

# Orthogonal Structure of Noether Identities

## Abstract

Any finite-order gauge-invariant curvature Lagrangian on an  $\mathrm{SO}(n+k)$  orthonormal frame bundle  $P \rightarrow M$  produces an Euler field  $\mathcal{E}(\omega^\oplus) \in \Omega^1(M, \mathrm{Ad}P)$  satisfying

$$(D^\oplus)^\dagger \mathcal{E}(\omega^\oplus) \equiv 0,$$

equivalently  $D^\oplus(\star \mathcal{E}(\omega^\oplus)) \equiv 0$ . The associated Hodge-dual current is covariantly closed, and the orthogonal splitting  $E = TM \oplus N$  yields tangent, mixed, and normal components of this covariant Noether identity. This formulation isolates the canonical orthogonal block decomposition and places the Euler field in the curvature-generated algebra determined by  $\omega^\oplus$ .

## 1 Geometric Data

### 1.1 Base and connection

Let  $M^m$  be an oriented manifold and  $E \rightarrow M$  a rank- $(n+k)$  Euclidean bundle. Let

$$P := P_{\mathrm{SO}(E)} \rightarrow M$$

be the orthonormal frame bundle. A principal  $\mathrm{SO}(n+k)$ -connection

$$\omega^\oplus \in \Omega^1(P, \mathfrak{so}(n+k)),$$

induces the covariant derivative

$$D^\oplus : \Omega^q(M, \mathrm{Ad}P) \rightarrow \Omega^{q+1}(M, \mathrm{Ad}P), \quad \mathrm{Ad}P := P \times_{\mathrm{Ad}} \mathfrak{so}(n+k).$$

Its curvature is

$$\mathcal{F}^\oplus = d\omega^\oplus + \omega^\oplus \wedge \omega^\oplus \in \Omega^2(M, \mathrm{Ad}P), \quad D^\oplus \mathcal{F}^\oplus = 0,$$

with standard conventions as in [2]. All objects are assumed smooth; variations and gauge parameters are compactly supported unless  $M$  is closed.

### 1.2 Pairing and gauge action

Fix an Ad-invariant nondegenerate pairing

$$\kappa : \mathfrak{so}(n+k) \otimes \mathfrak{so}(n+k) \rightarrow \mathbb{R}.$$

Fix background metric and orientation on  $M$ , and set

$$\langle U, V \rangle_\star := \kappa(U, \star V)$$

for equal-degree  $\mathrm{Ad}P$ -valued forms. Let

$$\mathcal{G} = \Gamma(M, P \times_{\mathrm{Ad}} \mathrm{SO}(n+k)), \quad \mathrm{Lie}(\mathcal{G}) = \Omega^0(M, \mathrm{Ad}P),$$

with infinitesimal action

$$\delta_\epsilon \omega^\oplus = D^\oplus \epsilon, \quad \delta_\epsilon \mathcal{F}^\oplus = [\mathcal{F}^\oplus, \epsilon].$$

## 2 Admissible Curvature Densities

**Definition 2.1.** A local density  $\mathcal{L}$  on  $\mathcal{A}(P)$  is admissible if:

- (i) finite order  $r$ ,
- (ii) local and horizontal,
- (iii) gauge-invariant up to horizontal exact term:  $\delta_\epsilon \mathcal{L} = d\alpha_\epsilon$ ,
- (iv) dependence on  $\omega^\oplus$  exclusively through the curvature jets

$$\mathfrak{j}_r(\omega^\oplus) = (\mathcal{F}^\oplus, D^\oplus \mathcal{F}^\oplus, \dots, (D^\oplus)^r \mathcal{F}^\oplus),$$

via invariant contractions built from  $\kappa$  and the Hodge operator.

Write

$$\mathcal{L}(\omega^\oplus) = L(\Phi_0, \dots, \Phi_r) \text{vol}_M, \quad \Phi_j = (D^\oplus)^j \mathcal{F}^\oplus,$$

and define momenta by

$$\delta L = \sum_{j=0}^r \kappa(P_j, \star \delta \Phi_j) + d\vartheta(\delta \omega^\oplus), \quad \deg P_j = \deg \Phi_j = 2 + j.$$

Here  $\vartheta(\delta \omega^\oplus) \in \Omega^{m-1}(M)$  depends linearly on  $\delta \omega^\oplus$ .

