

Dust in the Diffuse Universe

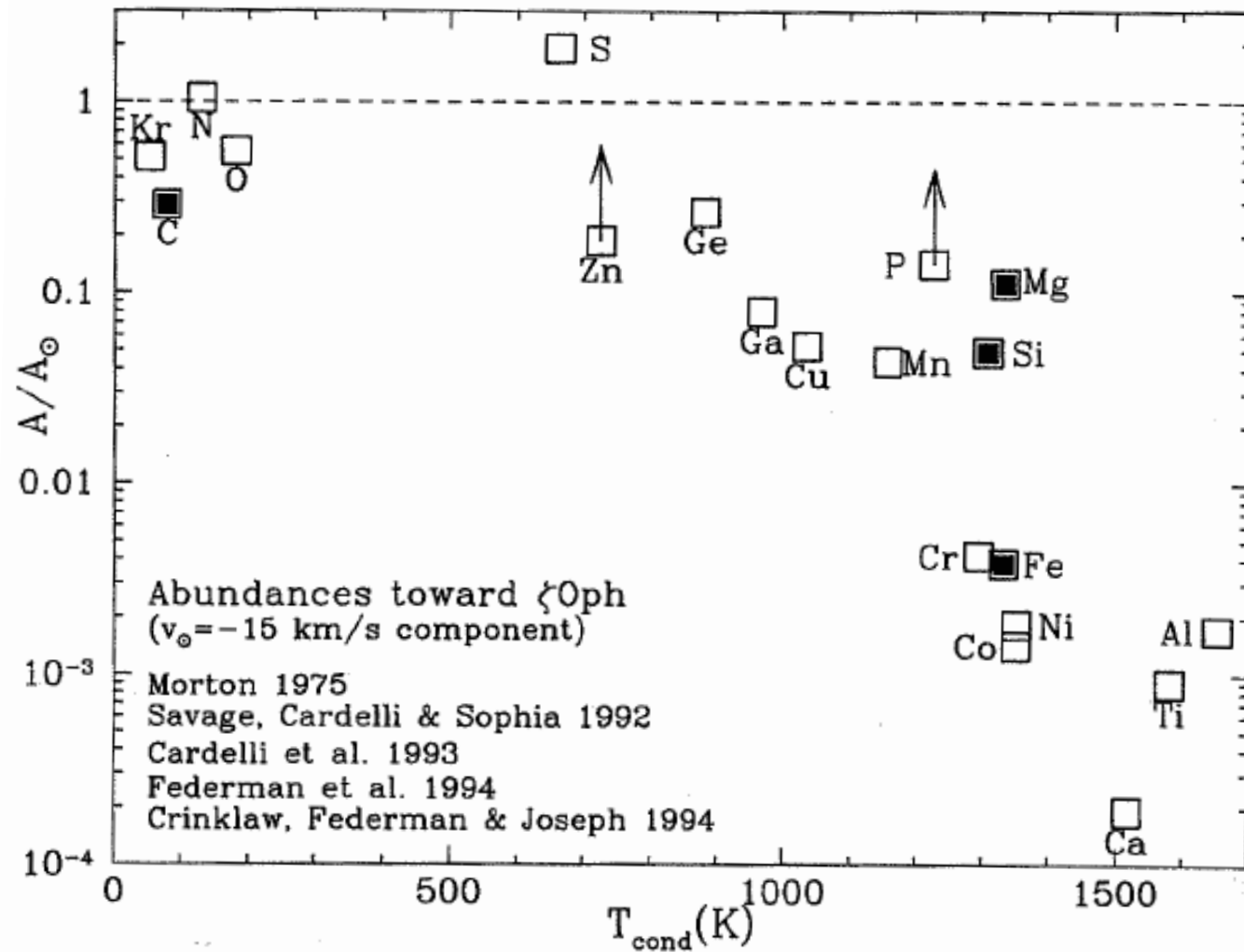
- Obscuring Effects
- Chemical Effects
- Thermal Effects
- Dynamical Effects
- Diagnostic Power

Evidence for Grains:

Chemical Effects

- Catalyzes molecular hydrogen formation.
- Depletion of elements (gas to solid)
 - The relative elemental abundances in interstellar clouds are quite different from that measured in the Sun.
 - Elements that can form refractory solids are among those which have high depletion.
 - Si and O can form silicates with Mg and Fe
 - Fe may form graphite, silicon carbide, or iron oxides
 - Suggests that as cooling gas moves away from a star, some elements condense into solid particles and are removed from the gas phase

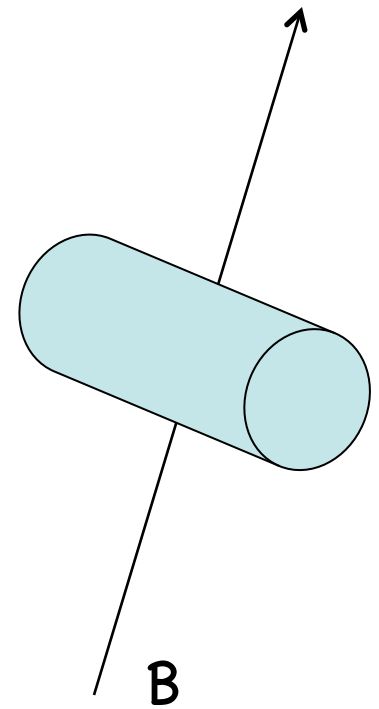
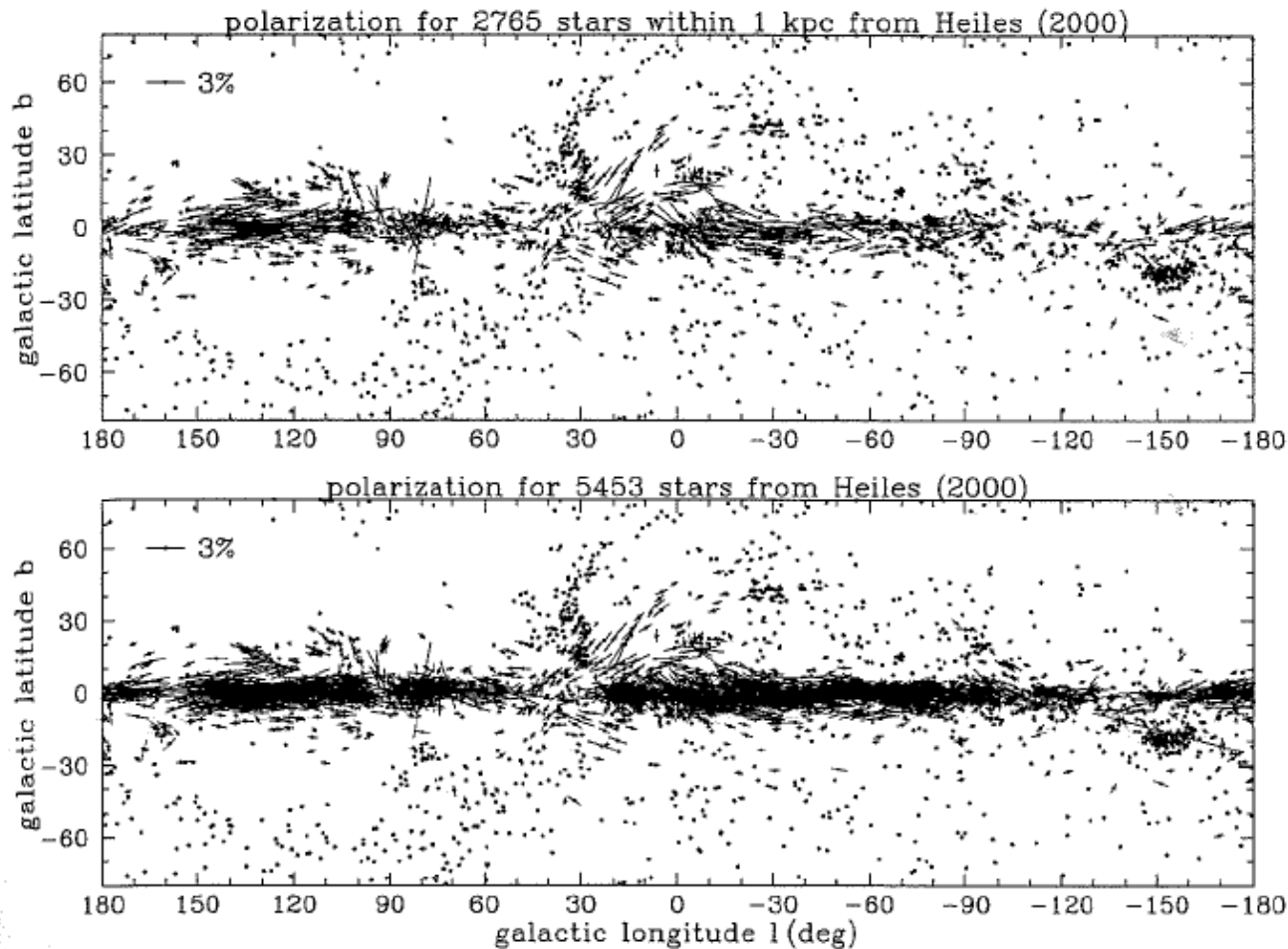
Gas-phase Abundances Relative to Solar in a Diffuse Cloud



