

Electric Drive Fundamentals

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Version: 0.1.0

MCI



B.Sc - Drive Systems



1. Introduction
2. Return to Fundamentals
3. Magnetic Circuits and Materials
4. Appendix

Introduction



First Steps

Introduction

Individual Assignment

Lecture Contents

Final Examination

Point Distribution

Point Distribution

Resources



- The goal of this lecture is to give you the fundamentals of electric machinery using theory and engineering practice.
- This lecture is a total of 2 SWS with a total of thirty (30) UE.
 - With 28 UE is devoted to lectures.
- There is a written exam at the end of the module worth two (2) UE.
- There is one (1) assignment for this course:
1st will be a pre-defined work which is individual based.



- The individual assignments focus on understanding electric machinery.
- The assignment is uploaded to SAKAI for you to work on along with what is required of you for submission.
 - The assignments contain electric drive problems to solve.
- The strict deadline is the end day of the **last lecture**.
- Any submission after this date will not be accepted.
- If all submissions are sent early we can do a statistical analysis and go through questions before the final exam.



- Lecture materials and all possible supplements will be present in its Github Repo.
 - You can easily access the link to the web-page from [here](#).

Github is chosen for easy access to material management and CI/CD capabilities and allowing hosting websites.

- In the lecture some exercises are solved using Python and other examples and can be accessed from the [Repo website](#).



- At the end of the lectures there will be a final examination which you will be tested.
- You will be asked three (3) questions related to electric drives.
- The exam will be ninety (90) minutes.
- You are NOT allowed a personal formula sheet or any kind of supporting material.
- You are allowed a calculator.



Assessment Type	Overall Points	Breakdown	%
Individual Assignment	40		
		Report	20
		Solution(s)	80
			20
Final Exam	60		
		Question 1	40
		Question 2	30
		Question 2	30

Table 1: Assessment Grade breakdown for the lecture.



Covered Topic	Appointment
Return to Fundamentals	1
Magnetic Circuits & Materials	1
Transformers	2
Electromechanical Energy Conversion	2-3
Rotating Magnetic Fields	3
DC Drives	4
Poly-phase Induction Drives	4-5
Single-phase Induction Drives	5-6
Linear Induction Drives	6
Poly-phase Synchronous Drives	6-7
Solid-state Commutation Drives	7

Table 2: Distribution of materials across the semester.



Return to Fundamentals

- Complex number notation,
- Multi-phase systems,
- Phasors and Wave-forms





Transformers

- Construction & Physical Properties
- Modelling
- Parameter Tests
- Connection Types





Rotating Magnetic Fields

- MMF in Winding
- Torque Generation
- Generated Voltage





DC Drives

- Construction
- Physical laws governing its operation
- Types of Connections used in industry and commercial applications.
- Applications in Industrial/Commercial venues.





Poly-phase Induction Drives

- Construction
- Physical laws governing its operation principle.
- Modelling an induction drive mathematically
- Methods of starting an induction drive.





Single-Phase Induction Drives

- Creating a Rotating Magnetic Field in a single phase
- Types of single-phase types
- Salient-pole





Poly-phase Synchronous Drives

- Construction and Rotor types,
- Operation principles,
- Regulatory behaviours,
- V-curves,





Solid-State Commutation Drives

- Solid State Commutation
- BLDC & PMSM Drives
- Switched Reluctance Motor (SRM)
- Stepper Drives





Books

- Mohan Ned. "*Advanced electric drives: analysis control and modeling using MATLAB/Simulink*" John Wiley & Sons 2014.
- Krause Paul C. et. al. "*Analysis of electric machinery and drive systems*" Vol. 2 IEEE Press 2002.
- Pyrhonen Juha et. al. "*Design of rotating electrical machines*" John Wiley & Sons 2013.
- Stephen J. Chapman. "*Electric Machinery Fundamentals (5th Edition)*" (2012).
- Fitzgerald A. E. et. al. "*Electric Machinery*" McGraw Hill (2003).



Books

- Hughes A. et. al. "*Electric Motors and Drives: Fundamentals Types and Applications*" Newnes 2019.
- Melkebeek A. "*Electrical Machines and Drives: Fundamentals and Advanced Modelling*" Springer 2018.
- Wildi T. "*Electrical machines, drives, and power systems*" Pearson Education 2006.
- Veltman A. et. al. "*Fundamentals of Electrical Drives*" Springer 2007.



White Papers

- Maddox Transformer "*Guide to transformer cores: types, construction, & purpose*"
- Control Engineering *Springtime for Switched-Reluctance Motors?* .



Lecture Notes

- Power Transformers "*ESE 470 Energy Distribution Systems*" Oregon State University,
- Principles of Electromechanical Energy Conversion "*Actuators & Sensors in Mechatronics Electromechanical Motion Fundamentals*" NYU,

Return to Fundamentals



Learning Outcomes

Introduction

Connection Types

Three-phase Waveform

Delta Connection

Wye Connection

Polar Coordinates

Power in AC



Learning Outcomes

- (LO1) An Overview of Poly-phase circuits,
- (LO2) Definitions on Active-Reactive power,
- (LO3) Polar Coordinate System.





- A rotating magnetic field is a magnetic field with **moving polarities**.
 - Which its opposite poles rotate about a **central point** or **axis**.
- To create a rotating magnetic field (RMF) you need at least a **2-phase** system.
- In the industry, RMFs are mostly produced using a 3-phase supply ¹.
 - There are also economic reasons, as it is cheaper to design a 3-phase system with minimal cost to wiring ².
- Before starting with electric drives, it is a good time to look at some **fundamental concepts** in power engineering.

¹3-phase supply produces smoother operation of motors compared to 2-phase. This is due to the power transfer in 3 phase supply being less pulsating than in 2 phase supply.

²However, there have been test and feasibility on 6-phase power in 1970s with minimal success.



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- In such a system, each voltage source consists of a group of voltages having related **magnitudes** and **phase angles**.
- Thus, an n -phase system employs voltage sources which typically consist of n voltages substantially equal in magnitude and successively displaced by a phase angle of $360^\circ/n$.
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- The three individual voltages of a three-phase source may each be connected to its own independent circuit.
- We would then have three separate single-phase systems.
- Alternatively, symmetrical electric connections can be made between the three voltages and the associated circuitry to form a three-phase system.
- It is the latter alternative that we are concerned.
- Note that the word phase now has two distinct meanings.
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Figure 1: A three phase waveform of currents **R**, **B**, and **Y**.



