Working Title

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1 Toy model: 0d GLSM

1.1 Setup

We start with a 0d GLSM toy model, i.e. we consider the source manifold $\Sigma = \{pt\}$ to be an abstract point and the target manifold to be

$$X=\mathbb{CP}^{N-1}\times\mathbb{CP}^{N-1}$$

The space of fields is then simply given by points on the target manifold X, namely

$$\mathcal{F} = \mathbb{CP}^{N-1} \times \mathbb{CP}^{N-1} \tag{1}$$

As the action of the model we consider

$$S(z,w) = -\beta |\langle \bar{z},w\rangle|^2 = -\beta \sum_{j,k=0}^{N-1} \bar{z}_j z_k \bar{w}_j w_k \tag{2} \qquad \text{eq:toy_S} \}$$

where $z, w \in \mathbb{C}^N$ subject to the condition

$$|z|^2 = \sum_j |z_j|^2 = 1$$
 , $|w|^2 = \sum_k |w_k|^2 = 1$ (3) {eq:toy_constr}

The action enjoys a $U(1) \times U(1)$ gauge freedom (which here is simply a global $U(1) \times U(1)$ symmetry, acting as

$$e^{i\theta} \times e^{i\varphi} \colon (z, w) \mapsto (e^{i\theta}z, e^{i\varphi}w)$$
 (4)

Under the assumption of the condition (3), the action (2) does define a function on $\mathbb{CP}^{N-1} \times \mathbb{CP}^{N-1}$.

The path integral of the model is defined by

$$Z = \int \prod_{i} dz_{j} d\bar{z}_{j} dw_{j} d\bar{w}_{j} \delta(|z|^{2} - 1) \delta(|w|^{2} - 1) e^{-S(z,w)}$$
 (5) {eq:toy_Z}

In order to evaluate (5), we want to embed the space of fields $\mathcal{F} = \mathbb{CP}^{N-1} \times \mathbb{CP}^{N-1}$ into a higher dimensional complex space such that

- 1. the new action \tilde{S} is holomorphic in the new variables (fields)
- 2. when we restrict to \mathcal{F} , \tilde{S} reduces to S

We think of this embedding as an analytic continuation of the space of fields in an appropriate sense.

1.2 Complexification of Real Analytic Manifolds

Let us first recall a basic construction of complexification of real analytic manifolds due to Bruhat and Whitney [1].

Let M be a real analytic manifold of dimension $\dim_{\mathbb{R}} M = m$. Moreover, let $\{U_i, \phi_i\}$ be a real analytic atlas of M, with $U_i \subset \mathbb{R}^m$ and charts $\phi_i \colon U_i \to M$ so that the transition functions

$$\phi_{ij} = \phi_j^{-1} \circ \phi_i \colon U_{ij} \to U_{ij} \tag{6}$$

are real analytic diffeomorphisms. The idea of complexifying M is to find a complex manifold $M^{\mathbb{C}}$ with $\dim_{\mathbb{C}} M^{\mathbb{C}} = m$ and a (real analytic) isomorphism $f \colon M \to \tilde{M} \subset M^{\mathbb{C}}$ of M onto a submanifold of $M^{\mathbb{C}}$. (Fancy way to say that M should be a real analytic submanifold of $M^{\mathbb{C}}$ up to isomorphism) Now, find opens $U_i^{\mathbb{C}} \subset \mathbb{C}^m$ such that $U_i^{\mathbb{C}} \cap \mathbb{R}^m = U_i$ and extend the charts ϕ_i charts $\phi_i^{\mathbb{C}}$ such that

(i) the transition functions $\phi_{ij}^{\mathbb{C}} \colon U_{ij}^{\mathbb{C}} \to U_{ij}^{\mathbb{C}}$ are biholomorphic

(ii)
$$\phi_{ji}^{\mathbb{C}} = \left(\phi_{ij}^{\mathbb{C}}\right)^{-1}$$
 and $\phi_{ii}^{\mathbb{C}} = id$

(iii) the transition functions $\phi_{ij}^{\mathbb{C}}$ satisfy the usual 2-cocycle condition (gluing condition) on $U_{ijk}^{\mathbb{C}}$: $\phi_{ij}^{\mathbb{C}} \circ \phi_{jk}^{\mathbb{C}} = \phi_{ik}^{\mathbb{C}}$

These conditions ensure that we can glue $M^{\mathbb{C}}$ from the local data $\{U_i^{\mathbb{C}}, \phi_i^{\mathbb{C}}\}$:

$$M^{\mathbb{C}} = \coprod_{i} U_{i}^{\mathbb{C}} / \sim \quad , \quad z_{i} \sim z_{j} \text{ iff } z_{j} = \phi_{ji}^{\mathbb{C}}(z_{i}) \text{ on } U_{ij}^{\mathbb{C}}$$
 (7)

For more details on this construction see Cieliebak and Eliashberga's book [2]