Evidence for Grains: Polarization

- Starlight is linearly polarized by a few percent
- The amount of polarization observed is correlated with the amount of extinction in that direction.
- Grains such as graphite are anisotropic.
- Polarization requires partial alignment of elongated grains.
- Anisotropic grains will normally be rotating due to collisions. Magnetic moment tends to align with the interstellar magnetic field.

Linear Polarization of Starlight



Evidence for Grains: Scattered Light

- Cross section for scattering by grains is several orders of magnitude larger than that for Rayleigh scattering by atoms & small molecules
- Albedo = scattering cross section / extinction cross section
- Scattering is about as important as absorption at optical wavelengths (0.5 μm)

Zodiacal Light



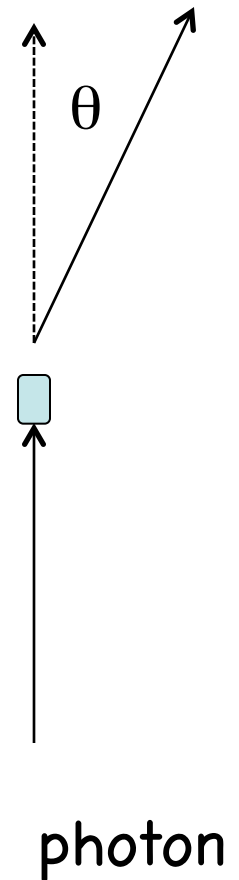
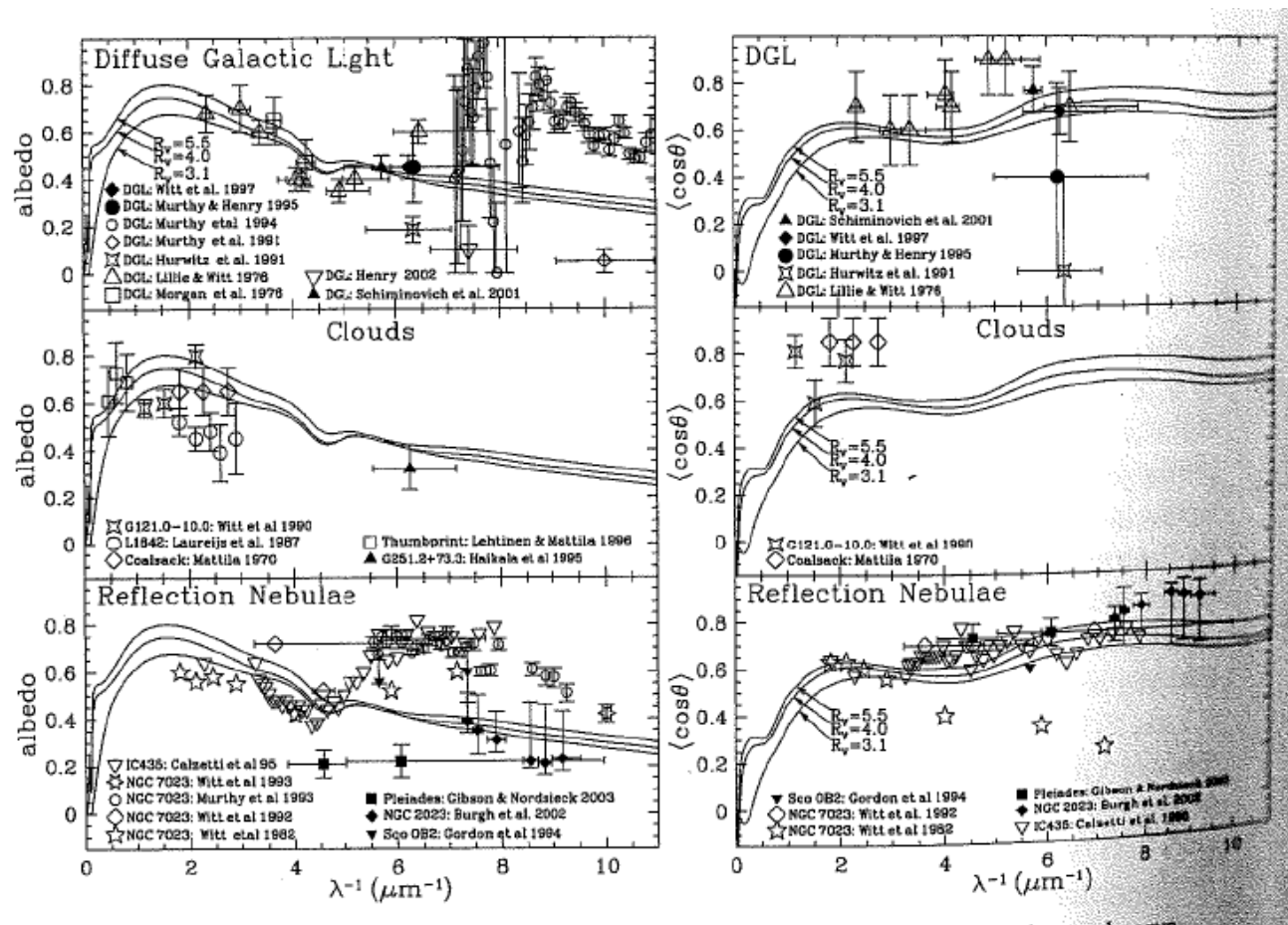
Winter 2016

