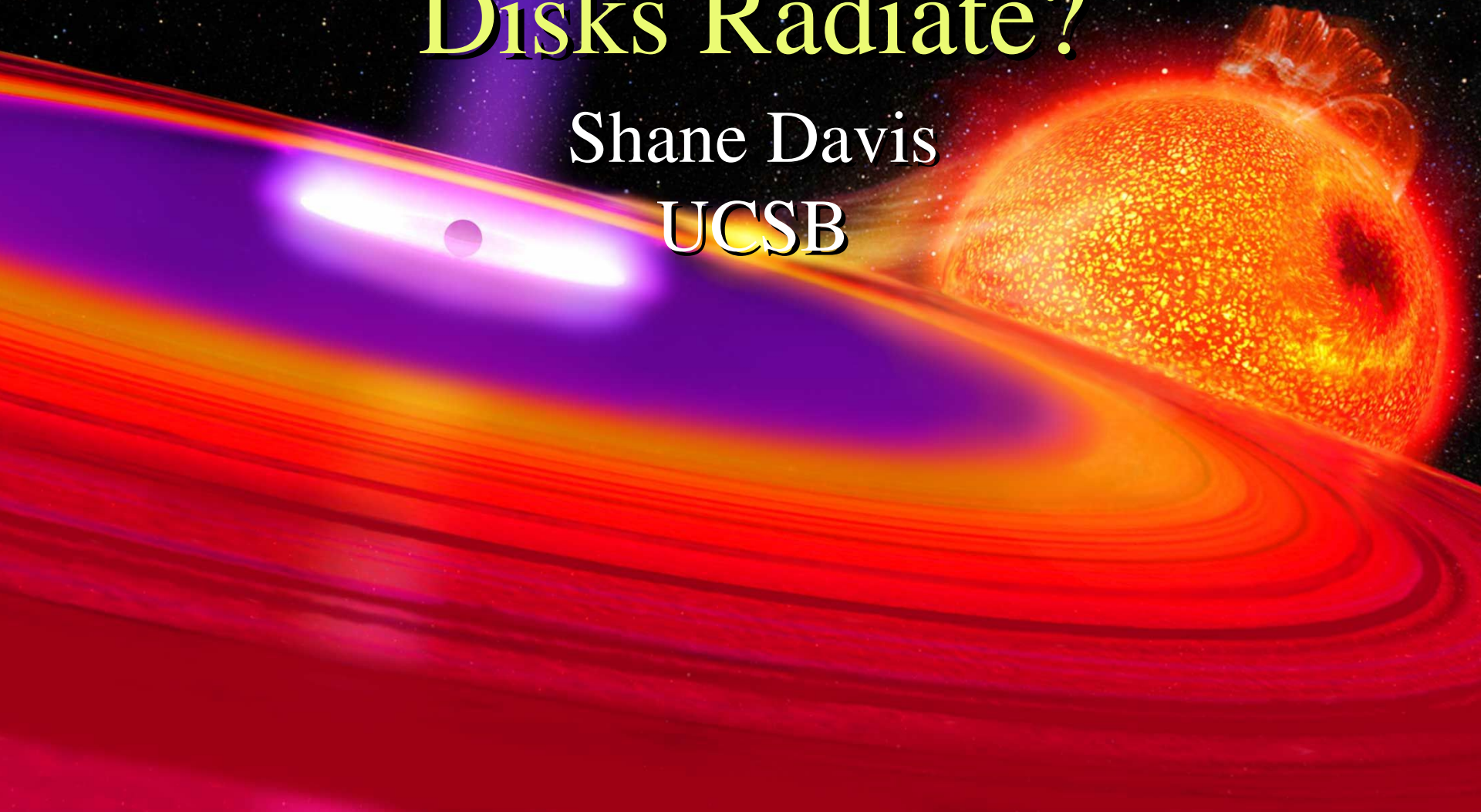
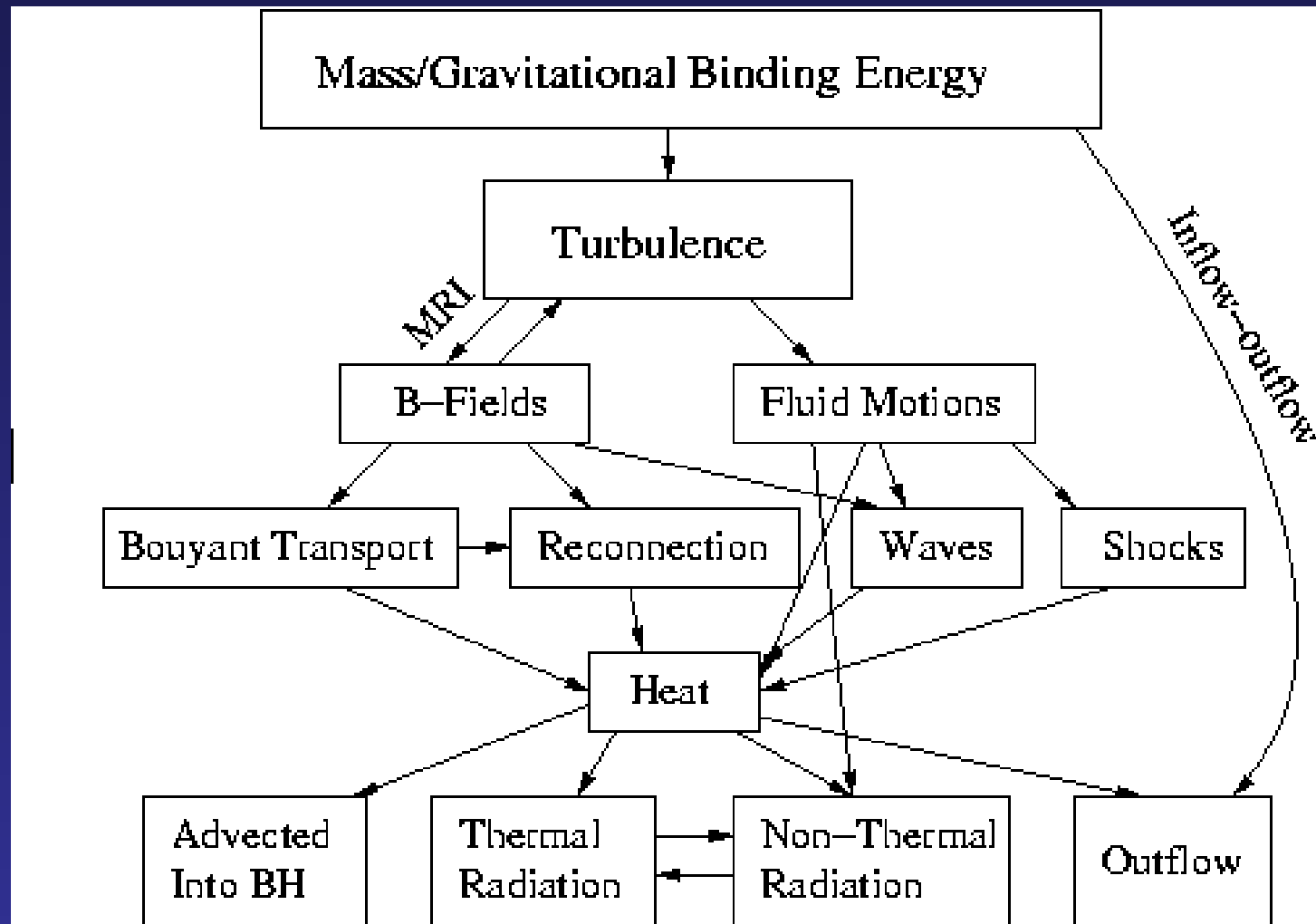


How do Black Hole Accretion Disks Radiate?

Shane Davis
UCSB



Standard Accretion Picture



Issues Addressed in This Talk

- How does the release of gravitational binding energy lead to \sim thermal radiation?
- Are thin disk (α -disk) models sufficient to explain observations? (yes and no)
- How do lifting the simple assumptions of the α -disk model alter the spectra?

I: Summary of Models and Observations

The thin disk model

- $H/R \ll 1$
- Constant accretion rate; time-steady
- Gravitational binding energy released locally

as radiation:

$$F \simeq -H \tau_{R\phi} R \frac{\partial \Omega}{\partial R}$$
$$\dot{M} \frac{\partial R^2 \Omega}{\partial R} = \frac{\partial}{\partial R} (4\pi R^2 H \tau_{R\phi})$$

The α -disk model

- $\tau_{R\phi} = \alpha P$ determines surface density

Σ

- Vertical dissipation distribution:

$$\frac{dF}{dz} = \frac{2\sigma T_{\text{eff}}^4}{\Sigma} \rho$$

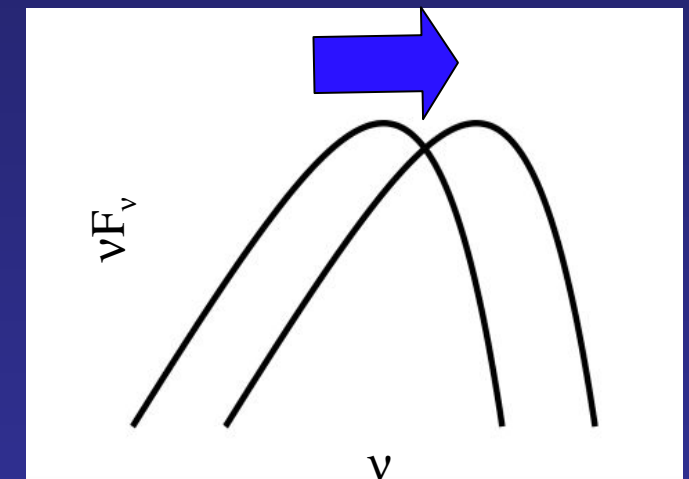
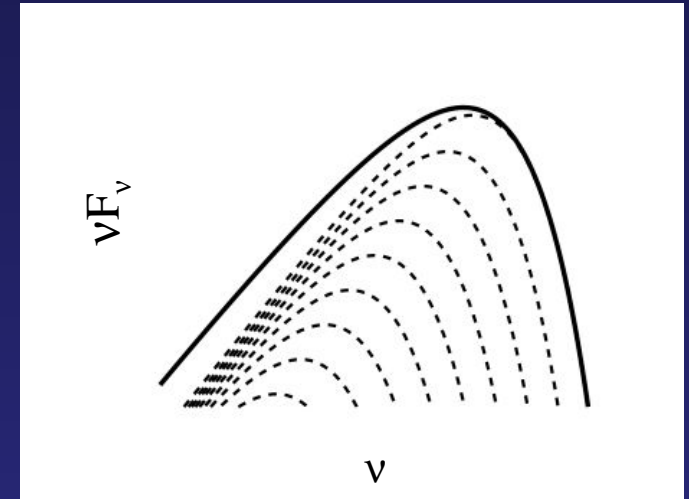
The Multicolor Disk Model (MCD)