1.2.1 Example: The N-sphere

Consider the N-sphere $S^N \subset \mathbb{R}^{N+1}$. First, consider the following atlas: let $p_{\pm} = (0, \dots, 0, \pm 1) \in S^N$ be the north and south pole respectively. We denote points on the sphere by $x = (x_1, \dots, x_{N+1})$, $||x||^2 = 1$ and points in \mathbb{R}^N by $X = (X_1, \dots, X_N)$. The atlas we consider is given by steoreographic projection through p_{\pm} : Let $U_{\pm} = \mathbb{R}^N$ and $V_{\pm} = S^N - \{p_{\pm}\}$. Then define charts

$$\phi_{\pm} \colon U_{\pm} \to V_{\pm} \subset S^N \quad , \quad X \mapsto \left(\frac{2X}{1 + \|X\|^2}, \pm \frac{\|X\|^2 - 1}{\|X\|^2 + 1}\right)$$
 (8)

with inverse

$$\phi_{\pm}^{-1} \colon x \mapsto \left\{ \frac{x_i}{1 \mp x_{N+1}} \right\} \tag{9}$$

{subsec:cmplxfy_spher

From this, one finds the transition functions

$$\phi_{+-} = (\phi_{-+})^{-1} = \phi_{-}^{-1} \circ \phi_{+} \colon X \mapsto \frac{X}{\|X\|^{2}}$$
 (10)

Remark 1. There is a nice geometric interpretation of these transition functions. Note that the map

$$X \mapsto \frac{X}{\|X\|^2} \tag{11}$$

describes an involution at the unit sphere S^{N-1} . On the sphere, the maps differ merely by a sign switch in the x_{N+1} component and indeed, if one computes the transition functions directly on the embedded sphere (that is as maps from $\mathbb{R}^{N+1} \to \mathbb{R}^{N+1}$) one finds

$$\phi_{+} \circ \phi_{-} \colon (x_i, x_{N+1}) \mapsto (x_i, -x_{N+1})$$
 (12)

which corresponds to a reflection of x about the equator! (It helps visualising this for the case 2-sphere) But under stereographic projection, a reflection about the equator becomes an inversion on the unit sphere S^{N-1} in \mathbb{R}^N (again, it helps working this out in the case N=2).

Now, since $U_{\pm} = \mathbb{R}^N$ there exist obvious candidates for a complexification, namely $U_{\pm}^{\mathbb{C}} = \mathbb{C}^N$. We thus promote every X to a complex variable Z = X + iY. Conversely, we can promote any $x \in \mathbb{R}^{N+1}$ satisfying $||x||^2 = \sum_j x_j^2 = 1$ to complex variables z satisfying

$$\sum_{j} z_{j}^{2} = 1 \tag{13} \quad \{eq:toy_quadric\}$$

The above equation defines a hypersurface (so-called quadric) inside \mathbb{C}^{N+1} .

There is a very interesting observation I found in this stackexchange post: the quadric Q defined by (13) is diffeomorphic to the tangent space TS^N . The diffeomorphism is realised by the following map:

$$\Psi \colon TS^N \to Q \quad , \quad (x,y) \mapsto z = \Psi(x,y) = x\sqrt{1 + ||y||^2} + iy$$
 (14)

with inverse

$$\Psi^{-1}(x+iy) = \left(\frac{x}{\sqrt{1+\|y\|^2}}, y\right)$$
 (15)

where $||y||^2 = \sum_i y_i^2$.

Remark 2. Verification that Ψ does indeed define a diffeomorphism goes by a straight forward calculation. (Recall in particular that the tangent space TS^N can be described by pairs $(x,y) \in \mathbb{R}^{N+1}$ such that $\langle x,y \rangle = \sum_i x_i y_i = 0$)

There exists another very interesting diffeomorphism (which I have discovered in this stackexchange post

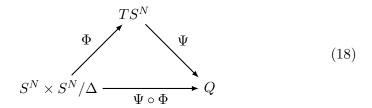
$$\Phi \colon S^N S^N / \Delta \to T S^N \quad , \quad (u,v) \mapsto \left(u, \frac{v - \langle u,v \rangle \, u}{1 - \langle u,v \rangle} \right) \tag{16} \quad \{ \texttt{eq:diff_SNSN_TSN} \}$$

Its inverse is given by

$$\Phi^{-1} \colon TS^N \to S^N \times S^N / \Delta \quad , \quad (x,y) \mapsto \left(x, \frac{x(\|y\|^2 - 1) + 2y}{\|y\|^2 + 1} \right)$$
 (17)

Remark 3. The map (16) is the sterepgraphic projection of $v \in S^N$ through the "pole" $u \in S^N$.

