are interested in the 14 VHS's of $(1, 1, 1, 1)$ -type on \mathbb{P}^1 found by Doran and Morgan [12]. However, it seems as if the authors of [13] were unaware of the fact that this result follows easily from the data in [23], in which the degrees of $\mathcal{E}^{3,0}$ and $\mathcal{E}^{2,1}$ are computed for all fourteen $(1, 1, 1, 1)$ -type variations of Hodge structure of hypergeometric type. We will see later that Theorem 2.7 can be used to perform computations that neither the authors of [3] nor [23] were able to perform.

This result is useful because, if we know the parabolic filtration of each bundle $\mathcal{E}^{p,q}$ and the Picard–Fuchs equation L_σ , then knowing the degrees of all maps θ_i allows us to easily compute the degrees of $\mathcal{E}^{p,q}$ when each $\mathcal{E}^{p,q}$ is a line bundle (i.e. when L_σ underlies a variation of Hodge structure of weight d and type $(1, 1, \dots, 1)$).

Next we will state several results due to del Angel, Müller-Stach, van Straten and Zuo, and Hollborn and Müller-Stach, which appear in [3, 23]. These results say that, in the situations that interest us, we can compute Hodge numbers of $H^1(\mathbb{P}^1, j_*\mathbb{V})$ (where $j: C \hookrightarrow \overline{C} = \mathbb{P}^1$ denotes the inclusion) using the degrees of $\mathcal{E}^{p,q}$ and the local monodromy data of \mathbb{V} . According to Theorem 2.7, all of this information can be obtained from L_σ .

Theorem 2.9 ([3, 23]). *Let \mathbb{V} be a local system on $\mathbb{P}^1 \setminus \Delta$ which supports a variation of Hodge structure of weight n and $(1, 1, \dots, 1)$ -type. Then the Hodge number $h^{0,n+1}$ of $H^1(\mathbb{P}^1, j_*\mathbb{V})$ is given by*

$$h^{0,n+1} = h^1(\mathbb{P}^1, \mathcal{E}^{0,n}) = h^0(\mathbb{P}^1, \mathcal{O}_{\mathbb{P}^1}(-2 - \deg \mathcal{E}^{0,n})).$$

In the case where $n = 1$ or 2 in the above theorem, then we can deduce the rest of the Hodge numbers of the parabolic cohomology $H^1(\mathbb{P}^1, j_*\mathbb{V})$ by Hodge symmetry and the following result, which generalizes the classical Riemann–Hurwitz formula (presumably this result is classically known, but the only reference that we know is [3]).

Proposition 2.10 ([3, Proposition 3.6]). *Let \mathbb{V} be a local system on a quasi-projective curve C and let $j: C \hookrightarrow \overline{C}$ be the smooth completion. Let γ_q be the local monodromy matrix around the point q in $\Delta := \overline{C} \setminus C$, acting on \mathbb{V}_p for some basepoint p in C , and define $R(q) := \dim \mathbb{V}_p - \dim \mathbb{V}_p^{\gamma_q}$, where $\mathbb{V}_p^{\gamma_q}$ denotes the part of \mathbb{V}_p fixed under γ_q . Then*

$$h^1(\overline{C}, j_*\mathbb{V}) = \sum_{q \in \Delta} R(q) + (2g(\overline{C}) - 2) \operatorname{rank} \mathbb{V}.$$

Finally, if \mathbb{V} is a VHS of $(1, 1, 1, 1)$ -type, then Hollborn and Müller-Stach show how to compute the Hodge number $h^{1,3}$ of the parabolic cohomology $H^1(\mathbb{P}^1, j_*\mathbb{V})$ from the degrees of its Hodge bundles.

Theorem 2.11 ([23, Theorem 4.3]). *If \mathbb{V} supports a variation of Hodge structure of type $(1, 1, 1, 1)$ over a Zariski open subset $j: C \hookrightarrow \mathbb{P}^1$, then the Hodge number $h^{1,3}$ of $H^1(\mathbb{P}^1, j_*\mathbb{V})$ is given by*

$$h^{1,3} = -2 + b - a + |\text{II}| + |\text{III}| + |\text{IV}|,$$

where $a = \deg(\mathcal{E}^{3,0})$ and $b = \deg(\mathcal{E}^{2,1})$. The sets II and III are the subsets of $\Delta = \mathbb{P}^1 \setminus C$ whose local monodromy matrices have Jordan normal forms

$$\begin{pmatrix} 1 & 1 & 0 & 0 \\ 0 & 1 & 1 & 0 \\ 0 & 0 & 1 & 1 \\ 0 & 0 & 0 & 1 \end{pmatrix}, \quad \text{and} \quad \begin{pmatrix} 1 & 1 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 1 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

respectively. The set IV is the subset of Δ for which the monodromy matrix is strictly quasi-unipotent.

2.5 Base change and characteristic exponents

The following instructive computation will be useful to us later. We want to understand how base change of a Fuchsian ODE affects its Riemann scheme. More precisely, let us take a Fuchsian ODE L and a map $g: \mathbb{P}^1 \rightarrow \mathbb{P}^1$. Pulling-back the solution sheaf of L by g yields a local system $g^*\text{Sol}(L)$ and, in particular, another ODE g^*L , of the same rank as L , whose solution sheaf is $g^*\text{Sol}(L)$. Our goal in this section is to analyze the local exponents of g^*L . Away from the ramification points of g , the differential equation g^*L will have the same behaviour as L at the corresponding point. So we can restrict ourselves to studying the local behaviour of L at ramification points, and this can be done in the power series ring $\mathbb{C}[[t]]$. Therefore, our question reduces to a question of base change of a Fuchsian ODE along the map $s^k = t$, for a positive integer k .

The following computation appears for rank 2 ODEs in [6].

Proposition 2.12 ([6, Lemma 3.21]). *Base change of order k at a point $p \in \mathbb{P}^1$ affects the characteristic exponents of L by multiplying everything by k .*

Proof. Assume that $w(t)$ is a local solution to $Lw = 0$. By the chain rule, we have that

$$\frac{dw}{ds}(s^k) = ks^{k-1} \frac{dw}{dt}(s^k),$$

therefore

$$\delta_t w(s^k) = s^k \frac{dw}{dt}(s^k) = \frac{s}{k} \frac{dw}{ds}(s^k) = \frac{\delta_s}{k} w(s^k).$$

Thus if w is a solution to the differential equation

$$\delta_t^n + \sum_{i=1}^n f_i(t) \delta_t^{n-i},$$

then $w(s^k)$ solves

$$\delta_s^n + \sum_{i=1}^n k^i f_i(s^k) \delta_s^{n-i}.$$

Suppose that the indicial equation of L at 0 is given by

$$\prod_{i=1}^n (T - \mu_i),$$

for μ_i the local exponents of L at 0. Then the indicial equation of the differential equation in terms of s is written as

$$\prod_{i=1}^n (T - k\mu_i).$$

This proves the proposition. ■

Corollary 2.13. *After base change, the ramification points of $g: \mathbb{P}^1 \rightarrow \mathbb{P}^1$ become apparent singularities of g^*L . If p is a point with characteristic exponents all 0, then for any point $q \in g^{-1}(p)$, the characteristic exponents of g^*L at q are all 0.*

Combining this with Theorem 2.7, we can say quite a bit about how the Hodge numbers of a variation of Hodge structure of type $(1, 1, \dots, 1)$ are altered by base change.