- There are **no neutral connection** available.
- Phase voltage appears **across the windings**.
- In a delta connection, Line voltage is equal to Phase voltage:

$$V_{\text{line}} = V_{\text{ph}}$$

- $i_{\text{YB}}, i_{\text{YR}}, i_{\text{BR}}$ are also called **Line current** (i.e., i_{line}).
- Line current is $\sqrt{3}$ times that of the phase current.

$$i_{\text{line}} = \sqrt{3} i_{\text{subs}}$$

- Power related definitions are:

$$P = \sqrt{3} v_{\text{line}} i_{\text{line}} \cos \varphi, \quad \text{and} \quad P = 3 v_{\text{ph}} i_{\text{ph}} \cos \varphi.$$



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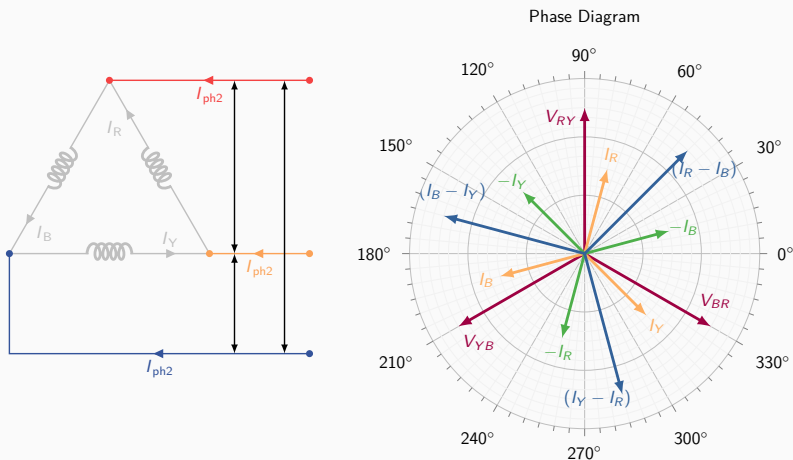


Figure 2: The connection and the phasor diagram of a 3-phase delta connection.



Example

Three impedance with a value $Z_{\Delta} = 12.00 + j9.00 = 15.00 \angle 36.9^{\circ} \Omega$ are connected in Δ .

For balanced line-to-line voltages of 208 V, find the line current, the power factor, and the total power.



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- Phase voltage appear across windings.

$$v_{\text{line}} = \sqrt{3} v_{\text{ph}},$$

- $i_{\text{YB}}, i_{\text{YR}}, i_{\text{BR}}$ are also known as line current (i_{line}).

- In a star connection, **line current** is equal to **phase current**.

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- Power related equations for 3-phase:

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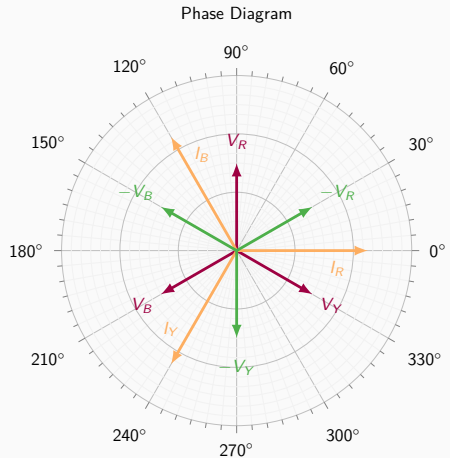
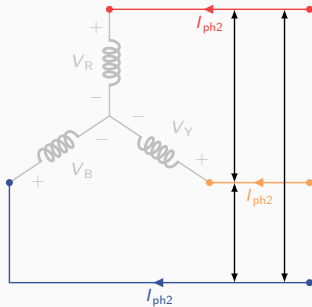


Figure 3: The connection and the phasor diagram of a 3-phase wye connection.



Example

A Y-connected 120 V source feeds a Δ -connected load through a distribution line having an impedance of $0.3 + j0.9 \Omega$. The Y-source impedance is $0.2 + j0.5 \Omega$

The load impedance is $118.5 + j85.8 \Omega/\phi$.

Use the a-phase internal voltage of the generator as the reference.

- (a) Construct a single-phase equivalent circuit of the three-phase system,
- (b) Calculate the line currents,
- (c) Calculate the phase voltages at the load terminals,
- (d) Calculate the phase currents of the load,
- (e) Calculate the line voltages at the source terminals.



$$Z = x + \mathbf{j}y = A\angle\theta,$$

where:

- Z is the complex vector,
- A is vector magnitude,
- x is real/active part,
- \mathbf{j} is defined as $\sqrt{-1}$.
- y is imag/reactive part,
- θ is the complex angle.

$$\theta = y/x.$$

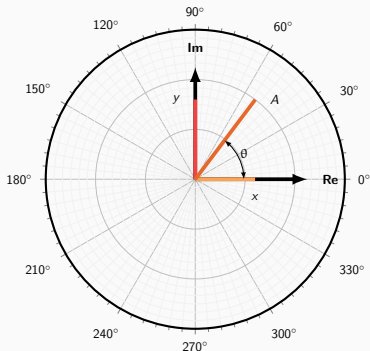


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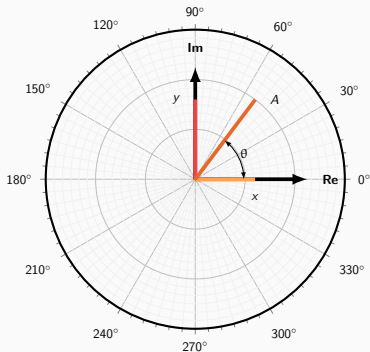


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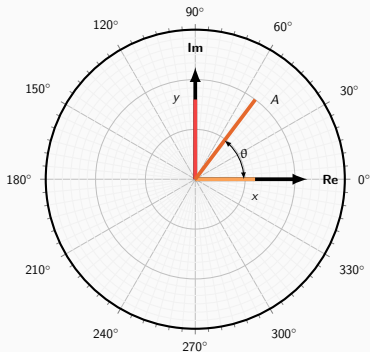


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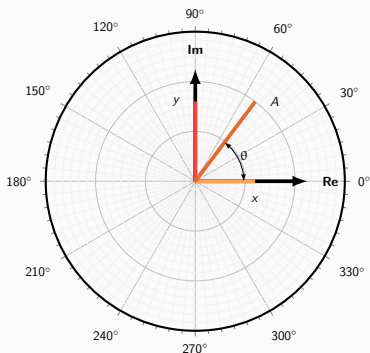