Finally we have the following commutative diagram



Remark 4. The original $S^N \subset Q$ was located at $S^N = Q \cap \mathbb{R}^{N+1}$. If we follow this through, we see that this S^N corresponds to the zero section $\{(x,0) \in TS^N\}$ inside TS^N and consequently is defined by $\Phi(u,v) = (u,0)$ inside $S^N \times S^N/\Delta$. Note that this is precisely the anti-diagonal

$$\bar{\Delta} = \{(u, -u) \mid u \in S^N\} \subset S^N \times S^N \tag{19}$$

Indeed, the stereographic projection of -u ("south pole") through u ("north pole") gives zero:

$$-u \mapsto \frac{(-u) - \langle (-u), u \rangle u}{1 - \langle (-u), u \rangle} = \frac{-u + ||u||^2 u}{1 + ||u||^2} = 0$$
 (20)

since $||u||^2 = 1$. Hence

$$\Phi(u, -u) = (u, 0) \tag{21}$$

and the original S^N is located along the anti-diagonal $\bar{\Delta}$.

1.3 Factor from δ -function Constraint

Recall that we are interested in integrals of the form

$$I = \int_{\mathbb{CP}^n \times \mathbb{CP}^n} d\mu(z) d\mu(w) e^{-S(z,w)} \mathcal{O}(z,w)$$
 (22)

where $d\mu(z) = \prod_i dz_i d\bar{z}_i$ and we have included some observable $\mathcal{O} \colon \mathbb{CP}^n \to \mathbb{C}$. For our toy model, we consider the action functional (2)

$$S(z, w) = -\beta |\langle \bar{z}, w \rangle|^2 \tag{23}$$

It is instructive pass to real coordinates. If $z_k = a_k + ib_k$, let x_k be the real vector $(a_k \ b_k)^t$. If we collect all components, we can form the real vector $x = (\ldots a_k \ b_k \ldots)^t \in \mathbb{R}^{2n+2}$. We denote the real vector associated to w by y. In terms of these, the path integral over can be written as

$$vol(S^{1})^{2} \int_{\mathbb{R}^{2n+2} \times \mathbb{R}^{2n+2}} d^{2n+2}x \ d^{2n+2y} \ \delta(\|x\|^{2} - 1)\delta(\|y\|^{2} - 1)e^{-S(x,y)}\mathcal{O}(x,y)$$
(24)

where the pre-factor stems from the compact part of the gauge group $\mathbb{C}^* \subset \mathbb{C}^{n+1}$. Let us focus on the following part of the integral:

$$I \sim \int_{\mathbb{R}^{2n+2}} \delta(\|x\|^2 - 1) = \int_{\mathbb{R}^{2n+2}} \delta(\langle x, x \rangle - 1)$$
 (25)

Following our idea, we complexify the space of integration, $x \to \zeta$

$$I \sim \int_{\Gamma_0 = \mathbb{R}^{2n+2} \cap \mathbb{C}^{2n+2}} \delta(\langle \zeta, \zeta \rangle - 1)$$
 (26)

where

$$\langle \zeta, \zeta \rangle = \sum_{k} \zeta_k^2 \tag{27}$$

The δ -function constraint restricts the support of the integrand to the quadric Q. Now, we would like to deform the "contour" (domain of integration) inside Q

$$I \sim \int_{\Gamma_a \subset \mathbb{C}^{2n+2}} \delta(\langle \zeta, \zeta \rangle - 1)$$
 (28)

Inspired by the discussion in 1.2.1, we would like to parametrise Q in terms of TS^{2n+1} . We thus choose a parametrisation of the following form: (by abuse of notation, we will use again x to denote a vector in \mathbb{R}^{2n+2})

$$\zeta(x) = x\sqrt{1 + \|y(x)\|^2} + iy(x) \tag{29} \quad \{\texttt{eq:parametrisation}\}$$

where y(x) is chosen in such a way that $\langle x, y \rangle = 0$. It follows that

$$I \sim \int_{\mathbb{R}^{2n+2}} d^{2n+2}x \det J(x)\delta(\langle \zeta(x), \zeta(x) \rangle - 1)$$
 (30)

where J(x) denotes the Jacobian of (29). Note that the δ -function constraint can be simplified as follows: Let

$$\lambda(x) = \sqrt{1 + ||y||^2} \tag{31}$$

Then

$$C(x) = \langle \zeta(x), \zeta(x) \rangle - 1$$

= $\lambda^{2}(x) ||x||^{2} + 2i\lambda(x) \langle x, y(x) \rangle - ||y||^{2} - 1$
= $\lambda^{2}(x) (||x||^{2} - 1)$ (32)

Importantly,

$$\lambda^2(x) > 0 \tag{33}$$

such that the δ -function constraint is of the following form:

$$\int \delta(C(x)) = \int d\mu(x) \delta(\underbrace{f(x)g(x)}_{\equiv \phi(x)})$$

$$= \int_{\phi^{-1}(0)} \frac{d\sigma}{\|\nabla\phi(x)\|}$$

$$= \int_{f^{-1}(0)} \frac{d\sigma}{\|f(x)\nabla g(x) + g(x)\nabla f(x)\|}$$

$$= \int_{f^{-1}(x)} \frac{d\sigma}{\|\nabla f(x)\|} \frac{1}{|g(x)|}$$

$$= \int d\mu \frac{\delta(f(x))}{|g(x)|}$$
(34)

where $f(x) = ||x||^2 - 1$ and $g(x) = \lambda^2(x) > 0$. Hence we find that with the chosen parametrisation, the integral nicely localises to an integral over S^{2n+1}

$$\int_{\mathbb{R}^{2n+2}} \delta(C(x)) = \int_{\mathbb{R}^{2n+2}} \frac{\delta(\|x\|^2 - 1)}{\lambda^2(x)} = \int_{S^{2n+1}} \frac{1}{\lambda^2(x)}$$
(35)

