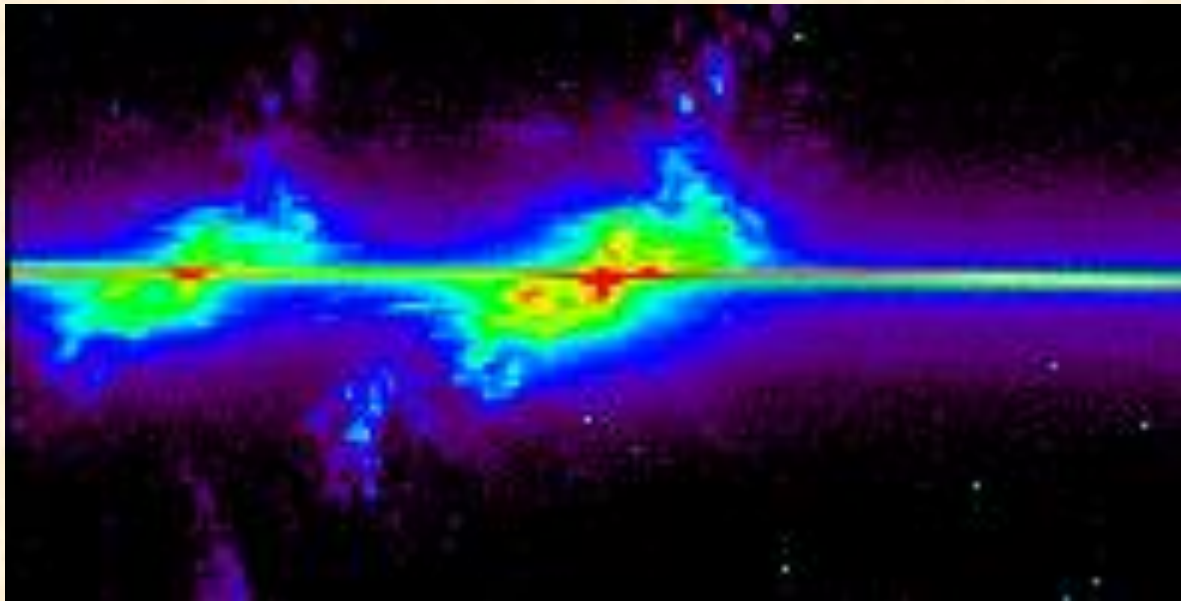
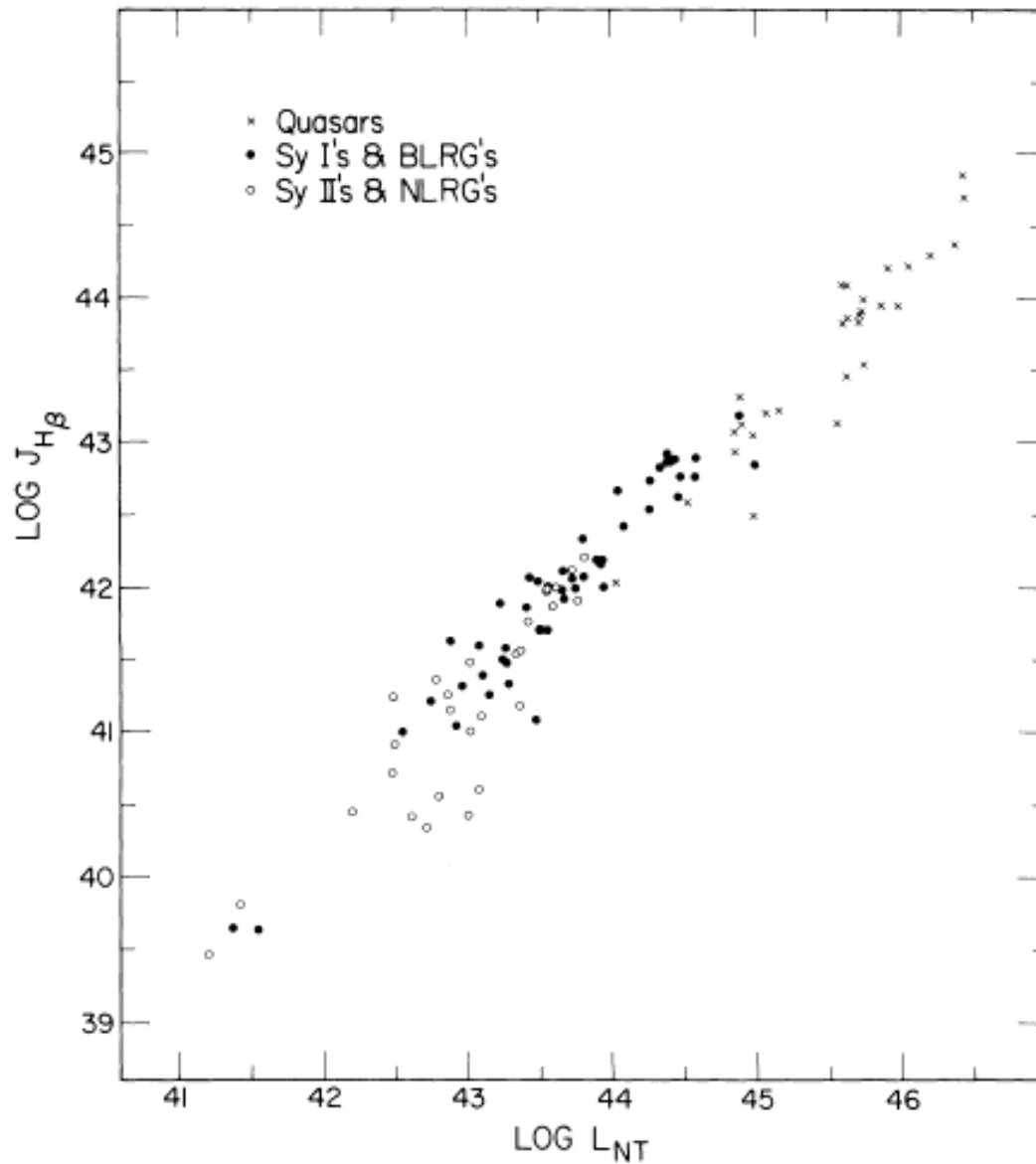


AGN – Physics of the Ionized Gas

- Physical conditions in the NLR
- Physical conditions in the BLR
- LINERs
- Emission-Line Diagnostics
- High-Energy Effects



Evidence for Photoionization



- continuum and
H β luminosity
correlated over a
huge range

(Yee, H. 1980, ApJ, 241, 894)

Emission-Line Diagnostics for Seyfert NLRs

- $T = 10,000 - 20,000$ K from [O III] lines → **photoionization** (shock heating gives temperatures $\approx 40,000$ K)
- Emission lines span a wide range in ionization potential (IP):
 - IP needed to create [O I]: 0 eV, [Fe X IV]: 361 eV
 - **Power-law SEDs with substantial X-ray contribution**
- UV radiation forms a classic H II region on the “front face”
- X-rays penetrate deep into the cloud to create a “partially-ionized zone” (PIZ): $N(\text{H II})/N(\text{H I}) \approx 0.1$ to 0.2
 - In the PIZ, elements are neutral or singly ionized
 - substantial emission from HI, [O I], [N II], [S II], Mg II
 - **Large column densities ($N_{\text{H}} = 10^{19} - 10^{21} \text{ cm}^{-2}$)**
- HST resolved spectroscopy shows wide range in number density.
 - **$n_{\text{H}} = 10^2 - 10^6 \text{ cm}^{-3}$ (from lines with a range in critical density)**