- Consider simple temperature distribution:

$$T_{\text{eff}} \propto R^{-3/4}$$

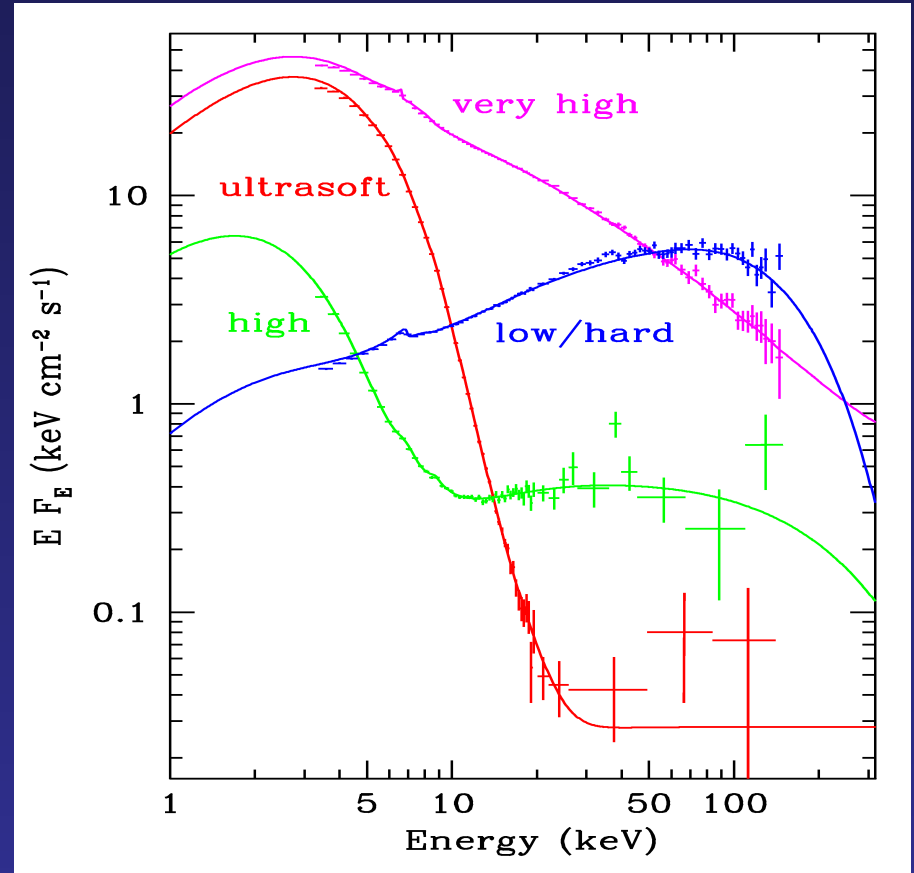
- Spectrum assumed to be color-corrected blackbody:

$$I_{\nu} = f_{\text{col}}^{-4} B_{\nu}(f_{\text{col}} T_{\text{eff}})$$

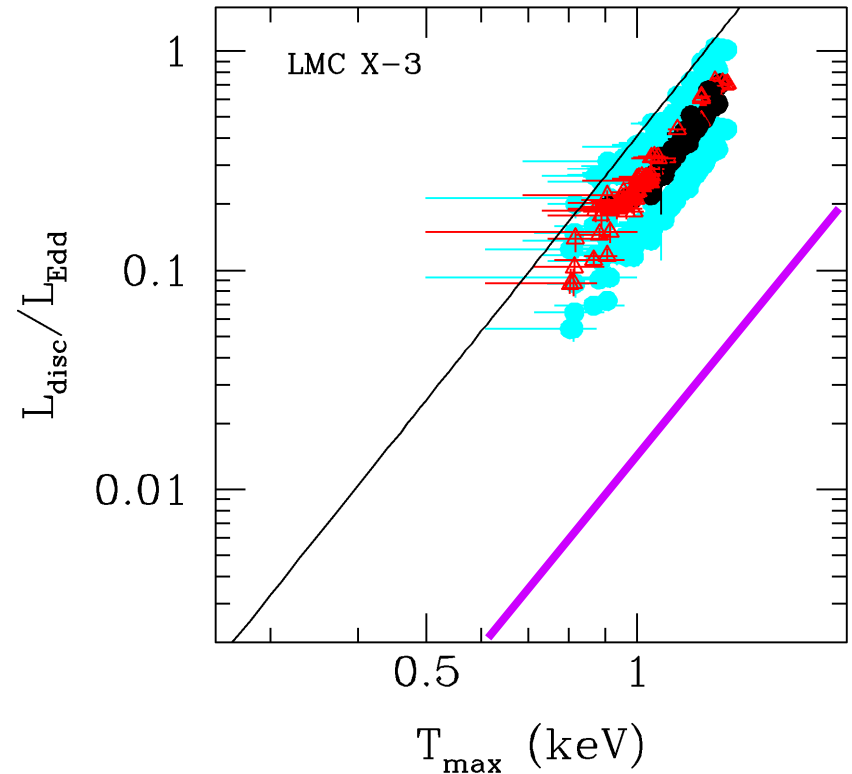
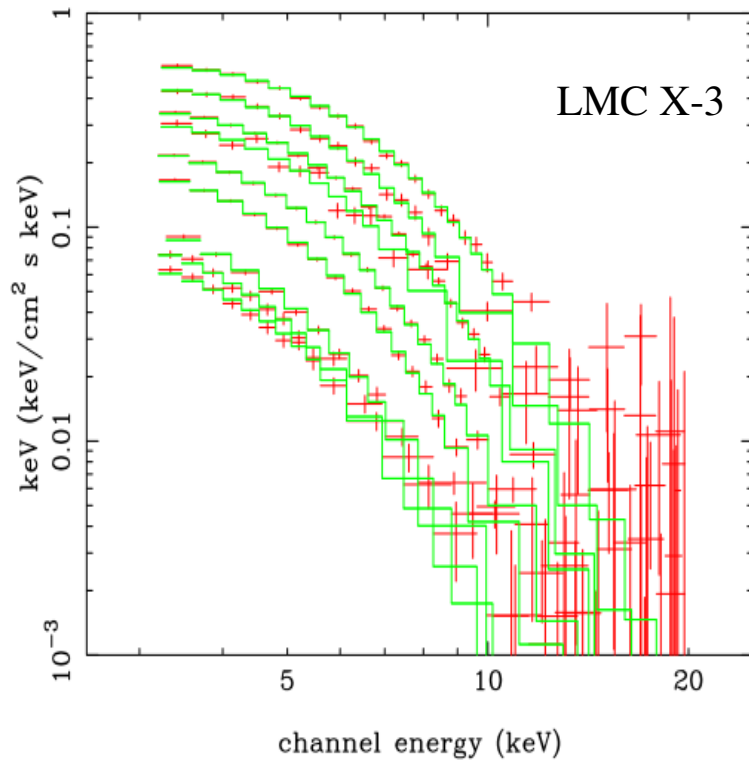


Spectral States of BHBs

- Spectral states specified by relative contributions of thermal (disk) and non-thermal emission (corona)
- High/Soft state is dominated by thermal disk-like component



Disk Dominated Spectra

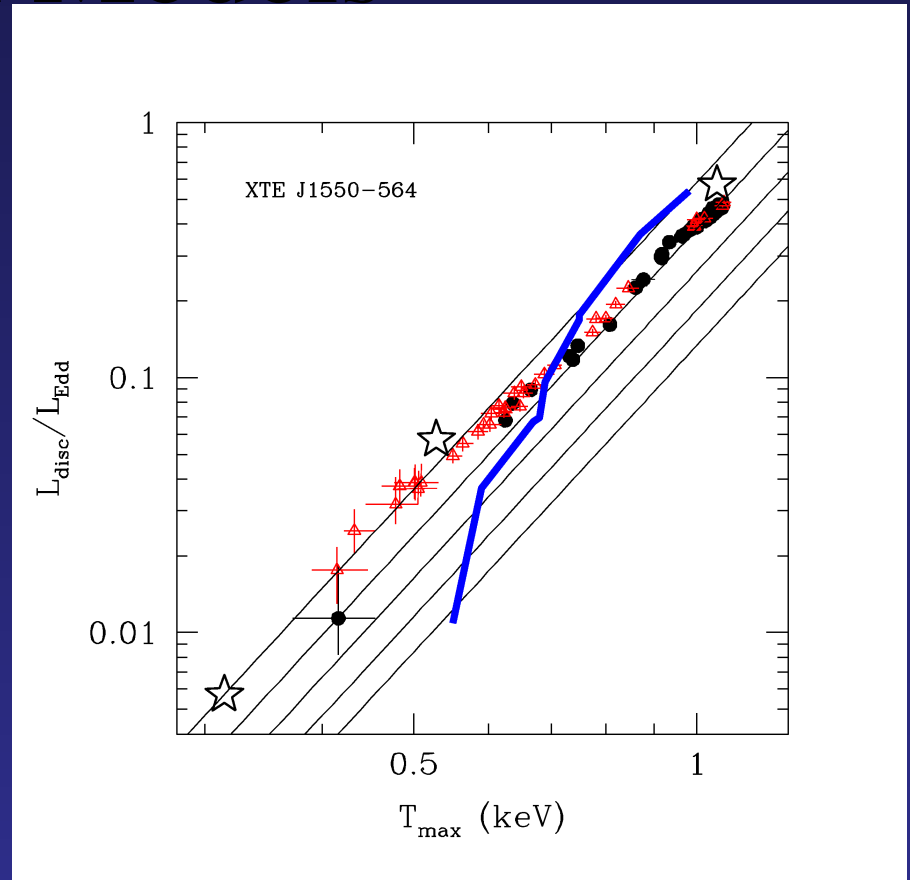


Gierlinski & Done 2004

- L prop. to T^4 suggests f_{col} and emitting area are constant

Comparison with Previous Spectral Models

- Previous α -disk models provide mixed results
- Shimura & Takahara (1995); full atmosphere calculation, free-free emission and Compton scattering: stars
- Merloni, Fabian & Ross (2000); relativistic effects, constant density, f-f emission and Compton scattering: blue curve
- Small scatter in L-T relation as L varies by over an order of magnitude



Gierlinski & Done 2004

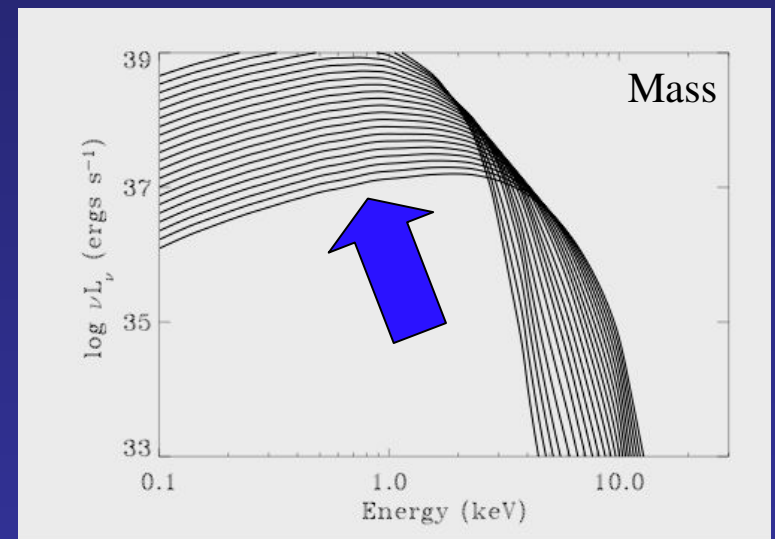
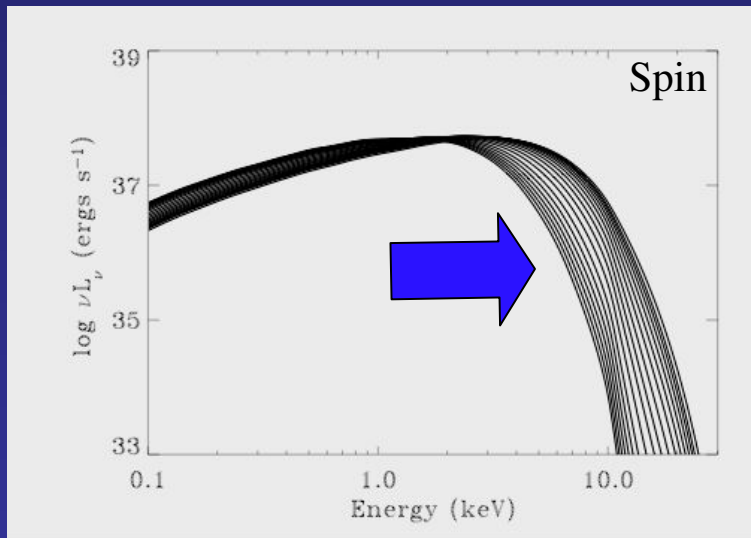
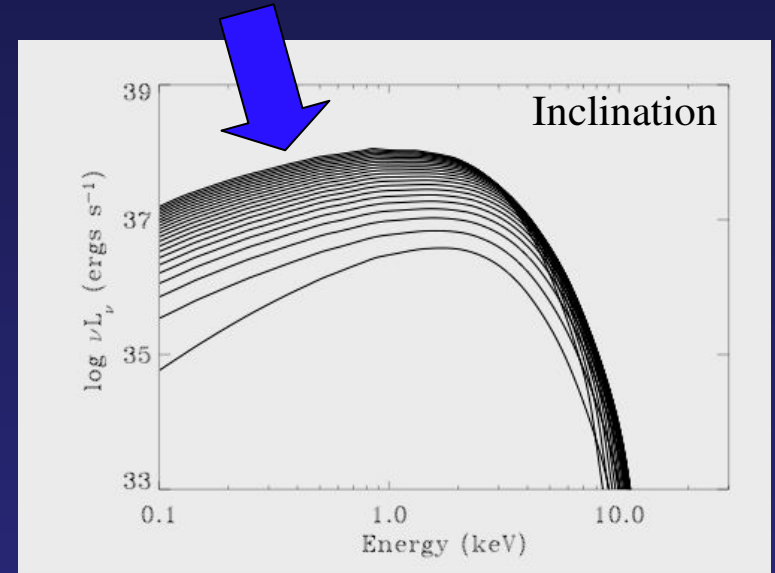
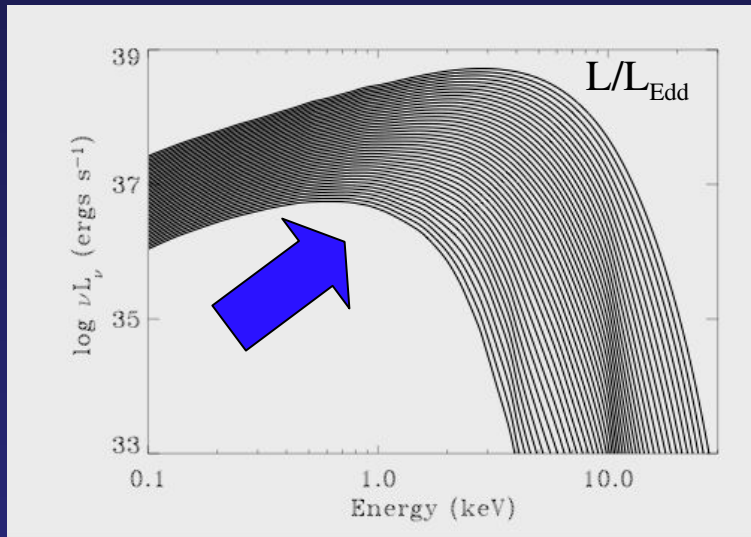
How do we Improve on these Models

