

Orthogonal Structure of Noether Identities

Abstract

Any finite-order gauge-invariant curvature Lagrangian on an $\mathrm{SO}(n+k)$ orthonormal frame bundle $P \rightarrow M$ produces an Euler field $\mathcal{E}(\omega^\oplus) \in \Omega^1(M, \mathrm{Ad}P)$ satisfying

$$(D^\oplus)^\dagger \mathcal{E}(\omega^\oplus) \equiv 0,$$

equivalently $D^\oplus(\star \mathcal{E}(\omega^\oplus)) \equiv 0$. The associated Hodge-dual current is covariantly closed, and the orthogonal splitting $E = TM \oplus N$ yields tangent, mixed, and normal components of this covariant Noether identity. This formulation isolates the canonical orthogonal block decomposition and places the Euler field in the curvature-generated algebra determined by ω^\oplus .

1 Geometric Data

1.1 Base and connection

Let M^m be an oriented manifold and $E \rightarrow M$ a rank- $(n+k)$ Euclidean bundle. Let

$$P := P_{\mathrm{SO}(E)} \rightarrow M$$

be the orthonormal frame bundle. A principal $\mathrm{SO}(n+k)$ -connection

$$\omega^\oplus \in \Omega^1(P, \mathfrak{so}(n+k)),$$

induces the covariant derivative

$$D^\oplus : \Omega^q(M, \mathrm{Ad}P) \rightarrow \Omega^{q+1}(M, \mathrm{Ad}P), \quad \mathrm{Ad}P := P \times_{\mathrm{Ad}} \mathfrak{so}(n+k).$$

Its curvature is

$$\mathcal{F}^\oplus = d\omega^\oplus + \omega^\oplus \wedge \omega^\oplus \in \Omega^2(M, \mathrm{Ad}P), \quad D^\oplus \mathcal{F}^\oplus = 0,$$

with standard conventions as in [2]. All objects are assumed smooth; variations and gauge parameters are compactly supported unless M is closed.

1.2 Pairing and gauge action

Fix an Ad-invariant nondegenerate pairing

$$\kappa : \mathfrak{so}(n+k) \otimes \mathfrak{so}(n+k) \rightarrow \mathbb{R}.$$

Fix background metric and orientation on M , and set

$$\langle U, V \rangle_\star := \kappa(U, \star V)$$

for equal-degree $\mathrm{Ad}P$ -valued forms. Let

$$\mathcal{G} = \Gamma(M, P \times_{\mathrm{Ad}} \mathrm{SO}(n+k)), \quad \mathrm{Lie}(\mathcal{G}) = \Omega^0(M, \mathrm{Ad}P),$$

with infinitesimal action

$$\delta_\epsilon \omega^\oplus = D^\oplus \epsilon, \quad \delta_\epsilon \mathcal{F}^\oplus = [\mathcal{F}^\oplus, \epsilon].$$

2 Admissible Curvature Densities

Definition 2.1. A local density \mathcal{L} on $\mathcal{A}(P)$ is admissible if:

- (i) finite order r ,
- (ii) local and horizontal,
- (iii) gauge-invariant up to horizontal exact term: $\delta_\epsilon \mathcal{L} = d\alpha_\epsilon$,
- (iv) dependence on ω^\oplus exclusively through the curvature jets

$$\mathbf{j}_r(\omega^\oplus) = (\mathcal{F}^\oplus, D^\oplus \mathcal{F}^\oplus, \dots, (D^\oplus)^r \mathcal{F}^\oplus),$$

via invariant contractions built from κ and the Hodge operator.

Write

$$\mathcal{L}(\omega^\oplus) = L(\Phi_0, \dots, \Phi_r) \text{vol}_M, \quad \Phi_j = (D^\oplus)^j \mathcal{F}^\oplus,$$

and define momenta by

$$\delta L = \sum_{j=0}^r \kappa(P_j, \star \delta \Phi_j) + d\vartheta(\delta \omega^\oplus), \quad \deg P_j = \deg \Phi_j = 2 + j.$$

Here $\vartheta(\delta \omega^\oplus) \in \Omega^{m-1}(M)$ depends linearly on $\delta \omega^\oplus$.

