

# Working Title

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## 1 Toy model: $0d$ GLSM

### 1.1 Setup

We start with a  $0d$  GLSM toy model, i.e. we consider the source manifold  $\Sigma = \{pt\}$  to be an abstract point and the target manifold to be  $X = \mathbb{CP}^1$ . The space of fields is then simply given by points on the target manifold  $X$ , namely

$$\mathcal{F} = \mathbb{CP}^1 \tag{1}$$

As the action of the model we consider

$$S(z_0, z_1) = \beta \left( |z_0|^2 - |z_1|^2 \right) \tag{2} \quad \{\text{eq:toy\_S}\}$$

where  $[z_0 : z_1] \in \mathbb{CP}^1$ . As written, the partition function explicitly shows the  $U(1)$  gauge freedom (which here is simply a global  $U(1)$  symmetry, acting as

$$e^{i\theta} : (z_0, z_1) \mapsto (e^{i\theta} z_0, e^{i\theta} z_1) \tag{3}$$

and defines indeed a function on  $\mathbb{CP}^1$  if we implicitly assume the constraint

$$|z_0|^2 + |z_1|^2 = 1 \tag{4} \quad \{\text{eq:toy\_constr}\}$$

The path integral of the model is thus defined by

$$Z = \int dz_0 dz_1 \delta(|z_0|^2 + |z_1|^2 - 1) e^{-S(z_0, z_1)} \tag{5} \quad \{\text{eq:toy\_Z}\}$$

In order to evaluate (5), we want to embed  $\mathbb{CP}^1$  into a higher dimensional complex space such that

1. the new action  $\tilde{S}$  is holomorphic in the new variables (fields)
2. when we restrict to  $\mathbb{CP}^1$ ,  $\tilde{S}$  reduces to  $S$

We think of this embedding as an analytic continuation of the space of fields in an appropriate sense.

Going to real fields,  $z_k = x_k + iy_k$  for  $k = 0, 1$ , the action (2) becomes

$$S = \beta (x_0^2 + y_0^2 - x_1^2 - y_1^2) \quad (6) \quad \{\text{eq:toy\_S\_real}\}$$

while the constraint becomes

$$x_0^2 + y_0^2 + x_1^2 + y_1^2 = 1 \quad (7) \quad \{\text{eq:toy\_constr\_real}\}$$

## 1.2 Manifold Deformation

Suppose that we now complexify in the sense that we consider  $x_k, y_k \in \mathbb{C}$ . For simplicity, let us rename the variables according to

$$x_0 = u_0 \quad , \quad x_1 = u_1 \quad , \quad y_0 = u_2 \quad , \quad y_1 = u_3 \quad (8) \quad \{\text{eq:toy\_us}\}$$

where now  $u_i \in \mathbb{C}$ . This slightly unusual renaming is done for later convenience.

The constraint (7) now becomes

$$\sum_{i=0}^3 u_i^2 = 1 \quad (9)$$

which can be seen as a complex hypersurface inside  $\mathbb{CP}^4$ , as follows: Let us introduce the complex variable  $t \in \mathbb{C}$  and consider the equation

$$\sum_{i=0}^3 u_i^2 = t^2 \quad (10)$$

This equation describes the zero set of a homogeneous quadratic polynomial

$$P(t, u_i) = t^2 - \sum_{i=0}^3 u_i^2 \quad (11)$$

Notice that for any  $\lambda \in \mathbb{C}^* = \mathbb{C} - \{0\}$ ,

$$P(\lambda t, \lambda u_i) = \lambda^2 P(t, u_i) \quad (12)$$

Therefore, the solution space of

$$P(u_i, t) = 0 \quad (13)$$

is invariant under the action of  $\mathbb{C}^*$  (by multiplication) and hence descends to an equation on  $\mathbb{CP}^4$  whose coordinates are  $[t : u_0 : u_1 : u_2 : u_3]$

The important observation is that the constraint surface (7) coincides with  $P(t, u_i) = 0$  for  $t = 1$ . But  $t = 1$  simply defines the  $U_{t \neq 0} \subset \mathbb{CP}^4$  whose coordinates are given by  $[1 : \frac{u_0}{t} : \frac{u_1}{t} : \frac{u_2}{t}]$ . Hence, the constraint surface can be embedded into  $\mathbb{CP}^4$ :

$$[u_0 : u_1 : u_2 : u_3] \mapsto [t(u_i) : u_0 : u_1 : u_2 : u_3] \quad , \quad t(u_i) = \sqrt{\sum_{i=0}^3 u_i^2} \quad (14) \quad \{\text{eq:toy\_emb}\}$$

Notice that  $[t(u_i) : u_0 : u_1 : u_2 : u_3]$  simply describes a point in  $P(t, u_i) = 0$  and the original surface is reproduced in the chart  $t(u_i) = 1$ .