Collisional Excitation of H Lines in the PIZ

- X-rays penetrate deep into the cloud to create high-energy (“suprathermal”) electrons, which cause multiple ionizations in the mostly neutral gas.
- Suprathermal electrons also collisionally excite the $n = 1$ level in hydrogen:
- $L\alpha$ is collisionally enhanced relative to the other H lines

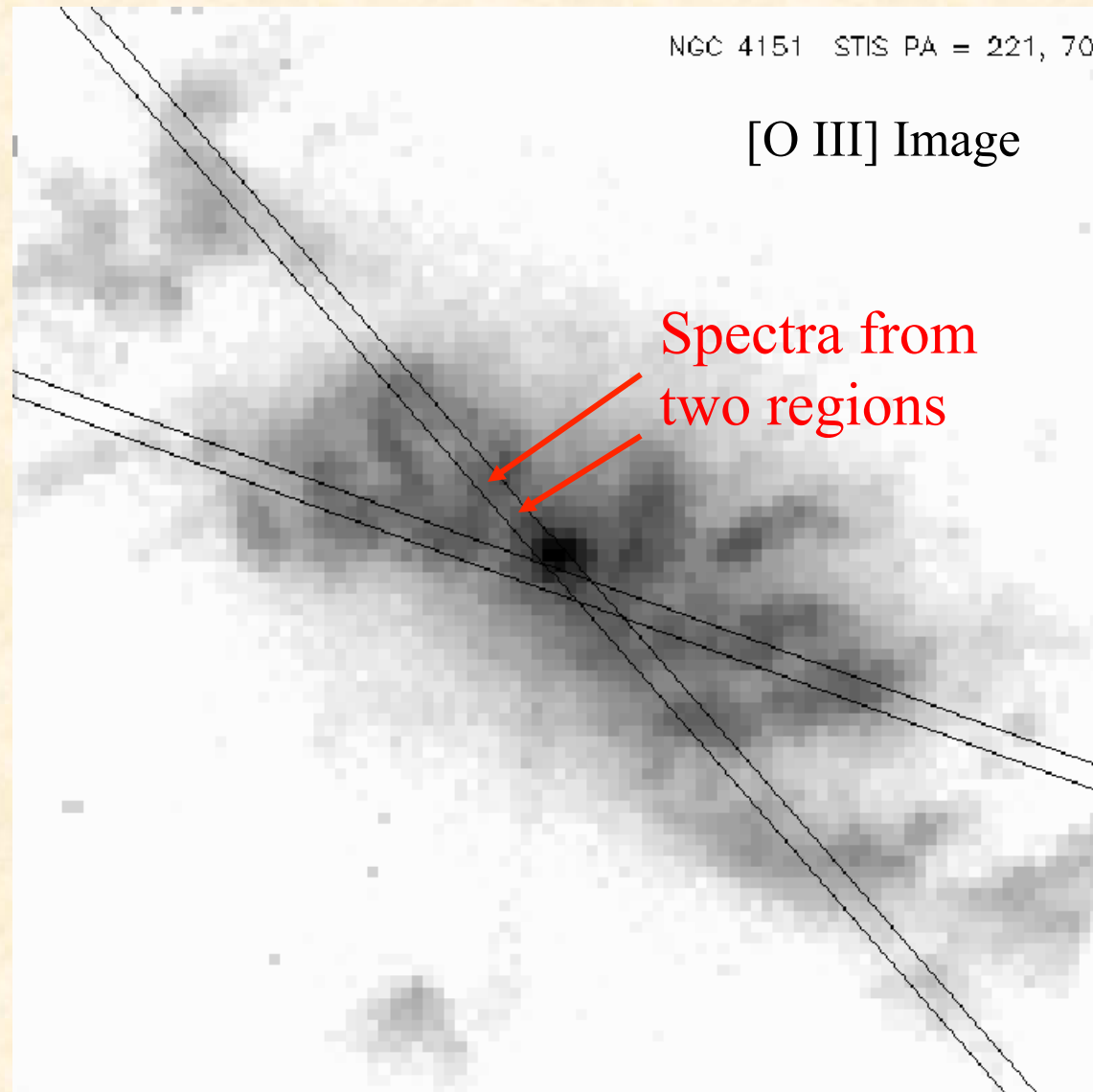
$$4\pi j_v = n_e n_{H^0} q_{12} h\nu_{L\alpha}, \text{ where } q \text{ is the collision rate}$$

- $L\alpha/H\beta$ can reach ~ 50 in the NLR, compared to recombination value of 33.
- $H\alpha$ is the next most collisionally enhanced line ($n = 1$ to $n = 3$)
- $H\alpha/H\beta$ can reach ~ 3.1 in the NLR, compared to the recombination value of 2.85

Results from NLR models

- Photoionization codes like CLOUDY contain all of the important physics (X-ray ionization of the PIZ, collisional excitation and ionization, Auger effect, charge exchange, etc.)
- Input parameters: U (or luminosity and distance for spatially resolved regions), continuum shape (SED), number density (n_H), abundances, column densities (N_H).
- Models indicate abundances are approximately solar
 - previous “low abundance” cases due to a high-density component, which suppresses the forbidden lines (CNO, etc.)
- Multiple components (with different U , n_H) are usually needed at each position.
- Power-law interpolation between UV and X-ray ($\alpha_v \approx 1.5$) works - no need for huge EUV bump (BBB)
- Dust within the clouds can suppress resonance lines (esp. Ly α)

Ex) STIS Long-Slit Spectra of the NLR in NGC 4151 (Kraemer et al. 2000, ApJ, 531, 278)