3 First Variation and Euler Field

Lemma 3.1 (Curvature variation).

$$\delta \mathcal{F}^\oplus = D^\oplus(\delta \omega^\oplus).$$

Lemma 3.2 (Jet recursion). For $\Phi_j = (D^\oplus)^j \mathcal{F}^\oplus$ and $j \geq 1$,

$$\delta \Phi_j = D^\oplus(\delta \Phi_{j-1}) + [\delta \omega^\oplus, \Phi_{j-1}].$$

Lemma 3.3 (Formal adjoint). Let

$$\mathcal{J}_{\omega^\oplus} : \Omega^1(M, \text{Ad}P) \rightarrow \bigoplus_{j=0}^r \Omega^{j+2}(M, \text{Ad}P), \quad \mathcal{J}_{\omega^\oplus}(\eta) = (\delta \Phi_0, \dots, \delta \Phi_r).$$

For compactly supported η , fix the formal adjoint

$$\mathcal{J}_{\omega^\oplus}^* : \bigoplus_{j=0}^r \Omega^{j+2}(M, \text{Ad}P) \rightarrow \Omega^1(M, \text{Ad}P)$$

such that

$$\int_M \langle \mathcal{J}_{\omega^\oplus}(\eta), Q \rangle_\star \text{vol}_M = \int_M \langle \eta, \mathcal{J}_{\omega^\oplus}^*(Q) \rangle_\star \text{vol}_M + \int_M d\mathcal{B}(\eta, Q).$$

Proposition 3.4 (First Variation).

$$\delta \mathcal{L} = \sum_{j=0}^r \langle P_j, \delta \Phi_j \rangle_\star \text{vol}_M = \langle \mathcal{E}(\omega^\oplus), \delta \omega^\oplus \rangle_\star \text{vol}_M + d\Theta, \quad \mathcal{E}(\omega^\oplus) := \mathcal{J}_{\omega^\oplus}^*(P_0, \dots, P_r) \in \Omega^1(M, \text{Ad}P).$$

Define

$$\tilde{\mathcal{E}}(\omega^\oplus) := \star \mathcal{E}(\omega^\oplus) \in \Omega^{m-1}(M, \text{Ad}P), \quad (D^\oplus)^\dagger := -\star D^\oplus \star.$$

4 Covariant Noether Identity

Theorem 4.1. *For admissible \mathcal{L} ,*

$$(D^\oplus)^\dagger \mathcal{E}(\omega^\oplus) \equiv 0 \iff D^\oplus \tilde{\mathcal{E}}(\omega^\oplus) \equiv 0.$$

Proof. With $\delta_\epsilon \omega^\oplus = D^\oplus \epsilon$,

$$\delta_\epsilon \mathcal{L} = \langle \mathcal{E}(\omega^\oplus), D^\oplus \epsilon \rangle_* \text{vol}_M + d\Theta = - \left\langle (D^\oplus)^\dagger \mathcal{E}(\omega^\oplus), \epsilon \right\rangle_* \text{vol}_M + d(\Theta + \Xi).$$

Since ϵ is arbitrary, the coefficient of ϵ vanishes identically. \square

5 Curvature-Generated Algebra

Define $\mathcal{I}_{\text{curv}} \subset \Omega^\bullet(M, \text{End}(E))$ as the graded subalgebra generated by the curvature jets and closed under wedge product, graded commutator, D^\oplus , and Hodge dual:

$$\mathcal{I}_{\text{curv}} := \langle (D^\oplus)^j \mathcal{F}^\oplus : j \geq 0 \rangle \subset \Omega^\bullet(M, \text{End}(E)),$$

i.e., the smallest graded subalgebra containing the curvature jets and closed under these operations.

