

# Working Title

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## 1 Toy model: $0d$ GLSM

### 1.1 Setup

We start with a  $0d$  GLSM toy model, i.e. we consider the source manifold  $\Sigma = \{pt\}$  to be an abstract point and the target manifold to be

$$X = \mathbb{CP}^{N-1} \times \mathbb{CP}^{N-1}$$

The space of fields is then simply given by points on the target manifold  $X$ , namely

$$\mathcal{F} = \mathbb{CP}^{N-1} \times \mathbb{CP}^{N-1} \tag{1}$$

As the action of the model we consider

$$S(z, w) = -\beta |\langle \bar{z}, w \rangle|^2 = -\beta \sum_{j,k=0}^{N-1} \bar{z}_j z_k \bar{w}_j w_k \tag{2} \quad \{\text{eq:toy\_S}\}$$

where  $z, w \in \mathbb{C}^N$  subject to the condition

$$|z|^2 = \sum_j |z_j|^2 = 1 \quad , \quad |w|^2 = \sum_k |w_k|^2 = 1 \quad (3) \quad \{\text{eq:toy\_constr}\}$$

The action enjoys a  $U(1) \times U(1)$  gauge freedom (which here is simply a global  $U(1) \times U(1)$  symmetry, acting as

$$e^{i\theta} \times e^{i\varphi}: (z, w) \mapsto (e^{i\theta} z, e^{i\varphi} w) \quad (4)$$

Under the assumption of the condition (3), the action (2) does define a function on  $\mathbb{CP}^{N-1} \times \mathbb{CP}^{N-1}$ .

The path integral of the model is defined by

$$Z = \int \prod_j dz_j d\bar{z}_j dw_j d\bar{w}_j \delta(|z|^2 - 1) \delta(|w|^2 - 1) e^{-S(z, w)} \quad (5) \quad \{\text{eq:toy\_Z}\}$$

In order to evaluate (5), we want to embed the space of fields  $\mathcal{F} = \mathbb{CP}^{N-1} \times \mathbb{CP}^{N-1}$  into a higher dimensional complex space such that

1. the new action  $\tilde{S}$  is holomorphic in the new variables (fields)
2. when we restrict to  $\mathcal{F}$ ,  $\tilde{S}$  reduces to  $S$

We think of this embedding as an analytic continuation of the space of fields in an appropriate sense.

## 1.2 Complexification of Real Analytic Manifolds

Let us first recall a basic construction of complexification of real analytic manifolds due to Bruhat and Whitney [1].

Let  $M$  be a real analytic manifold of dimension  $\dim_{\mathbb{R}} M = m$ . Moreover, let  $\{U_i, \phi_i\}$  be a real analytic atlas of  $M$ , with  $U_i \subset \mathbb{R}^m$  and charts  $\phi_i: U_i \rightarrow M$  so that the transition functions

$$\phi_{ij} = \phi_j^{-1} \circ \phi_i: U_{ij} \rightarrow U_{ij} \quad (6)$$

are real analytic diffeomorphisms. The idea of complexifying  $M$  is to find a complex manifold  $M^{\mathbb{C}}$  with  $\dim_{\mathbb{C}} M^{\mathbb{C}} = m$  and a (real analytic) isomorphism  $f: M \rightarrow \tilde{M} \subset M^{\mathbb{C}}$  of  $M$  onto a submanifold of  $M^{\mathbb{C}}$ . (Fancy way to say that  $M$  should be a real analytic submanifold of  $M^{\mathbb{C}}$  up to isomorphism) Now, find opens  $U_i^{\mathbb{C}} \subset \mathbb{C}^m$  such that  $U_i^{\mathbb{C}} \cap \mathbb{R}^m = U_i$  and extend the charts  $\phi_i$  charts  $\phi_i^{\mathbb{C}}$  such that

(i) the transition functions  $\phi_{ij}^{\mathbb{C}}: U_{ij}^{\mathbb{C}} \rightarrow U_{ij}^{\mathbb{C}}$  are biholomorphic

(ii)  $\phi_{ji}^{\mathbb{C}} = \left(\phi_{ij}^{\mathbb{C}}\right)^{-1}$  and  $\phi_{ii}^{\mathbb{C}} = id$

(iii) the transition functions  $\phi_{ij}^{\mathbb{C}}$  satisfy the usual 2-cocycle condition (gluing condition) on  $U_{ijk}^{\mathbb{C}}$ :  $\phi_{ij}^{\mathbb{C}} \circ \phi_{jk}^{\mathbb{C}} = \phi_{ik}^{\mathbb{C}}$

These conditions ensure that we can glue  $M^{\mathbb{C}}$  from the local data  $\{U_i^{\mathbb{C}}, \phi_i^{\mathbb{C}}\}$ :

$$M^{\mathbb{C}} = \coprod_i U_i^{\mathbb{C}} / \sim \quad , \quad z_i \sim z_j \text{ iff } z_j = \phi_{ji}^{\mathbb{C}}(z_i) \text{ on } U_{ij}^{\mathbb{C}} \quad (7)$$

For more details on this construction see [Cieliebak and Eliashberg's book \[2\]](#)