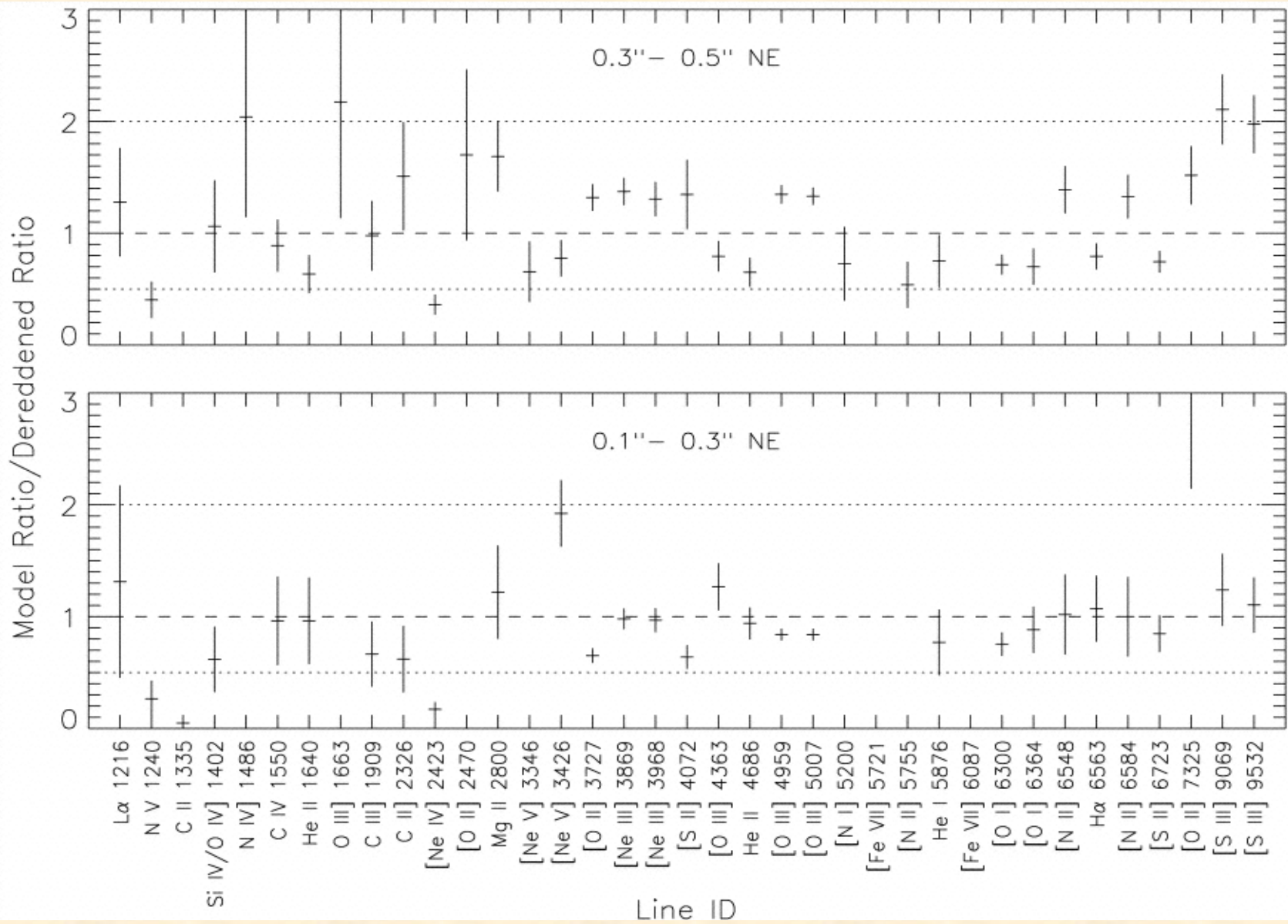
Model Results from Two Regions

Spectral Bin	Log U	n_e (cm ⁻³)	N_H (cm ⁻³)	% H β	Note
0.1-0.3 NE	-2.67	1.2 E4	1.6 E 21	50%	RB
	-3.0	1.0 E7	5.6 E 19	25%	MB
	-1.08	1.0 E5	5.6 E 20	25%	MB
0.3-0.5 NE	-2.67	1.2 E4	1.6 E21	90%	RB
	-1.36	6.0 E2	5.3 E 20	10%	MB

MB – matter bounded (optically thin)

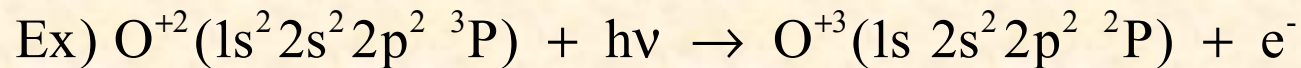
RB – radiation bounded (optically thick)

Comparison of Models and Observations

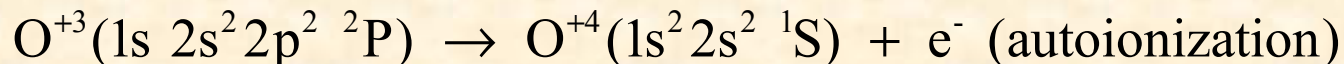


Physical Diagnostics of the BLR

- No forbidden lines, some semi-forbidden lines:
 - No broad [O III] $\lambda\lambda 4959, 5007 \rightarrow n_H \geq 10^8 \text{ cm}^{-3}$
 - Broad C III] $\lambda 1909: \rightarrow n_H \leq 10^{11} \text{ cm}^{-3}$
- Cooling is primarily done by recombination lines (H and He) and collisional excitation of permitted lines (e.g., C IV, N V in UV; Fe II in UV and optical)
- X-ray ionization (also important in NLR)
 - ejected outer shell (suprathermal) electrons causes ~ 6 collisional ionizations
 - Auger effect: X-ray photon can eject multiple electrons



– leaves O^{+3} in excited state

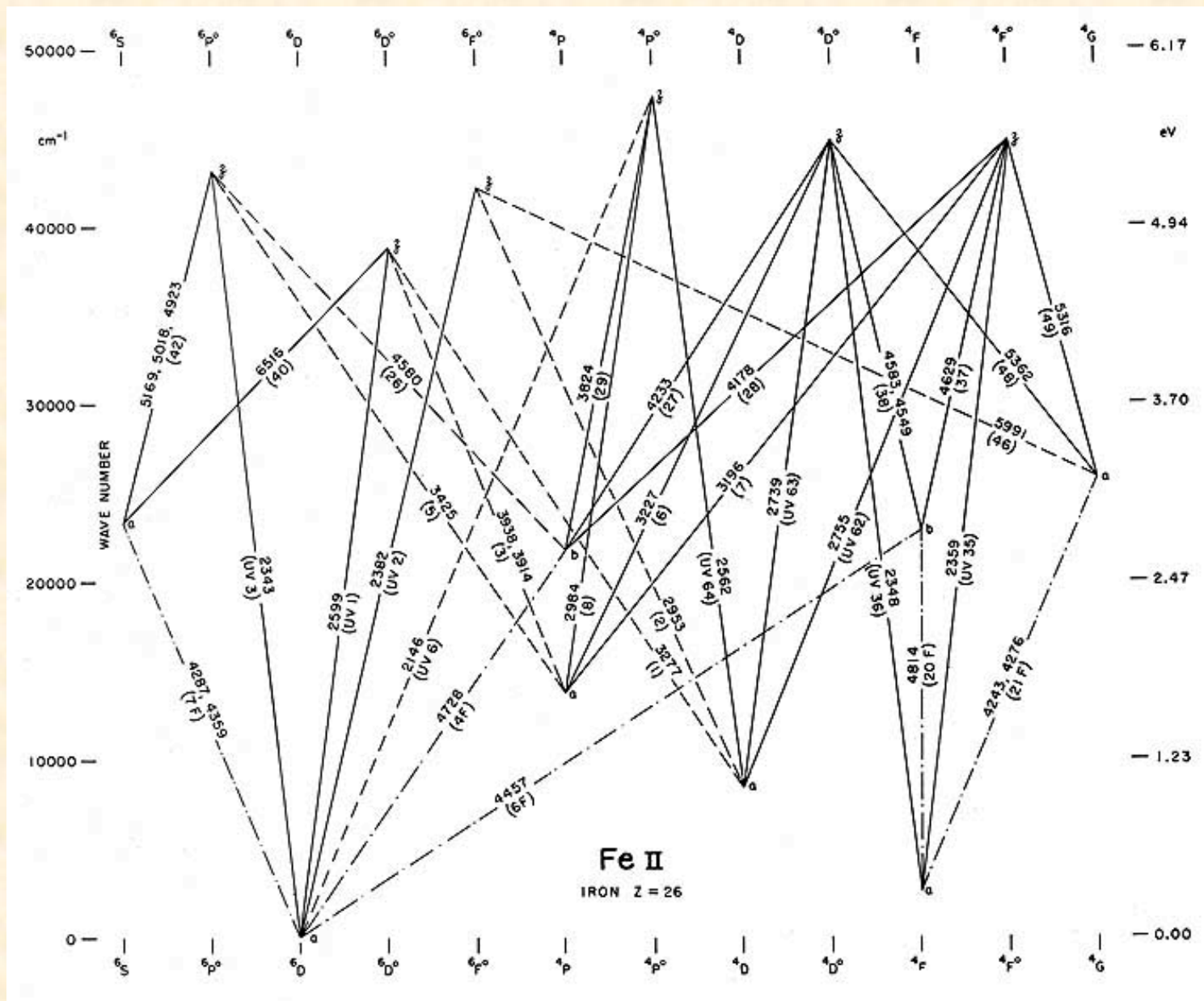


- Fe II, Mg II, C I, and O I are enhanced in the PIZ $\rightarrow N_H = 10^{22} - 10^{23} \text{ cm}^{-2}$
- BLR is not resolved: $U = 10^{-2} \text{ to } 10^{-1}$ from photoionization models
- Dust cannot survive in the BLR. Seyferts have “normal” abundances