Therefore, we are left with an integral of the form

$$I \sim \int_{S^{2n+1}} \frac{\det J(x)}{\lambda^2(x)}$$
 , $\lambda^2(x) = 1 + \|y(x)\|^2$ (36) {eq:toy_schema_I}

1.4 Homogeneous Deformations

1.4.1 The General Case

The n=0 example can be nicely generalised. Tej suggested to consider deformations of the form

eformations of the form
$$Y_a(X) = \Omega(a)X \quad , \quad \Omega(a) = \operatorname{diag}\left(\begin{pmatrix} 0 & -a_1 \\ a_1 & 0 \end{pmatrix}, \dots, \begin{pmatrix} 0 & -a_{n+1} \\ a_{n+1} & 0 \end{pmatrix}\right)$$

$$(37) \quad \{\text{eq:Tej_deformation}\}$$
or some $a = (a_1, \dots, a_{n+1}) \in \mathbb{R}^{n+1}$. This type of deformation have a very

for some $a = (a_1, \ldots, a_{n+1}) \in \mathbb{R}^{n+1}$. This type of deformation have a very nice geometric origin, as we now explain.

Remark 5. Under the isomorphism

$$\mathbb{R}^{2n+2} \cong \mathbb{C}^{n+1} \quad , \quad \begin{pmatrix} x_k \\ y_k \end{pmatrix} \mapsto z_k = x_k + iy_k$$
 (38)

the deformation (37) is acting on \mathbb{C}^{n+1} by multiplication with the diagonal matrix $\operatorname{diag}(ia_1,\ldots,ia_{n+1})$

$$\begin{pmatrix}
\vdots \\
z_k \\
\vdots
\end{pmatrix}
\mapsto
\begin{pmatrix}
\vdots \\
ia_k z_k \\
\vdots
\end{pmatrix}$$
(39)

The finite diffeomorphism generated by this vector field is given by

$$\begin{pmatrix} \vdots \\ z_k \\ \vdots \end{pmatrix} \mapsto \begin{pmatrix} \vdots \\ e^{ia_k} z_k \\ \vdots \end{pmatrix} \tag{40}$$

The generated finite action is thus the action of the torus

$$U(1)^{n+1} \subset SU(n+1) \tag{41}$$

Let G be a Lie group and H < G a subgroup with Lie algebras \mathfrak{g} and \mathfrak{h} respectively. The orbit space G/H is known as a homogeneous space. It is called reductive if \mathfrak{g} allows the following decomposition

$$\mathfrak{g} = \mathfrak{h} \oplus \mathfrak{m} \quad , \quad [\mathfrak{h}, \mathfrak{h}] \subset \mathfrak{h} \quad , \quad [\mathfrak{h}, \mathfrak{m}] \subset \mathfrak{m}$$
 (42)

If \mathfrak{g} is equipped with a Killing form B (or if G allows a G-invariant metric as a manifold), then G/H is automatically reductive with the choice

$$\mathfrak{m} = \mathfrak{h}^{\perp} \tag{43}$$

where \mathfrak{h}^{\perp} is the orthogonal to \mathfrak{h} with respect to B. Indeed, since B is G-invariant, schematically

$$0 = B(\mathfrak{h}, \mathfrak{m}) = B([\mathfrak{h}, \mathfrak{h}], \mathfrak{m}) = B(\mathfrak{h}, [\mathfrak{h}, \mathfrak{m}])$$
(44)

so that $[\mathfrak{h},\mathfrak{m}]\subset\mathfrak{m}$.

In summary, given a reductive homogeneous space G/H, we get a nice decomposition of the tangent space: first, let o = [e] be the class in G/H at the identity. Then

$$T_o(G/H) \cong \mathfrak{m}$$
 (45)

This isomorphism is given essentially by the pushforward of the natural projection $\pi\colon G\to G/H$. In fact, $\pi\colon G\to G/H$ defines a principal H-bundle. Its vertical vector fields are given by $\ker(\pi_*)=\mathfrak{h}$ and since π is a submersion (π and π_* are both surjective) we have

$$\pi_* : \mathfrak{g}/\ker(\pi_*) = \mathfrak{g}/\mathfrak{h} \cong T_o(G/H)$$
 (46)

Then, to compute the tangent space at any other point, we use translation by G.