### 1.2.1 Example: The $N$ -sphere

Consider the  $N$ -sphere  $S^N \subset \mathbb{R}^{N+1}$ . First, consider the following atlas: let  $p_{\pm} = (0, \dots, 0, \pm 1) \in S^N$  be the north and south pole respectively. We denote points on the sphere by  $x = (x_1, \dots, x_{N+1})$ ,  $\|x\|^2 = 1$  and points in  $\mathbb{R}^N$  by  $X = (X_1, \dots, X_N)$ . The atlas we consider is given by stereographic projection through  $p_{\pm}$ : Let  $U_{\pm} = \mathbb{R}^N$  and  $V_{\pm} = S^N - \{p_{\pm}\}$ . Then define charts

$$\phi_{\pm}: U_{\pm} \rightarrow V_{\pm} \subset S^N \quad , \quad X \mapsto \left( \frac{2X}{1 + \|X\|^2}, \pm \frac{\|X\|^2 - 1}{\|X\|^2 + 1} \right) \quad (8)$$

with inverse

$$\phi_{\pm}^{-1}: x \mapsto \left\{ \frac{x_i}{1 \mp x_{N+1}} \right\} \quad (9)$$

From this, one finds the transition functions

$$\phi_{+-} = (\phi_{-+})^{-1} = \phi_{-}^{-1} \circ \phi_{+}: X \mapsto \frac{X}{\|X\|^2} \quad (10)$$

**Remark 1.** There is a nice geometric interpretation of these transition functions. Note that the map

$$X \mapsto \frac{X}{\|X\|^2} \quad (11)$$

describes an involution at the unit sphere  $S^{N-1}$ . On the sphere, the maps differ merely by a sign switch in the  $x_{N+1}$  component and indeed, if one computes the transition functions directly on the embedded sphere (that is as maps from  $\mathbb{R}^{N+1} \rightarrow \mathbb{R}^{N+1}$ ) one finds

$$\phi_{+} \circ \phi_{-}: (x_i, x_{N+1}) \mapsto (x_i, -x_{N+1}) \quad (12)$$

which corresponds to a reflection of  $x$  about the equator! (It helps visualising this for the case 2-sphere) But under stereographic projection, a reflection about the equator becomes an inversion on the unit sphere  $S^{N-1}$  in  $\mathbb{R}^N$  (again, it helps working this out in the case  $N = 2$ ).

Now, since  $U_{\pm} = \mathbb{R}^N$  there exist obvious candidates for a complexification, namely  $U_{\pm}^{\mathbb{C}} = \mathbb{C}^N$ . We thus promote every  $X$  to a complex variable  $Z = X + iY$ . Conversely, we can promote any  $x \in \mathbb{R}^{N+1}$  satisfying  $\|x\|^2 = \sum_j x_j^2 = 1$  to complex variables  $z$  satisfying

$$\sum_j z_j^2 = 1 \quad (13) \quad \{\text{eq:toy\_quadric}\}$$

The above equation defines a hypersurface (so-called *quadric*) inside  $\mathbb{C}^{N+1}$ .

There is a very interesting observation I found in [this stackexchange post](#): the quadric  $Q$  defined by (13) is *diffeomorphic* to the tangent space  $TS^N$ . The diffeomorphism is realised by the following map:

$$\Psi: TS^N \rightarrow Q \quad , \quad (x, y) \mapsto z = \Psi(x, y) = x\sqrt{1 + \|y\|^2} + iy \quad (14)$$

with inverse

$$\Psi^{-1}(x + iy) = \left( \frac{x}{\sqrt{1 + \|y\|^2}}, y \right) \quad (15)$$

where  $\|y\|^2 = \sum_i y_i^2$ .

**Remark 2.** Verification that  $\Psi$  does indeed define a diffeomorphism goes by a straight forward calculation. (Recall in particular that the tangent space  $TS^N$  can be described by pairs  $(x, y) \in \mathbb{R}^{N+1}$  such that  $\langle x, y \rangle = \sum_i x_i y_i = 0$ )

There exists another very interesting diffeomorphism (which I have discovered in this [stackexchange post](#)

$$\Phi: S^N S^N / \Delta \rightarrow TS^N \quad , \quad (u, v) \mapsto \left( u, \frac{v - \langle u, v \rangle u}{1 - \langle u, v \rangle} \right) \quad (16) \quad \{\text{eq:diff\_SNSN\_TSN}\}$$

Its inverse is given by

$$\Phi^{-1}: TS^N \rightarrow S^N \times S^N / \Delta \quad , \quad (x, y) \mapsto \left( x, \frac{x(\|y\|^2 - 1) + 2y}{\|y\|^2 + 1} \right) \quad (17)$$

**Remark 3.** The map (16) is the stereographic projection of  $v \in S^N$  through the “pole”  $u \in S^N$ .

Finally we have the following commutative diagram

$$\begin{array}{ccc} & TS^N & \\ \Phi \nearrow & & \searrow \Psi \\ S^N \times S^N / \Delta & \xrightarrow{\Psi \circ \Phi} & Q \end{array} \quad (18)$$