## 3 First Variation and Euler Field

**Lemma 3.1** (Curvature variation).

$$\delta \mathcal{F}^\oplus = D^\oplus(\delta \omega^\oplus).$$

**Lemma 3.2** (Jet recursion). For  $\Phi_j = (D^\oplus)^j \mathcal{F}^\oplus$  and  $j \geq 1$ ,

$$\delta \Phi_j = D^\oplus(\delta \Phi_{j-1}) + [\delta \omega^\oplus, \Phi_{j-1}].$$

**Lemma 3.3** (Formal adjoint). Let

$$\mathcal{J}_{\omega^\oplus} : \Omega^1(M, \text{Ad}P) \rightarrow \bigoplus_{j=0}^r \Omega^{j+2}(M, \text{Ad}P), \quad \mathcal{J}_{\omega^\oplus}(\eta) = (\delta \Phi_0, \dots, \delta \Phi_r).$$

For compactly supported  $\eta$ , fix the formal adjoint

$$\mathcal{J}_{\omega^\oplus}^* : \bigoplus_{j=0}^r \Omega^{j+2}(M, \text{Ad}P) \rightarrow \Omega^1(M, \text{Ad}P)$$

such that

$$\int_M \langle \mathcal{J}_{\omega^\oplus}(\eta), Q \rangle_\star \text{vol}_M = \int_M \langle \eta, \mathcal{J}_{\omega^\oplus}^*(Q) \rangle_\star \text{vol}_M + \int_M d\mathcal{B}(\eta, Q).$$

**Proposition 3.4** (First Variation).

$$\delta \mathcal{L} = \sum_{j=0}^r \langle P_j, \delta \Phi_j \rangle_\star \text{vol}_M = \langle \mathcal{E}(\omega^\oplus), \delta \omega^\oplus \rangle_\star \text{vol}_M + d\Theta, \quad \mathcal{E}(\omega^\oplus) := \mathcal{J}_{\omega^\oplus}^*(P_0, \dots, P_r) \in \Omega^1(M, \text{Ad}P).$$

Define

$$\tilde{\mathcal{E}}(\omega^\oplus) := \star \mathcal{E}(\omega^\oplus) \in \Omega^{m-1}(M, \text{Ad}P), \quad (D^\oplus)^\dagger := -\star D^\oplus \star.$$

## 4 Covariant Noether Identity

**Theorem 4.1.** *For admissible  $\mathcal{L}$ ,*

$$(D^\oplus)^\dagger \mathcal{E}(\omega^\oplus) \equiv 0 \quad \Longleftrightarrow \quad D^\oplus \tilde{\mathcal{E}}(\omega^\oplus) \equiv 0.$$

*Proof.* With  $\delta_\epsilon \omega^\oplus = D^\oplus \epsilon$ ,

$$\delta_\epsilon \mathcal{L} = \langle \mathcal{E}(\omega^\oplus), D^\oplus \epsilon \rangle_\star \text{vol}_M + d\Theta = - \left\langle (D^\oplus)^\dagger \mathcal{E}(\omega^\oplus), \epsilon \right\rangle_\star \text{vol}_M + d(\Theta + \Xi).$$

Since  $\epsilon$  is arbitrary, the coefficient of  $\epsilon$  vanishes identically.  $\square$

## 5 Curvature-Generated Algebra

Define  $\mathcal{I}_{\text{curv}} \subset \Omega^\bullet(M, \text{End}(E))$  as the graded subalgebra generated by the curvature jets and closed under wedge product, graded commutator,  $D^\oplus$ , and Hodge dual:

$$\mathcal{I}_{\text{curv}} := \langle (D^\oplus)^j \mathcal{F}^\oplus : j \geq 0 \rangle \subset \Omega^\bullet(M, \text{End}(E)),$$

i.e., the smallest graded subalgebra containing the curvature jets and closed under these operations.