- Include metals; non-LTE populations
- Improved radiative transfer and better treatment of Compton Scattering
- Examine the effects of:
dissipation, magnetic stresses, and
inhomogeneities

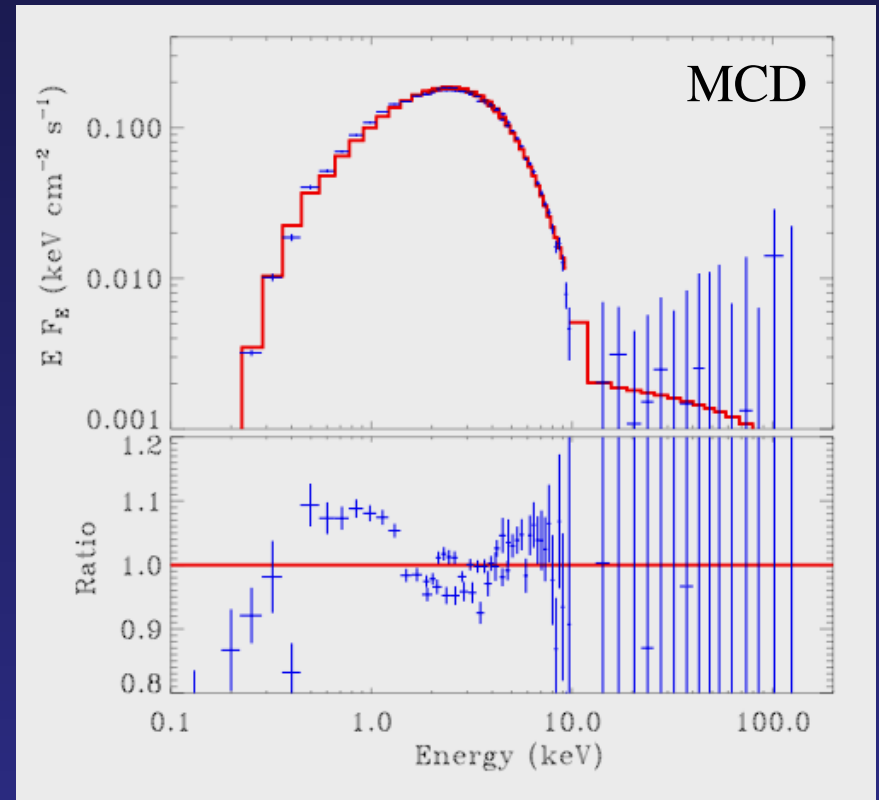
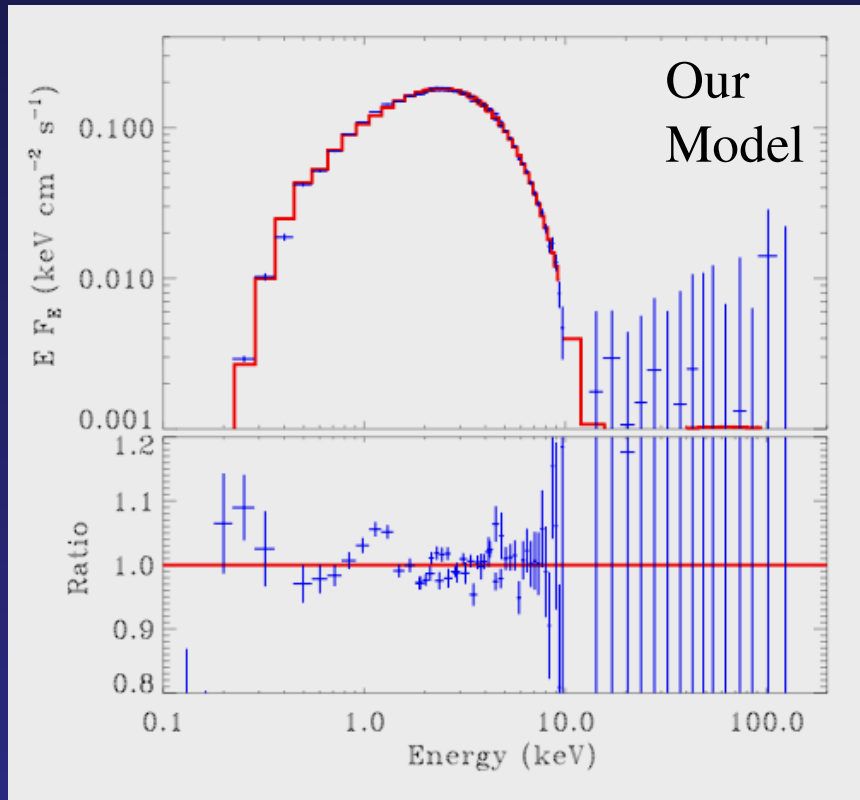
Our Models, Briefly

- Use conservation laws with Kerr metric to calculate one-zone model
- Full disk model is determined by ~ 4 parameters: $M, a, L/L_{\text{Edd}}, \alpha$
- Calculate self-consistent vertical structure and radiative transfer in a series of concentric annuli
- Each annulus is determined by 3 parameters (+ assumptions): $T_{\text{eff}}, Q, (g=Q z), \Sigma$
- Calculate photon geodesics in a relativistic spacetime for an observer at inclination i (ray tracing)

Model Parameters



'Broadband' Fits to LMC X-3



- MCD model is too narrow -- need relativistic broadening $\Delta\chi^2 \sim 100$
- Best fit find spinning black hole with $a_*=0.45$, consistent with *RXTE* fits

Spectral Formation

- Depth of formation τ_* : optical depth where $(\tau_{\text{es}} \tau_{\text{abs}})^{1/2} \sim 1$

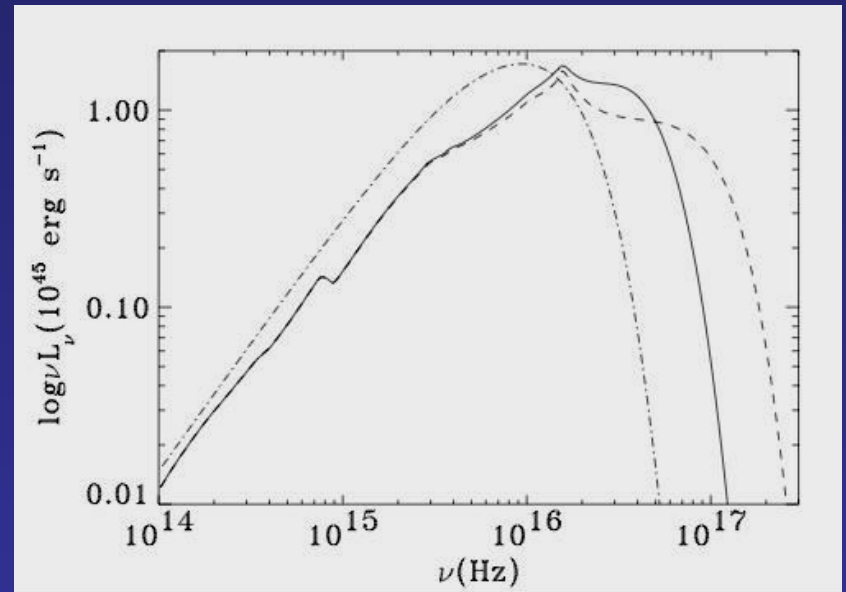
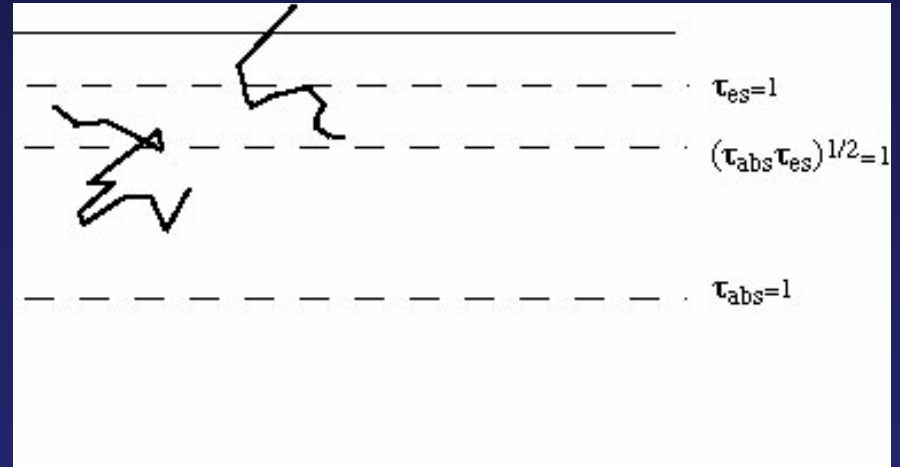
$\tau > \tau_*$: absorbed

$\tau < \tau_*$: escape

- Therefore Thomson scattering produces modified blackbody:

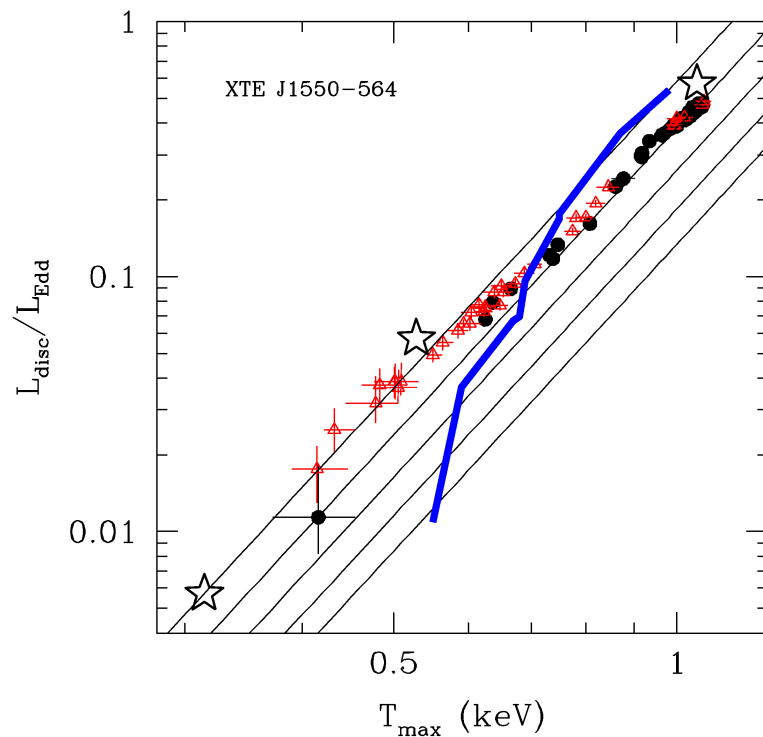
$$I_\nu \approx B_\nu \sqrt{\frac{\kappa_{\text{abs}}}{\kappa_{\text{es}}}}$$

- Compton scattering gives softer Wien spectrum
- From these considerations Shimura & Takahara (1995) find $f_{\text{col}} = 1.7 \pm 0.2$

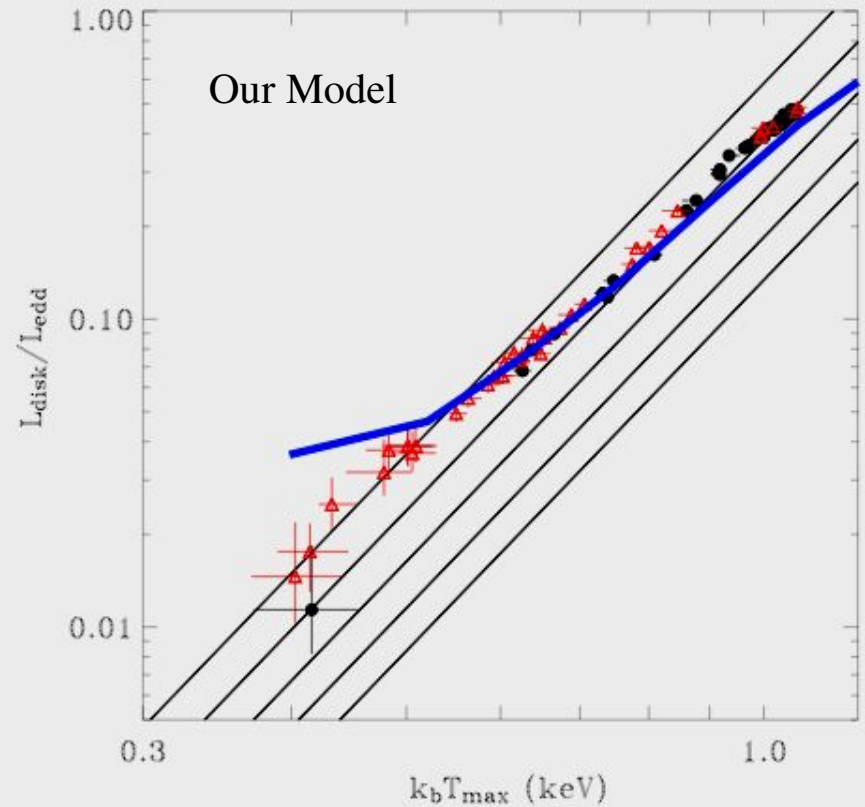


Luminosity vs. Temperature

Previous Models

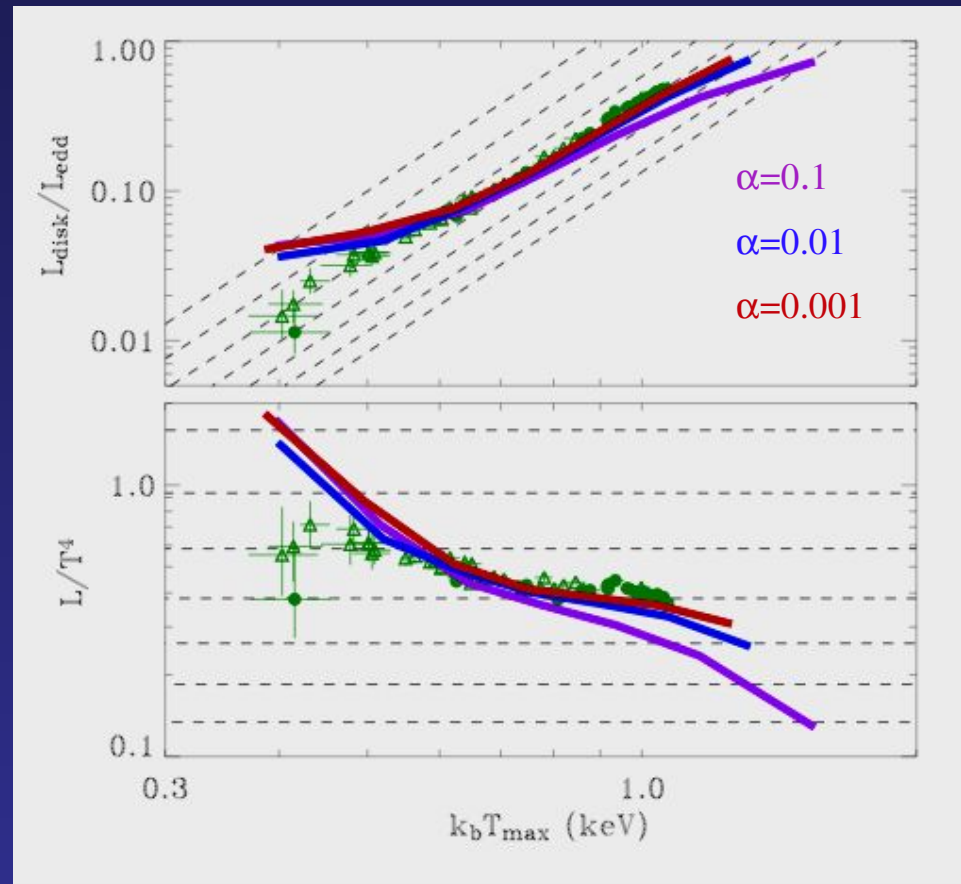


Our Model

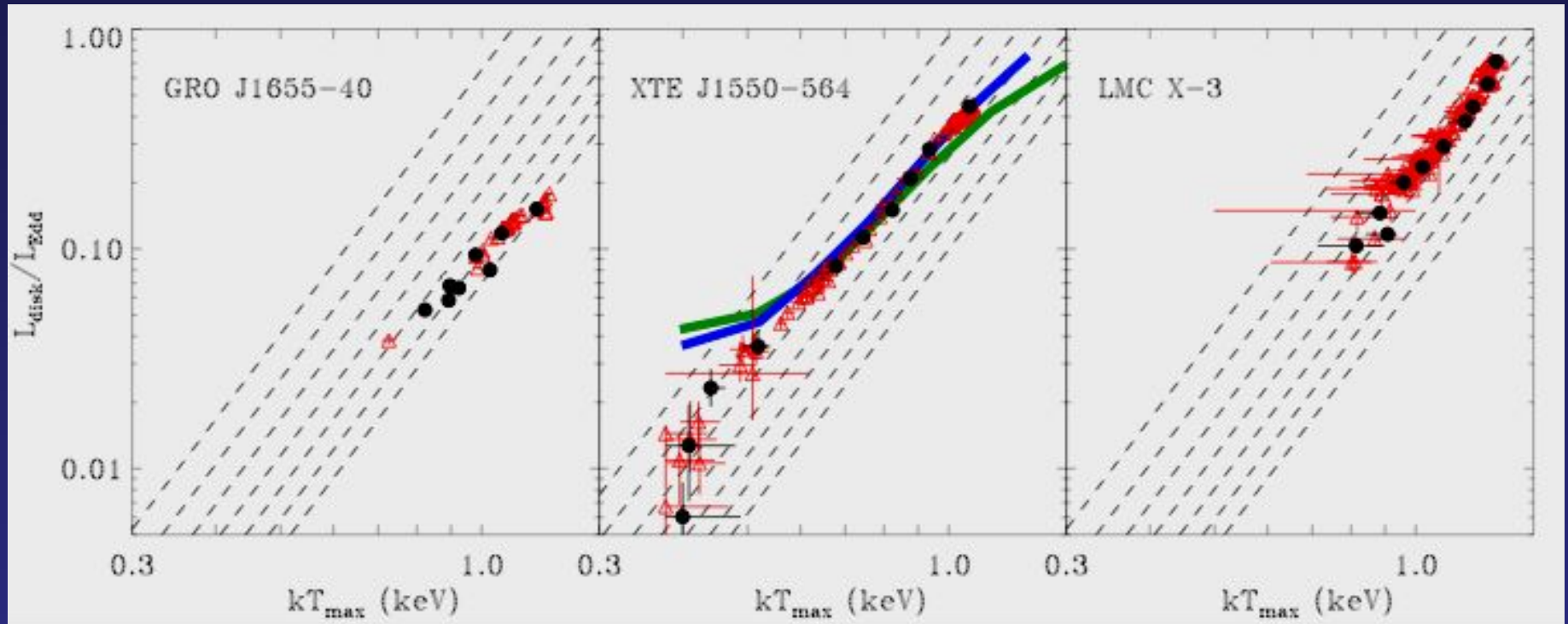


Luminosity vs. Temperature

- Slight hardening is consistent with some observations
- Allows one to constrain surface density and accretion stress
- Models become effectively optically thin at high accretion rates



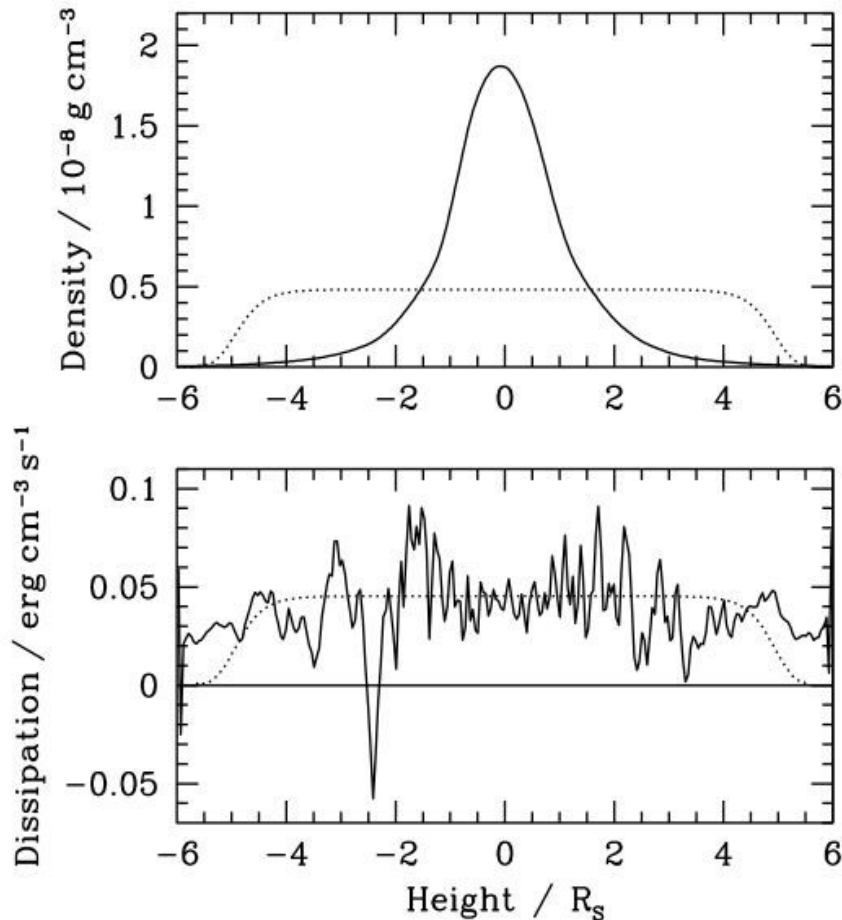
Summary of Fit Results



- Lack of hardening suggests hottest annuli remain effectively thick (especially LMC X-3)
- Models consistent with moderate spins, small torques for disk inclination = binary inclination ($a_* \sim 0.7, 0.1, 0.4$)

