

### Standard 1st-order systems

$$H(s) = \frac{K}{s\tau + 1}, \quad K, \tau > 0 \Rightarrow \text{pole} = -\frac{1}{\tau}$$

impulse response:

$$y(t) = \frac{K}{\tau} e^{-\frac{t}{\tau}} \mathbb{1}(t) \rightarrow \text{steady-state value} = 0$$

step response:

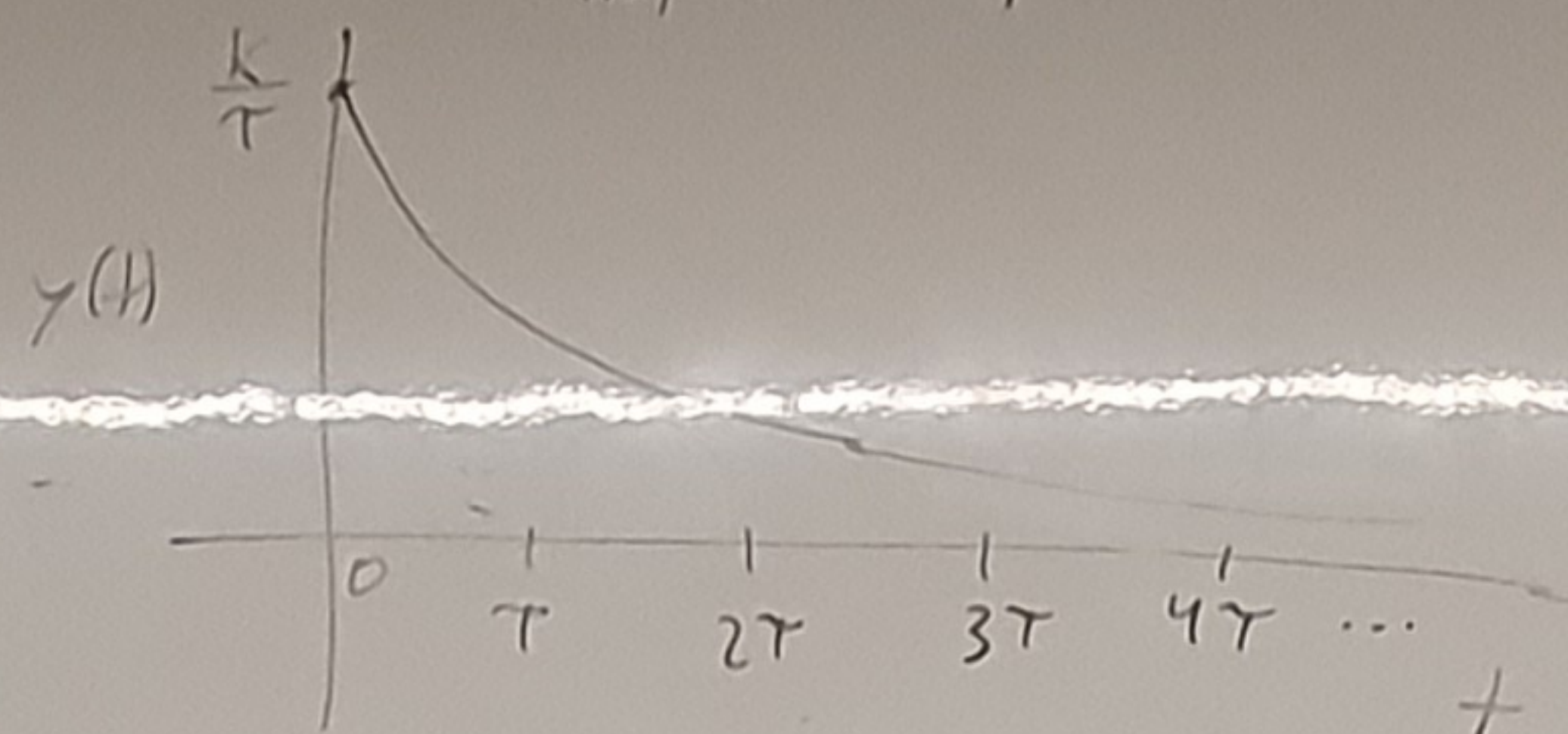
$$y(t) = K(1 - e^{-\frac{t}{\tau}}) \mathbb{1}(t) \rightarrow \text{steady-state value} = K$$