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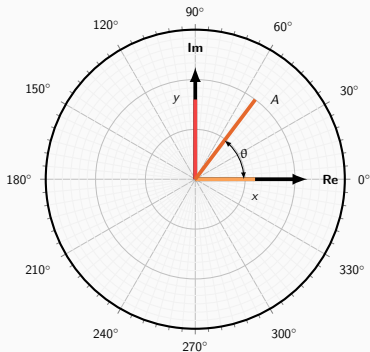


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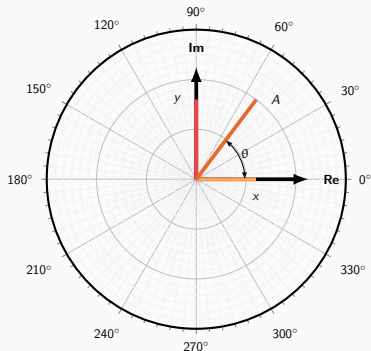


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When it comes to power in AC, there are three (3) definitions:

- If energy is used/generated by an **active** element it is called **Active** or **Real** (P) power and measured in W.
- If energy is used/generated by an **reactive** element it is called **Reactive** power (Q) and measured in $V \cdot A_r$.
- The combination of these two values is called **Apparent** power (S) and measured in $V \cdot A$.

$$S = \sqrt{P^2 + Q^2}$$

- Finally the angle difference between the voltage and current waveform is defined as the **phase** and tells how reactive/active a circuit is.

$$\varphi = \varphi_V - \varphi_I.$$



When it comes to power in AC, there are three (3) definitions:

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Figure 5: An animation showing the relations between Active, Reactive and Apparent power along with phase angle.



Example

A balanced three-phase load requires 480 kW at a lagging power factor of 0.8.

The load is fed from a line having an impedance of $0.005 + j0.025\Omega$.

The line voltage at the terminals of the load is 600 V.

- (a) Construct a single-phase equivalent circuit of the system.
- (b) Calculate the magnitude of the line current.
- (c) Calculate the magnitude of the line voltage at the sending end of the line.
- (d) Calculate the power factor at the sending end of the line.

Magnetic Circuits and Materials



Learning Outcomes

Introduction

Magnetic Materials

Magnetic Circuits

Maxwell's Equations

Flux Linkage, Inductance and Energy

Magnetic Materials

Introduction

Ferromagnetism

B-H Curve

AC Excitation

Hysteresis Losses



Learning Outcomes

- (LO1) An Overview of Maxwell's Equations,
- (LO2) Introduction to Magnetic Circuits,
- (LO3) Brief look on Magnetic Materials,
- (LO4) Magnetic Material Losses.





- Almost all electric drives use **ferromagnetic material** for shaping and directing **B**-fields.
 - These fields act as the medium for transferring and converting energy.

Permanent-magnets are also widely used in drive design.

- Without these materials, practical implementations of most familiar EEC devices would not be possible.
- Analysing and describing systems containing them is essential for designing effective drives.
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- Maxwell's equations are a set of coupled partial differential equations which form the foundation of electric and magnetic circuits.
- In their PDE form, they are written as [5]:

$$\begin{aligned}\nabla \cdot \mathbf{E} &= \frac{\rho}{\varepsilon_0}, \\ \nabla \cdot \mathbf{B} &= 0, \\ \nabla \times \mathbf{E} &= -\frac{\partial \mathbf{B}}{\partial t}, \\ \nabla \times \mathbf{B} &= \mu_0 \left(\mathbf{J} + \varepsilon_0 \frac{\partial \mathbf{E}}{\partial t} \right).\end{aligned}$$

where \mathbf{E} is the electric field, \mathbf{B} is the magnetic field, ρ is the electric charge density and \mathbf{J} the current density, ε_0 is the vacuum permittivity and μ_0 is the vacuum permeability.



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- Solution to Maxwell's equation can be hard if not attainable and usually simplifying assumptions are made to reach practical solutions.
- First, the displacement-current term ($\epsilon_0 \partial \mathbf{E} / \partial t$) can be neglected.
 - This term accounts for \mathbf{B} -fields produced in space by time-varying \mathbf{E} -fields (i.e., electromagnetic radiation) [10].
- Neglecting $\epsilon_0 \partial \mathbf{E} / \partial t$ results in the magneto-static form which relate \mathbf{B} -fields to the currents (\mathbf{J}) which produce them.
- In integral form [2]:

$$\oint_C \mathbf{H} \cdot d\mathbf{l} = \int_S \mathbf{J} \cdot d\mathbf{a}, \quad \text{and} \quad \oint_S \mathbf{B} \cdot d\mathbf{a} = 0.$$



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$$\oint_{\mathcal{C}} \mathbf{H} \, d\mathbf{l} = \int_{\mathcal{S}} \mathbf{J} \cdot d\mathbf{a}, \quad \text{and} \quad \oint_{\mathcal{S}} \mathbf{B} \cdot d\mathbf{a} = 0.$$

The line integral (\oint) of the tangential component of the magnetic field intensity (\mathbf{H}) around a closed contour (\mathcal{C}) is equal to the total current (\mathbf{J}) passing through any surface (\mathcal{S}) linking that contour.

- The second one states the magnetic flux density (\mathbf{B}) is conserved, i.e., no net flux enters or leaves a closed surface (i.e., $\nabla \cdot \mathbf{B} = 0$).

These simplifications have allowed us to remove the effect of the \mathbf{E} field on our calculations.



- A second simplifying assumption involves the concept of **magnetic circuits**.
- The general solution for the \mathbf{H} and the \mathbf{B} in a structure of complex geometry is extremely difficult if not practically impossible.
- However, a 3D field problem can often be reduced to what is essentially a circuit equivalent of magnetic elements,
 - with acceptable engineering accuracy [11].



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- A magnetic circuit is a structure composed of **highly permeable** ($\mu_r \gg 0$) material(s).
- The presence of high-permeability causes magnetic flux to be **confined** to the paths defined by the structure.

This is similar to how currents are confined to the conductors of an electric circuit.



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- The core is assumed to be composed of magnetic material with permeability much greater than of surrounding air (i.e., $\mu_r \gg \mu_0$).
- The core is of **uniform cross section** and is excited by a winding of N turns carrying a current of i amperes.
- This winding produces a **B-field** in the core.