**Remark 4.** The original  $S^N \subset Q$  was located at  $S^N = Q \cap \mathbb{R}^{N+1}$ . If we follow this through, we see that this  $S^N$  corresponds to the zero section  $\{(x, 0) \in TS^N\}$  inside  $TS^N$  and consequently is defined by  $\Phi(u, v) = (u, 0)$  inside  $S^N \times S^N / \Delta$ . Note that this is precisely the anti-diagonal

$$\bar{\Delta} = \{(u, -u) \mid u \in S^N\} \subset S^N \times S^N \quad (19)$$

Indeed, the stereographic projection of  $-u$  (“south pole”) through  $u$  (“north pole”) gives zero:

$$-u \mapsto \frac{(-u) - \langle (-u), u \rangle u}{1 - \langle (-u), u \rangle} = \frac{-u + \|u\|^2 u}{1 + \|u\|^2} = 0 \quad (20)$$

since  $\|u\|^2 = 1$ . Hence

$$\Phi(u, -u) = (u, 0) \quad (21)$$

and the original  $S^N$  is located along the anti-diagonal  $\bar{\Delta}$ .

### 1.3 Factor from $\delta$ -function Constraint

Recall that we are interested in integrals of the form

$$I = \int_{\mathbb{CP}^n \times \mathbb{CP}^n} d\mu(z) d\mu(w) e^{-S(z, w)} \mathcal{O}(z, w) \quad (22)$$

where  $d\mu(z) = \prod_i dz_i d\bar{z}_i$  and we have included some observable  $\mathcal{O}: \mathbb{CP}^n \rightarrow \mathbb{C}$ . For our toy model, we consider the action functional (2)

$$S(z, w) = -\beta |\langle \bar{z}, w \rangle|^2 \quad (23)$$

It is instructive pass to real coordinates. If  $z_k = a_k + ib_k$ , let  $x_k$  be the real vector  $(a_k \ b_k)^t$ . If we collect all components, we can form the real vector  $x = (\dots a_k \ b_k \dots)^t \in \mathbb{R}^{2n+2}$ . We denote the real vector associated to  $w$  by  $y$ . In terms of these, the path integral over can be written as

$$vol(S^1)^2 \int_{\mathbb{R}^{2n+2} \times \mathbb{R}^{2n+2}} d^{2n+2}x \ d^{2n+2}y \ \delta(\|x\|^2 - 1) \delta(\|y\|^2 - 1) e^{-S(x, y)} \mathcal{O}(x, y) \quad (24)$$

where the pre-factor stems from the compact part of the gauge group  $\mathbb{C}^* \mathbb{C}^{n+1}$ . Let us focus on the following part of the integral:

$$I \sim \int_{\mathbb{R}^{2n+2}} \delta(\|x\|^2 - 1) = \int_{\mathbb{R}^{2n+2}} \delta(\langle x, x \rangle - 1) \quad (25)$$

Following our idea, we complexify the space of integration,  $x \rightarrow \zeta$

$$I \sim \int_{\Gamma_0 = \mathbb{R}^{2n+2} \cap \mathbb{C}^{2n+2}} \delta(\langle \zeta, \zeta \rangle - 1) \quad (26)$$

where

$$\langle \zeta, \zeta \rangle = \sum_k \zeta_k^2 \quad (27)$$

The  $\delta$ -function constraint restricts the support of the integrand to the quadric  $Q$ . Now, we would like to deform the “contour” (domain of integration) inside  $Q$

$$I \sim \int_{\Gamma_a \subset \mathbb{C}^{2n+2}} \delta(\langle \zeta, \zeta \rangle - 1) \quad (28)$$

Inspired by the discussion in 1.2.1, we would like to parametrise  $Q$  in terms of  $TS^{2n+1}$ . We thus choose a parametrisation of the following form: (by abuse of notation, we will use again  $x$  to denote a vector in  $\mathbb{R}^{2n+2}$ )

$$\zeta(x) = x \sqrt{1 + \|y(x)\|^2} + iy(x) \quad (29) \quad \{\text{eq:parametrisation}\}$$

where  $y(x)$  is chosen in such a way that  $\langle x, y \rangle = 0$ . It follows that

$$I \sim \int_{\mathbb{R}^{2n+2}} d^{2n+2}x \det J(x) \delta(\langle \zeta(x), \zeta(x) \rangle - 1) \quad (30)$$

where  $J(x)$  denotes the Jacobian of (29). Note that the  $\delta$ -function constraint can be simplified as follows: Let

$$\lambda(x) = \sqrt{1 + \|y\|^2} \quad (31)$$

Then

$$\begin{aligned} C(x) &= \langle \zeta(x), \zeta(x) \rangle - 1 \\ &= \lambda^2(x) \|x\|^2 + 2i\lambda(x) \langle x, y(x) \rangle - \|y\|^2 - 1 \\ &= \lambda^2(x) (\|x\|^2 - 1) \end{aligned} \quad (32)$$