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Reflection Nebulae

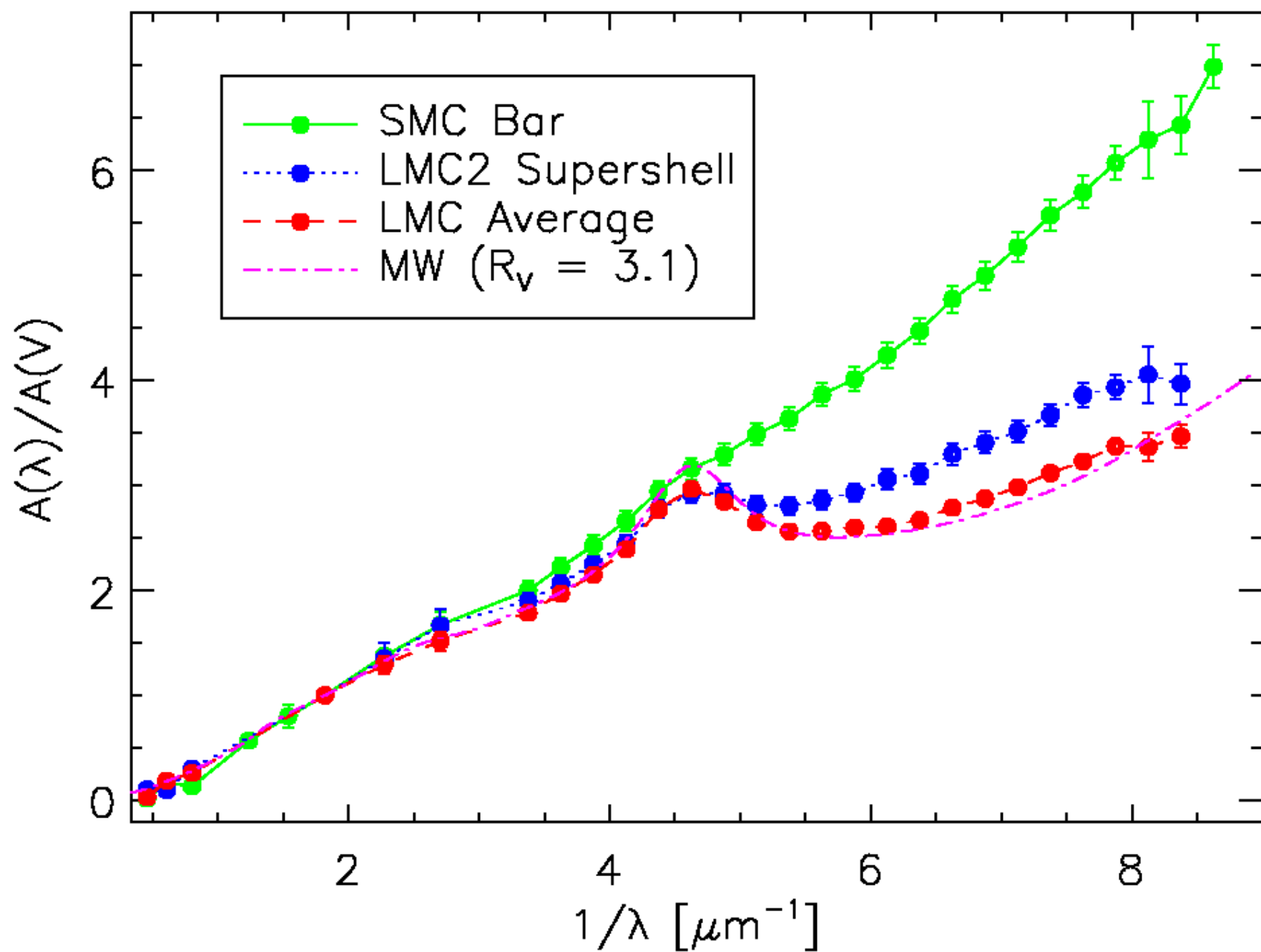


Albedo and Scattering Angle



Obscuring Effect of Grains

- Extinction reduces the observed intensity of starlight
 - Balmer decrement
- Measure brightness of two similar stars, only one of which is behind a gas cloud (correcting for distance differences).
 - Extinction curve
 - Can't measure it below the HI Lyman limit.
 - Hard to measure at very long wavelengths where it's small.
 - Ratio of selective to total extinction is sensitive to grain size.