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- Due to high permeability of the material, the magnetic flux is confined almost entirely to the core,
- the field lines follow the path defined by the core,
- and the flux density is essentially uniform over a cross section because the cross-sectional area is uniform.
- The magnetic field can be visualised in terms of flux lines which form closed loops interlinked with the winding.



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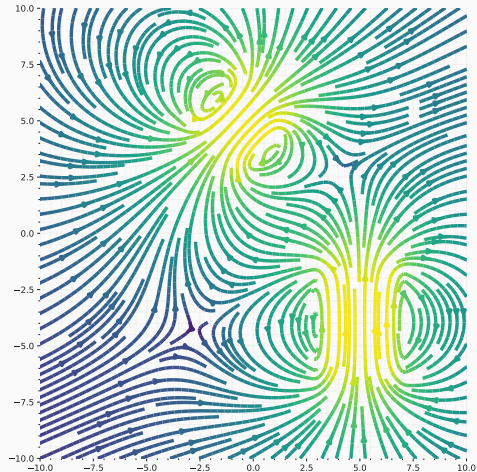


Figure 6: An example demonstration of flux paths passing highly permeable material.



- The source of the magnetic field in the core is the ampere-turn Ni .
- In magnetic circuit terminology Ni is the magnetomotive force (mmf) acting on the magnetic circuit.
- In systems with more than one winding, Ni must be replaced by the **algebraic sum** of the ampere-turns of all the windings.
- The magnetic flux ϕ crossing a surface S is the surface integral of the normal component of \mathbf{B} :

$$\phi = \int_S \mathbf{B} \cdot d\mathbf{a}$$

where the unit of flux (ϕ) is Wb.



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- This equation states that all flux entering the surface enclosing a volume must leave the volume as they form closed loops.
- This is enough to justify the uniformity of the magnetic flux density across the cross section of a magnetic circuit:

$$\phi_c = B_c A_c$$

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- Using the original magneto-static simplifications, we can come up to:

$$\mathcal{F} = Ni = \oint \mathbf{H} d\mathbf{l}.$$

- As the path of the flux line is close to the mean length of the core (l_c), this is simplified to:

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- The relationship between the **H** and **B** is a **property of the material**.
- It is common to assume a linear relationship shown as:

$$\mathbf{B} = \mu \mathbf{H},$$

where μ is the magnetic permeability.

- In SI units, **H** is measured in $\text{A} \cdot \text{m}^{-1}$ and **B** is in $\text{Wb} \cdot \text{m}^{-2}$ or T.



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- Transformers are wound on closed cores.
- However, EEC devices which incorporate a moving element must have **air gaps** in their magnetic circuits.
- When the air-gap (g) is much smaller than the dimensions of the adjacent core faces, the magnetic flux will follow the path defined by the core and the air gap and the techniques of magnetic-circuit analysis can be used.

If g becomes excessively large, the flux will be leak out of the sides of the air gap and the techniques of magnetic-circuit analysis will no longer be strictly applicable.



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- Provided the air-gap length g is sufficiently small, the configuration can be analysed as a magnetic circuit with two (2) series components:
 - a magnetic core of μ , cross-sectional area A_c , and mean length l_c ,
 - an air gap of permeability μ_0 , cross-sectional area A_g , and length l_g .
- In the core the flux density can be assumed uniform; thus

$$B_c = \frac{\phi}{A_c} \quad \text{in the air gap} \quad B_g = \frac{\phi}{A_g}$$



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- Applying this to the magnetic circuit gives:

$$\mathcal{F} = H_c l_c + H_g g,$$

using the **B-H** relationship:

$$\mathcal{F} = \frac{B_c}{\mu} l_c + \frac{B_g}{\mu_0} g.$$

- Here, \mathcal{F} is the mmf applied to the magnetic circuit where $\mathcal{F}_c = H_c l_c$ is the magnetic field in the core and $\mathcal{F}_g = H_g l_g$ is the magnetic field in the air-gap.
- Using ϕ we can re-write this equation as:

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$$\mathcal{F} = \phi \left(\frac{l_c}{\mu A_c} + \frac{g}{\mu_0 A_g} \right).$$



- The term which multiplies mmf to flux is called **Reluctance**.

$$\mathcal{F} = \phi \left(\underbrace{\frac{l_c}{\mu A_c}}_{\mathcal{R}_c} + \underbrace{\frac{g}{\mu_0 A_g}}_{\mathcal{R}_g} \right) \quad \text{or} \quad \mathcal{F} = \phi (\mathcal{R}_c + \mathcal{R}_g)$$

- Reluctance can be seen as analogous to resistance to electrical circuits.



Example

The magnetic circuit shown has dimensions:

$$A_c = A_g = 9 \text{ cm}^2,$$

$$g = 0.050 \text{ cm},$$

$$l_c = 30 \text{ cm},$$

$$N = 500 \text{ Turns}.$$

Assume the value $\mu_0 = 70,000 \text{ H} \cdot \text{m}^{-1}$ for core material.

- (a) Find the reluctance values \mathcal{R}_c and \mathcal{R}_g .
- (b) For the value of $B_c = 1 \text{ T}$, find the flux ϕ and the current i .



Example

The following magnetic structure has infinite permeability ($\mu \rightarrow \infty$). Find the air-gap flux ϕ and flux density B_g .

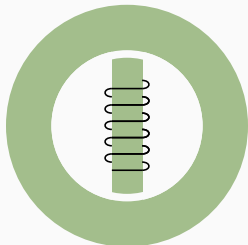


Figure 7: A simple synchronous drive.

Parameters: $I = 10$ A, $N = 1000$ Turns, $g = 1$ cm, and $A_g = 2000$ cm².



- When a **B**-field varies with time, an **E**-field is produced in space as determined by Faraday's law [5]:

$$\oint_C \mathbf{E} \cdot d\mathbf{s} = -\frac{d}{dt} \int_S \mathbf{B} \cdot d\mathbf{a}.$$

The line integral of the electric field intensity **E** around a closed contour **C** is equal to the time rate of change of the magnetic flux linking (i.e. passing through) that contour.

For working with electric drives, the effects of **E** field can be neglected.



- This simplification allows us to cancel the LHS of the equation to just the induced voltage (i.e., electromotive force)

$$\oint_C \mathbf{E} \cdot d\mathbf{s} = -\frac{d}{dt} \int_S \mathbf{B} \cdot d\mathbf{a}.$$

The equation is annotated with a red arrow pointing from the left-hand side towards the right-hand side, labeled with a red \mathcal{E} .