Remark 2.14. We will take this opportunity to emphasize that performing base change on a variation of Hodge structure and base change on a linear differential equation may not in fact be compatible. In other words, there is a natural variation of Hodge structure on \mathbb{P}^1 underlying $g^*\mathbb{V}$ with appropriate canonical extensions at the points in $g^{-1}(\Delta)$. The Hodge bundles $\mathcal{E}_g^{p,q}$ of this variation of Hodge structure may not agree with the pull-backs $g^*\mathcal{E}^{p,q}$ of the Hodge bundles associated to \mathbb{V} . However, the monodromy representation and period maps are correct, so if L_g is the Picard–Fuchs system of the pulled back variation of Hodge structure, then $g^*L = {}^hL_g$ for some meromorphic function h . Thus the differences between the characteristic exponents of g^*L are the same as for L_g , so we can apply Theorem 2.7 to compute the degrees of $\mathcal{E}_g^{p,q}$.

The underlying reason behind this incompatibility is that Deligne’s canonical extension does not commute with base change. Indeed, Deligne’s canonical extension of $g^*\mathbb{V}$ chooses a branch of the logarithm so that the eigenvalues of the maps N_α at each point of $g^{-1}(\Delta)$ lie in the interval $[0, 1)$ (see Section 2.1). However base change acts additively on these eigenvalues, so may take them out of this interval; the pulled-back variation of Hodge structure may thus correspond to a different choice of extension of $g^*\mathbb{V}$.

3 Families of elliptic curves

In the remainder of this paper we will study some applications of these results in various settings. We begin with the case of a family of elliptic curves.

Let L be a rank 2 ODE which underlies a real variation of Hodge structure of weight 1. We will show that if L is the Picard–Fuchs equation of a family of elliptic curves S , then we can determine the geometry of S from L .

First, we will restate Theorem 2.7 in the case of a $(1, 1)$ -VHS.

Corollary 3.1. *Let L be a rank 2 Fuchsian ODE which underlies an \mathbb{R} -VHS of weight 1, with regular singular points at Δ in \mathbb{P}^1 . Then the map $\theta: \mathcal{E}^{1,0} \rightarrow \mathcal{E}^{0,1} \otimes \Omega_{\mathbb{P}^1}(\Delta)$ has the following structure.*

- 1) if p is a nonsingular point of L , then θ is an isomorphism at p ,
- 2) if L has an apparent singularity at p , then θ has cokernel of length $\mu_2^p - \mu_1^p - 1$ at p ,
- 3) if $p \in \Delta$, then θ has cokernel of length $\lfloor \mu_2^p \rfloor - \lfloor \mu_1^p \rfloor$.

The following theorem follows easily.

Theorem 3.2. *Suppose that L underlies a family of elliptic curves over \mathbb{P}^1 . Let a_L denote the number of points in Δ at which the local monodromy of L is strictly quasi-unipotent, and let Δ_a denote the set of points at which L has an apparent singularity. Then*

$$\deg \mathcal{E}^{0,1} = \frac{1}{2} \left(2 - a_L + \sum_{p \in \Delta} (\lfloor \mu_2^p \rfloor - \lfloor \mu_1^p \rfloor - 1) + \sum_{p \in \Delta_a} (\mu_2^p - \mu_1^p - 1) \right).$$

Remark 3.3. If we note that the characteristic exponents μ_i^p are integers when p is nonsingular or an apparent singularity, and that $\mu_2^p = \mu_1^p + 1$ when p is a non-singular point, then we can simplify the formula in Theorem 3.2 to

$$\deg \mathcal{E}^{0,1} = \frac{1}{2} \left(2 - a_L + \sum_{p \in \mathbb{P}^1} (\lfloor \mu_2^p \rfloor - \lfloor \mu_1^p \rfloor - 1) \right).$$

However, the form given in Theorem 3.2 is simpler to work with from a computational standpoint, as both sums are clearly finite.

Proof of Theorem 3.2. By Example 2.2, we see that

$$\deg \mathcal{E}^{1,0} + \deg \mathcal{E}^{0,1} = -a_L,$$

and Corollary 3.1 gives that the map θ has cokernel of length

$$\sum_{p \in \Delta_a} (\mu_2^p - \mu_1^p - 1) + \sum_{p \in \Delta} (\lfloor \mu_2^p \rfloor - \lfloor \mu_1^p \rfloor).$$

Since $\deg(\mathcal{E}^{0,1} \otimes \Omega_{\mathbb{P}^1}(\Delta)) = |\Delta| - 2 + \deg \mathcal{E}^{0,1}$, it follows that

$$\deg \mathcal{E}^{0,1} + |\Delta| - 2 = -\deg \mathcal{E}^{0,1} - a_L + \sum_{p \in \Delta_a} (\mu_2^p - \mu_1^p - 1) + \sum_{p \in \Delta} (\lfloor \mu_2^p \rfloor - \lfloor \mu_1^p \rfloor).$$

Thus

$$\deg \mathcal{E}^{0,1} = \frac{1}{2} \left(2 - a_L + \sum_{p \in \Delta_a} (\mu_2^p - \mu_1^p - 1) + \sum_{p \in \Delta} (\lfloor \mu_2^p \rfloor - \lfloor \mu_1^p \rfloor - 1) \right)$$

as required. ■

The term $(\mu_2^p - \mu_1^p - 1)$ should be thought of as an invariant which counts the ramification of the j -function at $p \in \Delta_a$; Doran makes this precise in [6].

Remark 3.4. From this, we may use Theorem 2.9 to compute $h^{0,2}$ for the parabolic cohomology of our elliptic fibration. In [2], Cox and Zucker show that this $h^{0,2}$ computes the dimension of a particular space of modular forms. The computations in this section can be thought of as giving a way to compute the dimension of this space of modular forms using the Picard–Fuchs equation of the elliptic surface.

4 Families of K3 surfaces

Next we turn our attention to families of K3 surfaces, as studied in Example 2.3. Theorem 2.7 can be rephrased in the context of the corresponding $(1, 1, 1)$ -variations of Hodge structure, as follows. Recall that in this case we have the Kodaira–Spencer maps $\theta_0: \mathcal{E}^{2,0} \rightarrow \mathcal{E}^{1,1} \otimes \Omega_{\mathbb{P}^1}(\Delta)$ and $\theta_1: \mathcal{E}^{1,1} \rightarrow \mathcal{E}^{0,2} \otimes \Omega_{\mathbb{P}^1}(\Delta)$.