This embedding comes with a natural family of holomorphic deformations parametrised by a vector  $\omega \in \mathbb{C}^4 - \{0\}$ . Namely, instead of considering zeros of  $P(t, u_i)$ , one could consider zeros of a general homogeneous quadratic polynomial

$$P_\omega(t, u_i) = t^2 - \sum_{i=0}^3 \omega_i u_i^2 \quad (15)$$

which for  $\omega = (1, 1, 1, 1) \in \mathbb{C}^4$  coincides with  $P$ .

**Remark 1.** Note that the hypersurface  $\mathcal{C}_\omega$  defined by  $P_\omega(t, u_i) = 0$  defines a  $\mathbb{CP}^3 \subset \mathbb{CP}^4$ . Indeed, zeros of  $P_\omega(t, u_i)$  are given by points

$$[t(u_i) : u_0 : u_1 : u_2 : u_3] \quad , \quad t(u_i) = \sqrt{\sum_i u_i^2} \quad (16)$$

and hence parametrised by  $(u_0, \dots, u_3) \in \mathbb{C}^4 - \{0\}$ . However, we may freely scale the  $u_i$  by the same  $\lambda \in \mathbb{C}^*$  simultaneously, since  $t(\lambda u_i) = \lambda t(u_i)$ . Hence, the hypersurface  $\mathcal{C}_\omega$  is parametrised by  $(u_0, \dots, u_3) \in \mathbb{C}^4 - \{0\}$  only up to the action of  $\mathbb{C}^*$ , that is by a  $\mathbb{CP}^3$  with coordinates  $[u_0 : u_1 : u_2 : u_3]$ .

### 1.3 Action along the Deformed Manifold

In the chart  $t = 1$ , the action (6) becomes

$$S(u) = \beta (u_0^2 + u_1^2 - u_2^2 - u_3^2) \quad (17)$$

If we were to reinstate  $t$ , we would have to do it in a way that ensures that  $S(u)$  is a function on  $\mathbb{CP}^4$ , i.e. invariant under the  $\mathbb{C}^*$  action  $\lambda: (t, u_i) \mapsto (\lambda t, \lambda u_i)$ . An obvious candidate is

$$S(t, u) = \frac{\beta (u_0^2 + u_1^2 - u_2^2 - u_3^2)}{t^2} \quad (18)$$

On the constraint surface (inside  $\mathbb{CP}^4$ ) we now replace  $t^2$  by  $\sum_i \omega_i u_i^2$  for some non-zero  $\omega \in \mathbb{C}^4$ . We obtain

$$S(t, u)|_{\mathcal{C}_\omega} = \frac{\beta (u_0^2 + u_1^2 - u_2^2 - u_3^2)}{\sum_i \omega_i u_i^2} \quad (19) \quad \{\text{eq:toy\_S\_deformed}\}$$

where we recall that we denote the (generalised) constraint surface by

$$\mathcal{C}_\omega = \{P_\omega(t, u) = 0\} \quad (20)$$

Note that (19) is invariant under simultaneous scaling  $u_i \mapsto \lambda u_i$  for any  $\lambda \in \mathbb{C}^*$  and thus is indeed a function on  $\mathbb{CP}^3 \cong \mathcal{C}_\omega$ .

## 1.4 Integration Domain

So far we have just discussed how the “complexified” space of fields embeds nicely as a  $\mathbb{CP}^3$  into  $\mathbb{CP}^4$ . However, the original integration domain was a  $\mathbb{CP}^1$  and the correct deformation domain of the path integral should be a deformation of this  $\mathbb{CP}^1$  possibly embedded into  $\mathbb{CP}^3 \subset \mathbb{CP}^4$ . Here, we present *one way* to do it.