Proposition 5.1.

$$\mathcal{E}(\omega^\oplus) \in \mathcal{I}_{\text{curv}}, \quad D^\oplus \tilde{\mathcal{E}}(\omega^\oplus) \in \mathcal{I}_{\text{curv}}, \quad D^\oplus \tilde{\mathcal{E}}(\omega^\oplus) \equiv 0.$$

Proof. Admissibility gives P_j as invariant contractions of curvature jets. Construction of $\mathcal{E}(\omega^\oplus) = \mathcal{J}_{\omega^\oplus}^*(P_0, \dots, P_r)$ uses only D^\oplus , wedge, graded commutator, and $*$, so $\mathcal{E}(\omega^\oplus) \in \mathcal{I}_{\text{curv}}$. Therefore $\tilde{\mathcal{E}}(\omega^\oplus) = * \mathcal{E}(\omega^\oplus) \in \mathcal{I}_{\text{curv}}$ and $D^\oplus \tilde{\mathcal{E}}(\omega^\oplus) \in \mathcal{I}_{\text{curv}}$. The universal identity gives $D^\oplus \tilde{\mathcal{E}}(\omega^\oplus) \equiv 0$. \square

6 Orthogonal Block Projection

Under the orthogonal splitting $E = TM \oplus N$, the adjoint bundle $\text{Ad}P \subset \text{End}(E)$ inherits the block decomposition

$$\text{End}(E) = \text{End}(TM) \oplus \text{Hom}(N, TM) \oplus \text{End}(N).$$

Fix $E = TM \oplus N$ and write

$$\tilde{\mathcal{E}}(\omega^\oplus) = \begin{pmatrix} E_T & E_{\text{mix}} \\ E_{\text{mix}}^\top & E_N \end{pmatrix},$$

with $E_{\text{mix}} \in \text{Hom}(N, TM)$ and $E_{\text{mix}}^\top \in \text{Hom}(TM, N)$.

Proposition 6.1.

$$D^\oplus \tilde{\mathcal{E}}(\omega^\oplus) = 0 \iff \Pi_T(D^\oplus \tilde{\mathcal{E}}(\omega^\oplus)) = \Pi_{\text{mix}}(D^\oplus \tilde{\mathcal{E}}(\omega^\oplus)) = \Pi_N(D^\oplus \tilde{\mathcal{E}}(\omega^\oplus)) = 0.$$

Here D^{prod} denotes the block-diagonal covariant derivative induced by ω^\oplus . With

$$\Psi = \begin{pmatrix} 0 & -A \\ K & 0 \end{pmatrix}, \quad D^\oplus \tilde{\mathcal{E}}(\omega^\oplus) = D^{\text{prod}} \tilde{\mathcal{E}}(\omega^\oplus) + [\Psi, \tilde{\mathcal{E}}(\omega^\oplus)],$$

one projected equation is

$$D^h E_T - (A E_{\text{mix}}^\top + E_{\text{mix}} K) = 0.$$

Commutator channels.

$$[\Psi, \tilde{\mathcal{E}}(\omega^\oplus)]_{TT} = -(A E_{\text{mix}}^\top + E_{\text{mix}} K), \quad [\Psi, \tilde{\mathcal{E}}(\omega^\oplus)]_{TN} = -A E_N + E_T A, \quad [\Psi, \tilde{\mathcal{E}}(\omega^\oplus)]_{NN} = E_{\text{mix}}^\top A + K E_{\text{mix}}.$$

7 Yang–Mills Channel

$$S_{\text{YM}}[\omega^\oplus] := \frac{1}{2} \int_M \langle \mathcal{F}^\oplus, \mathcal{F}^\oplus \rangle_\star \text{vol}_M.$$

Then

$$\delta S_{\text{YM}} = \int_M \langle D^\oplus(\delta\omega^\oplus), \mathcal{F}^\oplus \rangle_\star \text{vol}_M = - \int_M \langle \delta\omega^\oplus, \star D^\oplus(\star\mathcal{F}^\oplus) \rangle_\star \text{vol}_M + \int_M d\Theta_{\text{YM}},$$

with

$$\begin{aligned} \Theta_{\text{YM}} &= \kappa(\delta\omega^\oplus \wedge \star\mathcal{F}^\oplus), & \mathcal{E}(\omega^\oplus)_{\text{YM}} &:= -\star D^\oplus(\star\mathcal{F}^\oplus), & \tilde{\mathcal{E}}(\omega^\oplus)_{\text{YM}} &:= \star\mathcal{E}(\omega^\oplus)_{\text{YM}}. \\ (D^\oplus)^\dagger \mathcal{E}(\omega^\oplus)_{\text{YM}} &\equiv 0 & \iff & & D^\oplus \tilde{\mathcal{E}}(\omega^\oplus)_{\text{YM}} &\equiv 0. \end{aligned}$$