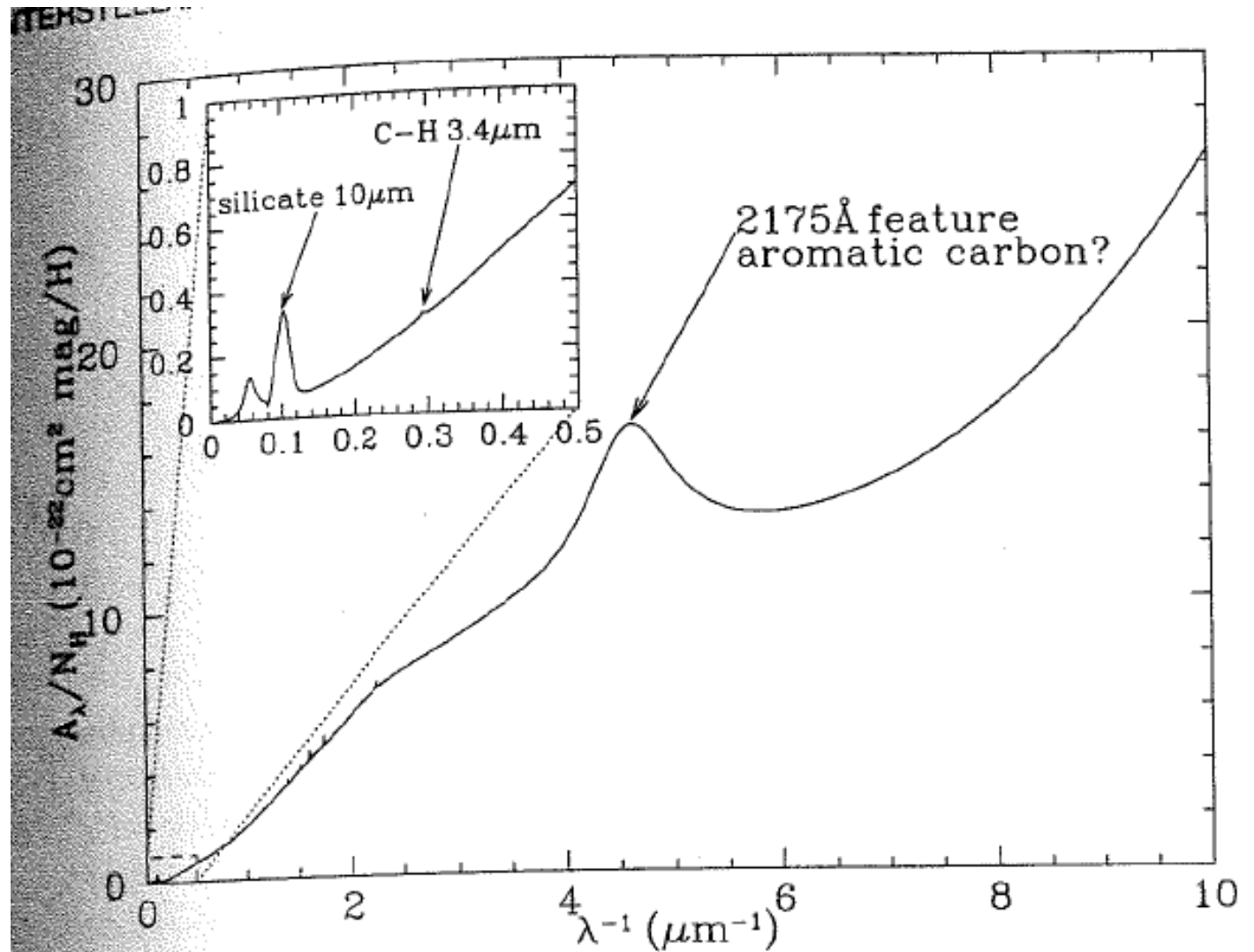
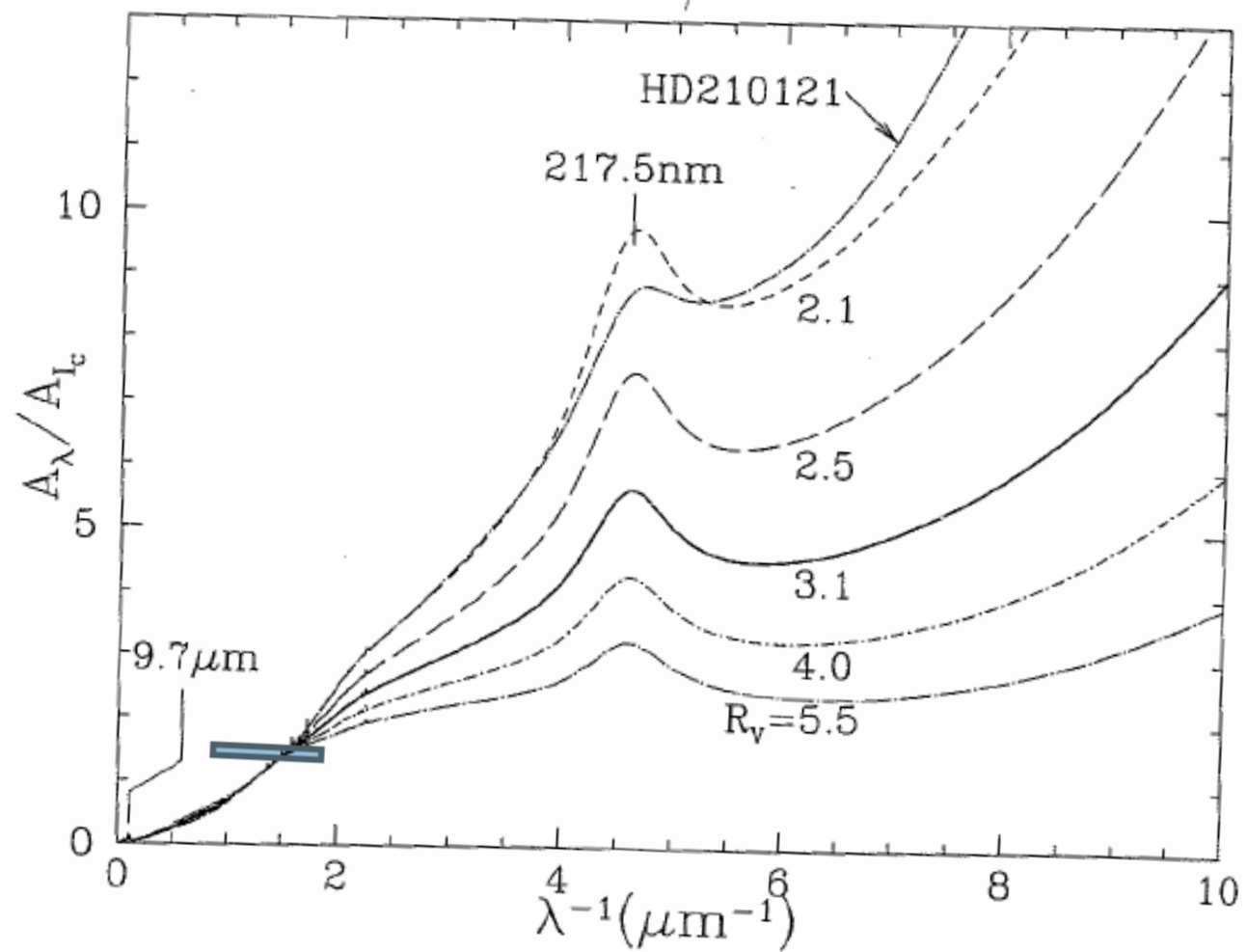


Figure 21.1 Extinction versus inverse wavelength λ^{-1} on a typical sightline in the local diffuse ISM. The inset shows the extinction at $\lambda > 2 \mu\text{m}$.

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Size & Shape of Grains

- Wavelength dependence of extinction (A_λ) in the optical is roughly λ^{-1} , even at 1000 Å.
 - Requires very small grains make a substantial contribution to extinction.
 - $2 \pi a \leq \lambda$, so $a \leq 0.015 \text{ } \mu\text{m}$.
- Starlight is significantly polarized at 5500, so the grains with $a \sim 0.1 \text{ } \mu\text{m}$ are non-spherical and substantially aligned with the B field.
- The polarization of starlight decreases towards bluer wavelengths. The grains with $a \leq 0.05 \text{ } \mu\text{m}$ which dominate the extinction at $\lambda < 0.3 \text{ } \mu\text{m}$ are either spherical or minimally aligned.
- Most of the mass is in the larger grains.
- Most of the surface area is in the smaller grains.

Composition of Grains

- Pronounced peak at 2175 Å in Milky Way is largely absent in Magellanic Clouds,
 - Attributed to C sheets
 - Drude profile
- 3.4 μm feature from C-H stretching
- Silicate features
 - 9.7 μm Si-O stretching
 - 18 μm Si-O-Si stretching