To find the steady-state value,

take  $\lim_{t \rightarrow \infty}$

\* Note: steady-state value is only well-defined for stable systems

impulse response



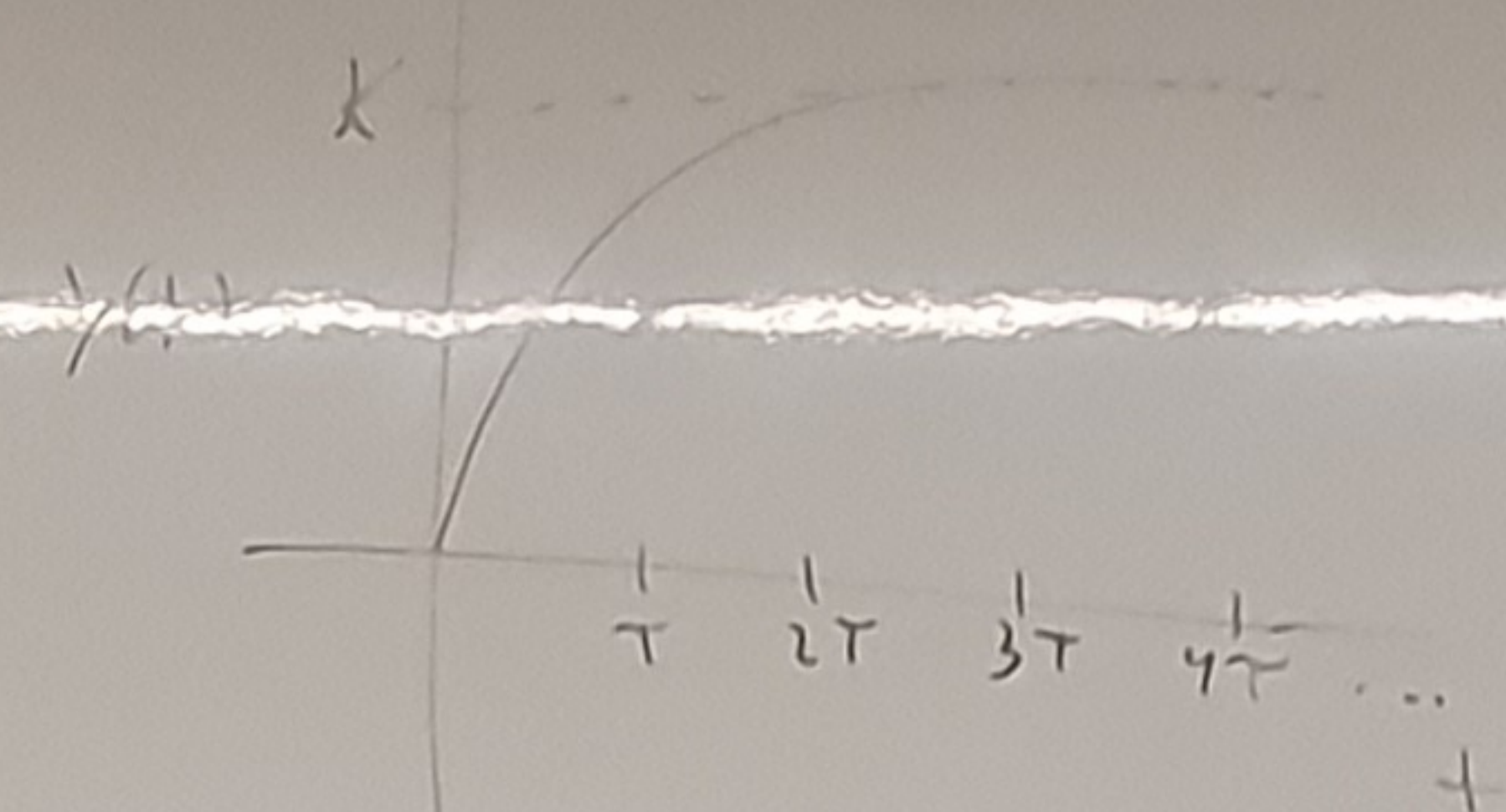
$\tau$  is called the time constant  $\rightarrow$  determines the rate of decay