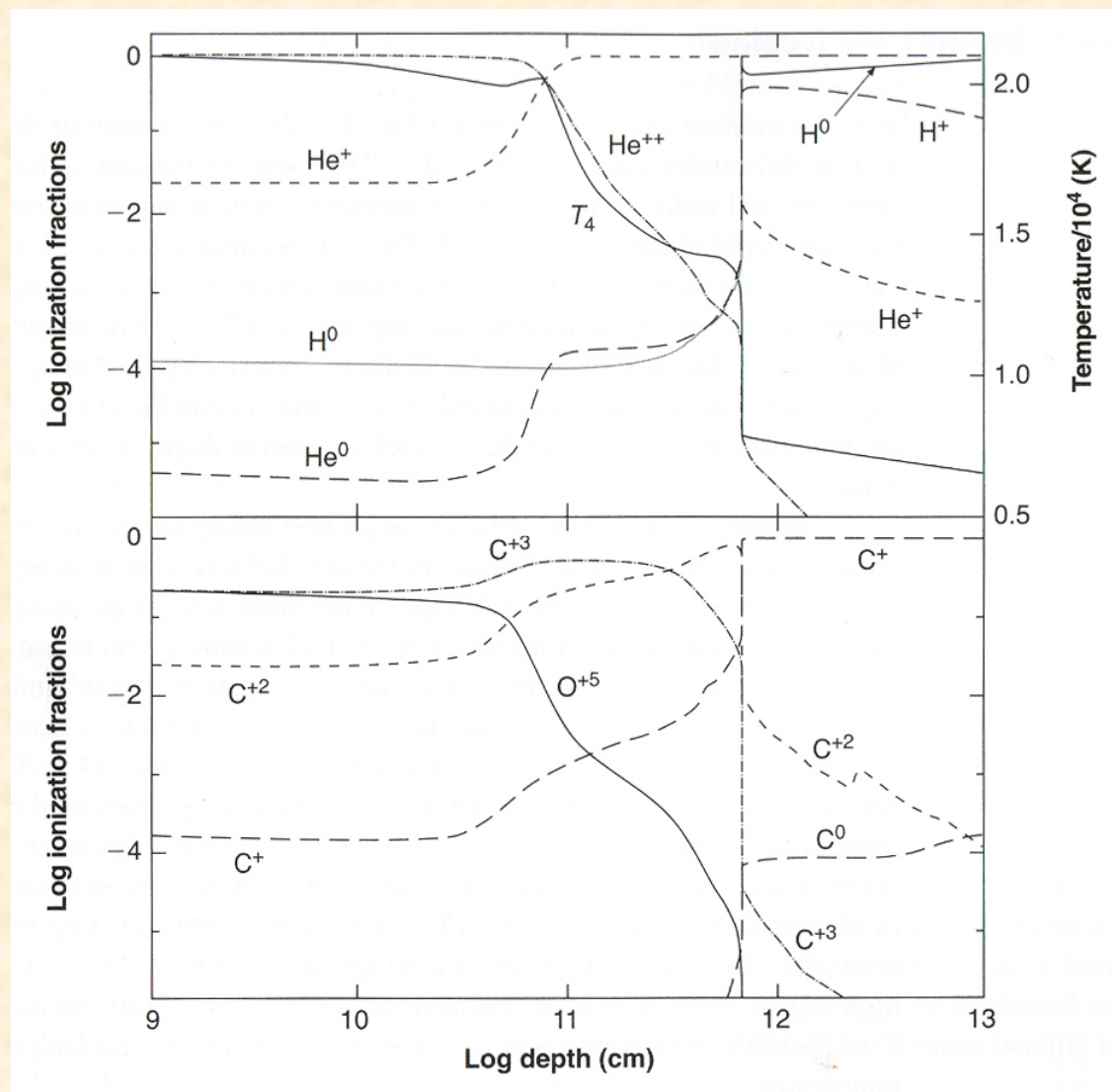
The “ $L\alpha/H\beta$ ’ ’ Problem

- Baldwin (1977) discovered the “ $L\alpha/H\beta$ ’ ’ problem by piecing together spectra of QSOs at different redshifts
 - $L\alpha/H\beta \approx 5 - 10$ for the BLR, whereas recombination gives ~ 33
 - What’s going on?
- BLR clouds have large column and number densities.
- $L\alpha$ scatters throughout the PIZ in BLR clouds, populating the $n = 2$ level
- $H\beta$ (and $H\alpha$) are collisionally excited in the PIZ, and therefore enhanced by factors of $3 - 6$ over recombination values
- $L\alpha$ is further reduced by ionization of electrons in $n = 2$ level
- Currently, there is still an “Fe II” problem: models underpredict the amount of Fe II emission
 - huge number of levels, so radiative transfer (radiative pumping, resonance fluorescence, and transition coincidences) and collisional excitations are complicated

Fe II Partial Grotrian Diagram



BLR “Cloud” Photoionization Model



(Osterbrock & Ferland, p. 364)

BLR Line Ratios

Observed and predicted relative BLR emission-line intensities

Ion	λ (Å)	Observed ^a	$U = 10^{-1.5}$ Model	Multi-component Model
O VI	1034	0.1–0.3	0.019	0.16
L α	1216	1.00	1.00	1.00
N V	1240	0.1–0.3	0.039	0.04
Si IV + O IV	~1400	0.08–0.24	0.091	0.06
C IV	1549	0.4–0.6	0.77	0.57
He II + O III]	1666	0.09–0.2	0.13	0.14
C III] + Si III]	1909	0.15–0.3	0.077	0.12
Mg II	2798	0.15–0.3	0.16	0.34
H β	4861	0.07–0.2	0.045	0.09

a. The-observed intensities from a sample of intermediate ($z \approx 2$) redshift quasars.

(Osterbrock & Ferland, p. 365)

- note: no prediction of Fe II emission ... hmmm

BLR Parameters from Photoionization Models

- Sizes: typically ~ 10 light days (in diameter) for Seyfert 1s
 - 1) Reverberation mapping – use time lag (τ) of emission lines with respect to continuum variations: $r = c\tau$
 - 2) Photoionization models: Determine ionization parameter and density from models. Determine Q_{ion} from luminosity and SED.

$$U = \frac{Q_{\text{ion}}}{4\pi r^2 c n_e} \rightarrow \text{solve for } r. \quad \text{To 1st order, } r \propto \sqrt{L}$$

- Mass of ionized gas in BLR:

$$L(\text{H}\beta) = n_e n_p \alpha_{\text{H}\beta}^{\text{eff}} h\nu_{\text{H}\beta} V \epsilon \quad \text{where } \epsilon = \text{filling factor}$$

$$M_{\text{BLR}} \approx V \epsilon n_p m_p = \frac{L(\text{H}\beta) m_p}{n_e \alpha_{\text{H}\beta}^{\text{eff}} h\nu_{\text{H}\beta}} \approx 0.7 L_{42}(\text{H}\beta) \frac{10^{10} \text{cm}^{-3}}{n_e} M_{\odot}$$