Now, recall that

$$S^{2n+1} = SU(n+1)/SU(n)$$
 (47)

is a homogeneous space where H = SU(n) sits inside G = SU(n+1) as the lower right corner

$$SU(n) \ni h \hookrightarrow \begin{pmatrix} 1 & 0 \\ 0 & h \end{pmatrix} \in SU(n+1)$$
 (48)

The reductive split is given by

$$\mathfrak{su}(n+1) = \mathfrak{su}(n) \oplus \mathfrak{m}$$
 (49) {eq:red_split}

Since $\mathfrak{su}(n+1)$ admits a Killing form $B(x,y) \propto tr(x,y)$, we can choose $\mathfrak{m} = \mathfrak{su}(n)^{\perp}$. In the split (49), we have

$$\mathfrak{su}(n) \equiv \begin{pmatrix} 0 & 0 \\ 0 & \mathfrak{su}(n) \end{pmatrix} \subset \mathfrak{su}(n+1) \tag{50}$$

We then find

$$\mathfrak{m} = \mathfrak{su}(n)^{\perp} = \left\{ \begin{pmatrix} 0 & -\zeta^* \\ \zeta & 0 \end{pmatrix} \right\} \quad , \quad \zeta \in \mathbb{C}^{n+1}$$
 (51)

Here $\zeta^* = \bar{\zeta}^t$ denotes the hermitian conjugate.

In order to describe the tangent space $T_{[g]}(G/K)$ at any point $[g] \in G/H$, we can use the tangential map of left translations. In fact, left translation by g is defined by the map

$$L_g \colon G \to G \quad , \quad g' \mapsto gg'$$
 (52)

so that its tangential map

$$\theta_g = (L_{q^{-1}})_* \colon T_g G \to T_e G = \mathfrak{g} \tag{53}$$

This map defines a \mathfrak{g} -valued 1-form on G known as the left-invariant Maurer-Cartan form. For a matrix group, it can be written as

$$\theta_q = g^{-1}dg \tag{54}$$

Note that the natural projection $\pi\colon G\to G/H$ is equivariant with respect to left translation: if

$$\tau_g \colon G/H \to G/H \quad , \quad [g'] \mapsto [gg']$$
 (55)

then we have

$$\pi \circ L_g = \tau_g \circ \pi \quad , \quad \tau_g \circ \tau_{g'} = \tau_{gg'}$$
 (56)

Moreover,

$$(\pi \circ L_{g^{-1}})_* = \pi_* \circ (L_{g^{-1}})_* = (\tau_{g^{-1}} \circ \pi)_* = (\tau_{g^{-1}})_* \circ \pi_*$$
 (57)

Hence, denoting $(\tau_{q^{-1}})_* = \vartheta_g$, we have

$$\pi_* \circ \theta_g = \vartheta_g \circ \pi_* \colon T_g G \to T_g(G/H) = \mathfrak{m} \tag{58}$$

For us, we will be mainly interested in the map

$$(\tau_q)_* : \mathfrak{m} = T_{[e]}(G/H) \to T_{[q]}(G/H)$$
 (59)

which parametrises the tangent space of G/H at a point [g] by \mathfrak{m} .

Now, if G is a matrix group the pushforward $(\tau_g)_*$ is essentially just given by left multiplication with g (matrix multiplication is linear) and so in we could describe the tangent space $T_p(G/H)$ at a point $p = g(p)p_0$ as follows

$$T_p(G/H) \cong \{ (p, g(p)\xi) \mid \xi \in \mathfrak{m} \}$$
(60)

where we implicitly identify g(p)X with $g(p)Xp_0$.

Remark 6. In theory, this is very nice. In practice, I believe it is computationally expensive: for a given $p \in S^{2n+1} = \mathrm{SU}(n+1)/\mathrm{SU}(n)$, we would need to find g(p) which on top of it is only defined up to a multiplication on the right of $\mathrm{SU}(n)$. We could, for example, choose $p_0 = (1, 0 \dots)^t$. There is a way how to algorithmically compute g(p): A matrix belongs to $\mathrm{SU}(n+1)$ iff its columns (rows) form an orthonormal basis of \mathbb{C}^{n+1} . If $p_0 = e_1$ is the first standard basis vector, then we start with the basis $\{p, e_2, \dots, e_{n+1}\}$ and by Gram-Schmidt produce an orthonormal basis $\{p, \tilde{e}_1, \dots, \tilde{e}_{n+1}\}$ and set

$$g(p) = (p \ \tilde{e}_2 \ \dots; \tilde{e}_n) \tag{61}$$

However, I believe that Gram-Schmidt is computationally expensive.