**Proposition 5.1.**

$$\mathcal{E}(\omega^\oplus) \in \mathcal{I}_{\text{curv}}, \quad D^\oplus \tilde{\mathcal{E}}(\omega^\oplus) \in \mathcal{I}_{\text{curv}}, \quad D^\oplus \tilde{\mathcal{E}}(\omega^\oplus) \equiv 0.$$

*Proof.* Admissibility gives  $P_j$  as invariant contractions of curvature jets. Construction of  $\mathcal{E}(\omega^\oplus) = \mathcal{J}_{\omega^\oplus}^*(P_0, \dots, P_r)$  uses only  $D^\oplus$ , wedge, graded commutator, and  $\star$ , so  $\mathcal{E}(\omega^\oplus) \in \mathcal{I}_{\text{curv}}$ . Therefore  $\tilde{\mathcal{E}}(\omega^\oplus) = \star \mathcal{E}(\omega^\oplus) \in \mathcal{I}_{\text{curv}}$  and  $D^\oplus \tilde{\mathcal{E}}(\omega^\oplus) \in \mathcal{I}_{\text{curv}}$ . The universal identity gives  $D^\oplus \tilde{\mathcal{E}}(\omega^\oplus) \equiv 0$ .  $\square$

## 6 Orthogonal Block Projection

Under the orthogonal splitting  $E = TM \oplus N$ , the adjoint bundle  $\text{Ad}P \subset \text{End}(E)$  inherits the block decomposition

$$\text{End}(E) = \text{End}(TM) \oplus \text{Hom}(N, TM) \oplus \text{End}(N).$$

Fix  $E = TM \oplus N$  and write

$$\tilde{\mathcal{E}}(\omega^\oplus) = \begin{pmatrix} E_T & E_{\text{mix}} \\ E_{\text{mix}}^\top & E_N \end{pmatrix},$$

with  $E_{\text{mix}} \in \text{Hom}(N, TM)$  and  $E_{\text{mix}}^\top \in \text{Hom}(TM, N)$ .

**Proposition 6.1.**

$$D^\oplus \tilde{\mathcal{E}}(\omega^\oplus) = 0 \quad \Longleftrightarrow \quad \Pi_T(D^\oplus \tilde{\mathcal{E}}(\omega^\oplus)) = \Pi_{\text{mix}}(D^\oplus \tilde{\mathcal{E}}(\omega^\oplus)) = \Pi_N(D^\oplus \tilde{\mathcal{E}}(\omega^\oplus)) = 0.$$

Here  $D^{\text{prod}}$  denotes the block-diagonal covariant derivative induced by  $\omega^\oplus$ . With

$$\Psi = \begin{pmatrix} 0 & -A \\ K & 0 \end{pmatrix}, \quad D^\oplus \tilde{\mathcal{E}}(\omega^\oplus) = D^{\text{prod}} \tilde{\mathcal{E}}(\omega^\oplus) + [\Psi, \tilde{\mathcal{E}}(\omega^\oplus)],$$

one projected equation is

$$D^h E_T - (A E_{\text{mix}}^\top + E_{\text{mix}} K) = 0.$$

**Commutator channels.**

$$[\Psi, \tilde{\mathcal{E}}(\omega^\oplus)]_{TT} = -(A E_{\text{mix}}^\top + E_{\text{mix}} K), \quad [\Psi, \tilde{\mathcal{E}}(\omega^\oplus)]_{TN} = -A E_N + E_T A, \quad [\Psi, \tilde{\mathcal{E}}(\omega^\oplus)]_{NN} = E_{\text{mix}}^\top A + K E_{\text{mix}}.$$

## 7 Yang–Mills Channel

$$S_{\text{YM}}[\omega^\oplus] := \frac{1}{2} \int_M \langle \mathcal{F}^\oplus, \mathcal{F}^\oplus \rangle_\star \text{vol}_M.$$

Then

$$\delta S_{\text{YM}} = \int_M \langle D^\oplus(\delta\omega^\oplus), \mathcal{F}^\oplus \rangle_\star \text{vol}_M = - \int_M \langle \delta\omega^\oplus, \star D^\oplus(\star \mathcal{F}^\oplus) \rangle_\star \text{vol}_M + \int_M d\Theta_{\text{YM}},$$

with

$$\begin{aligned} \Theta_{\text{YM}} &= \kappa(\delta\omega^\oplus \wedge \star \mathcal{F}^\oplus), & \mathcal{E}(\omega^\oplus)_{\text{YM}} &:= -\star D^\oplus(\star \mathcal{F}^\oplus), & \tilde{\mathcal{E}}(\omega^\oplus)_{\text{YM}} &:= \star \mathcal{E}(\omega^\oplus)_{\text{YM}}. \\ (D^\oplus)^\dagger \mathcal{E}(\omega^\oplus)_{\text{YM}} &\equiv 0 & \iff & D^\oplus \tilde{\mathcal{E}}(\omega^\oplus)_{\text{YM}} \equiv 0. \end{aligned}$$

