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August 3, 2019

Abstract

We are interested in the near flat regime of the radial basis function periodic interval $[0, 2\pi)$. Therefore we use radial basis functions on the sphere S^1

$$\phi(\theta; \theta_0) = \phi_{\epsilon}(t) = \psi_{\epsilon}(\sqrt{2 - 2t}), \quad t \in [-1, 1]$$
(1)

with $\mathbf{t} = \cos(\theta - \theta_0) = x^T x_0, x = e^{i\theta}, x_0 = e^{i\theta_0}, \theta, \theta_0 \in \mathbb{R}$ and $\psi_{\epsilon}(r) : [0, \infty) \to \mathbb{R}$ is the usual radial basis function. Hubbert and Baxter [1] give analytic expressions for the Fourier coefficients $c_{k,\epsilon}$ in

$$\phi_{\epsilon}(t) = c_{k,\epsilon} + \sum_{k=1}^{\infty} \epsilon^{2k} c_{k,\epsilon} \left(e^{i\theta} e^{-i\theta_0} + e^{-i\theta} e^{i\theta_0} \right)$$
 (2)

For Gaussians, i.e.,

$$\phi_{\epsilon}(t) = e^{-\epsilon^2(2-2t)},\tag{3}$$

these coefficients are

$$c_{0,\epsilon} = 0, \quad c_{k,\epsilon} = \frac{kI_k(2\epsilon^2)e^{-2\epsilon^2}}{\epsilon^{2k}}$$
 (4)

where I_k is the modified Bessel function of the first kind. These coefficients converge to finite numbers when ϵ grows small. Since $I_k(z) \sim \frac{(z/2)^k}{\Gamma(k+1)}$ for small z (HMF, ...), $\lim_{\epsilon \searrow 0} c_{k,\epsilon} = \frac{1}{\Gamma(k)}$. Keeping only the first K terms of (2) and filling in $x_m = 2\pi m/M$ and $x_n = 2\pi n/N$, the vector

of basis elements can be written as

$$\begin{bmatrix} \phi(\frac{2\pi m}{M}; \frac{2\pi 0}{N}) \\ \phi(\frac{2\pi m}{M}; \frac{2\pi 1}{N}) \\ \vdots \\ \phi(\frac{2\pi m}{M}; \frac{2\pi (M-2)}{N}) \\ \phi(\frac{2\pi m}{M}; \frac{2\pi (M-1)}{N}) \end{bmatrix} = \begin{bmatrix} W_N^0 & W_N^0 & \cdots & W^0 & W_N^0 & \cdots & W_N^0 & W_N^0 \\ W_N^0 & W_N^{-1} & \cdots & W_N^{-K} & W_N^K & \cdots & W_N^2 & W_N^1 \\ \vdots & \vdots & & \vdots & & \vdots & & \vdots \\ W_N^0 & W_N^{-1(N-2)} & \cdots & W_N^{-K(N-2)} & W_N^{K(N-2)} & \cdots & W_N^{2(N-2)} & W_N^{1(N-2)} \\ W_N^0 & W_N^{-1(N-2)} & \cdots & W_N^{-K(N-1)} & W_N^{K(N-1)} & \cdots & W_N^{2(N-1)} & W_N^{1(N-1)} \end{bmatrix} E_{K,\epsilon} C_{K,\epsilon} \begin{bmatrix} W_M^{\tilde{M}} \\ W_M^{\tilde{M}} \\ \vdots \\ W_N^{-K} \\ W_M^{\tilde{M}} \\ \vdots \\ W_M^{-2} \\ W_M^{\tilde{M}} \end{bmatrix}$$

$$(5)$$

where $W_N = e^{\frac{i2\pi}{N}}$, $W_M = e^{\frac{i2\pi}{M}}$ and

$$E_{K,\epsilon} = \operatorname{diag}\{\epsilon^0, \epsilon^2, \dots, \epsilon^K, \epsilon^K, \dots, \epsilon^4, \epsilon^2\}$$
 (6)

