

第一次作业答案

2. 写出静电场下单个氢气分子的精确哈密顿量（使用原子单位）。

$$\begin{aligned}\hat{H} &= \hat{T}_p + \hat{T}_e + \hat{U}_{pp} + \hat{U}_{ep} + \hat{U}_{ep} + \hat{U}_{Ee} \\ \hat{T}_p &= -\frac{1}{2M_a} \nabla_a^2 - \frac{1}{2M_b} \nabla_b^2 \\ \hat{T}_e &= -\frac{1}{2} \nabla_1^2 - \frac{1}{2} \nabla_2^2 \\ \hat{U}_{pp} &= \frac{1}{|R_a - R_b|}, \hat{U}_{ee} = \frac{1}{|r_1 - r_2|} \\ \hat{U}_{ep} &= -\frac{1}{|R_a - r_1|} - \frac{1}{|R_a - r_2|} - \frac{1}{|R_b - r_1|} - \frac{1}{|R_b - r_2|} \\ \hat{U}_{EP} &= -ER_a - ER_b, U_{Ee} = Er_1 + Er_2\end{aligned}$$

3. 把上一题中的哈密顿量，利用绝热近似分成原子核部分和电子部分。

$$\begin{aligned}H_e &= -\frac{1}{2} \nabla_1^2 - \frac{1}{2} \nabla_2^2 + \frac{1}{|r_1 - r_2|} + E_{r_1} + E_{r_2} \\ &\quad - \frac{1}{|R_a - r_1|} - \frac{1}{|R_a - r_2|} - \frac{1}{|R_b - r_1|} - \frac{1}{|R_b - r_2|} \\ \hat{H}_p &= -\frac{1}{2M_a} \nabla_a^2 - \frac{1}{2M_b} \nabla_b^2 + \frac{1}{|R_a - R_b|} - E \cdot R_a - E \cdot R_b\end{aligned}$$

4. 在Hartree – Fock近似下，推导 $\langle \Psi(r_1, r_2) | \frac{1}{|r_1 - r_2|} | \Psi(r_1, r_2) \rangle$

$$\begin{aligned}\psi(r_1, r_2) &= \frac{1}{\sqrt{2}} (\phi_1(r_1)\phi_2(r_2) - \phi_1(r_2)\phi_2(r_1)) \\ \iint \psi^*(r_1, r_2) \frac{1}{|r_1 - r_2|} \psi(r_1, r_2) dr_1 dr_2 \\ &= \frac{1}{2} \iint \phi_1^*(r_1)\phi_2^*(r_2) \frac{1}{|r_1 - r_2|} \phi_1(r_1)\phi_2(r_2) dr_1 dr_2 \\ &= \frac{1}{2} \iint \phi_1^*(r_1)\phi_2^*(r_2) \frac{1}{|r_1 - r_2|} \phi_1(r_2)\phi_2(r_1) dr_1 dr_2 \\ &= \frac{1}{2} \iint \phi_1^*(r_1)\phi_1^*(r_1) \frac{1}{|r_1 - r_2|} \phi_1(r_1)\phi_2(r_2) dr_1 dr_2 \\ &\quad + \frac{1}{2} \iint \phi_1^*(r_2)\phi_2^*(r_1) \frac{1}{|r_1 - r_2|} \phi_1(r_2)\phi_2(r_1) dr_1 dr_2 \\ &= \frac{1}{2} \iint \phi_1^*(r_1)\phi_1(r_1)\phi_2^*(r_2)\phi_2(r_2) \frac{1}{|r_1 - r_2|} dr_1 dr_2 \\ &\quad + \frac{1}{2} \iint \phi_1^*(r_1)\phi_1(r_1)\phi_2^*(r_2)\phi_2(r_2) \frac{1}{|r_1 - r_2|} dr_1 dr_2 \\ &\quad - \frac{1}{2} \iint \phi_1^*(r_1)\phi_1(r_2)\phi_2^*(r_2)\phi_2(r_1) \frac{1}{|r_1 - r_2|} dr_1 dr_2 \\ &\quad - \frac{1}{2} \iint \phi_1^*(r_1)\phi_1(r_2)\phi_2^*(r_2)\phi_2(r_1) \frac{1}{|r_1 - r_2|} dr_2 dr_1 dr_2 \\ &= \int \phi_1^*(r)\phi_1(r)\phi_2^*(r')\phi_2(r') \frac{1}{|r - r'|} dr dr' - \iint \phi_1^*(r)\phi_1(r')\phi_2^*(r')\phi_2(r) \frac{1}{|r - r'|} dr dr'\end{aligned}$$

5. 在 Hartree-Fock 近似下，说明关联能的来源及其定义，说明关联能不会大于零，并列举解决电子关联问题的常用方法

关联：电子波函数不能写成单粒子波函数乘积形式， $\Phi(r_1, r_2, \dots, r_n) \neq \varphi_1(r_1)\varphi_2(r_2) \dots \varphi_n(r_n)$

6. 利用变分原理推导 Hartree-Fock 方法交换积分的变分导数。

$$\begin{aligned}\delta\frac{1}{2}\sum_{i,j}k_{ij} &= \delta\frac{1}{2}\sum_{i,j}\int\phi_i^*(r_1)\phi_j^*(r_2)\frac{1}{|r_1-r_2|}\phi_i(r_2)\phi_j(r_1)dr_1dr_2 \\ &= \sum_j\int\delta\phi_a^*(r_1)\phi_j^*(r_2)\frac{1}{|r_1-r_2|}\phi_a(r_1)\phi_j(r_1)dr_1dr_2 \\ \frac{\delta\frac{1}{2}\sum_{ij}k_{ij}}{\delta\phi_a^*} &= \sum_j\phi_j(r_1)\int\phi_j^*(r_2)\frac{1}{|r_1-r_2|}\phi_a(r_2)dr_2\end{aligned}$$