Corollary 4.1. *If L is a Picard Fuchs equation corresponding to a $(1, 1, 1)$ -type variation of Hodge structure, then*

- 1) *if p is a nonsingular point of L , then θ_0 is an isomorphism,*
- 2) *if p is an apparent singularity of L , then θ_0 has cokernel of length $\mu_2^p - \mu_1^p - 1$ at p ,*
- 3) *if p is a regular singular point of L , then θ_0 has cokernel of length $\lfloor \mu_2^p \rfloor - \lfloor \mu_1^p \rfloor$.*

Let $U = \mathbb{P}^1 \setminus \Delta$ denote a Zariski open set and let $j: U \hookrightarrow \mathbb{P}^1$ denote the embedding. Suppose that $\mathcal{X} \rightarrow U$ is an M -polarized family of K3 surfaces, in the sense of [7, Definition 2.1], for some rank 19 lattice M . The transcendental lattices of the fibres of \mathcal{X} define a VHS \mathbb{V} of type $(1, 1, 1)$ over U ; we assume that all of the monodromy matrices of this VHS take the forms T_x considered in Example 2.3. Finally, let L denote the Picard–Fuchs equation underlying this $(1, 1, 1)$ -VHS.

By Lemma 2.1, it is easy to see that $\deg_{\text{par}} \mathcal{E}^{1,1} = 0$, so the degree of $\mathcal{E}^{1,1}$ can be computed easily using equation (2.1): indeed, if $a_{1/2}$ is the number of points p at which monodromy is of type T_{un} , T_i , T_ω or T_{nod} (see Example 2.3), then $\deg \mathcal{E}^{1,1} = -\frac{1}{2}a_{1/2}$. We let a_f denote the number of points whose monodromy matrices look like T_{un} , T_i , $-T_i$, T_i^2 , T_ω and T_ω^2 . Finally, let Δ_a denote the set of points at which L has an apparent singularity.

Proposition 4.2. *If L underlies a family of K3 surfaces $\mathcal{X} \rightarrow U$ as above, then we have*

$$\deg(\mathcal{E}^{0,2}) = 2 + \frac{1}{2}a_{1/2} - a_f + \sum_{p \in \Delta_a} (\mu_2^p - \mu_1^p - 1) + \sum_{p \in \Delta} (\lfloor \mu_2^p \rfloor - \lfloor \mu_1^p \rfloor - 1).$$

Proof. This is proved by a similar calculation to Theorem 3.2. ■

Remark 4.3. As in the case of elliptic curves, we may simplify this formula to the shorter, but less transparent, form

$$\deg(\mathcal{E}^{0,2}) = 2 + \frac{1}{2}a_{1/2} - a_f + \sum_{p \in \mathbb{P}^1} (\lfloor \mu_2^p \rfloor - \lfloor \mu_1^p \rfloor - 1).$$

Using this, one can easily use Theorem 2.9 to compute $h^{0,3}$ of the Hodge structure on $H^1(\mathbb{P}^1, j_* \mathbb{V})$.

Corollary 4.4. *The Hodge number $h^{0,3}$ of the Hodge structure on $H^1(\mathbb{P}^1, j_* \mathbb{V})$ is*

$$h^{0,3} = h^0 \left(\mathbb{P}^1, \mathcal{O}_{\mathbb{P}^1}^1 \left(a_f - 4 - \frac{1}{2}a_{1/2} - \sum_{p \in \Delta_a} (\mu_2^p - \mu_1^p - 1) - \sum_{p \in \Delta} (\lfloor \mu_2^p \rfloor - \lfloor \mu_1^p \rfloor - 1) \right) \right).$$

Now we may ask ourselves: when can such a fibration by K3 surfaces have a Calabi–Yau threefold total space?

Definition 4.5. A weight n Hodge structure with Hodge numbers $h^{p,q}$ is called *Calabi–Yau* if $h^{n,0} = 1$.

If a variety X is Calabi–Yau of dimension n , then the Hodge structure on $H^n(X, \mathbb{Q})$ is certainly Calabi–Yau, but the converse is not true. For example, look at any blow-up of an honest Calabi–Yau variety, or a Kulikov surface.

According to Corollary 4.4, a family of K3 surfaces $\mathcal{X} \rightarrow U$ has Calabi–Yau Hodge structure on $H^1(\mathbb{P}^1, j_* \mathbb{V})$ if and only if

$$a_f = 4 + \frac{1}{2}a_{1/2} + \sum_{p \in \Delta_a} (\mu_2^p - \mu_1^p - 1) + \sum_{p \in \Delta} (\lfloor \mu_2^p \rfloor - \lfloor \mu_1^p \rfloor - 1). \quad (4.1)$$

By the degeneration of the Leray spectral sequence [31, Section 15], for any compactification $\overline{\mathcal{X}}$ of \mathcal{X} , this parabolic cohomology group is a direct summand of the Hodge structure on $H^3(\overline{\mathcal{X}}, \mathbb{Q})$, and its complement has no $(3, 0)$ or $(0, 3)$ part (see [10]). Thus, if equation (4.1) does not hold, then $\mathcal{X} \rightarrow \mathbb{P}^1$ has no compactification which is a Calabi–Yau threefold. In [10] we will study the converse problem in specific examples.

4.1 Base change of a hypergeometric family

We illustrate these results with an example regarding families of K3 surfaces and their pullbacks. Let us take the family of mirror quartic K3 surfaces written as the compactifications in \mathbb{P}^3 of the fibres of the Laurent polynomial

$$f(x, y, z) = \frac{(x + y + z + 1)^4}{xyz}.$$

Narumiya and Shiga [27, Section 5] show that this family has Picard–Fuchs equation which is hypergeometric and may be written as

$$\delta^3 + t \left(\delta + \frac{1}{4} \right) \left(\delta + \frac{1}{2} \right) \left(\delta + \frac{3}{4} \right).$$

The local monodromies around its singular points are of type T_{un}^2 , T_i , and T_{nod} , over 0, ∞ , and 1 respectively, so the Riemann scheme of this differential operator is

$$\left\{ \begin{array}{ccc} 0 & \infty & 1 \\ 0 & 1/4 & 0 \\ 0 & 1/2 & 1/2 \\ 0 & 3/4 & 1 \end{array} \right\}.$$

Now consider the base change of this local system along a map $g: \mathbb{P}^1 \rightarrow \mathbb{P}^1$. Let $\mathcal{E}_g^{i,j}$ denote the appropriate Hodge bundles of the pulled-back local system $g^*\mathbb{V}$ and let $h_g^{i,j}$ be the Hodge numbers of its parabolic cohomology. Let d be the degree of g , and set k, ℓ, m to be the numbers of points over 0, ∞ and 1 respectively. Let

$$r = \sum_{p \in \mathbb{P}^1 \setminus \{0, 1, \infty\}} (e_p - 1)$$

denote the degree of ramification of g away from $\{0, 1, \infty\}$. We will write $[y_1, \dots, y_\ell]$ for the partition of d encoding the ramification profile over ∞ . Finally, define

$$D_g := a_f - a_{1/2} - 4 - \sum_{p \in \Delta_a} (\mu_2^p - \mu_1^p - 1) - \sum_{p \in \Delta} (\lfloor \mu_2^p \rfloor - \lfloor \mu_1^p \rfloor - 1),$$

where all terms in this expression are computed for the pulled-back local system $g^*\mathbb{V}$. Then we have the following proposition.