- Whereas on the RHS, the flux is dominated by the core flux (φ).

$$\mathcal{E} = -\frac{d}{dt} \int_S \mathbf{B} \cdot d\mathbf{a}.$$

The equation is annotated with a red arrow pointing from the right-hand side towards the left-hand side, labeled with a red φ .



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- As winding links the core flux N times, our equation is simplified to:

$$\mathcal{E} = N \frac{d\phi}{dt} = \frac{d\lambda}{dt} \quad \text{where} \quad \lambda = N\phi \quad \blacksquare$$

where λ is flux linkage, measured in Weber-turns,



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- For a magnetic circuit composed of magnetic material of constant magnetic permeability or which includes a dominating air gap, the relationship between ϕ and i will be **linear** and we can define the inductance L as:

$$L = \frac{\lambda}{i} \quad \text{which can be written as} \quad L = \frac{N^2}{\mathcal{R}_{\text{tot}}}$$

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Example

The magnetic circuit consists of an N -turn winding on a magnetic-core of infinite permeability with two parallel air gaps of lengths g_1 and g_2 and areas A_1 and A_2 respectively. Find:

- (a) The winding inductance,
- (b) The flux density B_1 in gap 1 when the winding is carrying current i .

Note: Neglect fringing effects at the air gap.



- The mmf (\mathcal{F}) is given by the total ampere-turns.
- The reference directions for the currents (i_1, i_2) have been chosen to produce flux in the same direction.
- The total mmf is therefore:

$$\mathcal{F} = N_1 i_1 + N_2 i_2.$$

- With the reluctance of the core neglected and assuming that $A_c = A_g$, the core flux ϕ is:

$$\phi = (N_1 i_1 + N_2 i_2) \frac{\mu_0 A_c}{g},$$

where ϕ is the **resultant core flux** produced by the total mmf of the two windings.



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- This resultant ϕ determines the operating point of the core material.
- If ϕ is broken up into terms attributable to the individual currents, the resultant flux linkages of coil 1 can be expressed as:

$$\lambda_1 = N_1 \phi = \overbrace{N_1^2 \left(\frac{\mu_0 A_c}{g} \right)}^{L_{11}} i_1 + \overbrace{N_1 N_2 \left(\frac{\mu_0 A_c}{g} \right)}^{L_{12}} i_2,$$

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- However, in many situations of practical interest, the reluctance of the system is dominated by that of an air gap (which is of course linear) and the nonlinear effects of the magnetic material can be ignored.
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 - obtain large magnetic flux densities
 - Requiring relatively low levels of magnetizing force.

Since magnetic forces and energy density increase with increasing flux density, this effect plays a large role in the performance of energy-conversion devices.

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- In a transformer they are used to maximize the coupling between the windings as well as to lower the excitation current required for transformer operation.
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Magnetic Circuits and Materials



1 IA																18 VIIIA																													
1	H Hydrogen 1.0079																														He Helium 4.0026														
2 IIA																																													
3	Li Lithium 6.941	4	Be Beryllium 9.0122														5	B Boron 10.811	6	C Carbon 12.011	7	N Nitrogen 14.007	8	O Oxygen 15.999	9	F Fluorine 18.998	10	Ne Neon 20.180																	
11	Na Sodium 22.990	12	Mg Magnesium 24.305														13	Al Aluminum 26.982	14	Si Silicon 28.086	15	P Phosphorus 30.974	16	S Sulfur 32.065	17	Cl Chlorine 35.453	18	Ar Argon 39.948																	
3 IIA		4 IIB		5 VB		6 VIB		7 VIIB		8 VIIIB		9 VIIIB		10 VIIIB		11 IB		12 IIB		13 IIIA		14 IVA		15 VA		16 VIA		17 VIIA																	
19	K Potassium 39.098	20	Ca Calcium 40.078	21	Sc Scandium 44.956	22	Ti Titanium 47.887	23	V Vanadium 50.942	24	Cr Chromium 51.996	25	Mn Manganese 54.938	26	Fe Iron 55.845	27	Co Cobalt 58.933	28	Ni Nickel 58.693	29	Cu Copper 63.546	30	Zn Zinc 65.39	31	Ga Gallium 69.723	32	Ge Germanium 72.64	33	As Arsenic 74.922	34	Se Selenium 78.96	35	Br Bromine 79.904	36	Kr Krypton 83.8										
37	Rb Rubidium 85.468	38	Sr Strontium 87.62	39	Y Yttrium 88.906	40	Zr Zirconium 91.224	41	Nb Niobium 92.906	42	Mo Molybdenum 95.94	43	Tc Technetium 98	44	Ru Ruthenium 101.07	45	Rh Rhodium 102.91	46	Pd Palladium 106.42	47	Ag Silver 107.87	48	Cd Cadmium 112.41	49	In Indium 114.82	50	Sn Tin 118.71	51	Sb Antimony 121.76	52	Te Tellurium 127.6	53	I Iodine 126.9	54	Xe Xenon 131.29										
55	Cs Cesium 132.91	56	Ba Barium 137.33	57-71 La-Lu Lanthanide		72	Hf Hafnium 178.49	73	Ta Tantalum 180.95	74	W Tungsten 183.84	75	Re Rhenium 186.21	76	Os Osmium 190.23	77	Ir Iridium 192.22	78	Pt Platinum 195.08	79	Au Gold 196.97	80	Hg Mercury 200.59	81	Tl Thallium 204.38	82	Pb Lead 207.2	83	Bi Bismuth 208.98	84	Po Polonium 209	85	At Astatine 210	86	Rn Radon 222										
87	Fr Francium 223	88	Ra Radium 226	89-103 Ac-Lr Actinide		104	Rf Rutherfordium 261	105	Db Dubnium 262	106	Sg Seaborgium 266	107	Bh Bohrium 264	108	Hs Hassium 277	109	Mt Meitnerium 268	110	Ds Darmstadtium 281	111	Rg Roentgenium 280	112	Uub Ununbium 285	113	Uut Ununtrium 284	114	Uuq Ununquadium 289	115	Uup Ununpentium 288	116	Uuh Ununhexium 293	117	Uus Ununseptium 292	118	Uuo Ununoctium 294										
																119	La Lanthanum 138.91	120	Ce Cerium 140.12	121	Pr Praseodymium 140.91	122	Nd Neodymium 144.24	123	Pm Promethium 145	124	Sm Samarium 150.36	125	Eu Europium 151.96	126	Gd Gadolinium 157.25	127	Tb Terbium 158.93	128	Dy Dysprosium 162.50	129	Ho Holmium 164.93	130	Er Erbium 167.26	131	Tm Thulium 168.93	132	Yb Ytterbium 173.04	133	Lu Lutetium 174.97
																134	Ac Actinium 227	135	Th Thorium 232.04	136	Pa Protactinium 231.04	137	U Uranium 238.03	138	Np Neptunium 237	139	Pu Plutonium 244	140	Am Americium 243	141	Cm Curium 247	142	Bk Berkelium 247	143	Cf Californium 251	144	Es Einsteinium 252	145	Fm Fermium 257	146	Md Mendelevium 258	147	No Nobelium 259	148	Lr Lawrencium 262
																2 max Symbol Name																													
																black: natural gray: man-made																													