Importantly,

$$\lambda^2(x) > 0 \quad (33)$$

such that the  $\delta$ -function constraint is of the following form:

$$\begin{aligned} \int \delta(C(x)) &= \int d\mu(x) \delta(\underbrace{f(x)g(x)}_{\equiv \phi(x)}) \\ &= \int_{\phi^{-1}(0)} \frac{d\sigma}{\|\nabla \phi(x)\|} \\ &= \int_{f^{-1}(0)} \frac{d\sigma}{\|f(x)\nabla g(x) + g(x)\nabla f(x)\|} \\ &= \int_{f^{-1}(x)} \frac{d\sigma}{\|\nabla f(x)\|} \frac{1}{|g(x)|} \\ &= \int d\mu \frac{\delta(f(x))}{|g(x)|} \end{aligned} \quad (34)$$

where  $f(x) = \|x\|^2 - 1$  and  $g(x) = \lambda^2(x) > 0$ . Hence we find that with the chosen parametrisation, the integral nicely localises to an integral over  $S^{2n+1}$

$$\int_{\mathbb{R}^{2n+2}} \delta(C(x)) = \int_{\mathbb{R}^{2n+2}} \frac{\delta(\|x\|^2 - 1)}{\lambda^2(x)} = \int_{S^{2n+1}} \frac{1}{\lambda^2(x)} \quad (35)$$

Therefore, we are left with an integral of the form

$$I \sim \int_{S^{2n+1}} \frac{\det J(x)}{\lambda^2(x)} \quad , \quad \lambda^2(x) = 1 + \|y(x)\|^2 \quad (36) \quad \{\text{eq:toy\_schema\_I}\}$$

## 1.4 Homogeneous Deformations

### 1.4.1 The General Case

The  $n = 0$  example can be nicely generalised. Tej suggested to consider deformations of the form

$$Y_a(X) = \Omega(a)X \quad , \quad \Omega(a) = \text{diag} \left( \begin{pmatrix} 0 & -a_1 \\ a_1 & 0 \end{pmatrix}, \dots, \begin{pmatrix} 0 & -a_{n+1} \\ a_{n+1} & 0 \end{pmatrix} \right) \quad (37) \quad \{\text{eq:Tej\_deformation}\}$$

for some  $a = (a_1, \dots, a_{n+1}) \in \mathbb{R}^{n+1}$ . This type of deformation have a very nice geometric origin, as we now explain.

**Remark 5.** Under the isomorphism

$$\mathbb{R}^{2n+2} \cong \mathbb{C}^{n+1} \quad , \quad \begin{pmatrix} x_k \\ y_k \end{pmatrix} \mapsto z_k = x_k + iy_k \quad (38)$$

the deformation (37) is acting on  $\mathbb{C}^{n+1}$  by multiplication with the diagonal matrix  $\text{diag}(ia_1, \dots, ia_{n+1})$

$$\begin{pmatrix} \vdots \\ z_k \\ \vdots \end{pmatrix} \mapsto \begin{pmatrix} \vdots \\ ia_k z_k \\ \vdots \end{pmatrix} \quad (39)$$

The finite diffeomorphism generated by this vector field is given by

$$\begin{pmatrix} \vdots \\ z_k \\ \vdots \end{pmatrix} \mapsto \begin{pmatrix} \vdots \\ e^{ia_k} z_k \\ \vdots \end{pmatrix} \quad (40)$$

The generated finite action is thus the action of the torus

$$U(1)^{n+1} \subset \text{SU}(n+1) \quad (41)$$

Let  $G$  be a Lie group and  $H < G$  a subgroup with Lie algebras  $\mathfrak{g}$  and  $\mathfrak{h}$  respectively. The orbit space  $G/H$  is known as a *homogeneous space*. It is called *reductive* if  $\mathfrak{g}$  allows the following decomposition

$$\mathfrak{g} = \mathfrak{h} \oplus \mathfrak{m} \quad , \quad [\mathfrak{h}, \mathfrak{h}] \subset \mathfrak{h} \quad , \quad [\mathfrak{h}, \mathfrak{m}] \subset \mathfrak{m} \quad (42)$$

If  $\mathfrak{g}$  is equipped with a Killing form  $B$  (or if  $G$  allows a  $G$ -invariant metric as a manifold), then  $G/H$  is automatically reductive with the choice

$$\mathfrak{m} = \mathfrak{h}^\perp \quad (43)$$

where  $\mathfrak{h}^\perp$  is the orthogonal to  $\mathfrak{h}$  with respect to  $B$ . Indeed, since  $B$  is  $G$ -invariant, schematically

$$0 = B(\mathfrak{h}, \mathfrak{m}) = B([\mathfrak{h}, \mathfrak{h}], \mathfrak{m}) = B(\mathfrak{h}, [\mathfrak{h}, \mathfrak{m}]) \quad (44)$$

so that  $[\mathfrak{h}, \mathfrak{m}] \subset \mathfrak{m}$ .

In summary, given a reductive homogeneous space  $G/H$ , we get a nice decomposition of the tangent space: first, let  $o = [e]$  be the class in  $G/H$  at the identity. Then

$$T_o(G/H) \cong \mathfrak{m} \quad (45)$$

This isomorphism is given essentially by the pushforward of the natural projection  $\pi: G \rightarrow G/H$ . In fact,  $\pi: G \rightarrow G/H$  defines a principal  $H$ -bundle. Its vertical vector fields are given by  $\ker(\pi_*) = \mathfrak{h}$  and since  $\pi$  is a submersion ( $\pi$  and  $\pi_*$  are both surjective) we have

$$\pi_*: \mathfrak{g} / \ker(\pi_*) = \mathfrak{g} / \mathfrak{h} \cong T_o(G/H) \quad (46)$$

Then, to compute the tangent space at any other point, we use translation by  $G$ .