Proposition 4.6. *With assumptions and notation as above,*

$$D_g = a_f - 2 + \sum_{i=1}^{\ell} \left\lfloor \frac{y_i}{4} \right\rfloor.$$

In particular, $h_g^{0,3} = h^0(\mathbb{P}^1, \mathcal{O}_{\mathbb{P}^1}(D_g))$.

Proof. It follows from the definition that a_f is equal to the number of points over 0 with ramification of order not divisible by 4. Let ℓ_{odd} and m_{odd} denote the number of points over 0 and 1, respectively, where the order of ramification is odd. Then

$$a_{1/2} = m_{\text{odd}} + \ell_{\text{odd}}.$$

Apparent singularities appear only at ramification points of g away from points over $\{0, 1, \infty\}$, and each ramification point has $\mu_2^p - \mu_1^p - 1$ equal to $e_p - 1$. Therefore

$$\sum_{\Delta_a} (\mu_2^p - \mu_1^p - 1) = \sum_{p \in \mathbb{P}^1 \setminus \{0, 1, \infty\}} (e_p - 1) = r.$$

We split the set Δ into three components, corresponding to points over 0, 1, and ∞ respectively, and denoted Δ_0 , Δ_1 , and Δ_∞ . The corresponding sums are

$$\begin{aligned} \sum_{p \in \Delta_0} (\lfloor \mu_2^p \rfloor - \lfloor \mu_1^p \rfloor - 1) &= \frac{1}{2}(d - \ell_{\text{odd}}) - \sum_{i=1}^{\ell} \left(\lfloor \frac{y_i}{4} \rfloor \right) - \ell, \\ \sum_{p \in \Delta_1} (\lfloor \mu_2^p \rfloor - \lfloor \mu_1^p \rfloor - 1) &= \frac{1}{2}(d - m_{\text{odd}}) - m, \\ \sum_{p \in \Delta_\infty} (\lfloor \mu_2^p \rfloor - \lfloor \mu_1^p \rfloor - 1) &= k. \end{aligned}$$

Substituting everything into D_g , we obtain

$$D_g = a_f - 4 - d - r + k + \ell + m + \sum_{i=1}^{\ell} \left\lfloor \frac{y_i}{4} \right\rfloor.$$

Now, since g is a map from \mathbb{P}^1 to \mathbb{P}^1 , the Riemann–Hurwitz formula gives

$$2d - 2 = \sum_p (e_p - 1) = r + (d - \ell) + (d - m) + (d - k),$$

which implies that $k + \ell + m - d - r = 2$. Making this substitution, we find that

$$D_g = a_f - 2 + \sum_{i=1}^{\ell} \left\lfloor \frac{y_i}{4} \right\rfloor,$$

which proves the proposition. ■

As a corollary, we can approximately recover [9, Proposition 2.4] (see also [10] for some further applications of this result).

Corollary 4.7. *The Hodge structure of $g^*\mathbb{V}$ is Calabi–Yau if and only if one of the following two conditions holds:*

- 1) *the number ℓ is equal to 2 and $1 \leq y_1, y_2 \leq 4$; or*
- 2) *the number ℓ is equal to 1 and $5 \leq y_1 \leq 8$.*

Proof. By Corollary 4.4 and Proposition 4.6, we may assume that

$$a_f + \sum_{i=1}^{\ell} \left\lfloor \frac{y_i}{4} \right\rfloor = 2.$$

Therefore, we have at most two points over 0. If we have two points over 0, then they must both have $1 \leq y_i \leq 4$. If there is a single point over 0, then $5 \leq y_1 \leq 8$. ■

5 Families of Calabi–Yau threefolds

In this section, we will first apply the results of Section 2.4 to complete the computations of [3] and [23]. Then we will determine the Hodge numbers of some specific fourfolds fibred by quintic mirror Calabi–Yau threefolds.

5.1 Inhomogeneous Picard–Fuchs equations and Calabi–Yau threefolds

In this section, we will perform some computations related to work of del Angel, Müller-Stach, van Straten, and Zuo [3]. These authors were interested in determining situations in which families of Calabi–Yau threefolds admit normal functions; here a *normal function* is an algebraic cycle on each fibre of a family of varieties, which is homologous to zero on all fibres.

In particular, if we have some algebraic $(2, 2)$ -class Z on a family of Calabi–Yau threefolds, whose restriction to each fibre is homologous to zero, then Z determines a multivalued function Φ_Z on the base of the fibration. From such a normal function, one obtains an inhomogeneous extension of the Picard–Fuchs equation L_σ which is satisfied by Φ_Z . In other words,

$$L_\sigma \Phi_Z = g_Z,$$

for some holomorphic function g_Z . Morrison and Walcher [26, 30] studied a specific such normal function on the Fermat family of quintics, and showed that Φ_Z can be interpreted as the domain wall tension of certain D -branes.

The computations in [3] aimed to determine whether such normal functions exist on other 1-parameter families of Calabi–Yau threefolds. Evidence for the existence of normal functions can be found by determining whether the parabolic cohomology of L_σ has classes of type $(2, 2)$ or not. They looked at the 14 hypergeometric examples of Doran and Morgan [12] and their pullbacks along maps $[t : s] \mapsto [t^d : s^d]$, for certain values of d , in an effort to determine whether, after such base change, the 14 families could admit interesting normal functions.

In both [3] and the subsequent work of Hollborn and Müller-Stach [23], the authors were not able to compute the degrees of the bundles $\mathcal{E}^{2,1}$ in many situations, and therefore were not able to determine the Hodge numbers of the parabolic cohomology groups that they were interested in. By applying the results of Section 2.4, we are now able to finish their computations.