II. Adding 'Real' Physics: Dissipation, Magnetic Stresses & Inhomogeneities

Local MRI Simulations (with radiation)



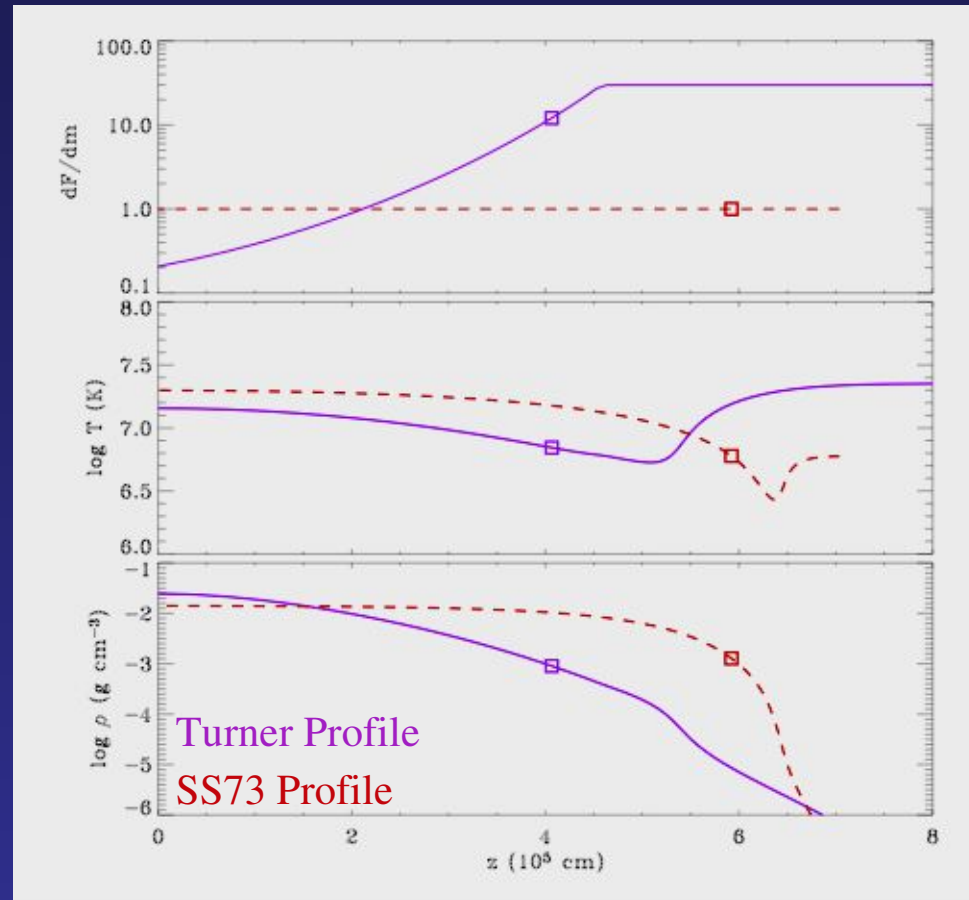
- Shakura & Sunyaev (1973):

$$-\frac{dF}{dm} = \frac{1}{\rho} \frac{dF}{dz} = \frac{2\sigma T_{\text{eff}}^4}{\Sigma}$$

- Turner (2004): Mass more centrally concentrated towards midplane in simulations.
- Magnetic fields produced near midplane are buoyant and dissipate near surface

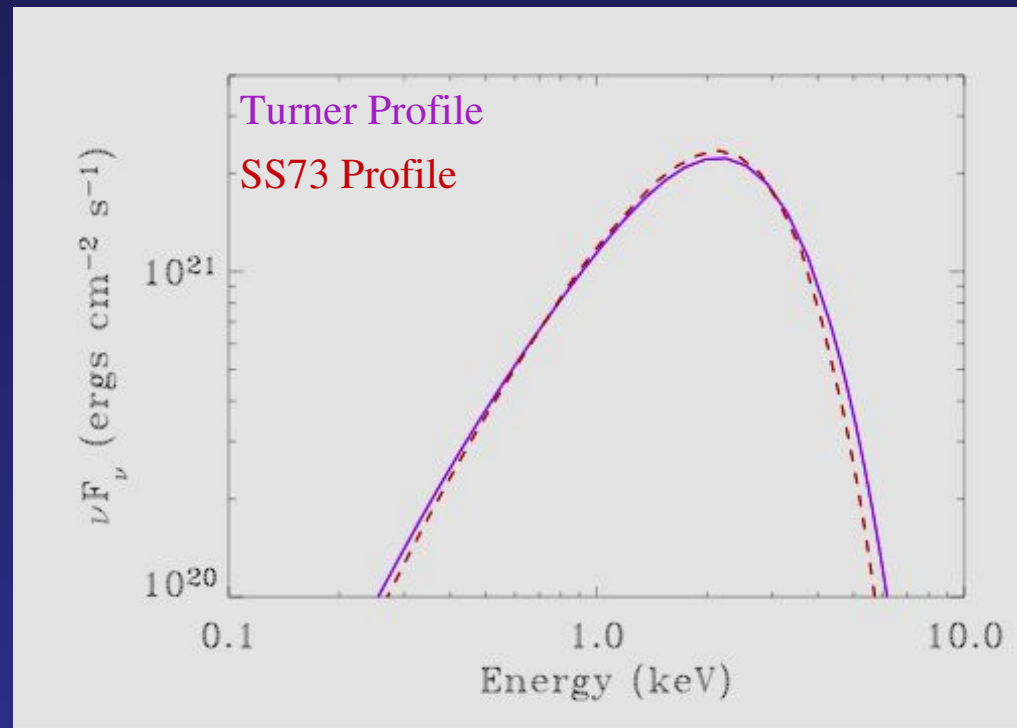
Dissipation Profile

- Modified dissipation profile changes structure significantly
- Where $-dF/dm$ is small (large), density is larger (smaller)
- Note that T and ρ near τ^* are similar

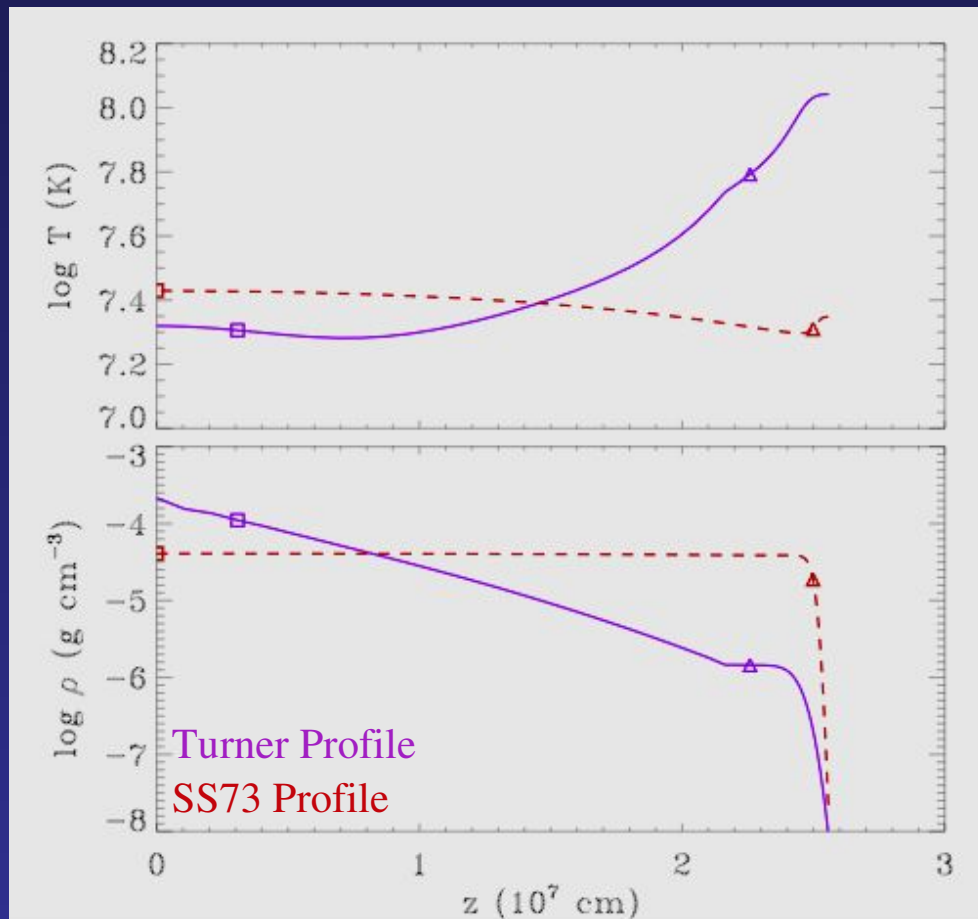


Dissipation Profile

- ... but modified dissipation profile has limited affect on the spectrum.
- This particular annulus is very effectively optically thick and τ^* is close to surface

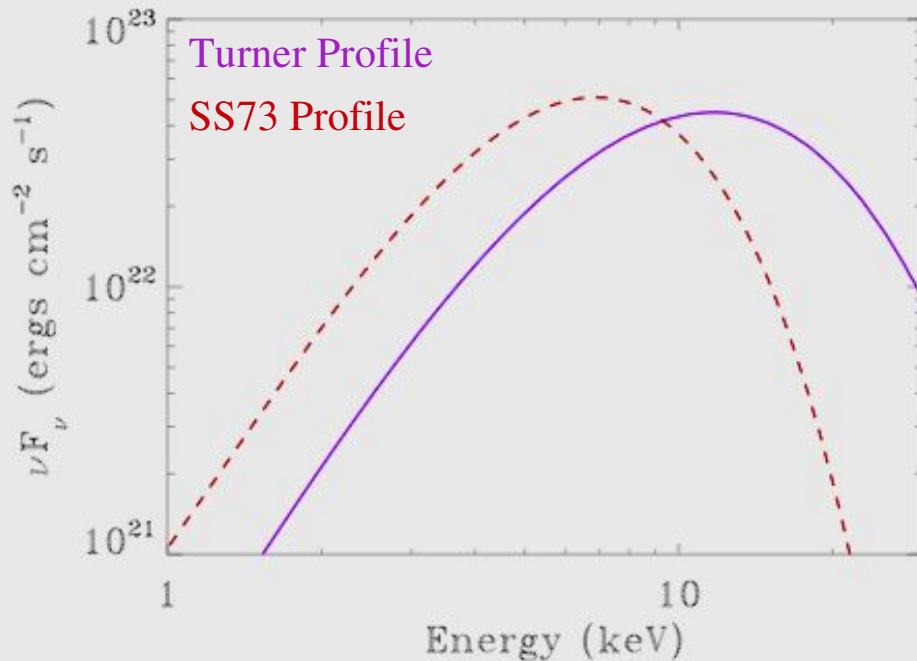


Dissipation Profile



- Consider same dissipation profile, but in an annulus that is not effectively optically thick.
- Spectrum with modified dissipation profile is effectively optically thick, but has a much greater surface temperature