Moreover, it might simply not be necessary. In order to define interesting deformations, we could simply consider any $X \in \mathfrak{g}$. Any such ξ will admit a split into $\xi = \xi_{\mathfrak{m}} + \xi_{\mathfrak{h}}$ and since $\mathfrak{h} = \text{Lie}(H) \cong \text{Stab}(p)$ is in the Lie algebra of the stablizer of p (note that we actually mean p here, not p_0), $\xi_h p = 0$ (as hp = p for any h in H). Thus, we could simply not care and just work with X as a whole and define deformations for any $\xi \in \mathfrak{g}$ by

$$Y_{\xi}(X) = \rho(\xi) \cdot X \quad , \quad \xi \in \mathfrak{g} \quad , \quad X \in S^{2n+1}$$
 (62)

Here, ρ denotes the appropriate representation of \mathfrak{g} on S^{2n+1} . I believe its easiest description is as follows: $\mathrm{SU}(n+1)$ and hence $\mathfrak{su}(n+1)$ naturally acts (by matrix multiplication) on S^{2n+1} when viewed as a subset of \mathbb{C}^{n+1} . Under the isomorphism $\mathbb{C} \cong \mathbb{R}^2$, where $z = x + iy \mapsto (x, y)$, multiplication by a complex number a + ib becomes matrix multiplication:

$$(a+ib)z \iff \begin{pmatrix} a & -b \\ b & a \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} \tag{63}$$

Hence, using our representation of \mathbb{C}^{n+1}

$$\begin{pmatrix} \vdots \\ z_k = x_k + iy_k \\ \vdots \end{pmatrix} \mapsto \begin{pmatrix} \vdots \\ x_k \\ y_k \\ \vdots \end{pmatrix}$$
 (64)

we have to simply replace any complex number a + ib in $\xi \in \mathfrak{su}(n+1)$ by the appropriate matrix:

$$\rho(\xi) \colon a + ib \leadsto \begin{pmatrix} a & -b \\ b & a \end{pmatrix} \tag{65}$$

In practice, it might be easier to have the following workflow:

$$S^{2n+1} \subset \mathbb{R}^{2n+2} \longrightarrow \mathbb{C}^{n+1} \xrightarrow{\xi} \mathbb{C}^{n+1} \longrightarrow \mathbb{R}^{2n+2}$$

$$X \longrightarrow Z \xrightarrow{\xi} \xi \cdot Z \longrightarrow Y_{\xi}(X)$$
(66)

1.5 Example: n = 0

1.5.1 Homogeneous deformation

A particular nice example, which can be computed explicitly and serves as a check that everything works nicely is given for the case n = 0. Consider the deformation

$$y(x) = \alpha \Omega x \quad , \quad \|x\|^2 = 1 \quad , \quad \Omega = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \in \mathrm{SO}(2) \quad , \quad \alpha \in \mathbb{R} \quad (67) \quad \{\mathrm{eq:toy_cst_def}\}$$

so that indeed

$$\langle x, y(x) \rangle = 0 \tag{68}$$

Moreover, note that

$$||y(x)||^2 = \langle y(x), y(x) \rangle = \alpha^2 \langle x, \Omega^t \Omega x \rangle = \alpha^2 ||x||^2 = \alpha^2$$
 (69)

and hence

$$\lambda^2(x) = 1 + \alpha^2 \tag{70}$$

We compute the Jacobian of this parametrisation as follows: let $M_i^{\ j} = \partial_i y^j(x)$. Then

$$J_i^{\ j} = \lambda(x)\delta_i^{\ j} + \frac{\langle y, \partial_i y \rangle x^j}{\lambda(x)} + iM_i^{\ j} = \lambda(x)\delta_i^{\ j} + \frac{M_{ik}y^k x^j}{\lambda(x)} + iM_i^{\ j}$$
 (71)

Using bra-ket notation and trivially lowering indices (considering the flat metric on \mathbb{R}^2), we can write J as follows:

$$J = \lambda(x) \cdot id + \lambda^{-1}(x)|My\rangle\langle x| + iM \tag{72}$$

Remark 7. The matrix M has some interesting properties.

1. From $\langle x, y(x) \rangle = 0$ it follows that

$$0 = \partial_i \langle x, y(x) \rangle = y_i(x) + M_{ik} x_k$$

so that

$$y_i(x) = M_{ik}x_k \implies y(x) = Mx$$

and hence

$$\langle x, y \rangle = \langle x, Mx \rangle = \langle M^t x, x \rangle = 0$$

from which follows (if $y(x) = Mx \neq 0$)

$$\langle x|M=0$$

Now, for our choice of

$$y(x) = \alpha \Omega x$$

we have

$$M = \alpha \Omega^t = -\alpha \Omega \tag{73}$$

and hence

$$J = \lambda(x) \cdot id + \lambda^{-1}(x)\alpha^{2}\Omega^{t}\Omega|x\rangle\langle x| - i\alpha\Omega$$

= $\lambda(x) \cdot id + \lambda^{-1}(x)\alpha^{2}|x\rangle\langle x| - i\alpha\Omega$ (74)

Now, here is a neat trick due to Tej. Consider the unitary matrix

$$U = (x \ \Omega x) \quad , \quad \Omega^* = \begin{pmatrix} x^t \\ x^t \Omega^t \end{pmatrix}$$
 (75)

where x is a 2×2 column vector with unit norm. Then

$$U^*U = \begin{pmatrix} \langle x, x \rangle & \langle x, \Omega x \rangle \\ \langle \Omega x, x \rangle & \langle x, x \rangle \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$
 (76)