Example 5.1. Let's compute the Hodge numbers of the VHS obtained from the second hypergeometric example in the list of Doran and Morgan [12]. This example is obtained in the following way. Begin with the hypergeometric operator

$$\delta^4 - t \left(\delta + \frac{1}{10} \right) \left(\delta + \frac{3}{10} \right) \left(\delta + \frac{7}{10} \right) \left(\delta + \frac{9}{10} \right).$$

This operator underlies a $(1, 1, 1, 1)$ -type rational VHS, and has local monodromy matrices

$$T_0 = \begin{pmatrix} 1 & 1 & 0 & 0 \\ 0 & 1 & 1 & 0 \\ 0 & 0 & 1 & 1 \\ 0 & 0 & 0 & 1 \end{pmatrix}, \quad T_1 = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 1 \\ 0 & 0 & 0 & 1 \end{pmatrix}, \quad T_\infty = \begin{pmatrix} \zeta_{10} & 0 & 0 & 0 \\ 0 & \zeta_{10}^3 & 0 & 0 \\ 0 & 0 & \zeta_{10}^7 & 0 \\ 0 & 0 & 0 & \zeta_{10}^9 \end{pmatrix},$$

where ζ_{10} is a primitive tenth root of unity. Its Riemann scheme is

$$\left\{ \begin{pmatrix} 0 & 1 & \infty \\ 0 & 0 & 1/10 \\ 0 & 1 & 3/10 \\ 0 & 1 & 7/10 \\ 0 & 2 & 9/10 \end{pmatrix} \right\}.$$

Let us take the base change of this differential equation under the map of degree 5 which is totally ramified over 0 and ∞ . The new differential equation has seven singular points and

Riemann scheme

$$\left\{ \begin{array}{ccc} 0 & t^5 - 1 = 0 & \infty \\ 0 & 0 & 1/2 \\ 0 & 1 & 3/2 \\ 0 & 1 & 7/2 \\ 0 & 2 & 9/2 \end{array} \right\}.$$

The points over 1 all have monodromy of type T_1 , the point over 0 has monodromy of type T_0 , and the point over ∞ has monodromy matrix equal to $-\text{Id}_4$. Therefore, the only point with nontrivial parabolic filtration is ∞ and the only filtrands are of weight $1/2$. Thus equation (2.1) and Lemma 2.1 give

$$\deg \mathcal{E}^{1,2} + \deg \mathcal{E}^{2,1} = -1.$$

By Theorem 2.7, the map $\theta_1: \mathcal{E}^{2,1} \rightarrow \mathcal{E}^{1,2} \otimes \Omega_{\mathbb{P}^1}(\Delta)$ has cokernel of length 2. Therefore,

$$\deg \mathcal{E}^{2,1} - \deg \mathcal{E}^{1,2} = 3.$$

Hence $\deg \mathcal{E}^{2,1} = 1$. This gives the correct b value in Theorem 2.11, which in turn allows us to complete the corresponding row of the table in [23, Section 4]; we obtain that the Hodge numbers of g^*L are $(0, 1, 2, 1, 0)$.

Similarly, if we take g to be the degree 10 map ramified completely at 0 and ∞ , then

$$\deg \mathcal{E}^{1,2} + \deg \mathcal{E}^{2,1} = 0.$$

In this case θ_1 has cokernel of length 4, so

$$\deg \mathcal{E}^{2,1} - \deg \mathcal{E}^{1,2} = 6,$$

and therefore $\deg \mathcal{E}^{2,1} = 3$. Once again, this allows us to complete the corresponding row of the table in [23, Section 4]; we obtain that the Hodge numbers of g^*L are $(0, 1, 3, 1, 0)$.

Using this, we may complete the calculations left unfinished by [3] and [23]. Table 3 is a reproduction of the table from [23, Section 4], with all missing pieces completed; our new results are highlighted in gray. Our notation is as follows. # is the case number from the table in [23, Section 4], and the numbers $(\alpha_1, \alpha_2, \alpha_3, \alpha_4)$ are the numbers determining the corresponding hypergeometric differential equations

$$L = \delta^4 - t(\delta + \alpha_1)(\delta + \alpha_2)(\delta + \alpha_3)(\delta + \alpha_4).$$

The number d indicates that we have pulled back by the covering $[t : s] \mapsto [t^d : s^d]$, to obtain a VHS over \mathbb{P}^1 with Hodge bundles $\mathcal{E}^{3,0}$ of degree a and $\mathcal{E}^{2,1}$ of degree b . We then apply the results of Section 2.4 to compute the Hodge numbers of the corresponding parabolic cohomology groups in the final column.

Remark 5.2. Recently there has been a great deal of interest in studying the monodromy subgroups of $\text{Sp}(4, \mathbb{R})$ in the 14 hypergeometric examples of Doran and Morgan [12]. These subgroups may be either arithmetic or non-arithmetic (more commonly called *thin*). Singh and Venkataramana [28, 29] have proved that the monodromy subgroup is arithmetic in seven of the 14 cases, and Brav and Thomas [1] have proved that it is thin in the remaining seven. The arithmetic/thin status of each of the examples is given in the “Monodromy” column of Table 3.

Computations carried out in [13] suggest that the arithmetic/thin dichotomy for monodromy groups of hypergeometric local systems is closely related to the value of the *Lyapunov exponents*

of the corresponding variations of Hodge structure. Lyapunov exponents are dynamical invariants associated to a flat vector bundle over a hyperbolic Riemann surface, that are computed using the monodromy of the flat bundle as well as the metric structure on the hyperbolic curve.

The main goal of [13] is to study these Lyapunov exponents; [13, Lemma 6.3] (restated above as Theorem 2.7) is a step towards this. In [13] it is conjectured that for the 14 hypergeometric variations of Hodge structure of Doran and Morgan [12], the associated Lyapunov exponents are closely related to the parabolic degrees of the bundles $\mathcal{E}^{3,0}$ and $\mathcal{E}^{2,1}$, provided that these degrees satisfy a certain relation (see [13, Conjecture 6.4] and [15] for the higher rank case).

It is not clear whether there is any relation between the computations presented here and Lyapunov exponents; however, it would be interesting to see whether there is any relation between Lyapunov exponents and the asymptotic growth of $\deg_{\text{par}} \mathcal{E}^{3,0}$ and $\deg_{\text{par}} \mathcal{E}^{2,1}$ as the degree d of the covering $[t : s] \mapsto [t^d : s^d]$ grows. Such growth has been investigated in the rank 2 case by Kappes [24].

Finally we note that, according to work of Doran and Malmendier [11], most of Doran's and Morgan's [12] 14 hypergeometric local systems can be obtained from hypergeometric local systems of rank 3. These rank 3 hypergeometric local systems underlie variations of Hodge structure corresponding to families of K3 surfaces, for which the Lyapunov exponents have been computed by Filip [14]. It seems plausible that the iterated construction of [11] along with the computations of [14] could allow one to compute Lyapunov exponents for the fourteen rank 4 hypergeometric local systems of [12].

Remark 5.3. We pose the question as to whether the above Hodge structures are reducible or not. It is possible that the $h^{2,2}$ value in each of the sets of Hodge numbers in the above table all correspond to rational Hodge classes, thus induce normal functions on the corresponding families of Calabi–Yau threefolds. The maps along which we have pulled back are rigid in the sense that any deformation of this map will change the Hodge numbers, so this indeed seems possible.

5.2 Calabi–Yau fourfolds fibred by mirror quintics

Finally, we will apply the same approach to classifying Calabi–Yau fourfolds fibred by Calabi–Yau threefolds as we did with Calabi–Yau threefolds fibred by K3 surfaces. We do not obtain concrete classification results, since modularity results for families of Calabi–Yau varieties of higher dimension are scarce, but we are at least able to place bounds on the space of possibilities.