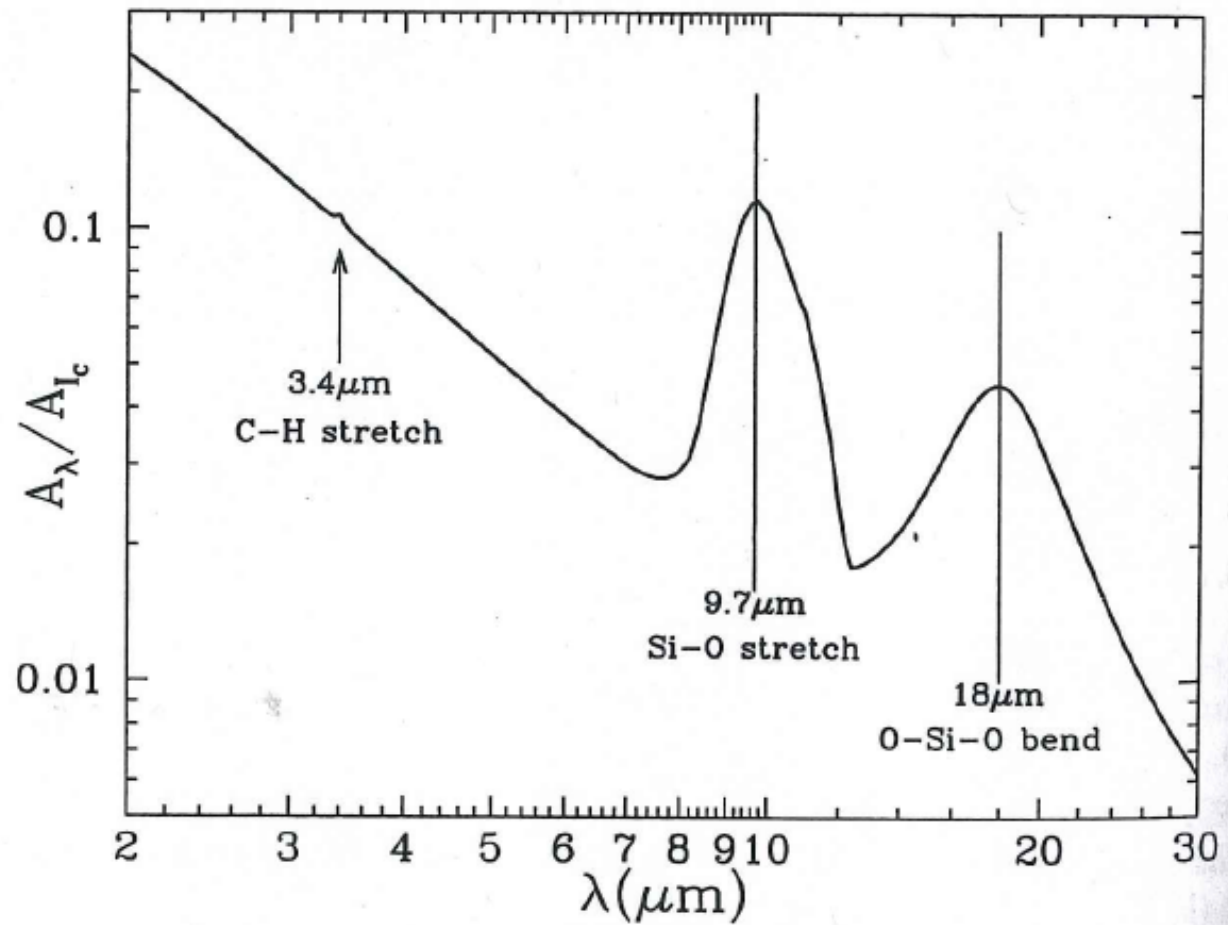


Figure 23.2 Infrared extinction curve. The 8 to 13 μm silicate profile is as observed toward the Galactic Center by Kemper et al. (2004), but with $A_V/\Delta\tau_{9.7\mu\text{m}} = 18.5$, as appropriate for sightlines through diffuse gas within a few kpc of the Sun (see Table 1 of Draine 2003a). The 3.4 μm C-H stretching feature is indicated.

Temperature of Grains

- Heated by starlight
- Grain temperature \ll Stellar temperature

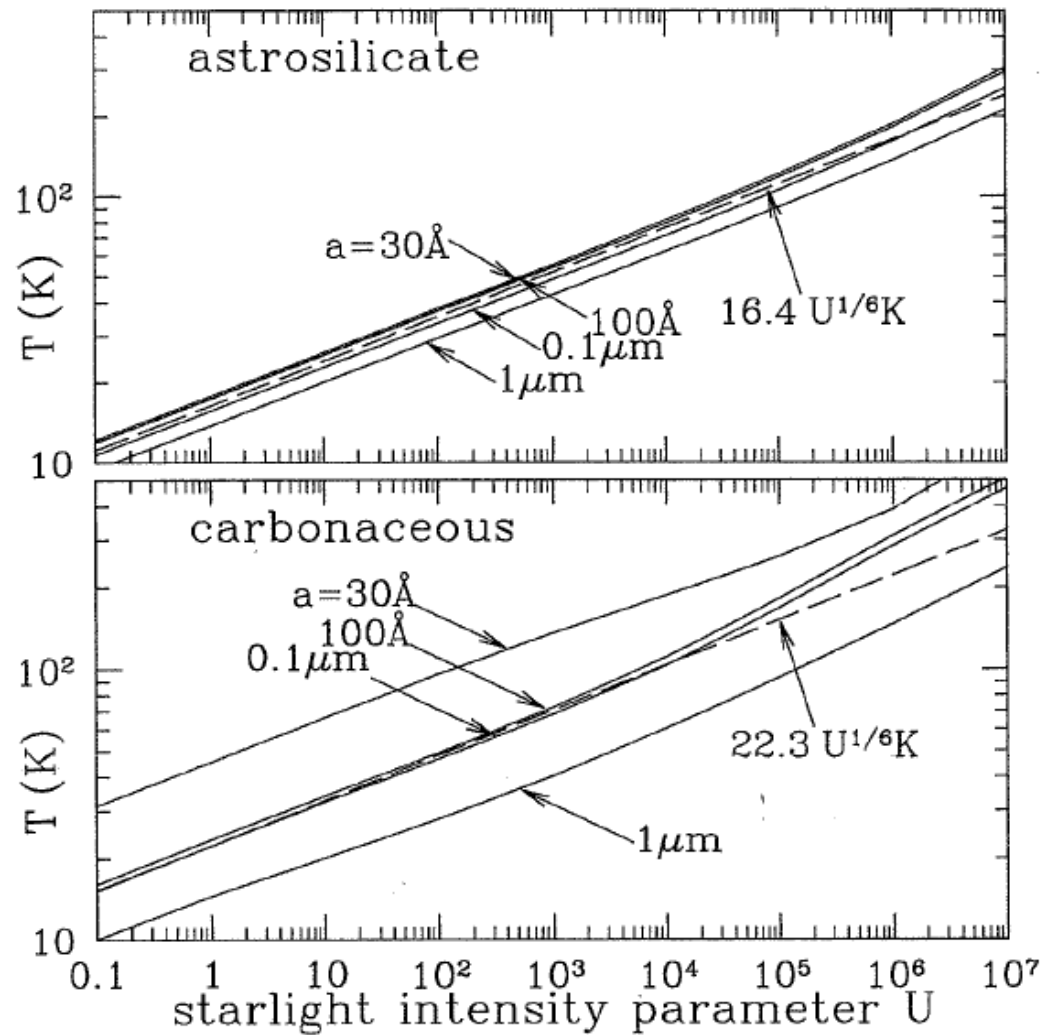
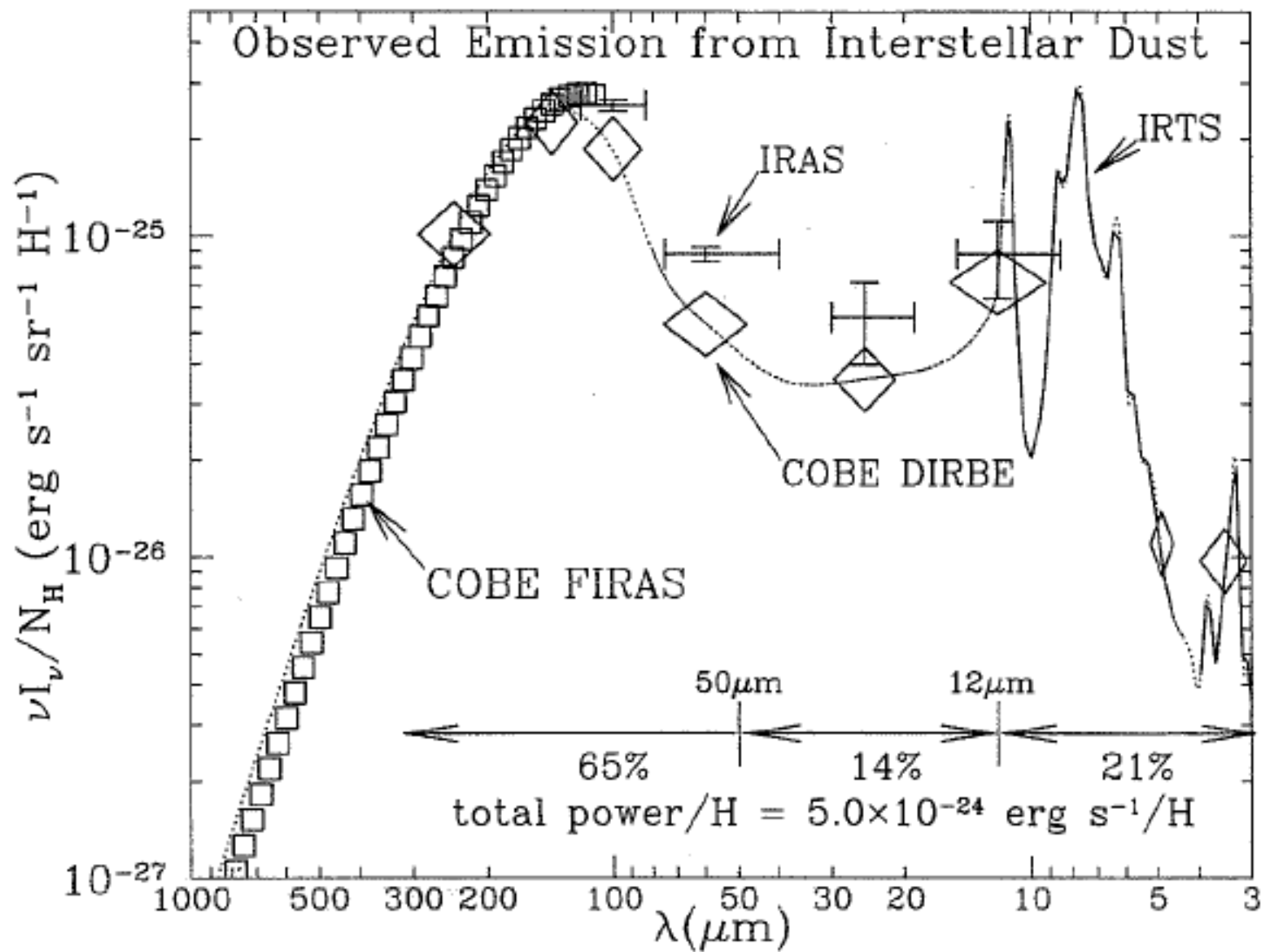