- Filling factor ϵ : assume a spherical BLR ($V = \frac{4}{3} \pi r^3$)

$$\text{From above: } \epsilon = \frac{L(\text{H}\beta)}{n_e n_p \alpha_{\text{H}\beta}^{\text{eff}} h\nu_{\text{H}\beta} V} \approx 0.01 - 0.1$$

- Covering factor – fraction of sky covered by BLR clouds
 - assume all ionizing photons are absorbed and use predicted equivalent width of emission line $W(\text{H}\beta)$

$$L_{\text{H}\beta} = h\nu_{\text{H}\beta} \frac{\alpha_{\text{H}\beta}^{\text{eff}}(\text{H}^0, T)}{\alpha_{\text{B}}(\text{H}^0, T)} \int_{\nu_0}^{\infty} \frac{L_{\nu}}{h\nu} d\nu$$

$$L_{\text{H}\beta} = L_{\lambda}(\lambda 4861) W_{\lambda}(\text{H}\beta) = L_{\nu}(\lambda 4861) \frac{d\nu}{d\lambda} W_{\lambda}(\text{H}\beta)$$

Assume power - law continuum : $L_{\nu} = C\nu^{-n}$

$$\text{Then } W_{\lambda}(\text{H}\beta) = \frac{\lambda_{\text{H}\beta}}{n} \frac{\alpha_{\text{H}\beta}^{\text{eff}}(\text{H}^0, T)}{\alpha_{\text{B}}(\text{H}^0, T)} \left(\frac{\nu_0}{\nu_{\text{H}\beta}} \right)^{-n} = \frac{568}{n} (5.33)^{-n}$$

(for a covering factor of 1)

So for $n = 1$, the predicted EW is $W_{\lambda}(\text{H}\beta) \approx 106 \text{ Ang.}$

$$\text{The covering factor is : } C_f = \frac{W_{\text{obs}}(\text{H}\beta)}{W_p(\text{H}\beta)} = \frac{20}{106} \approx 0.2$$

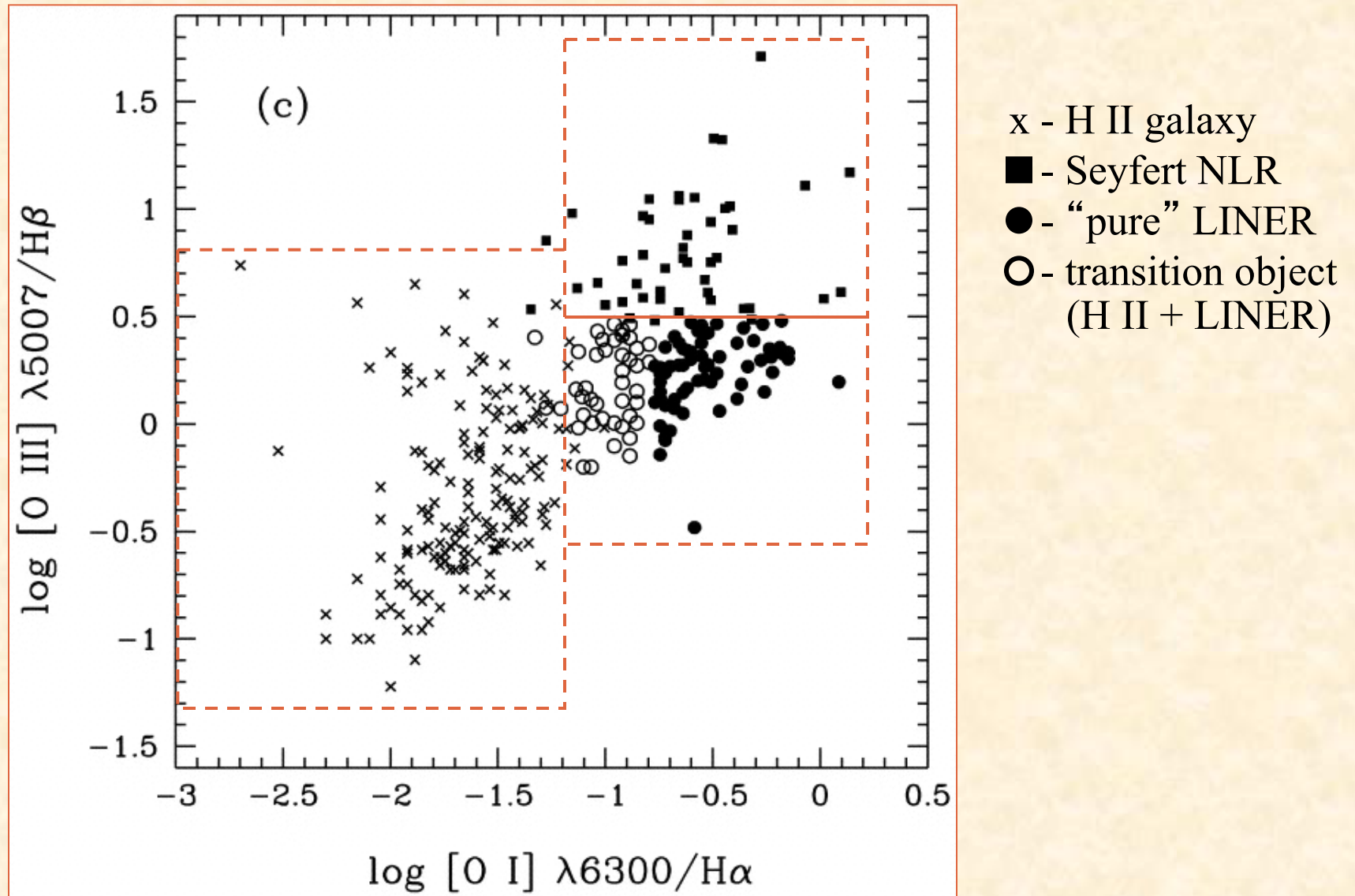
LINERs

- $[\text{O III}]/\text{H}\beta < 3$ (like galaxies with H II nuclei)
- $[\text{O I}]/\text{H}\alpha > 0.05$ (like the NLR in Seyferts)
- Original suggestion: shock heating or hot stars
- However, subsequent evidence indicates photoionization by AGN continuum (including X-rays) is likely for most
- $U = 10^{-3}$ to 10^{-5} for LINERs (rather than 10^{-1} to 10^{-2} for Seyferts)
- Probably due to low luminosity of continuum source, rather than higher density or greater distance

$$U = \frac{Q_{\text{ion}}}{4\pi r^2 c n_e}$$

- Further evidence for AGN: $\sim 20\%$ of LINERs show a mini BLR (type 1 LINERs)
- Transition objects: may be combination of starburst and AGN

Emission-Line Diagnostics (BPT Diagram)