Figure 8: The periodic table of elements.



- Typically composed of iron (Fe) and alloys of iron with:
 - cobalt, tungsten, nickel, aluminum, and other metals
- These materials are characterised by a wide range of properties,
 - but, the basic phenomena responsible for their properties are common to them all.
- They are found to be composed of a large number of domains, i.e., regions in which the magnetic moments of all the atoms are parallel, giving rise to a net magnetic moment for that domain.
- In an un-magnetised sample of material, the domain magnetic moments are **randomly** oriented, and the net resulting magnetic flux in the material is zero (0).



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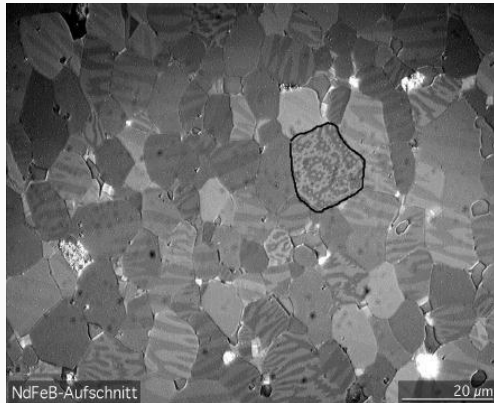


Figure 9: Microcrystalline grains within a piece of $\text{Nd}_2\text{Fe}_{14}\text{B}$ (the alloy used in neodymium magnets) with magnetic domains made visible with a Kerr microscope. The domains are the light and dark stripes visible within each grain. The outlined grain has its magnetocrystalline axis almost vertical, so the domains are seen end-on [4].



- When **B**-field is applied to the material, the magnetic domain moments tend to align with the applied magnetic field.
- As a result, the magnetic moments add to the applied field, producing a much larger value of flux density than would exist due to the magnetizing force alone.
- Therefore the effective permeability μ , equal to the ratio of the total magnetic flux density to the applied magnetic-field intensity, is large compared with the permeability of free space (μ_0).
- As the magnetising force is increased, this behavior continues until all the magnetic moments are aligned with the applied field; at this point they can no longer contribute to increasing the magnetic flux density
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- Without an external **B**-field, magnetic domains naturally align along certain directions associated with the crystal structure of the domain.
 - This is known as axes of easy magnetisation [6].
- Thus if the applied **B** is reduced, the domain magnetic moments relax to the direction of easy magnetism nearest to the applied field.
- As a result, **B** is reduced to zero (0), although they will tend to relax towards their initial orientation, the magnetic dipole moments will no longer be totally random in their orientation;
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- They will retain a net magnetisation along the applied field.
- It is this effect which is responsible for the phenomenon known as magnetic hysteresis [13].



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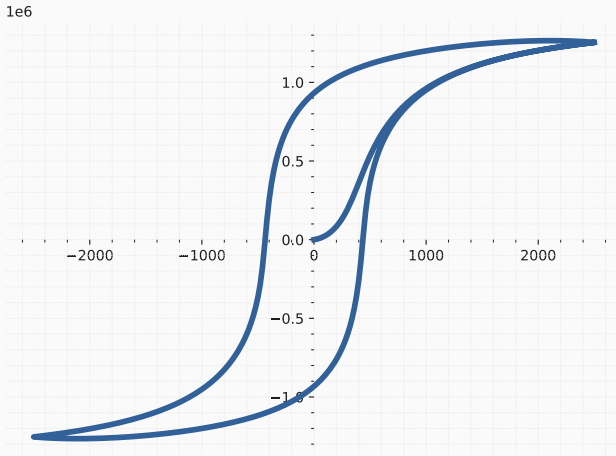


Figure 10: A Hysteresis Curve (B-H) Generated using the Jiles-Artherton model.



1st Quadrant

- The domains are very small, below the single domain size where there is resistance to demagnetization [1].
- The domains start to increase in size. For a small interval, the magnetization will be reversible.
- As the field increases, magnetization will no longer reverse to zero but move on a minor hysteresis loop.
- Eventually the curve starts to bend over to the right.
 - It will still increase as more magnetic domains reach their full size and their magnetizations become parallel to the external field [9].



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1st Quadrant - Saturation

- Eventually the **B**-field will become high enough where no more change in the magnetization occurs.
 - This is called **technical saturation**.
- technology is sufficient to reach 99+% of “technical saturation.”
- As the field is backed off from “saturation”, the magnetization declines very slightly to the B_r point.
 - This is the remanence, remanent induction, or residual induction [8].

All of the magnetic energy is now in the magnet and its field.



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2nd Quadrant - Demagnetisation

- In this quadrant the applied field **opposes** the materials **B**-field.
- As the external **B**-field increases in magnitude, some domains will reverse.
- At the knee of the demagnetisation curve, this increase has become rapid and the magnetisation will fall to the H_{ci} point.
- At H_{ci} , the number of domains aligned with the original magnetization is the same as the number aligned with the opposing magnetic field.

The net magnetisation is zero (0).



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3rd Quadrant - Re-magnetisation

- The total magnetization of the part will be reversed.
- If we go far enough, magnetization will reach the saturation level in the negative direction.



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4th Quadrant - Demagnetisation

- After fully reversing the magnet and removing the field in the third quadrant, magnetization will recoil to a point that is the negative of the B_r observed when in the first and second quadrant.
- If we apply additional field in the positive direction, we duplicate the second quadrant curve.