$$\text{If } t = \tau \Rightarrow e^{-t/\tau} = e^{-1}$$

$$\text{If } t = 3\tau \Rightarrow e^{-t/\tau} = e^{-3} \approx 0.05$$

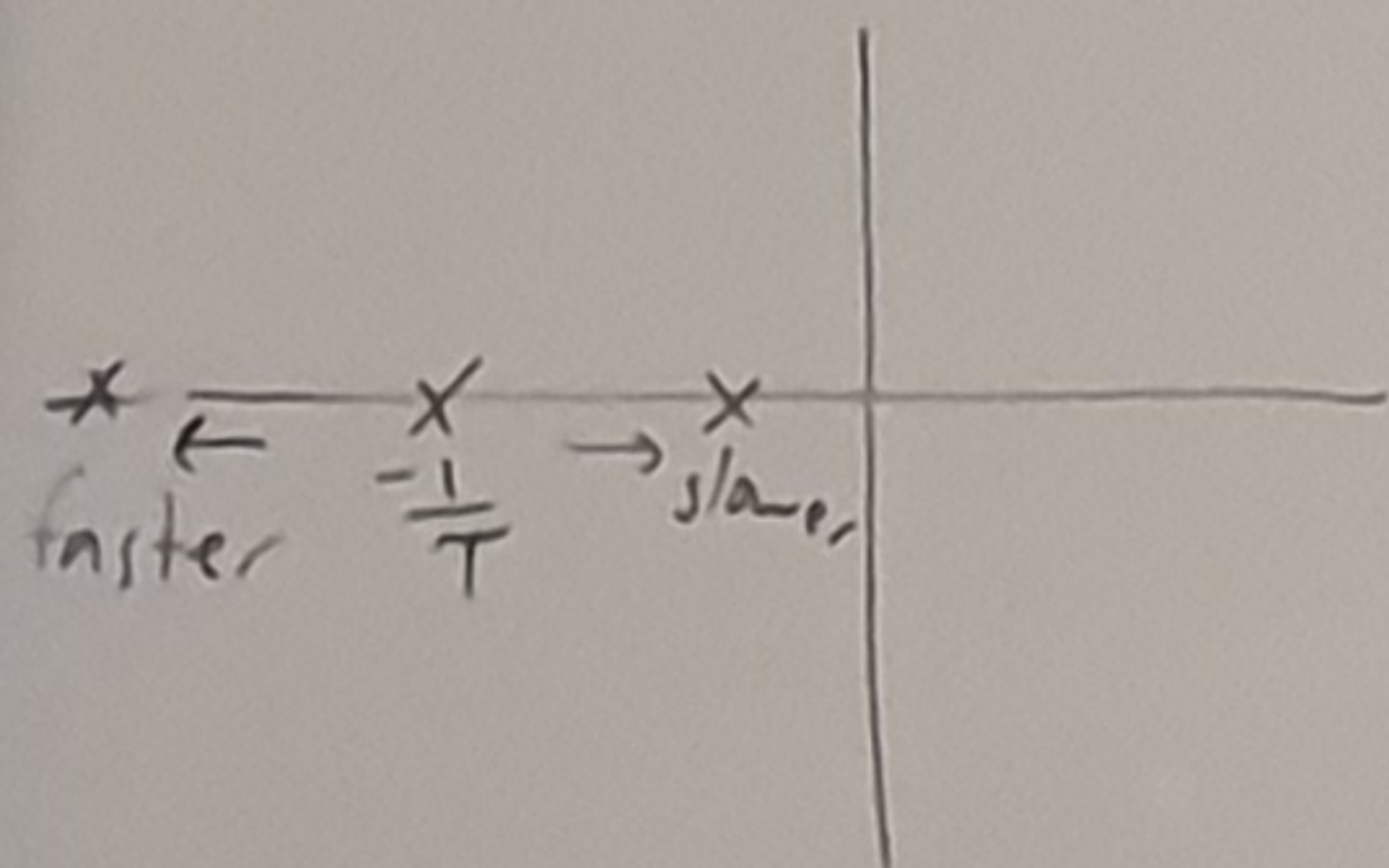
$$\text{If } t = 4\tau \Rightarrow e^{-t/\tau} = e^{-4} \approx 0.02$$