We start with the mirror quintic family of K3 surfaces. This family can be represented as a compactification of the fibres of the Laurent polynomial

$$\frac{(x + y + z + w + 1)^5}{xyzw}.$$

The periods of this family satisfy the hypergeometric differential equation L

$$\delta^4 - t \left(\delta + \frac{1}{5} \right) \left(\delta + \frac{2}{5} \right) \left(\delta + \frac{3}{5} \right) \left(\delta + \frac{4}{5} \right),$$

which has local monodromy matrices

$$T_0 = \begin{pmatrix} 1 & 1 & 0 & 0 \\ 0 & 1 & 1 & 0 \\ 0 & 0 & 1 & 1 \\ 0 & 0 & 0 & 1 \end{pmatrix}, \quad T_1 = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 1 \\ 0 & 0 & 0 & 1 \end{pmatrix}, \quad T_\infty = \begin{pmatrix} \zeta_5 & 0 & 0 & 0 \\ 0 & \zeta_5^2 & 0 & 0 \\ 0 & 0 & \zeta_5^3 & 0 \\ 0 & 0 & 0 & \zeta_5^4 \end{pmatrix},$$

#	$(\alpha_1, \alpha_2, \alpha_3, \alpha_4)$	d	a	b	Hodge numbers	Monodromy
1	$(1/5, 2/5, 3/5, 4/5)$	1	0	0	$(0, 0, 0, 0, 0)$	Thin
		2	0	0	$(0, 0, 1, 0, 0)$	
		5	1	2	$(0, 0, 0, 0, 0)$	
		10	2	4	$(1, 1, 1, 1, 1)$	
2	$(1/10, 3/10, 7/10, 9/10)$	1	0	0	$(0, 0, 0, 0, 0)$	Arithmetic
		2	0	0	$(0, 0, 1, 0, 0)$	
		5	0	1	$(0, 1, 2, 1, 0)$	
		10	1	3	$(0, 1, 3, 1, 0)$	
3	$(1/2, 1/2, 1/2, 1/2)$	1	0	0	$(0, 0, 0, 0, 0)$	Thin
		2	1	1	$(0, 0, 0, 0, 0)$	
		$2k$	k	k	$(k-1, 0, 0, 0, k-1)$	
4	$(1/3, 1/3, 2/3, 2/3)$	1	0	0	$(0, 0, 0, 0, 0)$	Arithmetic
		2	0	0	$(0, 0, 1, 0, 0)$	
		3	1	1	$(0, 0, 0, 0, 0)$	
		6	2	2	$(1, 0, 1, 0, 1)$	
5	$(1/3, 1/2, 1/2, 2/3)$	1	0	0	$(0, 0, 0, 0, 0)$	Thin
		6	1	2	$(0, 0, 2, 0, 0)$	
6	$(1/4, 1/2, 1/2, 3/4)$	1	0	0	$(0, 0, 0, 0, 0)$	Thin
		4	1	2	$(0, 0, 0, 0, 0)$	
		8	2	4	$(1, 1, 0, 1, 1)$	
7	$(1/8, 3/8, 5/8, 7/8)$	1	0	0	$(0, 0, 0, 0, 0)$	Thin
		2	0	0	$(0, 0, 1, 0, 0)$	
		4	0	1	$(0, 1, 1, 1, 0)$	
		8	1	3	$(0, 1, 1, 1, 0)$	
8	$(1/6, 1/3, 2/3, 5/6)$	1	0	0	$(0, 0, 0, 0, 0)$	Arithmetic
		2	0	0	$(0, 0, 1, 0, 0)$	
		6	1	2	$(0, 0, 1, 0, 0)$	
9	$(1/12, 5/12, 7/12, 11/12)$	1	0	0	$(0, 0, 0, 0, 0)$	Thin
		2	0	0	$(0, 0, 1, 0, 0)$	
		3	0	1	$(0, 1, 0, 1, 0)$	
		4	0	1	$(0, 1, 1, 1, 0)$	
		6	0	2	$(0, 2, 1, 2, 0)$	
		12	1	5	$(0, 3, 1, 3, 0)$	
10	$(1/4, 1/4, 3/4, 3/4)$	1	0	0	$(0, 0, 0, 0, 0)$	Arithmetic
		2	0	0	$(0, 0, 1, 0, 0)$	
		4	1	1	$(0, 0, 1, 0, 0)$	
		8	2	2	$(1, 0, 3, 0, 1)$	
11	$(1/6, 1/4, 3/4, 5/6)$	1	0	0	$(0, 0, 0, 0, 0)$	Arithmetic
		2	0	0	$(0, 0, 1, 0, 0)$	
		12	1	3	$(0, 1, 5, 1, 0)$	
12	$(1/4, 1/3, 2/3, 3/4)$	1	0	0	$(0, 0, 0, 0, 0)$	Arithmetic
		2	0	0	$(0, 0, 1, 0, 0)$	
		3	0	1	$(0, 1, 0, 1, 0)$	
		12	1	4	$(0, 2, 3, 2, 0)$	
13	$(1/6, 1/6, 5/6, 5/6)$	1	0	0	$(0, 0, 0, 0, 0)$	Arithmetic
		2	0	0	$(0, 0, 1, 0, 0)$	
		3	0	0	$(0, 0, 2, 0, 0)$	
		6	1	1	$(0, 0, 3, 0, 0)$	
14	$(1/6, 1/2, 1/2, 5/6)$	1	0	0	$(0, 0, 0, 0, 0)$	Thin
		3	0	1	$(0, 1, 0, 1, 0)$	
		6	1	3	$(0, 1, 0, 1, 0)$	

Table 3. Invariants for families of Calabi–Yau threefolds.

and Riemann scheme

$$\left\{ \begin{array}{ccc} 0 & 1 & \infty \\ 0 & 0 & 1/5 \\ 0 & 1 & 2/5 \\ 0 & 1 & 3/5 \\ 0 & 2 & 4/5 \end{array} \right\}.$$

We can ask: along which maps $g: \mathbb{P}^1 \rightarrow \mathbb{P}^1$ does the pull-back g^*L have Hodge number $h^{0,4} = 1$? In other words, when can the fourfold total space obtained from the mirror quintic family be a Calabi–Yau fourfold?

We have all the tools to perform this calculation. Let us take a map $g: \mathbb{P}^1 \rightarrow \mathbb{P}^1$ of degree d . Let k , ℓ , and m denote the numbers of points over 0 , ∞ , and 1 respectively and let $[y_1, \dots, y_\ell]$ denote the partition of d encoding the ramification profile over ∞ . Finally, let ℓ_0 be the number of values of i such that $5 \nmid y_i$. We compute the Hodge numbers of the parabolic cohomology associated to g^*L .