4th Quadrant - Demagnetisation

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- In AC power systems, voltage and flux wave-forms closely approximate sinusoidal functions over time.
- Time to describe the excitation characteristics and losses associated with steady-state ac operation of magnetic materials.
- Assume a closed-core magnetic circuit (i.e., with no air gap),
- And a sinusoidal variation of the core flux (ϕ) with the following:

$$\phi(t) = \phi_{\max} \sin(\omega t) = A_c B_{\max} \sin(\omega t)$$

where:

ϕ_{\max} amplitude of core flux in Wb,

B_{\max} Amplitude of flux density B_c in T,

ω Angular frequency in rad/s,

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- From, the voltage induced in the N -turn winding is:

$$e(t) = \omega N \phi_{\max} \cos(\omega t) = E_{\max} \cos(\omega t)$$

where:

$$E_{\max} = \omega N \phi_{\max} = 2\pi f N A_c B_{\max}$$

- In steady-state AC, it is more important to use RMS rather than instant values.
- Generally, the rms value of a periodic function of time $f(t)$, of T is:

$$F_{\text{rms}} = \sqrt{\left(\frac{1}{T} \int_0^T f^2(t) dt \right)}$$



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- The rms value of a sine wave can be shown to be $1/\sqrt{2}$ times its peak value. Therefore the induced voltage rms is:

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- To produce magnetic flux in the core requires current in the exciting winding known as the exciting current, i_ϕ .
- The nonlinear magnetic properties of the core require that the waveform of the exciting current differs from the sinusoidal waveform of the flux.
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- Part of this energy is dissipated as losses and results in heating of the core.
- The rest appears as reactive power associated with energy storage in the magnetic field.

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- There are two (2) mechanisms associated with time varying **B**-fields.

Ohmic Heating

- Associated with induced currents in the core material.
- From Faraday's law, we see that **B**-fields give rise to **E**-fields.
- In magnetic materials these **E**-fields result in induced currents,
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- These laminations, which are aligned in the direction of the field lines, are insulated from each other by an oxide layer on their surfaces or by a thin coat of insulating enamel or varnish.
- This greatly reduces the magnitude of the eddy currents since the layers of insulation interrupt the current paths; the thinner the laminations, the lower the losses.
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- Due to the **hysteretic nature** of magnetic material.
- In a magnetic circuit or the transformer, a time-varying excitation (i_ϕ) will cause the magnetic material to undergo a cyclic variation described by a hysteresis loop.
- The energy input W to the magnetic core as the material undergoes a single cycle is shown to be:

$$W = \oint i_\phi d\lambda = \oint \left(\frac{H_c l_c}{N} \right) (A_c N dB_c) = A_c l_c \oint H_c dB_c$$

- Notice $A_c l_c$ is the **core volume** and the integral is the area of the ac hysteresis loop, we see that each time the magnetic material undergoes a cycle, there is a net energy input into the material.



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- Therefore for a given flux level, the corresponding hysteresis losses are **proportional to the area of the hysteresis loop** and to the **total volume of material**.
- As there is an energy loss per cycle, hysteresis power loss is **proportional to the frequency of the applied excitation**. (Data Sheet)



- In general, both losses depend on the metallurgy of the material as well as the flux density and frequency.
- Information on core loss is typically presented in graphical form.
- It is plotted in terms of watts per unit weight as a function of flux density;
 - Often a family of curves for different frequencies are given.
 - Generally it is either 50 or 60 Hz.



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- Nearly all transformers and certain sections of electric machines use sheet-steel material that has highly favorable directions of magnetization along which the core loss is low and the permeability is high.
- This material is termed **grain-oriented steel** [12].
- The reason lies in the atomic structure of a crystal of the silicon-iron alloy,
 - which is a body-centred cube.



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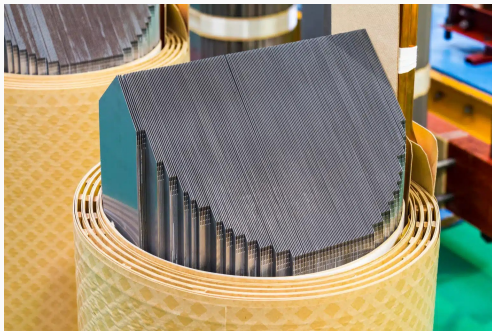


Figure 11: An example of Grain Oriented Steel used in industry [7].

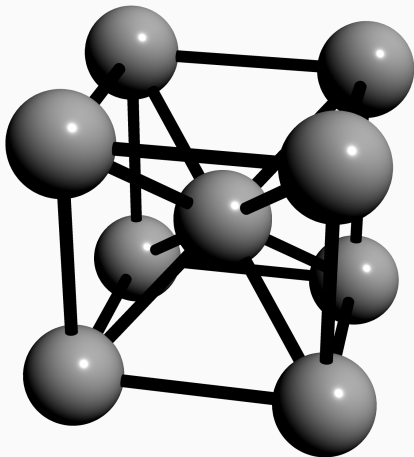


Figure 12: The atomic structure of grain oriented steel.



- Each crystalline cube has an atom at each corner as well as one in the center of the cube.
- In the cube, the easiest axis of magnetization is the cube edge.
- The diagonal across the cube face is more difficult
- By suitable manufacturing techniques most of the crystalline cube edges are aligned in the rolling direction to make it the favorable direction of magnetisation.



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- The behaviour in this direction is superior in core loss and permeability to non-oriented steels in which the crystals are randomly oriented to produce a material with characteristics which are uniform in all directions.
- As a result, oriented steels can be operated at higher flux densities than the nonoriented grades.
- Non-oriented electrical steels are used in applications where the flux does not follow a path which can be oriented with the high-Temperature rolling direction or where low cost is of importance.
- In these steels the losses are somewhat higher and the permeability is very much lower than in grain-oriented steels.



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Appendix



- Assume the following **simplified** induction drive equivalent circuit.

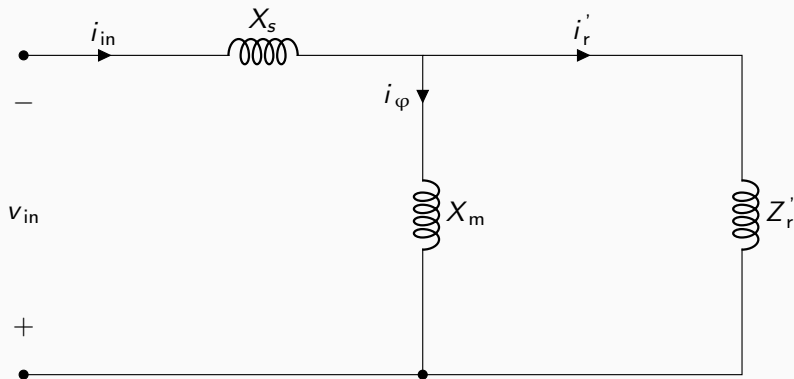


Figure 13: An abstract rendition of the **simplified** induction circuit equivalent circuit.