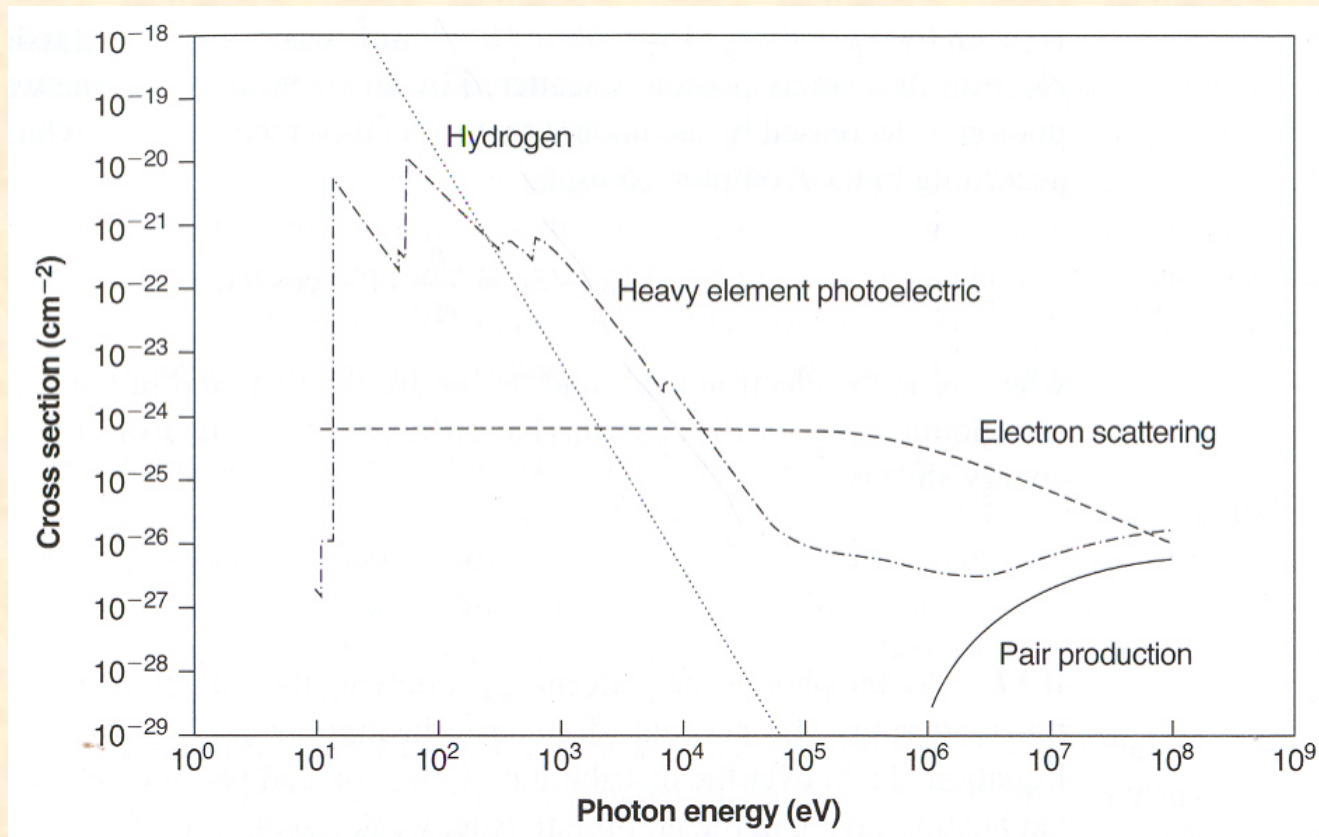


(Ho, Filippenko, & Sargent, 1997, ApJS 112, 315)

High Energy Processes/ X-ray Spectra of AGN

- X-ray spectra of AGN show evidence for hot photoionized gas ($T = 30,000 - 100,000 \text{ K}$; $U = 1 - 10$)
- Heating:
 - Photoionization of inner and outer shell electrons
 - Collisional ionization from ejected outer-shell (suprathermal) electrons
- Cooling:
 - Recombination lines: dominant in X-ray spectra (transitions to inner shells: $n = 1, 2, 3$ corresponding to K, L, M)
 - Fluorescence after ejection of inner-shell electrons: competes with Auger effect
 - Radiative recombination continuum (RRC), e.g., Lyman continuum (LC) : narrow, since $kT \ll \text{I.P.}$
 - Two-photon: significant, since critical density for 2s in H-like heavy elements is $> 10^{14} \text{ cm}^{-3}$.
 - Photoexcitation important due to many lines in spectra.
 - Collisional excitation: not so important in X-ray spectra: $kT \ll \chi$

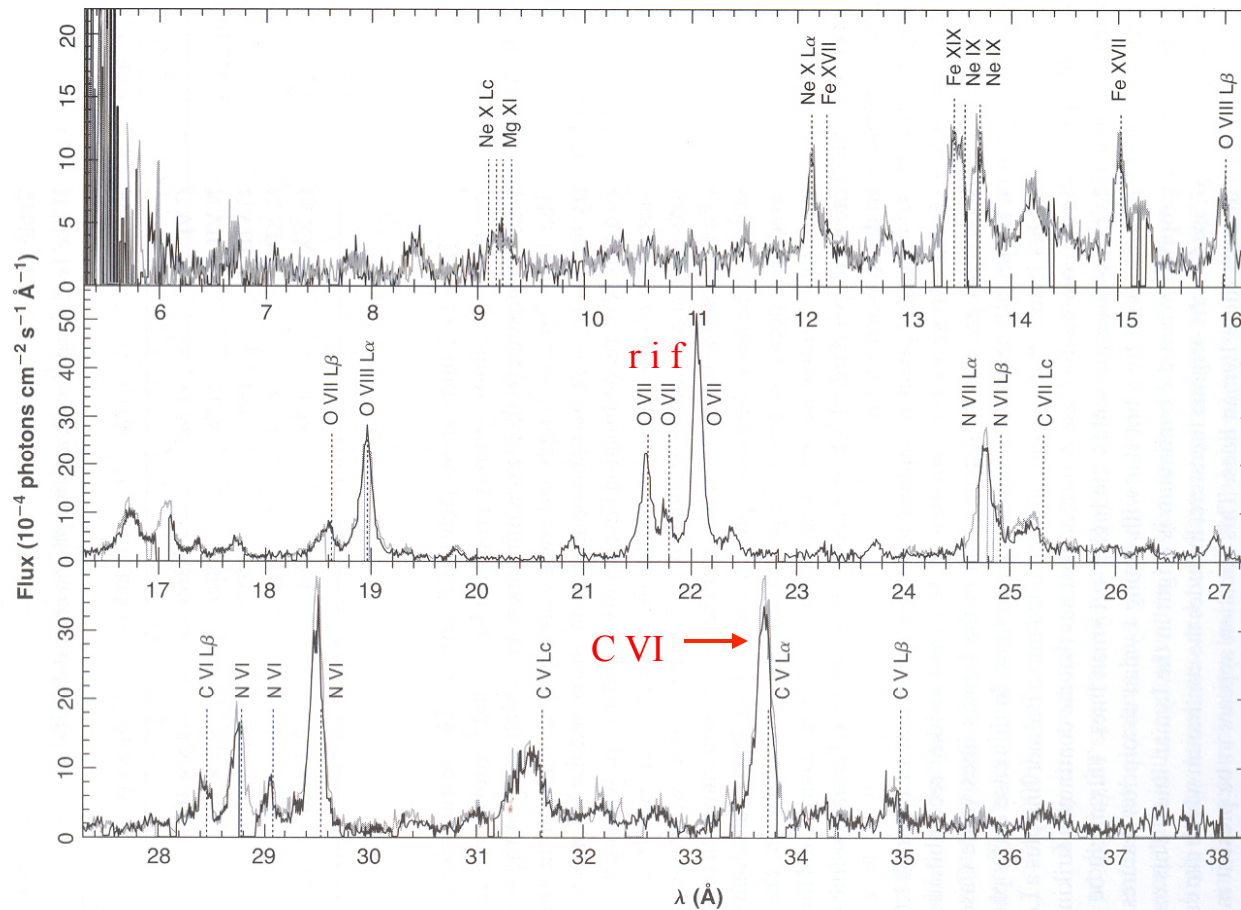
X-ray opacities



(Osterbrock & Ferland, p. 283)