step response



Position of the pole in  $\mathbb{C}$ :

- 1 real pole at  $s = -\frac{1}{\tau}$



- as pole  $\rightarrow 0$ ,  $\tau \rightarrow \infty \Rightarrow$  response slows (transient decays more slowly)
- as pole  $\rightarrow -\infty$ ,  $\tau \rightarrow 0 \Rightarrow$  response speeds up (transient decays faster)

potential issues with fast responses:

- larger overshoot
- larger oscillations
- larger control effort
- more vulnerable to disturbances and noise (less inertia)

When designing controllers, we'll try to keep poles away from the imaginary axis to ensure acceptable speed  
- but also we will try to avoid potential issues with being too fast

### Standard 2nd order system

$$H(s) = \frac{\omega_n^2}{s^2 + 2\zeta\omega_n s + \omega_n^2}, \quad \omega_n, \zeta > 0$$

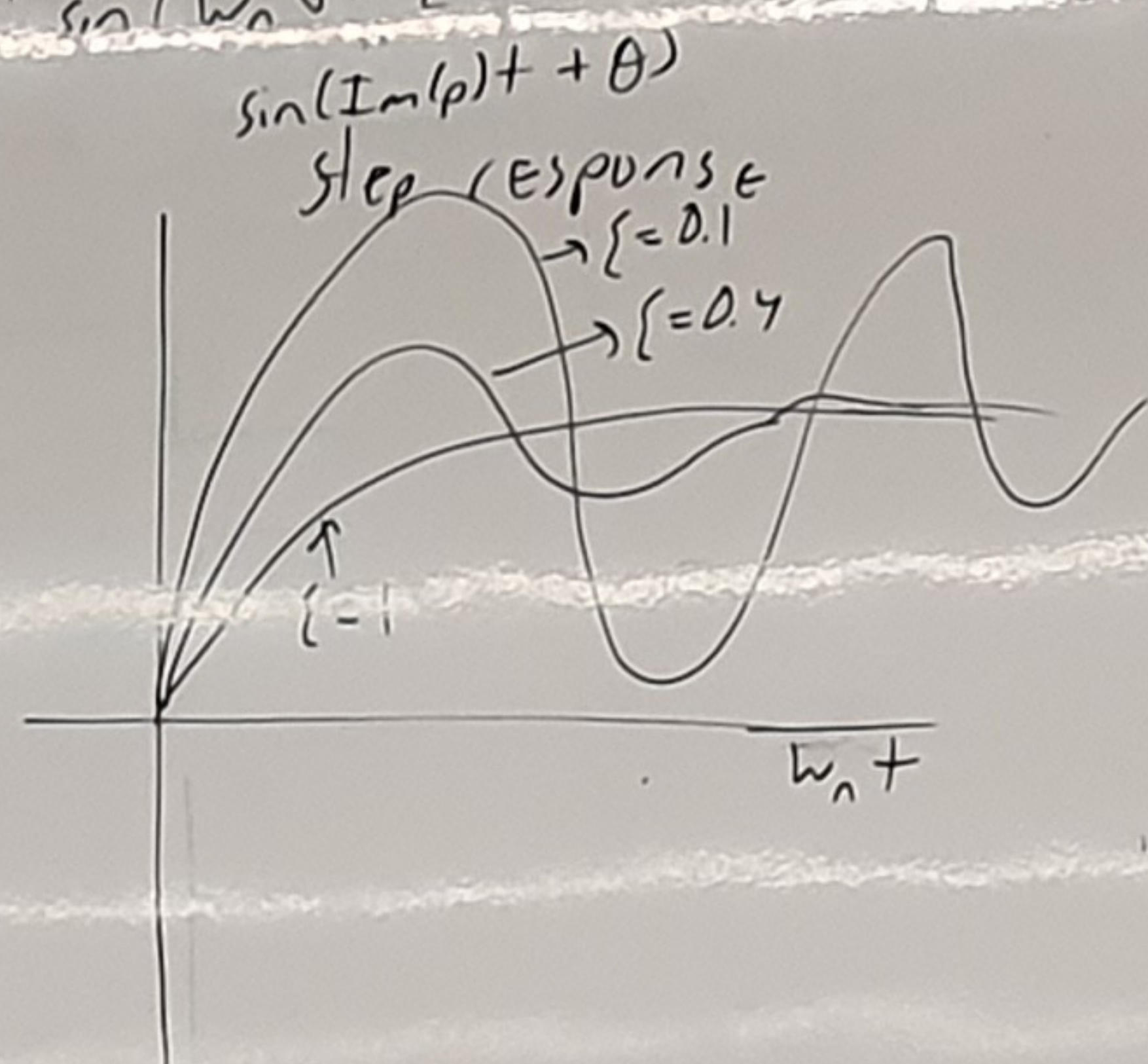
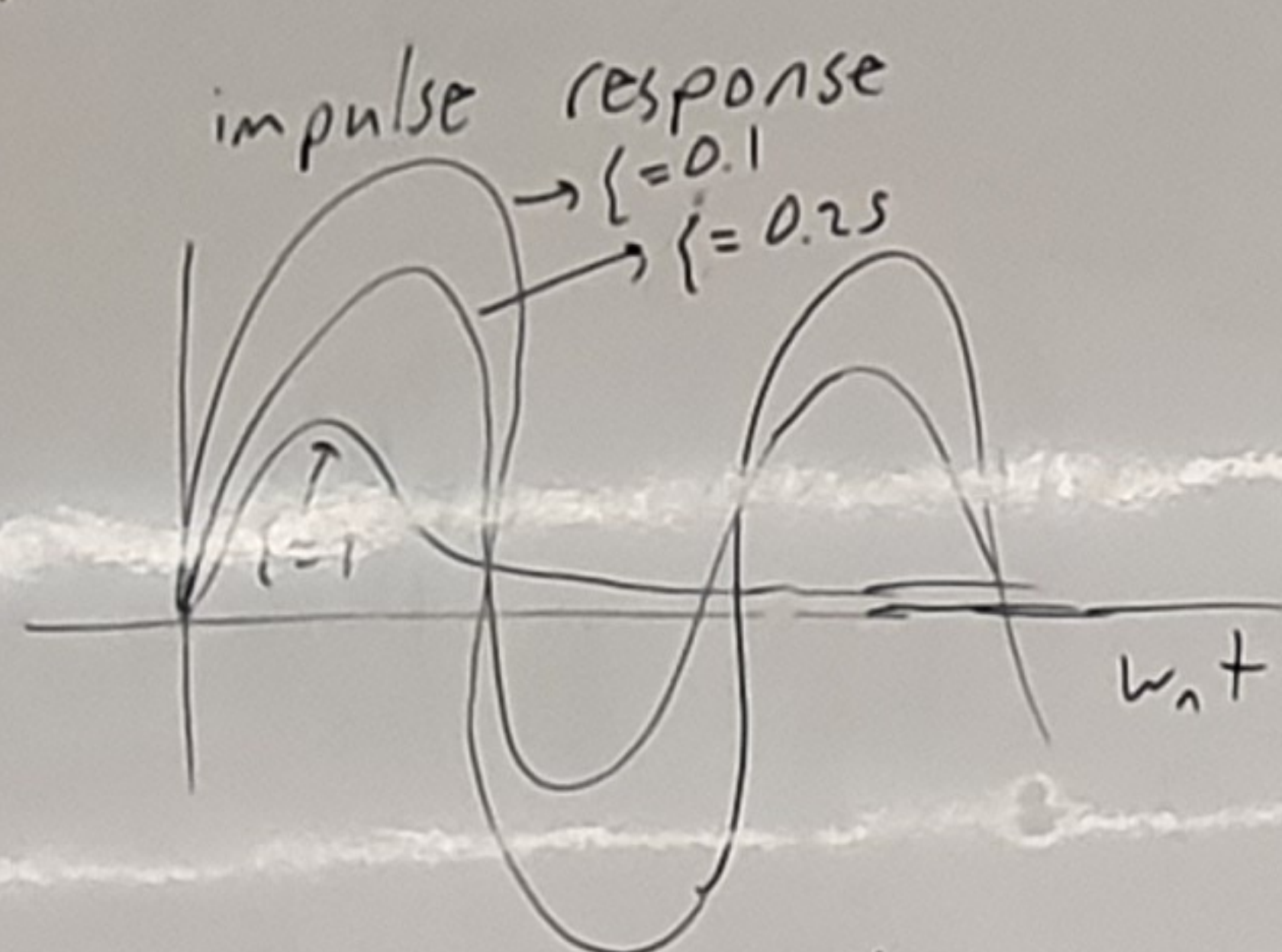
we'll consider only this case