Now, recall that

$$S^{2n+1} = \mathrm{SU}(n+1) / \mathrm{SU}(n) \quad (47)$$

is a homogeneous space where  $H = \mathrm{SU}(n)$  sits inside  $G = \mathrm{SU}(n+1)$  as the lower right corner

$$\mathrm{SU}(n) \ni h \hookrightarrow \begin{pmatrix} 1 & 0 \\ 0 & h \end{pmatrix} \in \mathrm{SU}(n+1) \quad (48)$$

The reductive split is given by

$$\mathfrak{su}(n+1) = \mathfrak{su}(n) \oplus \mathfrak{m} \quad (49) \quad \{\text{eq:red\_split}\}$$

Since  $\mathfrak{su}(n+1)$  admits a Killing form  $B(x, y) \propto \mathrm{tr}(x, y)$ , we can choose  $\mathfrak{m} = \mathfrak{su}(n)^\perp$ . In the split (49), we have

$$\mathfrak{su}(n) \equiv \begin{pmatrix} 0 & 0 \\ 0 & \mathfrak{su}(n) \end{pmatrix} \subset \mathfrak{su}(n+1) \quad (50)$$



We then find

$$\mathfrak{m} = \mathfrak{su}(n)^\perp = \left\{ \begin{pmatrix} 0 & -\zeta^* \\ \zeta & 0 \end{pmatrix} \right\}, \quad \zeta \in \mathbb{C}^{n+1} \quad (51)$$

Here  $\zeta^* = \bar{\zeta}^t$  denotes the hermitian conjugate.

In order to describe the tangent space  $T_{[g]}(G/H)$  at any point  $[g] \in G/H$ , we can use the tangential map of left translations. In fact, left translation by  $g$  is defined by the map

$$L_g: G \rightarrow G, \quad g' \mapsto gg' \quad (52)$$

so that its tangential map

$$\theta_g = (L_{g^{-1}})_*: T_g G \rightarrow T_e G = \mathfrak{g} \quad (53)$$

This map defines a  $\mathfrak{g}$ -valued 1-form on  $G$  known as the left-invariant Maurer-Cartan form. For a matrix group, it can be written as

$$\theta_g = g^{-1} dg \quad (54)$$

Note that the natural projection  $\pi: G \rightarrow G/H$  is equivariant with respect to left translation: if

$$\tau_g: G/H \rightarrow G/H, \quad [g'] \mapsto [gg'] \quad (55)$$

then we have

$$\pi \circ L_g = \tau_g \circ \pi, \quad \tau_g \circ \tau_{g'} = \tau_{gg'} \quad (56)$$

Moreover,

$$(\pi \circ L_{g^{-1}})_* = \pi_* \circ (L_{g^{-1}})_* = (\tau_{g^{-1}} \circ \pi)_* = (\tau_{g^{-1}})_* \circ \pi_* \quad (57)$$

Hence, denoting  $(\tau_{g^{-1}})_* = \vartheta_g$ , we have

$$\pi_* \circ \theta_g = \vartheta_g \circ \pi_*: T_g G \rightarrow T_o(G/H) = \mathfrak{m} \quad (58)$$

For us, we will be mainly interested in the map

$$(\tau_g)_*: \mathfrak{m} = T_{[e]}(G/H) \rightarrow T_{[g]}(G/H) \quad (59)$$

which parametrises the tangent space of  $G/H$  at a point  $[g]$  by  $\mathfrak{m}$ .

Now, if  $G$  is a matrix group the pushforward  $(\tau_g)_*$  is essentially just given by left multiplication with  $g$  (matrix multiplication is linear) and so in we could describe the tangent space  $T_p(G/H)$  at a point  $p = g(p)p_0$  as follows

$$T_p(G/H) \cong \{(p, g(p)\xi) \mid \xi \in \mathfrak{m}\} \quad (60)$$

where we implicitly identify  $g(p)X$  with  $g(p)Xp_0$ .

**Remark 6.** In theory, this is very nice. In practice, I believe it is computationally expensive: for a given  $p \in S^{2n+1} = \text{SU}(n+1)/\text{SU}(n)$ , we would need to find  $g(p)$  which on top of it is only defined up to a multiplication on the right of  $\text{SU}(n)$ . We could, for example, choose  $p_0 = (1, 0 \dots)^t$ . There is a way how to algorithmically compute  $g(p)$ : A matrix belongs to  $\text{SU}(n+1)$  iff its columns (rows) form an orthonormal basis of  $\mathbb{C}^{n+1}$ . If  $p_0 = e_1$  is the first standard basis vector, then we start with the basis  $\{p, e_2, \dots, e_{n+1}\}$  and by Gram-Schmidt produce an orthonormal basis  $\{p, \tilde{e}_1, \dots, \tilde{e}_n\}$  and set

$$g(p) = (p \ \tilde{e}_2 \ \dots; \tilde{e}_n) \quad (61)$$

However, I believe that Gram-Schmidt is computationally expensive.