Let us briefly recall the first steps of “complexifying” the space of fields: Starting with  $[z_0 : z_1] \in \mathbb{CP}^1$ , we rewrite it into real fields,  $z_k = x_k + iy_k$  for  $k = 0, 1$  and then promoted  $x_k, y_k$  to complex fields  $u_i$  for  $i = 0, \dots, 3$  and ultimately we found that the resulting space is forms a  $\mathbb{CP}^3 \subset \mathbb{CP}^4$  with coordinates  $[u_0 : u_1 : u_2 : u_3]$ .

It is left to discuss how the  $\mathbb{CP}^1$  embeds into this  $\mathbb{CP}^3$ , i.e. how we have to extend the coordinates  $x_k, y_k$  to  $u_i$ . One way to do this is by *scaling*: for some  $\alpha = (\alpha_1, \dots, \alpha_3) \in \mathbb{C}^4$  we set

$$u_0 = \alpha_0 x_0 \quad , \quad u_1 = \alpha_1 x_1 \quad , \quad u_2 = \alpha_2 y_0 \quad , \quad u_3 = \alpha_3 y_1 \quad (21) \quad \{\text{eq:toy\_deform\_contour}\}$$

This scaling has the advantage that it preserves the polynomial degree of  $x_k, y_k$ , and hence of the original complex coordinates  $z_0, z_1$ , as well as the homogeneity. In turn, this implies that we preserve the action of  $\mathbb{C}^*$ . Put differently, the deformation (or analytical continuation) of  $x_k, y_k$  to  $u_i$  is *equivariant*, namely for any  $\lambda \in \mathbb{C}^*$

$$u_i(\lambda x, \lambda y) = \lambda u_i(x, y) \quad (22)$$

Due to this equivariance, the deformation  $z_k \rightsquigarrow u_i$  can be seen as a deformation of  $\mathbb{CP}^1$  *inside*  $\mathbb{CP}^3 \subset \mathbb{CP}^4$ , since

$$\begin{aligned} u: \mathbb{C}^2 &\rightarrow \mathbb{C}^4 \\ (z_0, z_1) &\mapsto u(z_0, z_1) \equiv (u_0(z_0, z_1), u_1(z_0, z_1), u_2(z_0, z_1), u_3(z_0, z_1)) \end{aligned} \quad (23)$$

descends to a map from  $\mathbb{CP}^1$  to  $\mathbb{CP}^3$  due to its equivariance property

$$u(\lambda z_0, \lambda z_1) = \lambda u(z_0, z_1) \quad (24)$$

which simply means that we can consider the quotient by  $\mathbb{C}^*$  on both sides (domain and codomain).

The full path integral domain would then be deformed as follows:

$$\mathcal{E}_{\alpha,\omega}: \mathbb{CP}^1 \xrightarrow{u_\alpha} \mathbb{CP}^3 \xrightarrow{P_\omega} \mathbb{CP}^4 \quad (25) \quad \{\text{eq:toy\_E}\}$$

The first map is given by

$$u_\alpha: [z_0 : z_1] = [x_0 + iy_0 : x_1 + iy_1] \mapsto [\alpha_0 u_0 : \cdots : \alpha_3 u_3] \quad (26) \quad \{\text{eq:toy\_cont\_deform1}\}$$

with  $u_i = u_i(z_0, z_1)$  given by (8). The second map is given by

$$P_\omega: [u_0 : u_1 : u_2 : u_3] \mapsto [P_\omega(u_i) : u_0 : u_1 : u_2 : u_3] \quad (27)$$

where  $P_\omega(u_i) = \sqrt{\sum_i \omega_i u_i^2}$ . The full embedding  $\mathcal{E}_{\alpha,\omega}$  is hence parametrised by two complex vectors,  $\alpha \in \mathbb{C}^4$  and  $\omega \in \mathbb{C}^4$ , which gives 8 complex or equivalently 16 real parameters which can be tuned to maximize the StN.

**Remark 2.** For  $x_k + iy_k$  for  $k = 0, 1$ , the embedding  $\mathcal{E}_{\alpha,\omega}$  can be explicitly expressed as

$$\mathcal{E}_{\alpha,\omega}(x, y) = [\sqrt{\omega_0 \alpha_0^2 x_0^2 + \omega_1 \alpha_1^2 x_1^2 + \omega_2 \alpha_2^2 y_0^2 + \omega_3 \alpha_3^2 y_1^2} : \alpha_0 x_0 : \alpha_1 x_1 : \alpha_2 y_0 : \alpha_3 y_1] \quad (28)$$