Figure 24.4 Equilibrium temperature for astrosilicate and carbonaceous grains heated by starlight with the spectrum of the local radiation field, and intensity U times the local intensity. Also shown are the power-laws $T = 16.4U^{1/6} \text{ K}$ and $T = 22.3U^{1/6}$ for $a = 0.1 \mu\text{m}$ from Eqs. (24.19 and 24.20).



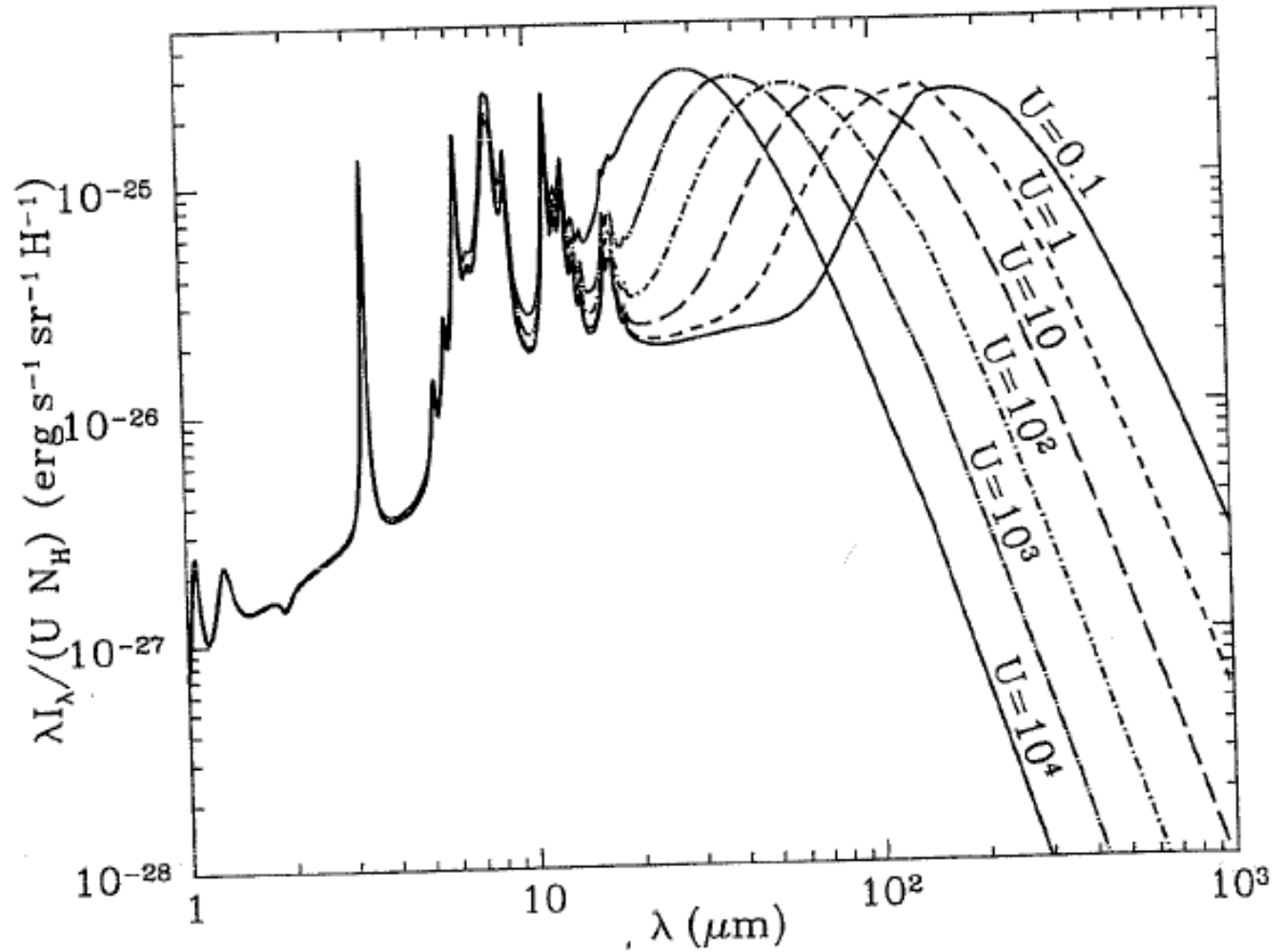
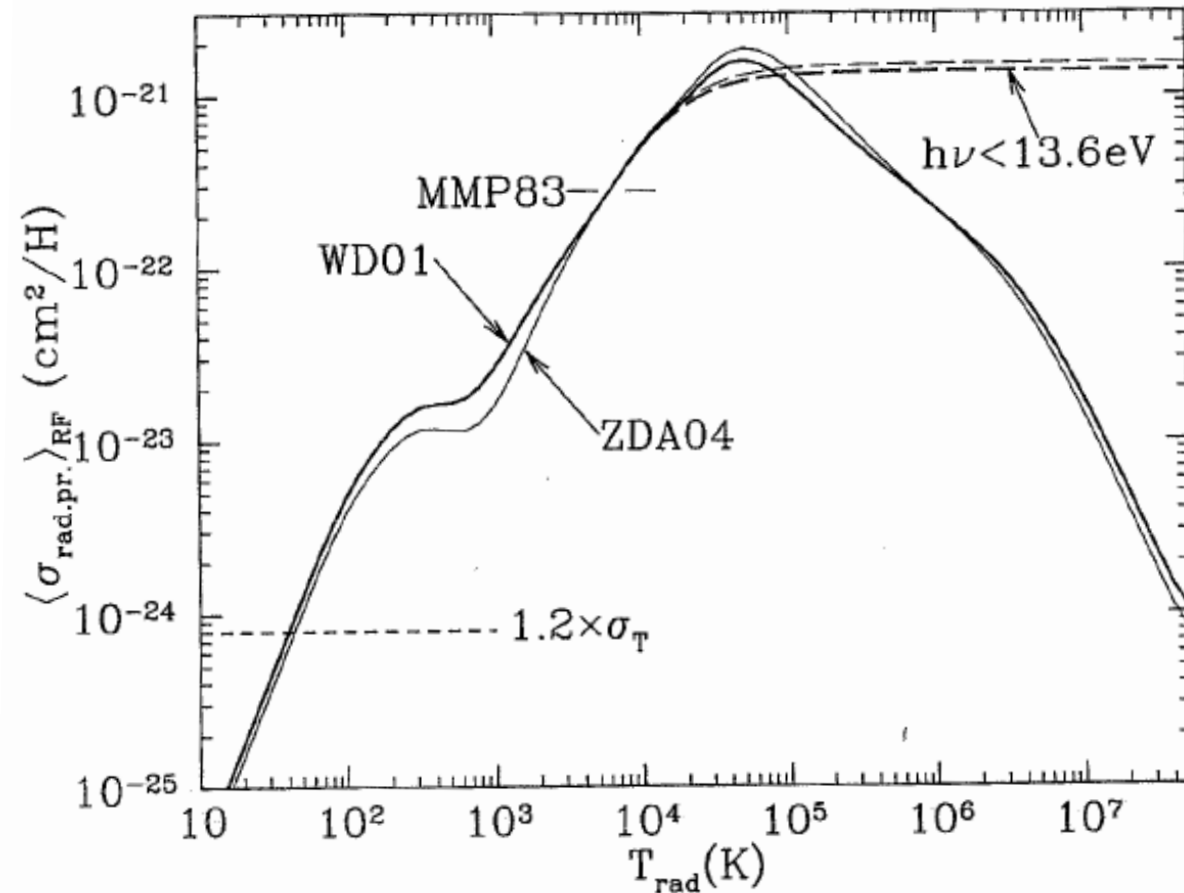