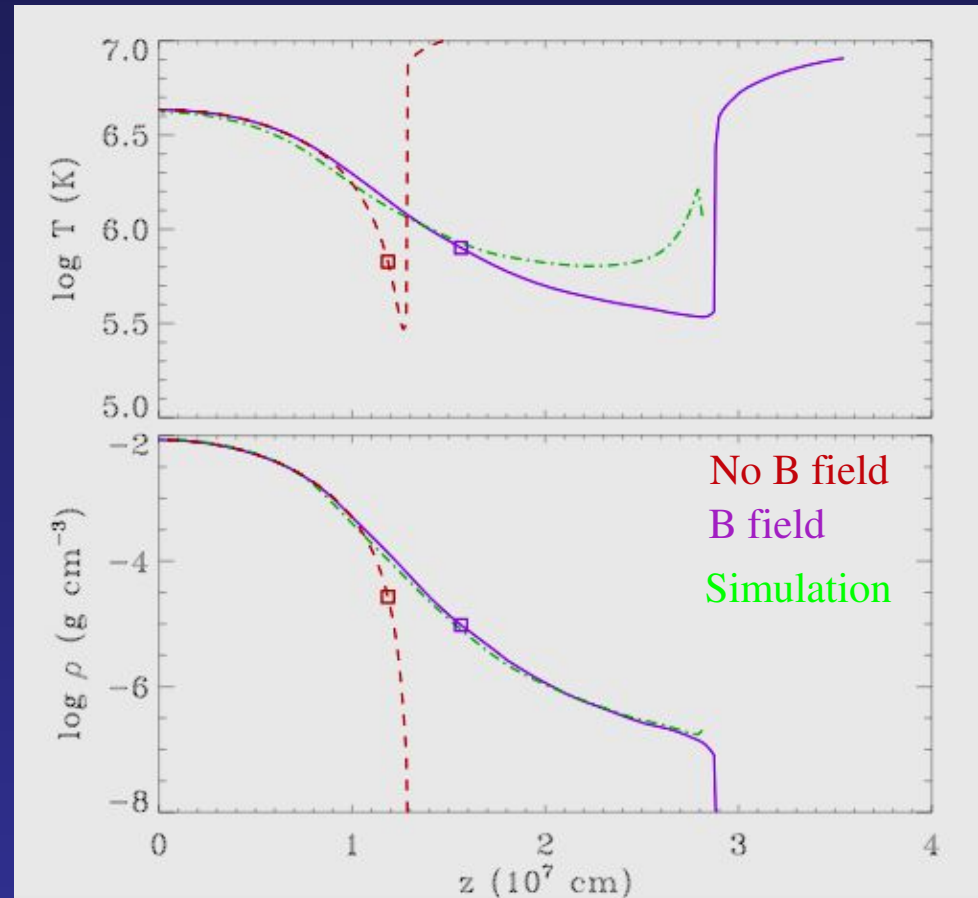
Dissipation Profile



- For this annulus, the spectra from the modified dissipation profile is considerably harder
- This will exacerbate the discrepancy between observations and models unless disks stay optically thick

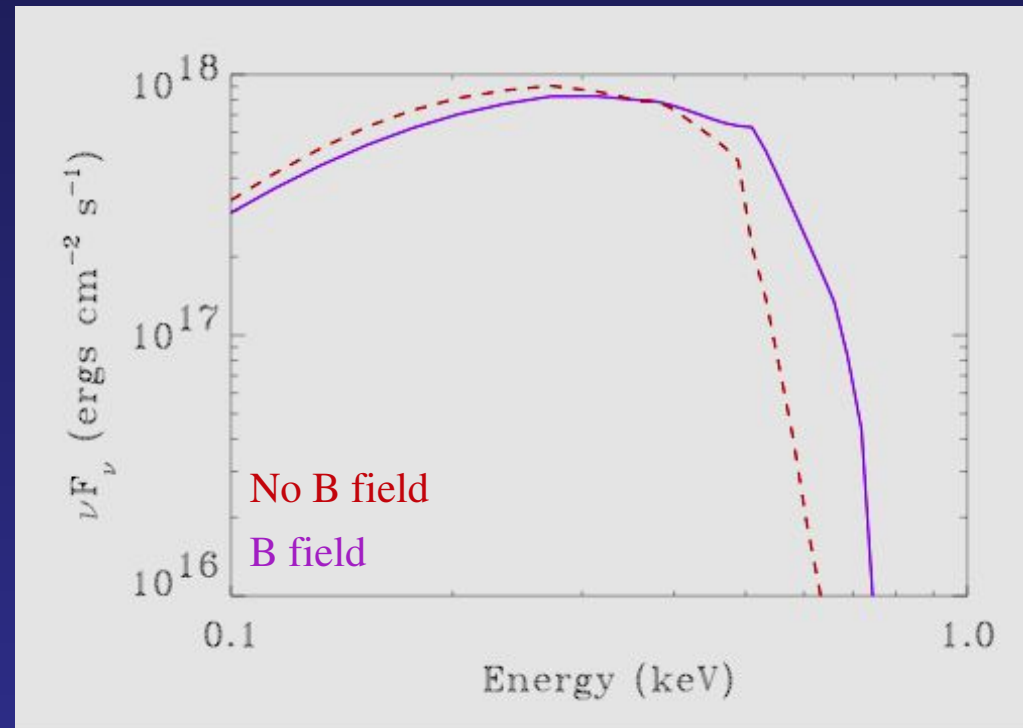
Magnetic Stresses

- $B^2/8\pi$ & dF/dm taken from simulations of Hirose, Krolik, & Stone
- Gas pressure dominated -- dF/dm has little effect on the structure
- Extra magnetic pressure support lengthens scale height h_ρ
- $\rho_* \sim \tau_*/(\kappa_{es} h_\rho)$

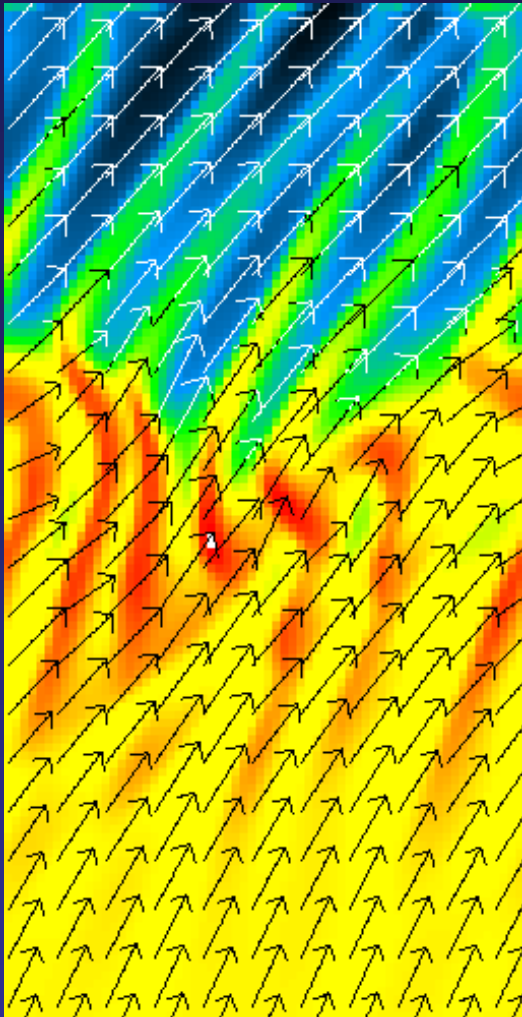


Magnetic Stresses

- Lower ρ_* and higher T_* combine to give somewhat harder spectrum
- In this case lower ρ_* alters statistical equilibrium -- lower recombination rate relative to photoionization rate give more highly ionized matter with lower bound-free opacity



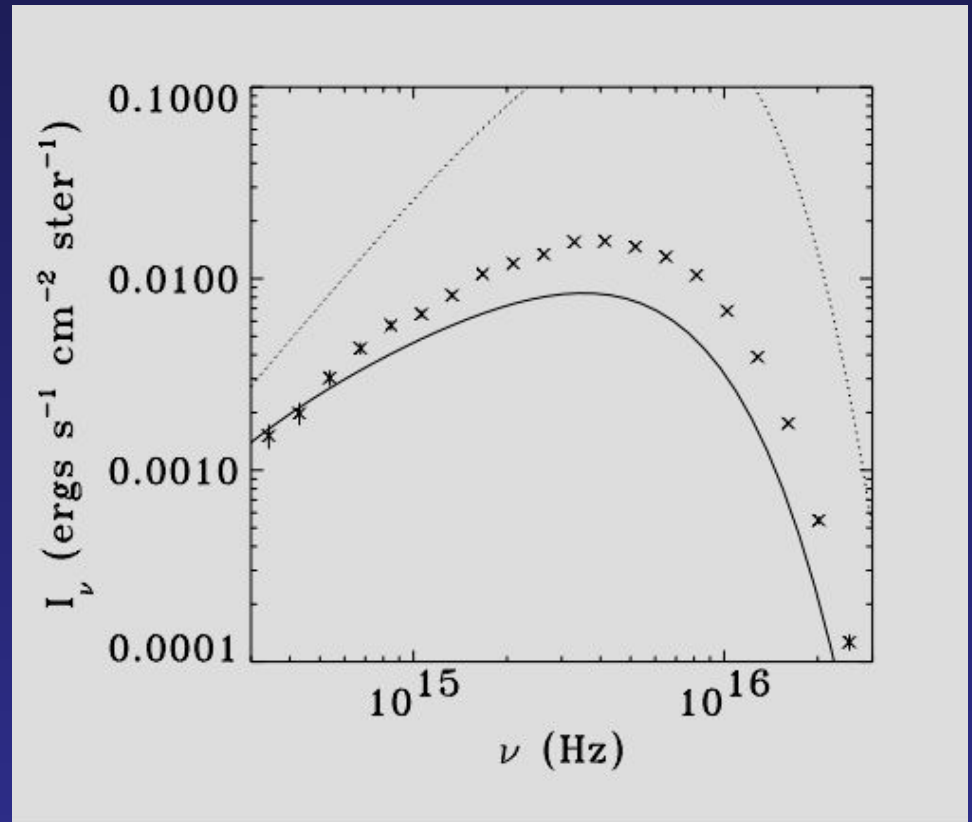
Inhomogeneities



- Radiation pressure dominated accretion disks expected to be homogeneous due to photon bubble instability and/or compressible MRI
- May make disk thinner and (therefore) denser (Begelman 2001, Turner et al. 2005)
- May also affect the radiative transfer...

Inhomogeneities

- 2-D Monte Carlo calculations: photon bubbles help thermalize the spectrum, making it softer
- Photon emission and absorption dominated by denser regions.



What Have We Learned?

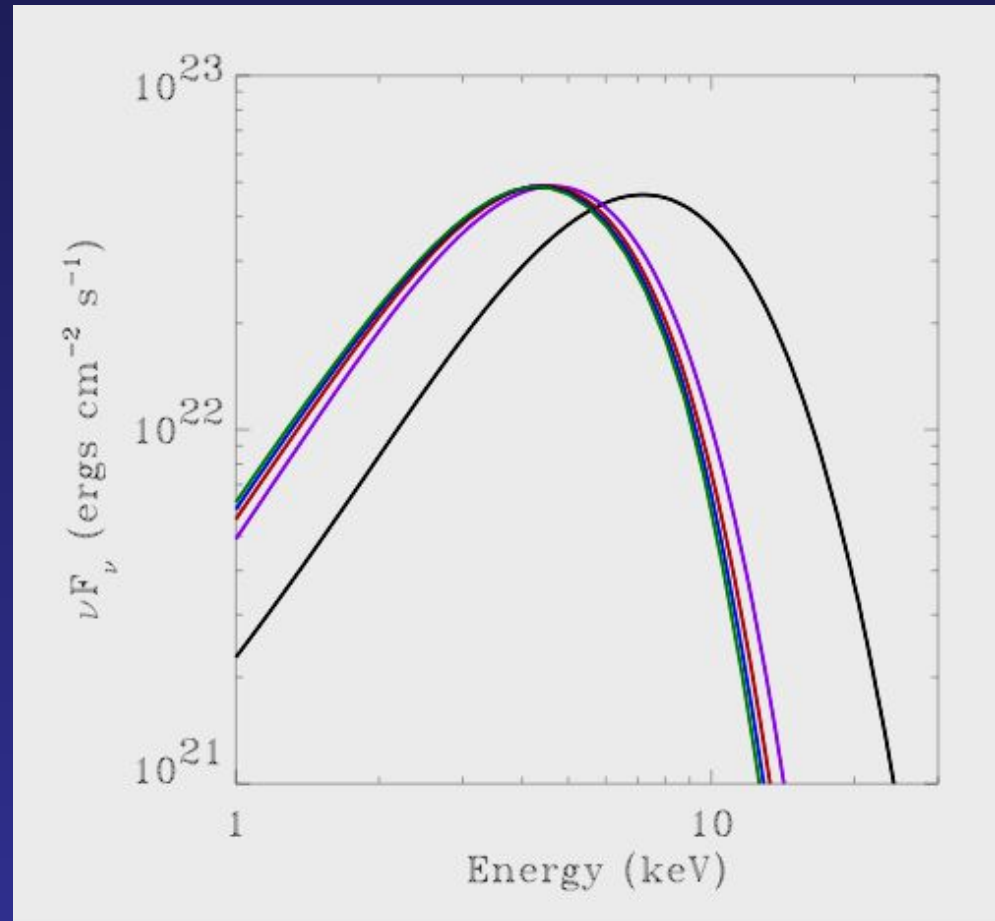
- Our detailed accretion disk models can reproduce the spectra of disk dominated BHBs (better than the MCD in LMC X-3)
- Models reproduce the observed evolution with luminosity; spectral modeling may constrain the nature of angular momentum transport and black hole spin (torques?)
- Dissipation, magnetic stresses, and inhomogeneities may affect the spectra -- more simulations or analytical progress needed.