## 8 First Higher-Jet Channel

$$S_1[\omega^\oplus] := \frac{1}{2} \int_M \langle D^\oplus(\star \mathcal{F}^\oplus), D^\oplus(\star \mathcal{F}^\oplus) \rangle_\star \text{vol}_M.$$

Set  $X := \star \mathcal{F}^\oplus$  and  $\Psi := D^\oplus X$ . Then

$$\delta X = \star D^\oplus \eta, \quad \delta \Psi = D^\oplus(\star D^\oplus \eta) + [\eta, \star \mathcal{F}^\oplus].$$

So

$$\delta S_1 = \int_M \langle \eta, \mathcal{E}(\omega^\oplus)_1 \rangle_\star \text{vol}_M + \int_M d\Theta_1, \quad \mathcal{E}(\omega^\oplus)_1 := \mathcal{J}_1^*(\Psi), \quad \mathcal{J}_1(\eta) := D^\oplus(\star D^\oplus \eta) + [\eta, \star \mathcal{F}^\oplus].$$

Substituting  $\delta_\epsilon \omega^\oplus = D^\oplus \epsilon$  and integrating by parts yields

$$(D^\oplus)^\dagger \mathcal{E}(\omega^\oplus)_1 \equiv 0.$$

## 9 Comparison with Classical Noether Theory

The identity  $(D^\oplus)^\dagger \mathcal{E}(\omega^\oplus) \equiv 0$  is the gauge Noether identity in curvature form [3, 4, 5]. The orthogonal splitting expresses this identity as tangent, mixed, and normal block components. While gauge-natural Noether identities are classical, the present formulation isolates their orthogonal block structure within the compatibility-connection framework.

## A Sign Conventions

$$\begin{aligned} D^\oplus \alpha &= d\alpha + [\omega^\oplus, \alpha], & \mathcal{F}^\oplus &= d\omega^\oplus + \omega^\oplus \wedge \omega^\oplus, \\ [\alpha, \beta] &= \alpha \wedge \beta - (-1)^{pq} \beta \wedge \alpha, & \alpha &\in \Omega^p, \beta \in \Omega^q. \end{aligned}$$

## B Covariant Integration by Parts

For compactly supported  $\text{Ad}P$ -valued forms  $U, V$ ,

$$\kappa(U, \star D^\oplus V) \text{vol}_M = d\mathcal{B}(U, V) - (-1)^{\deg U} \kappa(D^\oplus U, \star V) \text{vol}_M,$$

with

$$\mathcal{B}(U, V) := \kappa(U \wedge \star V), \quad \kappa(U, \star V) \text{vol}_M = \kappa(U \wedge \star V).$$

## C Graded Derivation Identities

$$D^\oplus(\alpha \wedge \beta) = (D^\oplus \alpha) \wedge \beta + (-1)^{\deg \alpha} \alpha \wedge (D^\oplus \beta), \tag{1}$$

$$D^\oplus[\alpha, \beta] = [D^\oplus \alpha, \beta] + (-1)^{\deg \alpha} [\alpha, D^\oplus \beta]. \tag{2}$$

## References

- [1] I. Kolář, P. W. Michor, and J. Slovák, *Natural Operations in Differential Geometry*, Springer, 1993.
- [2] S. Kobayashi and K. Nomizu, *Foundations of Differential Geometry, Vol. I*, Wiley, 1963.
- [3] R. Utiyama, Invariant theoretical interpretation of interaction, *Phys. Rev.* **101** (1956), 1597–1607.
- [4] M. Fatibene and M. Francaviglia, *Natural and Gauge Natural Formalism for Classical Field Theories*, Kluwer, 2003.
- [5] I. M. Anderson, The variational bicomplex, *Contemp. Math.* **132** (1992), 51–73.