Proposition 5.4. *The degrees of the Hodge bundles of the variation of Hodge structure on g^*L are given by*

$$\begin{aligned} \deg \mathcal{E}^{3,0} &= \frac{1}{2} \left(d - \ell_0 - \sum_{i=1}^{\ell} \left(\lfloor \frac{3y_i}{5} \rfloor - \lfloor \frac{2y_i}{5} \rfloor \right) \right) - \sum_{i=1}^{\ell} \left(\lfloor \frac{2y_i}{5} \rfloor - \lfloor \frac{y_i}{5} \rfloor \right) \\ \deg \mathcal{E}^{2,1} &= \frac{1}{2} \left(d - \ell_0 - \sum_{i=1}^{\ell} \left(\lfloor \frac{3y_i}{5} \rfloor - \lfloor \frac{2y_i}{5} \rfloor \right) \right), \end{aligned}$$

and the equations $\deg \mathcal{E}^{3,0} + \deg \mathcal{E}^{0,3} = -\ell_0$ and $\deg \mathcal{E}^{2,1} + \deg \mathcal{E}^{1,2} = -\ell_0$.

Remark 5.5. From this data, one can compute all Hodge numbers of the parabolic cohomology groups, by applying the results of Section 2.4.

Proof. By Theorem 2.7, the map $\theta_1: \mathcal{E}^{2,1} \rightarrow \mathcal{E}^{1,2} \otimes \Omega_{\mathbb{P}^1}(\Delta)$ has cokernel of length

$$r + \sum_{i=1}^{\ell} \left(\lfloor \frac{3y_i}{5} \rfloor - \lfloor \frac{2y_i}{5} \rfloor \right) - (\ell - \ell_0)$$

(note that $(\ell - \ell_0)$ counts the number of values of i so that $5 \mid y_i$). Thus

$$\deg \mathcal{E}^{2,1} - \deg \mathcal{E}^{1,2} = k + \ell + m - 2 - r - \sum_{i=1}^{\ell} \left(\lfloor \frac{3y_i}{5} \rfloor - \lfloor \frac{2y_i}{5} \rfloor \right).$$

Applying the Riemann–Hurwitz formula, we see that

$$2d - 2 = \sum_p (e_p - 1) = r + (d - k) + (d - \ell) + (d - m), \quad (5.1)$$

which implies that $d = k + \ell + m - r - 2$. Therefore,

$$\deg \mathcal{E}^{2,1} - \deg \mathcal{E}^{1,2} = d - \sum_{i=1}^{\ell} \left(\lfloor \frac{3y_i}{5} \rfloor - \lfloor \frac{2y_i}{5} \rfloor \right).$$

Equation (2.1) and Lemma 2.1 also give that $\deg \mathcal{E}^{3-i,i} + \deg \mathcal{E}^{i,3-i} = -\ell_0$. So we conclude that

$$\deg \mathcal{E}^{2,1} = \frac{1}{2} \left(d - \ell_0 - \sum_{i=1}^{\ell} \left(\lfloor \frac{3y_i}{5} \rfloor - \lfloor \frac{2y_i}{5} \rfloor \right) \right).$$

We also have that the cokernel of the map $\theta_0: \mathcal{E}^{3,0} \rightarrow \mathcal{E}^{2,1} \otimes \Omega_{\mathbb{P}^1}(\Delta)$ has length

$$r + d + \sum_{i=1}^{\ell} \left(\lfloor \frac{2y_i}{5} \rfloor - \lfloor \frac{y_i}{5} \rfloor \right) - (\ell - \ell_0),$$

thus

$$\deg \mathcal{E}^{3,0} - \deg \mathcal{E}^{2,1} = k + \ell + m - 2 - r - d - \sum_{i=1}^{\ell} \left(\lfloor \frac{2y_i}{5} \rfloor - \lfloor \frac{y_i}{5} \rfloor \right) = - \sum_{i=1}^{\ell} \left(\lfloor \frac{2y_i}{5} \rfloor - \lfloor \frac{y_i}{5} \rfloor \right),$$

where the second line follows by applying equation (5.1). Adding the expression for $\deg \mathcal{E}^{2,1}$ obtained above, we get

$$\deg \mathcal{E}^{3,0} = \frac{1}{2} \left(d - \ell_0 - \sum_{i=1}^{\ell} \left(\lfloor \frac{3y_i}{5} \rfloor - \lfloor \frac{2y_i}{5} \rfloor \right) \right) - \sum_{i=1}^{\ell} \left(\lfloor \frac{2y_i}{5} \rfloor - \lfloor \frac{y_i}{5} \rfloor \right),$$

which completes the proof of the theorem. ■

Since we have that $\deg \mathcal{E}^{3,0} + \deg \mathcal{E}^{0,3} = -\ell_0$, it follows that

$$\deg \mathcal{E}^{0,3} = -\frac{1}{2} \left(d + \ell_0 - \sum_{i=1}^{\ell} \left(\lfloor \frac{3y_i}{5} \rfloor - \lfloor \frac{2y_i}{5} \rfloor \right) \right) + \sum_{i=1}^{\ell} \left(\lfloor \frac{2y_i}{5} \rfloor - \lfloor \frac{y_i}{5} \rfloor \right).$$

Therefore,

Corollary 5.6. *The Hodge structure on g^*L is Calabi–Yau if and only if*

$$\frac{1}{2} \left(d + \ell_0 - \sum_{i=1}^{\ell} \left(\lfloor \frac{3y_i}{5} \rfloor - \lfloor \frac{2y_i}{5} \rfloor \right) \right) - \sum_{i=1}^{\ell} \left(\lfloor \frac{2y_i}{5} \rfloor - \lfloor \frac{y_i}{5} \rfloor \right) = 2.$$

Remark 5.7. Note that the value of $h^{0,3}$ of the parabolic cohomology only depends on the structure of ramification over the point ∞ , just as in the case of quartic threefolds.

A consequence of this corollary is that there is a finite and enumerable set of values for $[y_1, \dots, y_{\ell}]$ so that g^*L is Calabi–Yau. This set strictly contains the possible ramification data over ∞ under which the total space of the pullback of the family of mirror quintics is Calabi–Yau. The list of all values of $[y_1, \dots, y_{\ell}]$ are as follows:

$$\begin{aligned} &[1, 1], [1, 2], [1, 3], [1, 4], [1, 5], [2, 2], [2, 3], [2, 4], [2, 5], \\ &[3, 3], [3, 4], [3, 5], [4, 4], [4, 5], [5, 5], [6], [7], [8], [9], [10]. \end{aligned}$$

We expect that a proper subset of these ramification indices should correspond to maps along which we can pull back the family of mirror quintics and get a variety that admits a smooth (or terminal) Calabi–Yau resolution.

One can also check that the only cases in which the Hodge number $h^{3,0}$ of the parabolic cohomology is 0 are the cases where $\ell = 1$ and $y_1 \in \{1, 2, 3, 4, 5\}$.

Remark 5.8. Similar calculations can presumably be completed for any family of Calabi–Yau threefolds realising one of the fourteen hypergeometric variations of Hodge structure of type $(1, 1, 1, 1)$ classified by Doran and Morgan [12]. This includes the so-called “mirror twin” families [12, Section 3], which have the same Picard–Fuchs equations and \mathbb{R} -VHS’s as some well-known toric examples, but have different underlying \mathbb{Z} -VHS’s and Hodge number $h^{1,1}$. In particular, the calculations above apply equally to the “quintic mirror twin” family, which shares its Picard–Fuchs equation and \mathbb{R} -VHS with the quintic mirror example above.