Figure 24.7 Infrared emission spectrum for model with silicate and graphite/PAH grains in ISRF intensity scale factor U from 0.1 to 10^4 ($U = 1$ is the local ISRF). Spectra are scaled to give power per H nucleon per unit U , calculated using the model of Draine & Li (2007).

Radiation Pressure Cross Section per H (averaged over a Planck Spectrum)



Photoelectric Heating of Gas

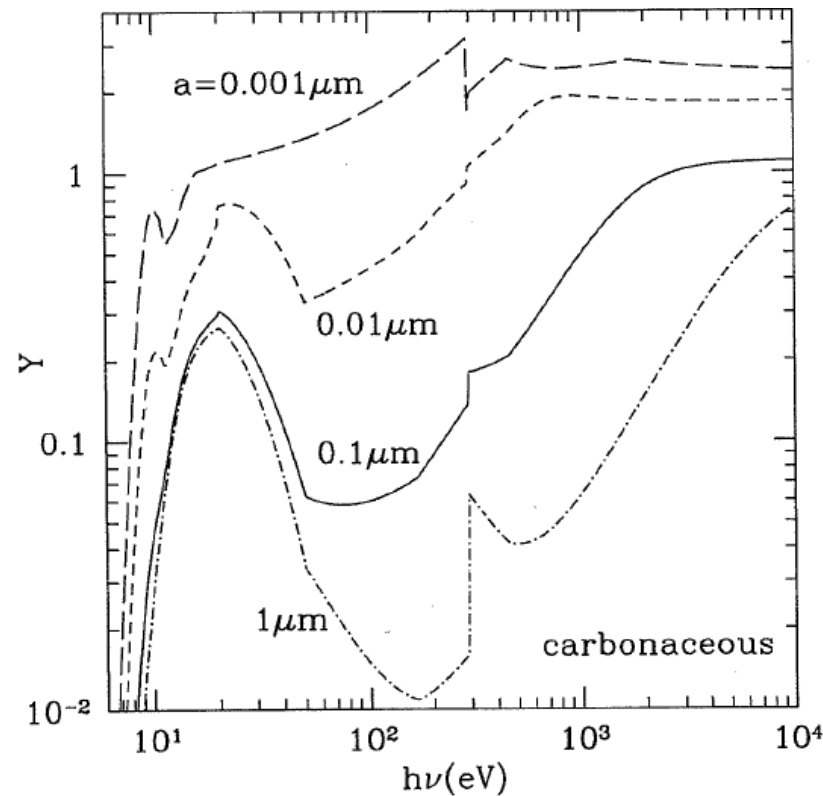
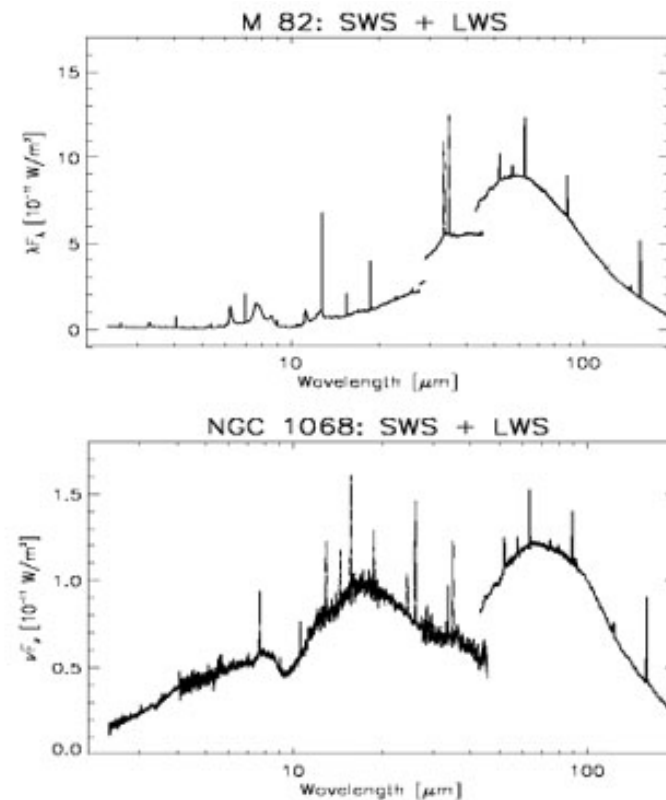
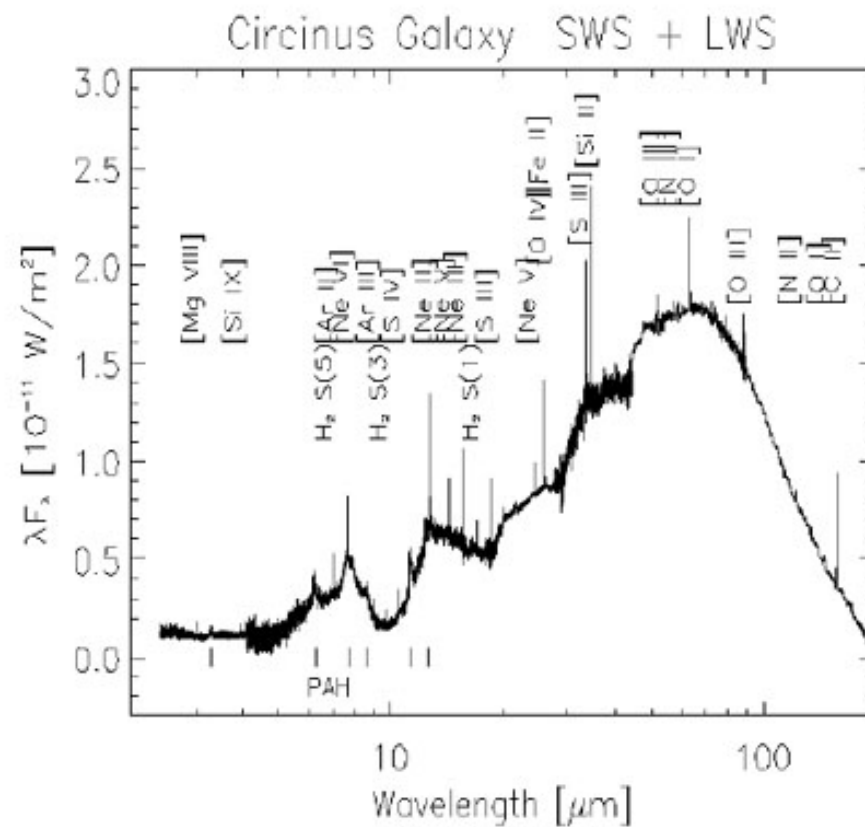