Moreover,

$$\langle x|U = x^t U = (\langle x, x \rangle \ \langle x, \Omega x \rangle) = (1\ 0) \tag{77}$$

as well as

$$U^*\Omega U = U^*(\Omega x \ \Omega^2 x) = U^*(\Omega x \ -x) = \begin{pmatrix} \langle x, \Omega x \rangle & -\langle x, x \rangle \\ \langle x, \Omega^t \Omega x \rangle & -\langle x, \Omega^t x \rangle \end{pmatrix} = -\Omega$$
(78)

Then

$$U^*JU = \lambda(x) \cdot id + \lambda^{-1}(x)\alpha^2 U^* |x\rangle \langle x|U - i\alpha U^*\Omega U$$

$$= \lambda(x) \cdot id + \lambda^{-1}(x)\alpha^2 U^* \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} + i\alpha\Omega$$

$$= \begin{pmatrix} \lambda(x) + \lambda^{-1}(x)\alpha^2 & i\alpha \\ -i\alpha & \lambda(x) \end{pmatrix}$$
(79)

It follows that

$$\det J(x) = \det(U^*J(x)U)$$

$$= \det\begin{pmatrix} \lambda(x) + \lambda^{-1}(x)\alpha^2 & i\alpha \\ -i\alpha & \lambda(x) \end{pmatrix}$$

$$= \lambda^2(x) + \alpha^2 - \alpha^2 = \lambda^2(x)$$
(80)

Finally, we obtain that

$$I \sim \int_{S^{2n+1}} \frac{\det J(x)}{\lambda^2(x)} = \int_{S^{2n+1}} \frac{\lambda^2(x)}{\lambda^2(x)} = \int_{S^{2n+1}}$$
(81)

This means that the constant deformation (67) comes with a trivial total Jacobian (Jacobian + factor from the φ -function constraint).

1.5.2 What it really means

Recall that for n=0 we consider the path integral over S^1 . Before, we have considered $x \in S^1 \subset \mathbb{R}^2$. Now, we might think of its complex formulation: let us consider $z = e^{i\theta} \in S^1 \subset \mathbb{C}$. In [?], the complexification was chosen in terms of $\theta \to \theta + i\tau$. This means that

$$z \mapsto \tilde{z} = e^{i(\theta + i\tau)} = e^{-\tau}z \tag{82}$$

On the other hand, we can parametrize x by

$$x = \begin{pmatrix} \sin(\theta) \\ \cos(\theta) \end{pmatrix} \tag{83}$$

so that

$$\tilde{z} = \begin{pmatrix} \sin(\theta + i\tau) \\ \cos(\theta + i\tau) \end{pmatrix} = \begin{pmatrix} \sin(\theta)\cosh(\tau) + i\cos(\theta)\sinh(\tau) \\ \cos(\theta)\cosh(\tau) - i\sin(\theta)\sinh(\tau) \end{pmatrix} \\
= \begin{pmatrix} \cosh(\tau) & i\sinh(\tau) \\ -i\sinh(\tau) & \cosh(\tau) \end{pmatrix} \begin{pmatrix} \sin(\theta) \\ \cos(\theta) \end{pmatrix} = \lambda x + i\underbrace{\Omega x}_{y(x)}$$
(84)

where

$$\Omega = \begin{pmatrix} 0 & -\sinh(\tau) \\ \sinh(\tau) & 0 \end{pmatrix} \tag{85}$$

and

$$\lambda^2 = 1 + ||y(x)||^2 = 1 + \sinh^2(\tau) = \cosh^2(\tau)$$
(86)

To summarise, if we consider the homogeneous deformation (67),

$$y(x) = \alpha \Omega x$$
 , $||x||^2 = 1$, $\Omega = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$ (87)

then choosing

$$\boxed{\alpha = \sinh(\tau)} \tag{88}$$

we find that the homogeneous deformation simply scales the complex variable $z=e^{i\theta}\in S^1\subset \mathbb{C},$

$$z \to \tilde{z} = e^{-\tau} z \tag{89}$$

{subsubsec:Torus_defo

1.5.3 Generalisation: Torus deformations

Let us now think about n > 0. Let $z = [z_0 : \cdots : z_n] \in \mathbb{CP}^n$. Again, we factor out the U(1) $\subset \mathbb{C}^*$ and fix the $\mathbb{R}_+ \subset \mathbb{C}^*$ by restricting the path integral to S^{2n+1} , i.e.