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References

- [1] Brav C., Thomas H., Thin monodromy in $\mathrm{Sp}(4)$, *Compos. Math.* **150** (2014), 333–343, [arXiv:1210.0523](#).
- [2] Cox D.A., Zucker S., Intersection numbers of sections of elliptic surfaces, *Invent. Math.* **53** (1979), 1–44.
- [3] Del Angel P.L., Müller-Stach S., Van Straten D., Zuo K., Hodge classes associated to 1-parameter families of Calabi–Yau 3-folds, *Acta Math. Vietnam.* **35** (2010), 7–22, [arXiv:0911.0277](#).
- [4] Deligne P., Équations différentielles à points singuliers réguliers, *Lecture Notes in Math.*, Vol. 163, Springer-Verlag, Berlin – New York, 1970.
- [5] Doran B., Doran C.F., Harder A., Picard–Fuchs uniformization of modular subvarieties, in Uniformization, Riemann–Hilbert Correspondence, Calabi–Yau Manifolds, and Picard–Fuchs Equations (Institut Mittag-Leffler, 2017), Editors L. Ji, S.-T. Yau, to appear.
- [6] Doran C.F., Picard–Fuchs uniformization: modularity of the mirror map and mirror-moonshine, in The Arithmetic and Geometry of Algebraic Cycles (Banff, AB, 1998), *CRM Proc. Lecture Notes*, Vol. 24, Amer. Math. Soc., Providence, RI, 2000, 257–281, [math.AG/9812162](#).
- [7] Doran C.F., Harder A., Novoseltsev A.Y., Thompson A., Families of lattice polarized K3 surfaces with monodromy, *Int. Math. Res. Not.* **2015** (2015), 12265–12318, [arXiv:1312.6434](#).
- [8] Doran C.F., Harder A., Novoseltsev A.Y., Thompson A., Calabi–Yau threefolds fibred by Kummer surfaces associated to products of elliptic curves, in String-Math 2014, *Proc. Sympos. Pure Math.*, Vol. 93, Amer. Math. Soc., Providence, RI, 2016, 263–287, [arXiv:1501.04024](#).
- [9] Doran C.F., Harder A., Novoseltsev A.Y., Thompson A., Calabi–Yau threefolds fibred by mirror quartic K3 surfaces, *Adv. Math.* **298** (2016), 369–392, [arXiv:1501.04019](#).
- [10] Doran C.F., Harder A., Novoseltsev A.Y., Thompson A., Calabi–Yau threefolds fibred by high rank lattice Polarized K3 surfaces, [arXiv:1701.03279](#).
- [11] Doran C.F., Malmendier A., Calabi–Yau manifolds realizing symplectically rigid monodromy tuples, [arXiv:1503.07500](#).
- [12] Doran C.F., Morgan J.W., Mirror symmetry and integral variations of Hodge structure underlying one-parameter families of Calabi–Yau threefolds, in Mirror Symmetry. V, *AMS/IP Stud. Adv. Math.*, Vol. 38, Editors J.D. Lewis, S.-T. Yau, N. Yui, Amer. Math. Soc., Providence, RI, 2006, 517–537, [math.AG/0505272](#).
- [13] Eskin A., Kontsevich M., Moeller M., Zorich A., Lower bounds for Lyapunov exponents of flat bundles on curves, [arXiv:1609.01170](#).
- [14] Filip S., Families of K3 surfaces and Lyapunov exponents, [arXiv:1412.1779](#).
- [15] Fougeron C., Parabolic degrees and Lyapunov exponents for hypergeometric local systems, [arXiv:1701.08387](#).

- [16] Fujino O., A canonical bundle formula for certain algebraic fiber spaces and its applications, *Nagoya Math. J.* **172** (2003), 129–171.
- [17] Green M., Griffiths P., Kerr M., Some enumerative global properties of variations of Hodge structures, *Mosc. Math. J.* **9** (2009), 469–530.
- [18] Griffiths P.A., Periods of integrals on algebraic manifolds. I. Construction and properties of the modular varieties, *Amer. J. Math.* **90** (1968), 568–626.
- [19] Griffiths P.A., Periods of integrals on algebraic manifolds. II. Local study of the period mapping, *Amer. J. Math.* **90** (1968), 805–865.
- [20] Griffiths P.A., On the periods of certain rational integrals. I, *Ann. of Math.* **90** (1969), 460–495.
- [21] Griffiths P.A., On the periods of certain rational integrals. II, *Ann. of Math.* **90** (1969), 496–541.
- [22] Griffiths P.A., Periods of integrals on algebraic manifolds. III. Some global differential-geometric properties of the period mapping, *Inst. Hautes Études Sci. Publ. Math.* **38** (1970), 125–180.
- [23] Hollborn H., Müller-Stach S., Hodge numbers for the cohomology of Calabi–Yau type local systems, in Algebraic and Complex Geometry, *Springer Proc. Math. Stat.*, Vol. 71, Springer, Cham, 2014, 225–240, [arXiv:1302.3047](#).
- [24] Kappes A., Lyapunov exponents of rank 2-variations of Hodge structures and modular embeddings, *Ann. Inst. Fourier (Grenoble)* **64** (2014), 2037–2066, [arXiv:1303.1088](#).
- [25] Miranda R., The basic theory of elliptic surfaces, Dottorato di Ricerca in Matematica, ETS Editrice, Pisa, 1989.
- [26] Morrison D.R., Walcher J., D-branes and normal functions, *Adv. Theor. Math. Phys.* **13** (2009), 553–598, [arXiv:0709.4028](#).
- [27] Narumiya N., Shiga H., The mirror map for a family of $K3$ surfaces induced from the simplest 3-dimensional reflexive polytope, in Proceedings on Moonshine and Related Topics (Montréal, QC, 1999), *CRM Proc. Lecture Notes*, Vol. 30, Amer. Math. Soc., Providence, RI, 2001, 139–161.
- [28] Singh S., Arithmeticity of four hypergeometric monodromy groups associated to Calabi–Yau threefolds, *Int. Math. Res. Not.* **2015** (2015), 8874–8889, [arXiv:1308.4039](#).
- [29] Singh S., Venkataramana T.N., Arithmeticity of certain symplectic hypergeometric groups, *Duke Math. J.* **163** (2014), 591–617, [arXiv:1208.6460](#).
- [30] Walcher J., Extended holomorphic anomaly and loop amplitudes in open topological string, *Nuclear Phys. B* **817** (2009), 167–207, [arXiv:0705.4098](#).
- [31] Zucker S., Hodge theory with degenerating coefficients. L_2 cohomology in the Poincaré metric, *Ann. of Math.* **109** (1979), 415–476.