10. 以氢原子为例，根据量子力学求解出来的波函数，计算其在 LDA 下的动能、外势场能、库伦能、交换能和关联能 (PW92)

交换能

$$\begin{aligned}\Phi &= \frac{1}{\sqrt{\pi}}e^{-r}, \rho = \frac{1}{\pi}e^{-2r} \\ E_x &= \int \rho \epsilon_x(p) dr \\ \epsilon_x &= \epsilon^0(n) + [\epsilon^1(n) - \epsilon^0(n)]f(\zeta) \\ \text{对于氢原子, } \zeta &= 0, f(\zeta) = 0 \\ \epsilon_x &= \epsilon^0(n) = -\frac{3}{4}\left(\frac{9}{4\pi^2}\right)^{1/3}\frac{1}{r_s} = -\frac{3}{4}\left(\frac{9}{4\pi^2}\right)^{1/3}\frac{1}{\left(\frac{3}{4\pi\rho}\right)^{1/3}} = -\frac{3}{4}\left(\frac{3\rho}{\pi}\right)^{1/3} \\ E_x &= \int \rho \left(-\frac{3}{4}\right)\left(\frac{3\rho}{\pi}\right)^{1/3} dr \\ &= -3\left(\frac{3}{\pi^2}\right)^{\frac{1}{3}}\int r^2 e^{-\frac{2}{3}r} dr = -0.2127\end{aligned}$$

关联能

$$\begin{aligned}E_c &= 4\pi \int \rho \epsilon_c(p) r^2 dr \\ \epsilon^0(p) &= -2A(1 + \alpha_1 r_s) \ln \left(1 + \frac{1}{2A(\beta_1 r_s^{1/2} + \beta_2 r_s + \beta_3 r_s^{3/2} + \beta_4 r_s^2)}\right) \\ A &= 0.031091, \alpha_1 = 0.2137, \beta_1 = 7.5957, \\ \beta_2 &= 3.5876, \beta_3 = 1.6382, \beta_4 = 0.49294 \\ E_c &= -0.044\end{aligned}$$

```
import numpy as np
from scipy.integrate import quad

def rho(r):
    return (1/np.pi) * np.exp(-2 * r)

def r_s(r):
    # 限制r的值以避免溢出
    r = np.minimum(r, 20) # 你可以根据需要调整这个值
    return (3/4) ** (1/3) * np.exp(2*r/3)

def epsilon_c(r):
    rs = r_s(r)
    x = np.sqrt(rs)

    A = 0.031091
    alpha1 = 0.2137
    beta1 = 7.5957
    beta2 = 3.5876
```

```

beta3 = 1.6382
beta4 = 0.49294

# 确保对数函数的参数为正且不会太小
term = 1 + 1/(2*A * (beta1*x + beta2*x**2 + beta3*x**3 + beta4*x**4))
if term <= 1:
    term = 1 + 1e-12 # 添加一个小的正数以避免对数为负或零

ep = -2*A * (1 + alpha1*x**2) * np.log(term)
return ep

def integrand(r):
    ep = epsilon_c(r)
    rho_val = rho(r)
    if np.isfinite(ep) and np.isfinite(rho_val):
        return ep * rho_val * 4 * np.pi * r**2
    else:
        return 0 # 如果epsilon_c或rho不是有限值, 则返回0

E_c, error = quad(integrand, 0, np.inf, epsabs=1e-12, epsrel=1e-12)

print(f'关联能 E_c = {E_c}')

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库伦能

$$\begin{aligned}
 E_{cl} &= \frac{1}{2} \iint \frac{\rho(r')\rho(r)}{|r-r'|} dr dr' \\
 |r-r'| &= \sqrt{r^2 + r'^2 - 2rr' \cos \theta} \\
 E_{cl} &= \frac{1}{2} \int dr \rho(r) \int \frac{\rho(r') dr'}{|r-r'|} \\
 &= \int \frac{\rho(r') dr'}{|r-r'|} \\
 &= \int \frac{1}{\pi} \frac{e^{-2r'}}{\sqrt{r^2 + r'^2 - 2rr' \cos \theta}} r'^2 \sin \theta dr' d\theta d\varphi \\
 &= \int \frac{2dr'}{\sqrt{2rr'}} e^{-2r'} r'^2 2\sqrt{\frac{r^2 + r'^2}{2rr'} - \cos \theta} \Big|_0^\pi \\
 &= 4 \int \frac{e^{-2r'} dr'}{2rr'} r'^2 (r + r' - |r - r'|) \\
 &= \frac{2}{r} \left(\int_0^r e^{-2r'} dr' r' 2r' + \int_r^\infty e^{-2r'} dr' r' 2r \right) \\
 &= -e^{-2r} - \frac{1}{r} e^{-2r} + \frac{1}{r} \\
 E_c &= \frac{1}{2} \int \frac{1}{\pi} e^{-2r} r^2 dr (4\pi) \left(-e^{-2r} - \frac{1}{r} e^{-2r} + \frac{1}{r} \right) \\
 &= 2 \int dr e^{-2r} (-r^2 e^{-2r} - r e^{-2r} + r) \\
 &= 2 \int_0^\infty dr (-r^2 e^{-4r} - r e^{-4r} + r e^{-2r}) \\
 &= 2 \left(-\frac{2}{4^3} - \frac{1}{4^2} + \frac{1}{2^2} \right) \\
 &= \frac{5}{16}
 \end{aligned}$$