Moreover, it might simply not be necessary. In order to define interesting deformations, we could simply consider any  $X \in \mathfrak{g}$ . Any such  $\xi$  will admit a split into  $\xi = \xi_{\mathfrak{m}} + \xi_{\mathfrak{h}}$  and since  $\mathfrak{h} = \text{Lie}(H) \cong \text{Stab}(p)$  is in the Lie algebra of the stablizer of  $p$  (note that we actually mean  $p$  here, not  $p_0$ ),  $\xi_{\mathfrak{h}} p = 0$  (as  $h p = p$  for any  $h$  in  $H$ ). Thus, we could simply not care and just work with  $X$  as a whole and define deformations for any  $\xi \in \mathfrak{g}$  by

$$Y_{\xi}(X) = \rho(\xi) \cdot X \quad , \quad \xi \in \mathfrak{g} \quad , \quad X \in S^{2n+1} \quad (62)$$

Here,  $\rho$  denotes the appropriate representation of  $\mathfrak{g}$  on  $S^{2n+1}$ . I believe its easiest description is as follows:  $\text{SU}(n+1)$  and hence  $\mathfrak{su}(n+1)$  naturally acts (by matrix multiplication) on  $S^{2n+1}$  when viewed as a subset of  $\mathbb{C}^{n+1}$ . Under the isomorphism  $\mathbb{C} \cong \mathbb{R}^2$ , where  $z = x + iy \mapsto (x, y)$ , multiplication by a complex number  $a + ib$  becomes matrix multiplication:

$$(a + ib)z \iff \begin{pmatrix} a & -b \\ b & a \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} \quad (63)$$

Hence, using our representation of  $\mathbb{C}^{n+1}$

$$\begin{pmatrix} \vdots \\ z_k = x_k + iy_k \\ \vdots \end{pmatrix} \mapsto \begin{pmatrix} \vdots \\ x_k \\ y_k \\ \vdots \end{pmatrix} \quad (64)$$

we have to simply replace any complex number  $a + ib$  in  $\xi \in \mathfrak{su}(n+1)$  by the appropriate matrix:

$$\rho(\xi): a + ib \rightsquigarrow \begin{pmatrix} a & -b \\ b & a \end{pmatrix} \quad (65)$$

In practice, it might be easier to have the following workflow:

$$\begin{array}{ccccccc} S^{2n+1} \subset \mathbb{R}^{2n+2} & \longrightarrow & \mathbb{C}^{n+1} & \xrightarrow{\xi} & \mathbb{C}^{n+1} & \longrightarrow & \mathbb{R}^{2n+2} \\ X & \longrightarrow & Z & \xrightarrow{\xi} & \xi \cdot Z & \longrightarrow & Y_{\xi}(X) \end{array} \quad (66)$$

## 1.5 Example: $n = 0$

### 1.5.1 Homogeneous deformation

A particular nice example, which can be computed explicitly and serves as a check that everything works nicely is given for the case  $n = 0$ . Consider the deformation

$$y(x) = \alpha \Omega x \quad , \quad \|x\|^2 = 1 \quad , \quad \Omega = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \in \text{SO}(2) \quad , \quad \alpha \in \mathbb{R} \quad (67) \quad \{\text{eq:toy\_cst\_def}\}$$

so that indeed

$$\langle x, y(x) \rangle = 0 \quad (68)$$

Moreover, note that

$$\|y(x)\|^2 = \langle y(x), y(x) \rangle = \alpha^2 \langle x, \Omega^t \Omega x \rangle = \alpha^2 \|x\|^2 = \alpha^2 \quad (69)$$

and hence

$$\lambda^2(x) = 1 + \alpha^2 \quad (70)$$

We compute the Jacobian of this parametrisation as follows: let  $M_i^j = \partial_i y^j(x)$ . Then

$$J_i^j = \lambda(x) \delta_i^j + \frac{\langle y, \partial_i y \rangle x^j}{\lambda(x)} + i M_i^j = \lambda(x) \delta_i^j + \frac{M_{ik} y^k x^j}{\lambda(x)} + i M_i^j \quad (71)$$

Using bra-ket notation and trivially lowering indices (considering the flat metric on  $\mathbb{R}^2$ ), we can write  $J$  as follows:

$$J = \lambda(x) \cdot id + \lambda^{-1}(x) |My\rangle \langle x| + iM \quad (72)$$

**Remark 7.** The matrix  $M$  has some interesting properties.

1. From  $\langle x, y(x) \rangle = 0$  it follows that

$$0 = \partial_i \langle x, y(x) \rangle = y_i(x) + M_{ik} x_k$$

so that

$$y_i(x) = M_{ik} x_k \implies y(x) = Mx$$

and hence

$$\langle x, y \rangle = \langle x, Mx \rangle = \langle M^t x, x \rangle = 0$$

from which follows (if  $y(x) = Mx \neq 0$ )

$$\langle x | M = 0$$

Now, for our choice of

$$y(x) = \alpha \Omega x$$

we have

$$M = \alpha \Omega^t = -\alpha \Omega \quad (73)$$

and hence

$$\begin{aligned} J &= \lambda(x) \cdot id + \lambda^{-1}(x) \alpha^2 \Omega^t \Omega |x\rangle \langle x| - i\alpha \Omega \\ &= \lambda(x) \cdot id + \lambda^{-1}(x) \alpha^2 |x\rangle \langle x| - i\alpha \Omega \end{aligned} \quad (74)$$