- where:

$$\begin{aligned}Z_s &= R_s + \mathbf{j} X_s, \\Z_m &= \frac{R_c \times \mathbf{j} X_m}{R_c + \mathbf{j} X_m}, \\Z_r &= \frac{R_r'}{s} + jX_r' .\end{aligned}$$

- From Kirchhoff's current law, the following statement holds true:

$$i_{\text{in}} = i_{\varphi} + i_r'.$$

- The second identity can be derived from parallel circuit principles;

$$Z_m i_{\varphi} = Z_r' i_r \quad \rightarrow \quad i_m = i_r' \times \left(\frac{Z_r}{Z_m} \right) .$$



- We can isolate the magnetizing current (i_φ).

$$i_\varphi = i_r' \times \left(\frac{Z_r}{Z_m} \right) + i_r'.$$

- Isolating the rotor current (i_r') gives us the final expression:

$$i_r' = \frac{i_s}{\left(1 + \frac{Z_r}{Z_m} \right)} \quad \blacksquare$$



- Consider the following equivalent circuit of an IM.

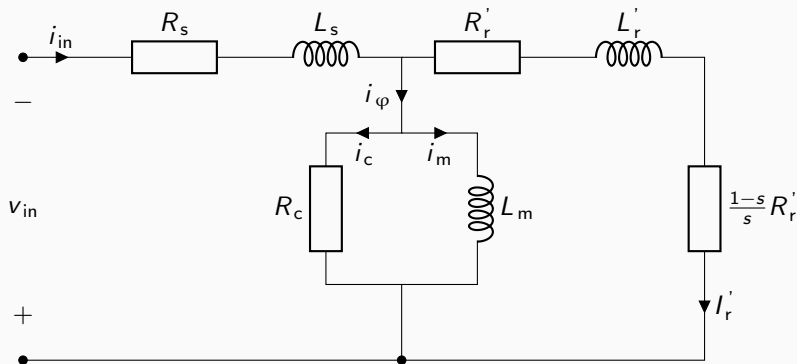


Figure 14: An abstract rendition of the induction circuit equivalent circuit.



- We can start by summing up the **rotor side**.

$$Z_r = \frac{R_r'}{s} + \mathbf{j} w_r L_r' \rightarrow Z_r = \frac{R_r'}{s} + X_r'.$$

- The **stator side** is calculated as:

$$Z_s = R_s + \mathbf{j} w_s L_s \rightarrow Z_s = \frac{R_s}{s} + X_s.$$

- The **magnetising side** is calculated as:

$$Z_m = \frac{R_c \times \mathbf{j} w_s L_m}{R_c + \mathbf{j} w_s L_m} \rightarrow Z_m = \frac{R_c \times X_m}{R_c + X_m}.$$



- Without substitution, the input impedance is calculate to be

$$Z_{in} = Z_s + \frac{Z_m \times Z_r}{Z_m + Z_r}.$$

- From here, we can derive the stator current (I_{in}).

$$V_{ph} = \frac{Z_{in}}{I_{in}} \rightarrow I_{in} = \frac{V_{in}}{Z_{in}}.$$

- Now, the Torque of an IM is related to the air-gap power by:

$$T_{mech} = \frac{P_{gap}}{w_s} \quad \text{and} \quad P_{gap} = n_{ph} \frac{I_r^2}{w_s} \left(\frac{R_2'}{s} \right)$$



- Using the stator-rotor relationship, we can obtain our final result.

$$I_r' = \frac{I_s}{\left(1 + \frac{Z_r}{Z_m}\right)} \quad \blacksquare$$



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an empirical relationship or phenomenological relationship is a relationship or correlation that is supported by experiment or observation but not necessarily supported by theory.



- [1] William Fuller Brown Jr. **“Rigorous approach to the theory of ferromagnetic microstructure”**. In: *Journal of Applied Physics* 29.3 (1958), pp. 470–471.
- [2] Richard P. Feynman. *Feynman lectures on physics. Volume 2: Mainly electromagnetism and matter*. 1964.
- [3] Fausto Fiorillo. *Measurement and characterization of magnetic materials (electromagnetism)*. Elsevier Academic Press, 2005.
- [4] Gorchy. *Photomicrograph of NdFeB*. 2005. URL: <https://en.wikipedia.org/wiki/File:NdFeB-Domains.jpg>.
- [5] David J Griffiths. *Introduction to electrodynamics*. Cambridge University Press, 2023.



- [6] Hiroaki Mamiya and Balachandran Jeyadevan. **“Design criteria of thermal seeds for magnetic fluid hyperthermia-from magnetic physics point of view”**. In: *Nanomaterials for magnetic and optical hyperthermia applications*. Elsevier, 2019, pp. 13–39.
- [7] Metal Miner. ***Grain Oriented Electrical Steel Demand Expected to Keep Growing***. 2024. URL: <https://agmetalminer.com/2023/01/30/grain-oriented-electrical-steel-demand-expected-to-keep-growing/>.
- [8] Terunobu Miyazaki and Hanmin Jin. ***The physics of ferromagnetism***. Vol. 158. Springer Science & Business Media, 2012.
- [9] Alliance Org. **“Understanding the Hysteresis (B-H) Curve”**. In: (2024).



- [10] Hsueh-Yuan Pao, Steven L Dvorak, and Donald G Dudley. “**The effects of neglecting displacement currents when studying transient wave propagation in the Earth**”. In: *IEEE Transactions on Antennas and Propagation* 44.9 (1996), pp. 1259–1265.
- [11] MIT Staff. ***Magnetic Circuits and Transformers: A First Course for Power and Communication Engineers***. MIT Press, 1943.
- [12] Colin Tong. ***Introduction to materials for advanced energy systems***. Springer, 2019.
- [13] Peter J Wasilewski. “**Magnetic hysteresis in natural materials**”. In: *Earth and Planetary Science Letters* 20.1 (1973), pp. 67–72.