- If seen in absorption, we can see the effects of absorption edges and scattering on the ionizing continuum: $\tau_v = a_v N_{\text{ion}}$
- The gas starts to become “Compton thick” at $N_e \sim 1/a_v \sim 10^{24} \text{ cm}^{-2}$

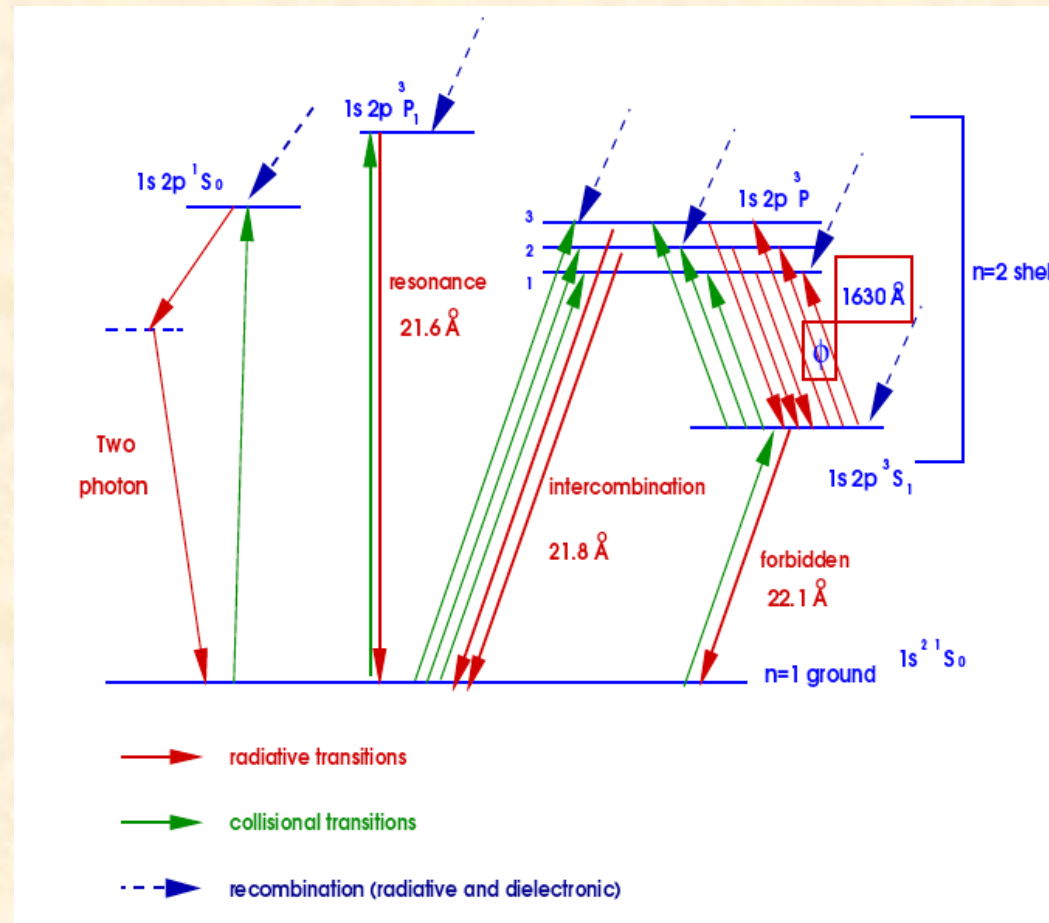
Soft X-ray Emission-Line Spectra



(Osterbrock & Ferland, p. 287)

- Chandra images reveal extended X-ray gas in the NLR of NGC 1068
- Chandra spectra reveal emission lines - mostly H and He-like.
- Observed lines can be matched by photoionization models with $U \approx 1 - 10$

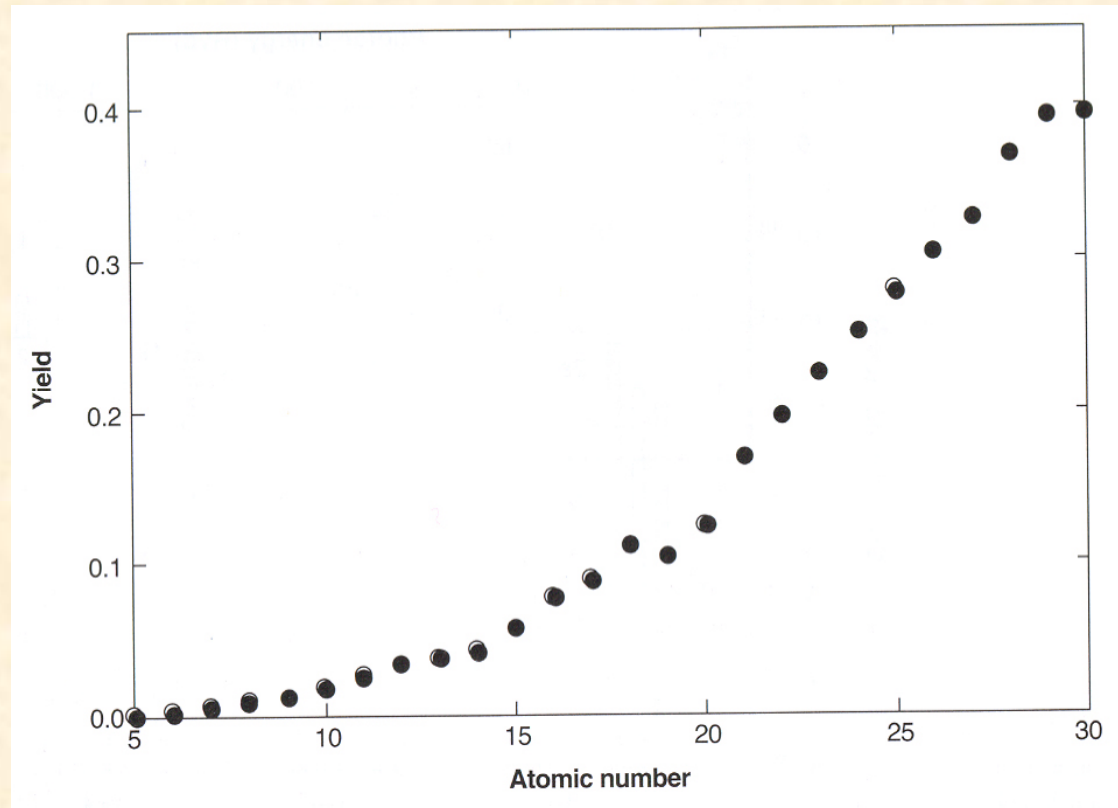
He-like Triplet Lines (rif): O VII



(Morales 2002)