Now, here is a neat trick due to Tej. Consider the unitary matrix

$$U = (x \ \Omega x) \quad , \quad \Omega^* = \begin{pmatrix} x^t \\ x^t \Omega^t \end{pmatrix} \quad (75)$$

where  $x$  is a  $2 \times 2$  column vector with unit norm. Then

$$U^* U = \begin{pmatrix} \langle x, x \rangle & \langle x, \Omega x \rangle \\ \langle \Omega x, x \rangle & \langle x, x \rangle \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \quad (76)$$

Moreover,

$$\langle x | U = x^t U = (\langle x, x \rangle \ \langle x, \Omega x \rangle) = (1 \ 0) \quad (77)$$

as well as

$$U^* \Omega U = U^* (\Omega x \ \Omega^2 x) = U^* (\Omega x \ -x) = \begin{pmatrix} \langle x, \Omega x \rangle & -\langle x, x \rangle \\ \langle x, \Omega^t \Omega x \rangle & -\langle x, \Omega^t x \rangle \end{pmatrix} = -\Omega \quad (78)$$

Then

$$\begin{aligned} U^* J U &= \lambda(x) \cdot id + \lambda^{-1}(x) \alpha^2 U^* |x\rangle \langle x| U - i\alpha U^* \Omega U \\ &= \lambda(x) \cdot id + \lambda^{-1}(x) \alpha^2 U^* \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} + i\alpha \Omega \\ &= \begin{pmatrix} \lambda(x) + \lambda^{-1}(x) \alpha^2 & i\alpha \\ -i\alpha & \lambda(x) \end{pmatrix} \end{aligned} \quad (79)$$

It follows that

$$\begin{aligned} \det J(x) &= \det(U^* J(x) U) \\ &= \det \begin{pmatrix} \lambda(x) + \lambda^{-1}(x) \alpha^2 & i\alpha \\ -i\alpha & \lambda(x) \end{pmatrix} \\ &= \lambda^2(x) + \alpha^2 - \alpha^2 = \lambda^2(x) \end{aligned} \quad (80)$$

Finally, we obtain that

$$I \sim \int_{S^{2n+1}} \frac{\det J(x)}{\lambda^2(x)} = \int_{S^{2n+1}} \frac{\lambda^2(x)}{\lambda^2(x)} = \int_{S^{2n+1}} \quad (81)$$

This means that the *constant deformation* (67) comes with a trivial *total* Jacobian (Jacobian + factor from the  $\varphi$ -function constraint).

### 1.5.2 What it really means

Recall that for  $n = 0$  we consider the path integral over  $S^1$ . Before, we have considered  $x \in S^1 \subset \mathbb{R}^2$ . Now, we might think of its complex formulation: let us consider  $z = e^{i\theta} \in S^1 \subset \mathbb{C}$ . In [?], the complexification was chosen in terms of  $\theta \rightarrow \theta + i\tau$ . This means that

$$z \mapsto \tilde{z} = e^{i(\theta+i\tau)} = e^{-\tau} z \quad (82)$$

On the other hand, we can parametrize  $x$  by

$$x = \begin{pmatrix} \sin(\theta) \\ \cos(\theta) \end{pmatrix} \quad (83)$$

so that

$$\begin{aligned} \tilde{z} &= \begin{pmatrix} \sin(\theta + i\tau) \\ \cos(\theta + i\tau) \end{pmatrix} = \begin{pmatrix} \sin(\theta) \cosh(\tau) + i \cos(\theta) \sinh(\tau) \\ \cos(\theta) \cosh(\tau) - i \sin(\theta) \sinh(\tau) \end{pmatrix} \\ &= \begin{pmatrix} \cosh(\tau) & i \sinh(\tau) \\ -i \sinh(\tau) & \cosh(\tau) \end{pmatrix} \begin{pmatrix} \sin(\theta) \\ \cos(\theta) \end{pmatrix} = \lambda x + i \underbrace{\Omega x}_{y(x)} \end{aligned} \quad (84)$$

where

$$\Omega = \begin{pmatrix} 0 & -\sinh(\tau) \\ \sinh(\tau) & 0 \end{pmatrix} \quad (85)$$

and

$$\lambda^2 = 1 + \|y(x)\|^2 = 1 + \sinh^2(\tau) = \cosh^2(\tau) \quad (86)$$

To summarise, if we consider the homogeneous deformation (67),

$$y(x) = \alpha \Omega x \quad , \quad \|x\|^2 = 1 \quad , \quad \Omega = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \quad (87)$$

then choosing

$$\boxed{\alpha = \sinh(\tau)} \quad (88)$$

we find that the homogeneous deformation simply scales the complex variable  $z = e^{i\theta} \in S^1 \subset \mathbb{C}$ ,

$$z \rightarrow \tilde{z} = e^{-\tau} z \quad (89)$$

### 1.5.3 Generalisation: Torus deformations

Let us now think about  $n > 0$ . Let  $z = [z_0 : \dots : z_n] \in \mathbb{CP}^n$ . Again, we factor out the  $U(1) \subset \mathbb{C}^*$  and fix the  $\mathbb{R}_+ \subset \mathbb{C}^*$  by restricting the path integral to  $S^{2n+1}$ , i.e.