Figure 25.1 Photoelectric yield Y for uncharged carbonaceous grains for selected radii, as a function of photon energy. For sufficiently small grains, Y can exceed unity because of secondary electron emission for $h\nu \gtrsim 14$ eV, and Auger electron emission for $h\nu > 291$ eV. From Weingartner et al. (2006), reproduced by permission of the AAS.

Diagnostic Power

- Total heating rate often reveals 'obscured' star formation.
- Infrared spectroscopy of solids
- Absorption features from particular bonds in solid particles

Solid state spectral lines



Optics of Grains

- How do grains of a given size, shape, and refractive index modify the radiation which falls on them?
- Just apply Maxwell's equations with appropriate boundary conditions on the grain surface.
- Produce extinction curves for solids of various size, shape, and refractive index.
- Juggle these to build the observed extinction curve. For the optical, need
 - Size $a=3e-5$ cm
 - $n_{\text{grains}}=1e-12 n_{\text{H}}$
 - Refractive index ~ 1.3

End Lecture 13

- Draine
 - Ch 21, 24, 25.22 (assigned)
 - 22, 23, 25, 26 (as time permits)
- DS:
 - Ch 12
- Osterbrock
 - Ch 8
- Tielens:
 - Ch 4, 6, 9

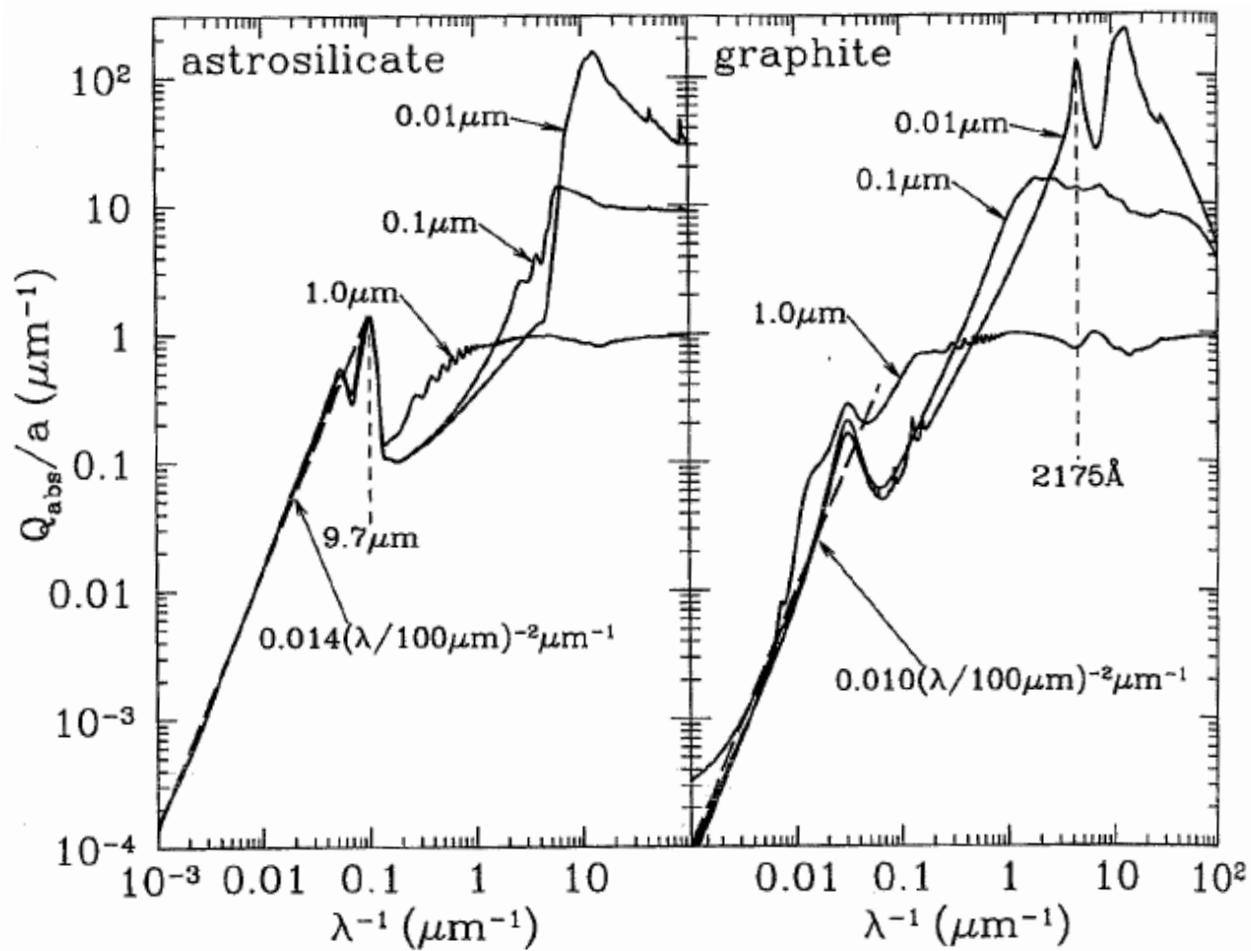


Figure 24.1 Absorption efficiency $Q_{\text{abs}}(\lambda)$ divided by grain radius a for spheres of amorphous silicate (left) and graphite (right). Also shown are power-laws that provide a reasonable approximation to the opacity for $\lambda \gtrsim 20\mu\text{m}$.