$$\sum_{k} |z_k|^2 = 1 \tag{90}$$

Following the logic of the homogeneous deformation, let us consider a deformation parametrised by a generic element in the Cartan subalgebra of $\mathfrak{su}(n+1)$

$$Y = \Omega(\alpha)X = \sum_{k=1}^{n} \rho(i\alpha_k H_k)$$
(91)

where we choose the following hermitian (hence the factor of i) basis of the Cartan $\mathfrak{h} \subset \mathfrak{su}(n+1)$ subalgebra

$$H_k = E_{kk} - E_{(k+1)(k+1)} (92)$$

where E_{ij} is the $n \times n$ matrix with $(E_{ij})_{ab} = \delta_{ai}\delta_{bj}$. Here, ρ is again the representation of complex multiplication in terms of 2×2 matrices,

$$\rho(a+ib) = \begin{pmatrix} a & -b \\ b & a \end{pmatrix} \tag{93}$$

and defined element-wise on a generic complex matrix as

$$\rho(M) = \begin{pmatrix} \rho(m_{11}) & \rho(m_{12}) & \dots \\ \rho(m_{21} & \rho(m_{22}) & \dots \\ \vdots & \vdots & \end{pmatrix}$$
(94)

In particular,

$$\rho(iH_k) = \rho(iE_{kk}) - \rho(iE_{(k+1)(k+1)}) = \begin{pmatrix} \ddots & & & & \\ & 0 & -1 & & \\ & 1 & 0 & & \\ & & & 0 & 1 & \\ & & & -1 & 0 & \\ & & & & \ddots \end{pmatrix}$$
(95)

Let us choose

$$\alpha_k = \sinh(\tau_k) \quad , \quad \tau_0 \equiv \tau_{n+1} \equiv 0$$
 (96)

then, acting on the complex vector $z = (z_0, \ldots, z_n)$, we find

$$[\tilde{z} = (\tilde{z}_0, \dots, \tilde{z}_n) = (e^{-\Delta_{01}} z_0, \dots, e^{-\Delta_{nn+1}} z_n) , \Delta_{kk+1} = \tau_k - \tau_{k+1}]$$
 (97)

1.6 Correlators

1.6.1 One-point function

It is time to look at interesting operators we would like to compute. We start with the one-point function. Let

$$\mathcal{O}_{ij}(z) = z_i \bar{z}_j \tag{98}$$

Clearly, \mathcal{O}_{ij} is invariant under the U(1) action $z_k \to e^{i\varphi} z_k$ for all k. It's expectation value is given by

$$\langle \mathcal{O}_{ij}(z)\rangle = \frac{1}{(2\pi)^2} \int_{S^{2n+1}\times S^{2n+1}} \prod_k d^n z_k d^n w_k \ z_i \bar{z}_j e^{-S(z,w)} \qquad (99) \quad \{\texttt{eq:toy_one_pt}\}$$

where

$$S(z,w) = -\beta |\langle \bar{z}, w \rangle|^2 = -\beta \sum_{k,\ell} \bar{z}_{\ell} z_k w_{\ell} \bar{w}_k$$
 (100)

Note that both, the action and the measure is invariant under the diagonal U(1) action

$$(z_k, w_k) \mapsto (e^{i\alpha} z_k, e^{i\alpha} w_k) \quad , \quad \forall k$$
 (101)

If $i \neq j$, by this change of coordinates, for k = i, it follows that

$$\langle \mathcal{O}_{ij}(z) \rangle = e^{-i\alpha} \langle \mathcal{O}_{ij}(z) \rangle$$
 (102)

for any α and hence we must have

$$\boxed{\langle \mathcal{O}_{ij}(z)\rangle = 0 \quad , \quad \forall z, i \neq j} \tag{103} \quad \{\texttt{eq:toy_var_one_pt}\}$$

Its variance is given by

$$\operatorname{Var}(\mathcal{O}_{ij}(z)) = \left\langle |\mathcal{O}_{ij}(z)|^2 \right\rangle - |\langle \mathcal{O}_{ij}(z) \rangle|^2 = \left\langle \mathcal{O}_{ij}(z) \mathcal{O}_{ij}^*(z) \right\rangle - \left\langle \mathcal{O}_{ij}(z) \right\rangle \left\langle \mathcal{O}_{ij}^*(z) \right\rangle$$
(104)

so that by (103)

$$\operatorname{Var}(\mathcal{O}_{ij}(z)) = \langle |\mathcal{O}_{ij}(z)|^2 \rangle = \langle |z_i|^2 |z_j|^2 \rangle \tag{105}$$

If we consider a torus / Cartan deformation as outlined in Section 1.5.3, then we find

$$\operatorname{Var}(\widetilde{\mathcal{O}}_{ij}(\tilde{z})) = \left\langle e^{-2\Delta_{ii+1}} |z_i|^2 e^{-2\Delta_{jj+1}} |z_j|^2 \right\rangle_{deformed}$$

$$= e^{-2\Delta_{ii+1} - 2\Delta_{jj+1}} \left\langle |z_i|^2 |z_j|^2 \det(J_{tot}) e^{-(\tilde{S} - S)} \right\rangle$$
(106)

1.6.2 Two-point function

References

- [1] F. Bruhat and H. Whitney, Quelques propriétés fondamentales des ensembles analytiques-réels, Comment. Math. Helv. 33, 132-160 (1959).
- [2] K. Cieliebak and Y. Eliashberg. From Stein to Weinstein and back: symplectic geometry of affine complex manifolds. Vol. 59. American Mathematical Soc., 2012.