$$\sum_k |z_k|^2 = 1 \quad (90)$$

Following the logic of the homogeneous deformation, let us consider a deformation parametrised by a generic element in the Cartan subalgebra of  $\mathfrak{su}(n+1)$

$$Y = \Omega(\alpha)X = \sum_{k=1}^n \rho(i\alpha_k H_k) \quad (91)$$

where we choose the following *hermitian* (hence the factor of  $i$ ) basis of the Cartan  $\mathfrak{h} \subset \mathfrak{su}(n+1)$  subalgebra

$$H_k = E_{kk} - E_{(k+1)(k+1)} \quad (92)$$

where  $E_{ij}$  is the  $n \times n$  matrix with  $(E_{ij})_{ab} = \delta_{ai}\delta_{bj}$ . Here,  $\rho$  is again the representation of complex multiplication in terms of  $2 \times 2$  matrices,

$$\rho(a + ib) = \begin{pmatrix} a & -b \\ b & a \end{pmatrix} \quad (93)$$

and defined element-wise on a generic complex matrix as

$$\rho(M) = \begin{pmatrix} \rho(m_{11}) & \rho(m_{12}) & \dots \\ \rho(m_{21}) & \rho(m_{22}) & \dots \\ \vdots & \vdots & \ddots \end{pmatrix} \quad (94)$$

In particular,

$$\rho(iH_k) = \rho(iE_{kk}) - \rho(iE_{(k+1)(k+1)}) = \begin{pmatrix} \ddots & & & & \\ & 0 & -1 & & \\ & 1 & 0 & & \\ & & & 0 & 1 \\ & & & -1 & 0 \\ & & & & \ddots \end{pmatrix} \quad (95)$$

Let us choose

$$\alpha_k = \sinh(\tau_k) \quad , \quad \tau_0 \equiv \tau_{n+1} \equiv 0 \quad (96)$$

then, acting on the complex vector  $z = (z_0, \dots, z_n)$ , we find

$$\boxed{\tilde{z} = (\tilde{z}_0, \dots, \tilde{z}_n) = (e^{-\Delta_{01}} z_0, \dots, e^{-\Delta_{nn+1}} z_n) \quad , \quad \Delta_k = \tau_k - \tau_{k+1}} \quad (97)$$

## 1.6 Operators

### 1.6.1 One-point function

It is time to look at interesting operators we would like to compute. We start with the one-point function. Let

$$\mathcal{O}_{ij}(z) = z_i \bar{z}_j \quad (98)$$

Clearly,  $\mathcal{O}_{ij}$  is invariant under the  $U(1)$  action  $z_k \rightarrow e^{i\varphi} z_k$  for all  $k$ .

It's expectation value is given by

$$\langle \mathcal{O}_{ij}(z) \rangle = \frac{1}{(2\pi)^2} \int_{S^{2n+1} \times S^{2n+1}} \prod_k d^n z_k d^n w_k z_i \bar{z}_j e^{-S(z,w)} \quad (99) \quad \{\text{eq:toy\_one\_pt}\}$$

where

$$S(z, w) = -\beta |\langle \bar{z}, w \rangle|^2 = -\beta \sum_{k,\ell} \bar{z}_\ell z_k w_\ell \bar{w}_k \quad (100)$$

Note that both, the action and the measure is invariant under the diagonal  $U(1)$  action

$$(z_k, w_k) \mapsto (e^{i\alpha} z_k, e^{i\alpha} w_k) \quad , \quad \forall k \quad (101)$$

By this change of coordinates, for  $k = i$ , it follows that

$$\langle \mathcal{O}_{ij}(z) \rangle = e^{-i\alpha} \langle \mathcal{O}_{ij}(z) \rangle \quad (102)$$

for any  $\alpha$  and hence we must have

$$\boxed{\langle \mathcal{O}_{ij}(z) \rangle = 0 \quad , \quad \forall z, i, j} \quad (103) \quad \{\text{eq:toy\_var\_one\_pt}\}$$

Its variance is given by

$$\begin{aligned} \text{Var}(\mathcal{O}_{ij}(z)) &= \langle |\mathcal{O}_{ij}(z)|^2 \rangle - |\langle \mathcal{O}_{ij}(z) \rangle|^2 \\ &= \langle \mathcal{O}_{ij}(z) \mathcal{O}_{ij}^*(z) \rangle - \langle \mathcal{O}_{ij}(z) \rangle \langle \mathcal{O}_{ij}^*(z) \rangle \end{aligned} \quad (104)$$

so that by (103)

$$\text{Var}(\mathcal{O}_{ij}(z)) = \langle |\mathcal{O}_{ij}(z)|^2 \rangle = \langle |z_i|^2 |z_j|^2 \rangle \quad (105)$$

## 1.7 Two-point function

## References

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