e.g. mass-spring-damper system, RLC circuit

- system has 2 poles
- $0 < \zeta < 1 \Rightarrow$  1 pair of complex conjugate poles  $\rightarrow$  underdamped
  - $\zeta = 1 \Rightarrow$  2 real poles at the same location  $\rightarrow$  critically damped
  - $\zeta > 1 \Rightarrow$  2 distinct real poles  $\rightarrow$  overdamped

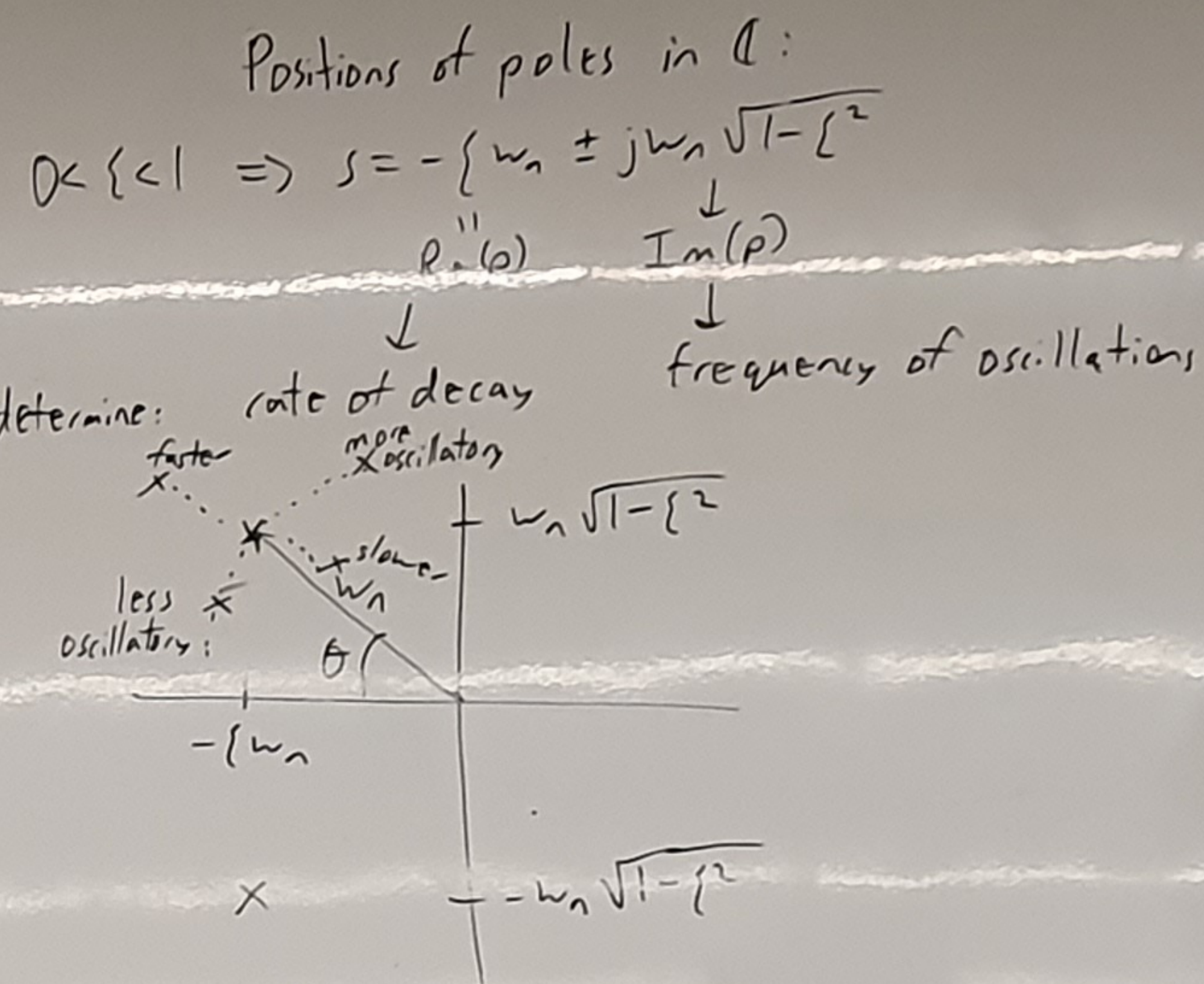
impulse response:  $y(t) = L^{-1}(H(s)) = \frac{\omega_n}{\sqrt{1-\zeta^2}} e^{-\zeta\omega_n t} \sin(\omega_n \sqrt{1-\zeta^2} t) \mathbb{1}(t) \rightarrow \text{steady-state value} = 0$

step response  $y(t) = L^{-1}(H(s) \frac{1}{s}) = 1 - \frac{1}{\sqrt{1-\zeta^2}} e^{-\zeta\omega_n t} \sin(\omega_n \sqrt{1-\zeta^2} t + \theta) \mathbb{1}(t) \rightarrow \text{steady-state value} = 1$



The plots show that as  $\zeta \uparrow \Rightarrow$  response is less oscillatory

We call  $\omega_n$  the natural frequency and  $\zeta$  the damping ratio



It is convenient to plot the responses as functions of  $\omega_n t \Rightarrow \omega_n$  acts as a time-scale factor:  
 $\omega_n \uparrow \Rightarrow$  response speeds up,  $\omega_n \downarrow \Rightarrow$  response slows down

The magnitude of the pole is  $\sqrt{(-\zeta\omega_n)^2 + (\pm\omega_n \sqrt{1-\zeta^2})^2} = \sqrt{\omega_n^2(\zeta^2 + 1 - \zeta^2)} = \sqrt{\omega_n^2} = \omega_n$

The phase of the pole is  $\theta = \cos^{-1} \frac{\text{Re}(p)}{\omega_n} = \cos^{-1} \zeta$

So we have the following relationship:

magnitude  $\uparrow \Rightarrow \omega_n \uparrow \Rightarrow$  response speeds up

$\theta \downarrow \Rightarrow \zeta \uparrow \Rightarrow$  response less oscillatory

Again, we find that the further the poles from the imaginary axis, the faster the response

Also, the closer the poles are to the real axis, the less oscillatory the response

To quantify the effects of  $\zeta$  and  $\omega_n$  on the speed and oscillation of the transient response, let's apply some standard control engineering performance specifications (specs).