$$C_{K,\epsilon} = \operatorname{diag}\{c_{0,\epsilon}, c_{1,\epsilon}, \dots, c_{K,\epsilon}, c_{K,\epsilon}, \dots, c_{2,\epsilon}, c_{1,\epsilon}\}.$$

$$(7)$$

The collocation matrix $A_{m,n} = \phi(x_m; x_n)$ for m = 1, ..., M, n = 1, ..., N becomes

$$A = \begin{bmatrix} W_{M}^{0} & W_{M}^{0} & \cdots & W_{M}^{0} & W_{M}^{0} \\ W_{M}^{0} & W_{M}^{-1} & \cdots & W_{M}^{-(M-2)} & W_{M}^{-(M-1)} \\ \vdots & \vdots & & \vdots & & \vdots \\ W_{M}^{0} & W_{M}^{-K} & \cdots & W_{M}^{-K(M-2)} & W_{M}^{-K(M-1)} \\ W_{M}^{0} & W_{M}^{K} & \cdots & W_{M}^{K(M-2)} & W_{M}^{K(M-1)} \\ \vdots & \vdots & & \vdots & & \vdots \\ W_{M}^{0} & W_{M}^{2} & \cdots & W_{M}^{2(M-2)} & W_{M}^{2(M-1)} \\ W_{M}^{0} & W_{M}^{1} & \cdots & W_{M}^{(M-2)} & W_{M}^{M-1} \end{bmatrix} \times$$

$$\begin{bmatrix} W_{N}^{0} & W_{N}^{0} & \cdots & W_{N}^{0} \\ \end{bmatrix}$$

$$\begin{bmatrix} W_{N}^{0} & W_{N}^{0} & \cdots & W_{N}^{0} \\ \end{bmatrix}$$

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$$E_{K,\epsilon} \times C_{K,\epsilon} \times \begin{bmatrix} W_N^0 & W_N^0 & \cdots & W_N^{M-1} \\ W_M^0 & W_M^1 & \cdots & W_M^{(M-2)} & W_M^{M-1} \end{bmatrix}$$

$$E_{K,\epsilon} \times C_{K,\epsilon} \times \begin{bmatrix} W_N^0 & W_N^0 & \cdots & W_N^0 & W_N^0 \\ W_N^0 & W_N^{-1} & \cdots & W_N^{-(N-2)} & W_N^{-(N-1)} \\ \vdots & \vdots & & \vdots & & \vdots \\ W_N^0 & W_N^{-K} & \cdots & W_N^{-K(N-2)} & W_N^{-K(N-1)} \\ W_N^0 & W_N^{K} & \cdots & W_N^{K(N-2)} & W_N^{K(N-1)} \\ \vdots & \vdots & & \vdots & & \vdots \\ W_N^0 & W_N^2 & \cdots & W_N^{2(N-2)} & W_N^{2(N-1)} \\ W_N^0 & W_N^1 & \cdots & W_N^{1(N-2)} & W_N^{1(N-1)} \end{bmatrix}$$

$$(9)$$

Note the adjoint applied at the first matrix. The first and last matrix are very similar to the DFT matrix of size $N \times N$

$$F_{N} = \begin{bmatrix} 1 & 1 & 1 & 1 & \cdots & 1\\ 1 & W_{N}^{-1} & W_{N}^{-2} & W_{N}^{-3} & \cdots & W_{N}^{-(N-1)}\\ 1 & W_{N}^{-2} & W_{N}^{-4} & W_{N}^{-6} & \cdots & W_{N}^{-2(N-1)}\\ 1 & W_{N}^{-3} & W_{N}^{-6} & W_{N}^{-9} & \cdots & W_{N}^{-3(N-1)}\\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots\\ 1 & W_{N}^{-(N-1)} & W_{N}^{-2(N-1)} & W_{N}^{-3(N-1)} & \cdots & W_{N}^{-(N-1)(N-1)} \end{bmatrix}$$

$$(10)$$

References

[1] S. Hubbert and B. Baxter, *Radial basis functions for the sphere*, Recent Progress in Multivariate Approximation. 4th International Conference, September 2000, (2001), pp. 33–47.