Non-aligned Jet

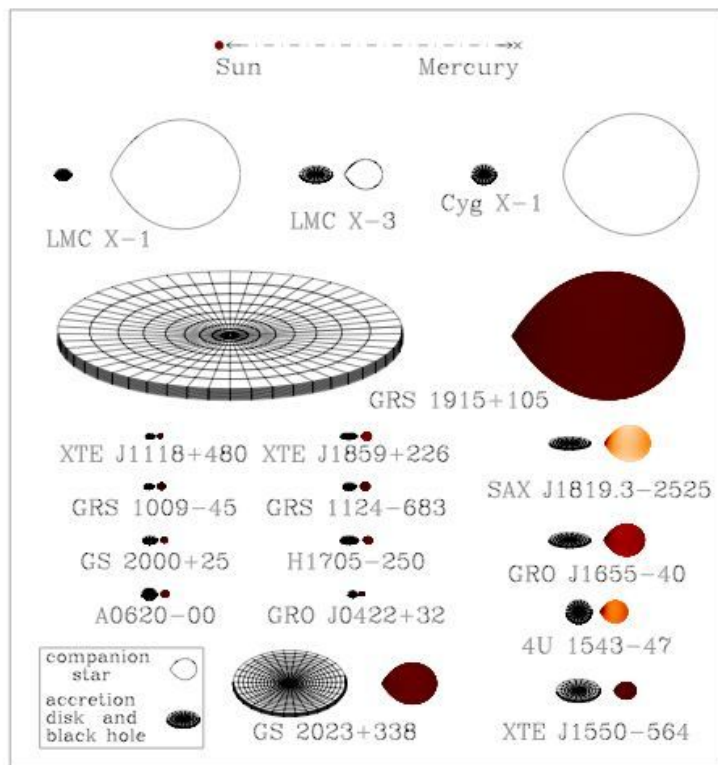
- XTE J1550-564 is a microquasar
- Hannikainen et al. (2001) observe superluminal ejections with $v > 2c$
- Ballistic model:
$$\beta_{\text{ap}} = \frac{\beta \sin \theta}{1 - \beta \cos \theta} \Rightarrow \theta < 53^\circ$$
- Orosz et al. (2002) found $i=72^\circ$
- Non-aligned jets not uncommon -- usually assumed that BH spin differs from binary orbit and inner disk aligns with BH -- Bardeen-Petterson effect
- Best fit inclination, spin: $i=43^\circ$, $a=0.44$

Spectral Dependence on Surface Density

- Spectra largely independent of S for large surface density ($\Sigma > \sim 10^3 \text{ g/cm}^2$)
- As disk becomes marginally effectively thin, spectra become sensitive to Σ and harden rapidly with decreasing Σ



Binaries Provide Independent Constraints on Models



- Orosz and collaborators derive reasonably precise estimates from modeling the light curve of secondary
- e.g. XTE J1550-564:

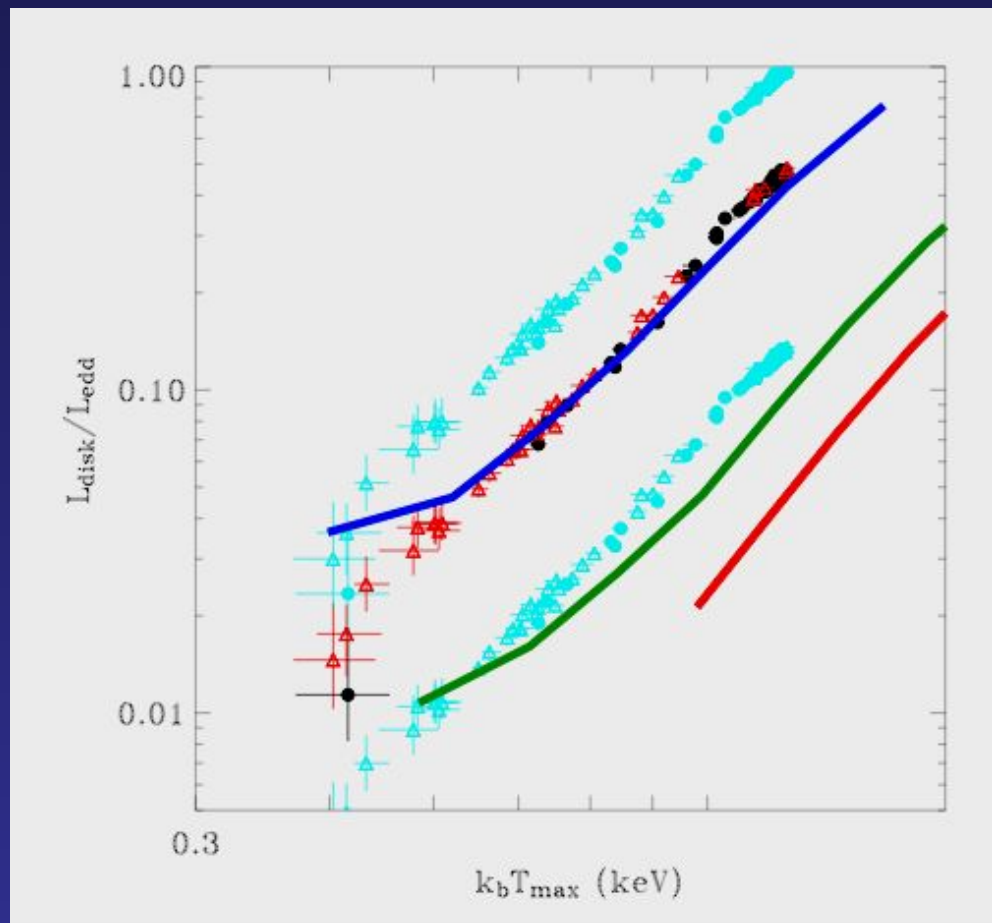
$$M = 10 (9.7 - 11.6) M_{\odot}$$

$$i = 72^{\circ} (70.8^{\circ} - 75.4^{\circ})$$

$$D = 5.3 (2.8 - 7.6) \text{ kpc}$$

Luminosity vs. Temperature

- Measured binary properties limit parameter space of fits
- Simultaneous fits to multiple observations of same source constrain spin/torque
- Spectra are too soft to allow for extreme spin/large torques

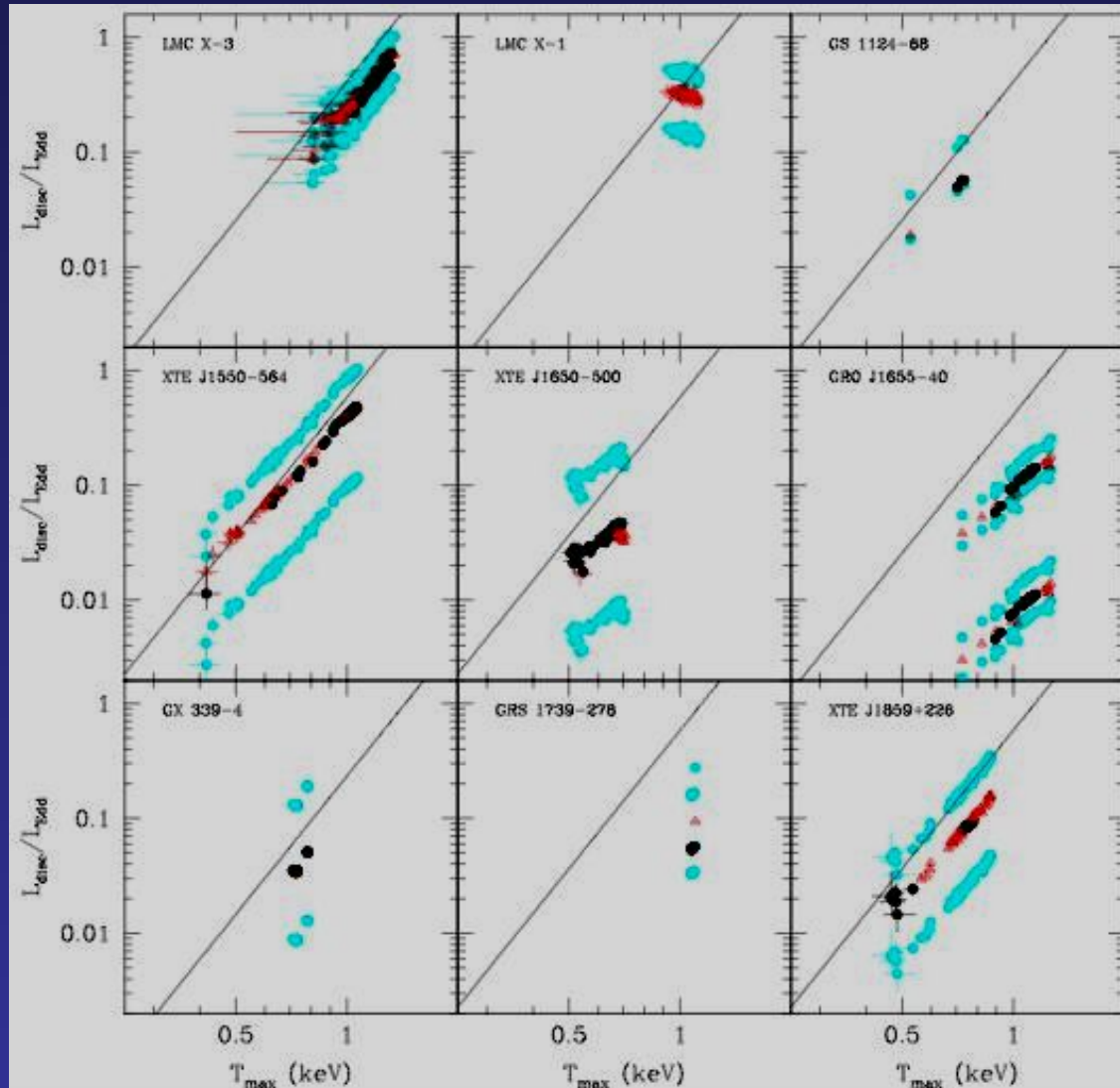


Stellar Atmospheres: Disk Annuli vs. Stellar Envelopes

- The spectra of stars are determined by T_{eff} , g , and the composition
- Annuli are determined by, T_{eff} , Q (where $g=Q z$), S , the composition, and the vertical dissipation profile: $F(m)$
- T_{eff} , Q , and S can be derived from radial disk structure equations
- Standard assumption is:

$$-\frac{dF}{dm} = \frac{2\sigma T_{\text{eff}}^4}{\Sigma}$$

Luminosity vs. Temperature



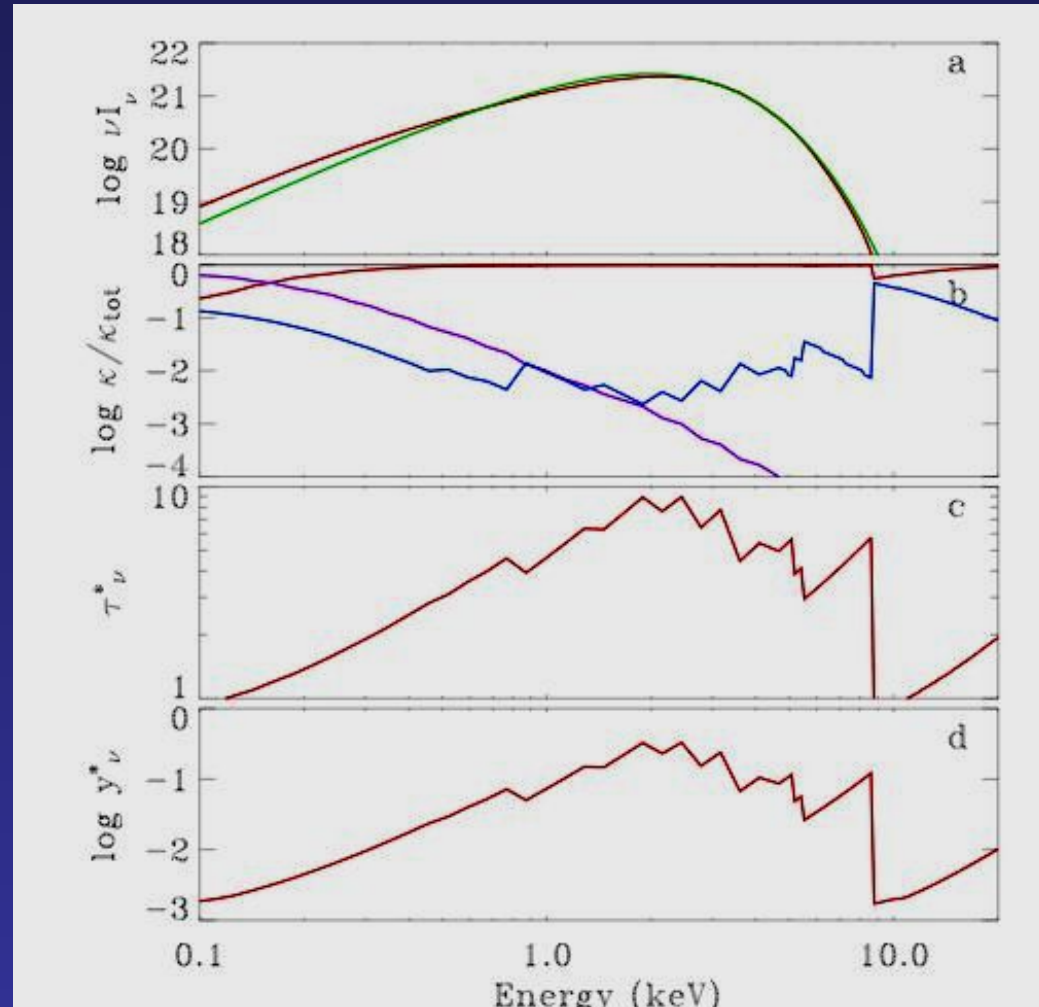
Effect of bound-free opacity

- Bound-free opacity decreases depth of formation: τ^*

$$y_\nu^* = \frac{4kT^*}{m_e c^2} \tau_\nu^{*2} < 1$$

- Absorption opacity approximately grey
- Spectrum still approximated by diluted blackbody:

$$I_\nu \approx B_\nu \sqrt{\frac{\kappa_{\text{abs}}}{\kappa_{\text{es}}}} \approx B_\nu \times \text{const.}$$



The Multicolor Disk Model (MCD)

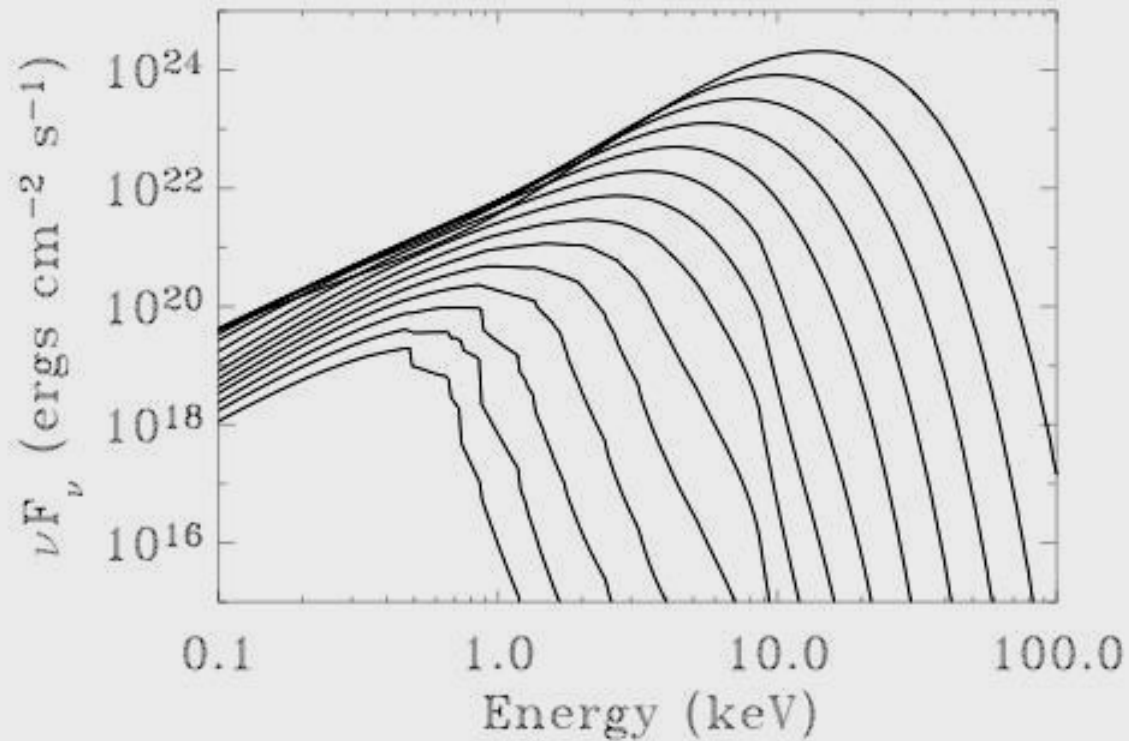
- Consider simplest temperature distribution:

$$\pi R^2 \sigma_{\text{sb}} T_{\text{eff}}^4 \sim \frac{GM\dot{M}}{R} \Rightarrow T_{\text{eff}}^4 \propto R^{-3/4}$$

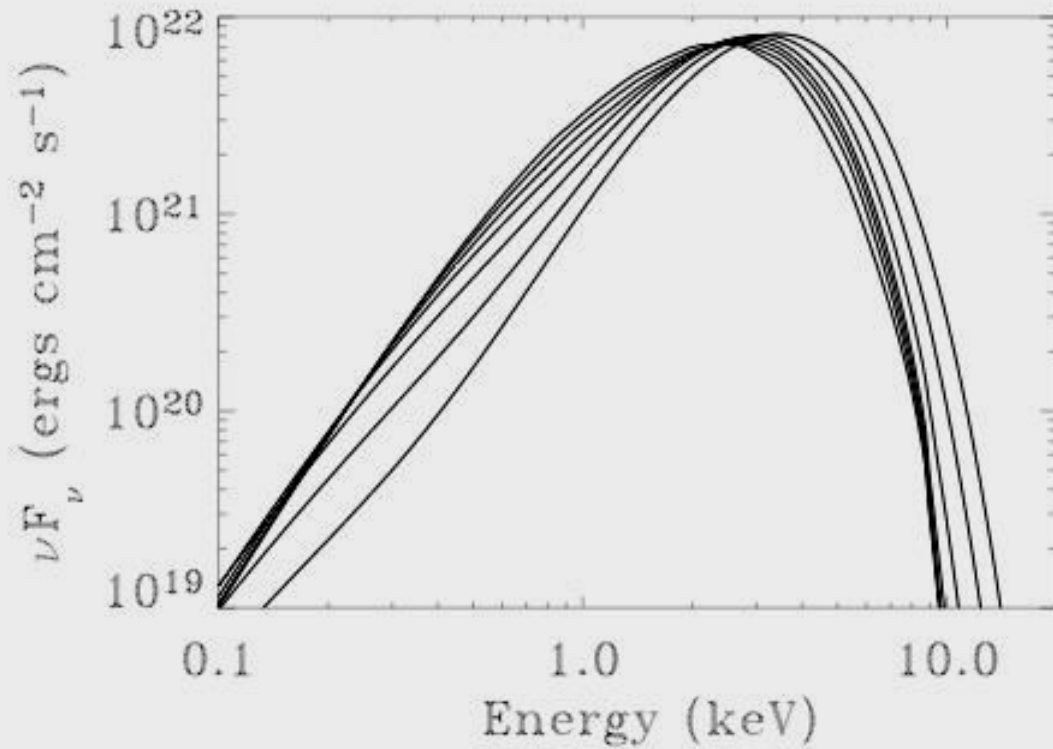
- Assume blackbody and integrate over R replacing R with T ($T_{\text{max}} \sim f_{\text{col}} T_{\text{eff}}$):

$$I_{\nu} = \frac{8\pi R_{\text{in}}^2 \cos i}{3D^2 f_{\text{col}}^4} \int_0^{T_{\text{max}}} B_{\nu}(T) \left(\frac{T}{T_{\text{max}}} \right)^{-11/3} \frac{dT}{T_{\text{max}}}$$

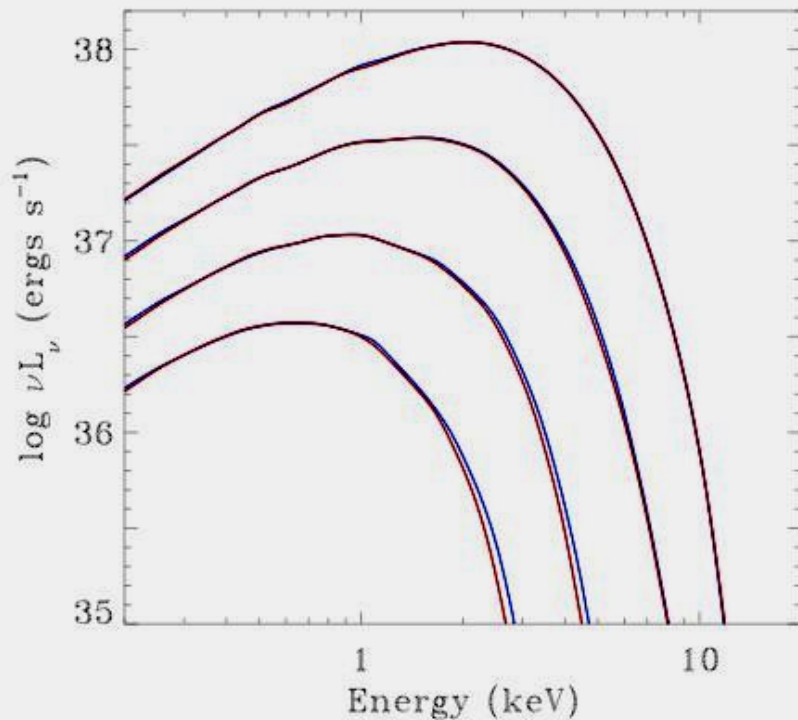
Effective Temperature: T_{eff}



Gravity Parameter: Q



Comparison Between Interpolation and Exact Models



- Interpolation best at high L/L_{edd}
- Exact: Blue curve
Interpolation: Red Curve