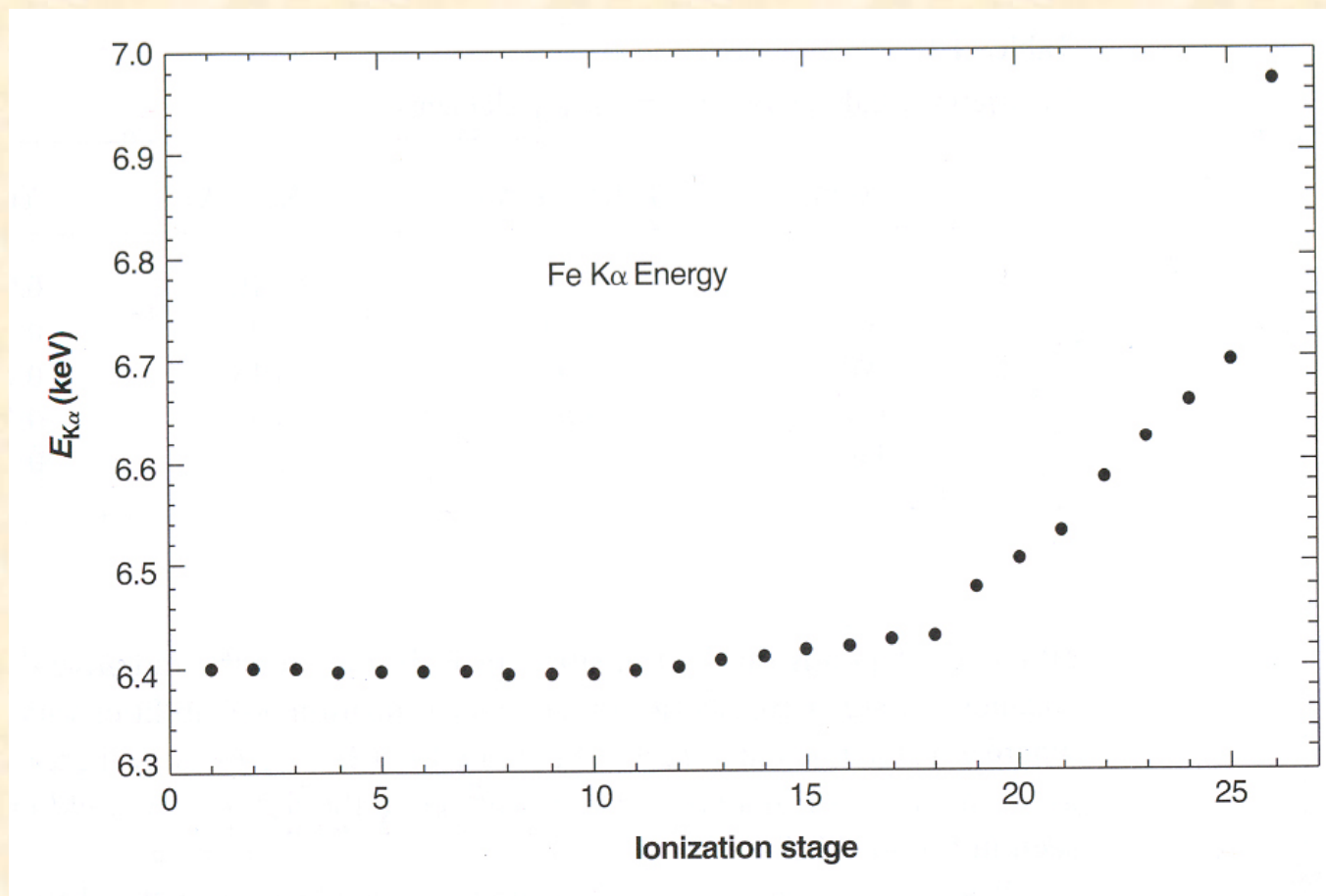
- Triplet lines are sensitive to density, since the intercombination and forbidden lines can be collisionally de-excited.
- Unfortunately, the critical densities are rather high: $n_e \approx 10^{10}\ \text{cm}^{-3}$, so not much help for the NLR (useful for higher density gas).

High-Energy Processes - Inner Shell Ionization



(Osterbrock & Ferland, p. 280)

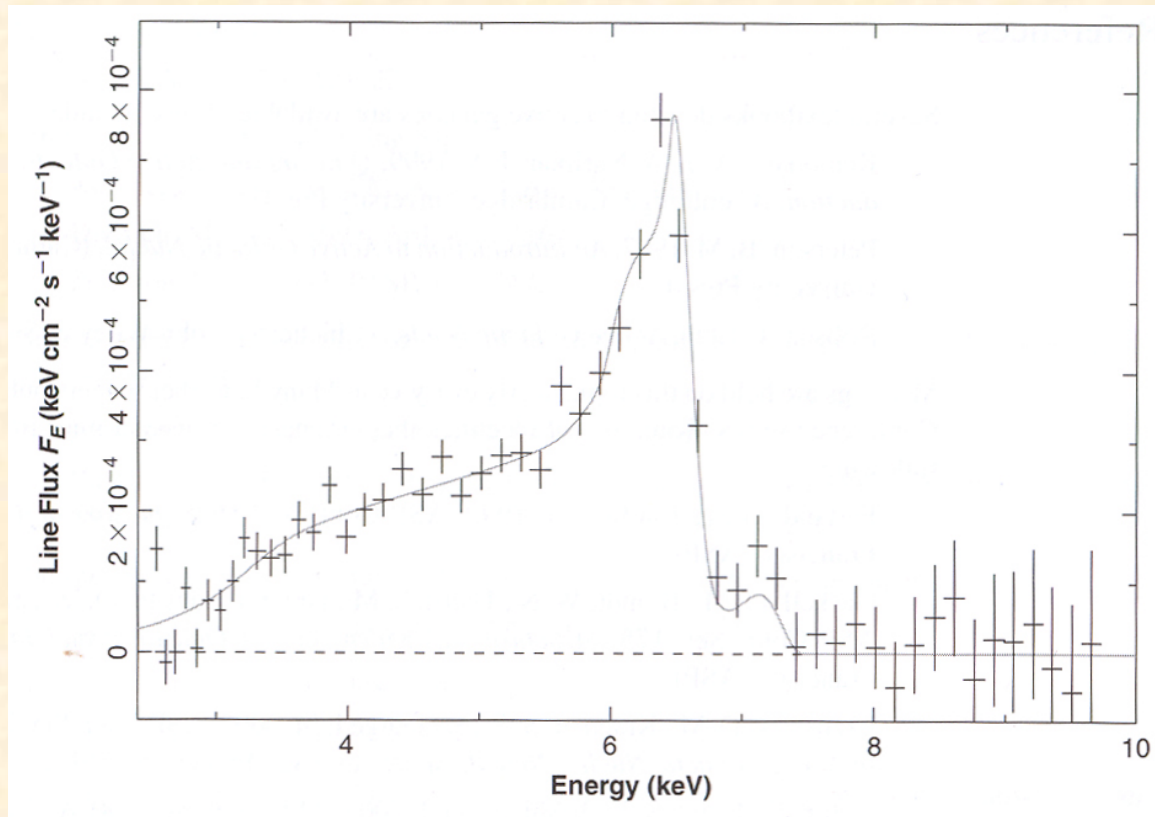
- Inner shell ionization - vacancy filled by ejection of electrons (Auger effect) and/or fluorescent emission
- Yield = probability of filling K-shell vacancy by emission of $K\alpha$ line.
- Fe ($Z = 26$) is abundant and has a high yield: Fe $K\alpha$ is strong in hard X-ray region.



(Osterbrock & Ferland, p. 282)

- Energy of Fe K α increases with ionization state, as there is less “screening” of nucleus by outer electrons with non-zero wave functions close in.
- With high-resolution spectroscopy, one can determine ionization state of gas from Fe peak (Fe XVIII and lower often known as “cold iron” by X-ray astronomers).

Fe K α Emission from MCG-6-30-15



(Osterbrock & Ferland, p. 349)

- Relativistic disk fit to Fe K α profile, velocity up to $\sim 0.4c$
- Rest-frame line center at 6.4 keV - consistent with emission from cold accretion disk.
- Peak slightly blueshifted due to Doppler boosting of approaching gas.
- Long red tail due to GR: can measure BH mass and spin.