8 First Higher-Jet Channel

$$S_1[\omega^\oplus] := \frac{1}{2} \int_M \langle D^\oplus(\star\mathcal{F}^\oplus), D^\oplus(\star\mathcal{F}^\oplus) \rangle_\star \text{vol}_M.$$

Set $X := \star\mathcal{F}^\oplus$ and $\Psi := D^\oplus X$. Then

$$\delta X = \star D^\oplus \eta, \quad \delta\Psi = D^\oplus(\star D^\oplus \eta) + [\eta, \star\mathcal{F}^\oplus].$$

So

$$\delta S_1 = \int_M \langle \eta, \mathcal{E}(\omega^\oplus)_1 \rangle_\star \text{vol}_M + \int_M d\Theta_1, \quad \mathcal{E}(\omega^\oplus)_1 := \mathcal{J}_1^*(\Psi), \quad \mathcal{J}_1(\eta) := D^\oplus(\star D^\oplus \eta) + [\eta, \star\mathcal{F}^\oplus].$$

Substituting $\delta_\epsilon \omega^\oplus = D^\oplus \epsilon$ and integrating by parts yields

$$(D^\oplus)^\dagger \mathcal{E}(\omega^\oplus)_1 \equiv 0.$$

9 Comparison with Classical Noether Theory

The identity $(D^\oplus)^\dagger \mathcal{E}(\omega^\oplus) \equiv 0$ is the gauge Noether identity in curvature form [3, 4, 5]. The orthogonal splitting expresses this identity as tangent, mixed, and normal block components. While gauge-natural Noether identities are classical, the present formulation isolates their orthogonal block structure within the compatibility-connection framework.

A Sign Conventions

$$\begin{aligned} D^\oplus \alpha &= d\alpha + [\omega^\oplus, \alpha], & \mathcal{F}^\oplus &= d\omega^\oplus + \omega^\oplus \wedge \omega^\oplus, \\ [\alpha, \beta] &= \alpha \wedge \beta - (-1)^{pq} \beta \wedge \alpha, & \alpha \in \Omega^p, \beta \in \Omega^q. \end{aligned}$$

B Covariant Integration by Parts

For compactly supported $\text{Ad}P$ -valued forms U, V ,

$$\kappa(U, \star D^\oplus V) \text{vol}_M = d\mathcal{B}(U, V) - (-1)^{\deg U} \kappa(D^\oplus U, \star V) \text{vol}_M,$$

with

$$\mathcal{B}(U, V) := \kappa(U \wedge \star V), \quad \kappa(U, \star V) \text{vol}_M = \kappa(U \wedge \star V).$$

C Graded Derivation Identities

$$D^\oplus(\alpha \wedge \beta) = (D^\oplus \alpha) \wedge \beta + (-1)^{\deg \alpha} \alpha \wedge (D^\oplus \beta), \quad (1)$$

$$D^\oplus[\alpha, \beta] = [D^\oplus \alpha, \beta] + (-1)^{\deg \alpha} [\alpha, D^\oplus \beta]. \quad (2)$$

References

- [1] I. Kolář, P. W. Michor, and J. Slovák, *Natural Operations in Differential Geometry*, Springer, 1993.
- [2] S. Kobayashi and K. Nomizu, *Foundations of Differential Geometry, Vol. I*, Wiley, 1963.
- [3] R. Utiyama, Invariant theoretical interpretation of interaction, *Phys. Rev.* **101** (1956), 1597–1607.
- [4] M. Fatibene and M. Francaviglia, *Natural and Gauge Natural Formalism for Classical Field Theories*, Kluwer, 2003.
- [5] I. M. Anderson, The variational bicomplex, *Contemp. Math.* **132** (1992), 51–73.