

Magnetism_Formulae

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Part 1

Bio Savarts law

Differential

$$d\vec{H} = \frac{I}{4\pi} \frac{d\vec{L} \times \vec{r}}{|\vec{r}|^3}$$

 Note

Remember

- Cartesian

$$d\vec{L} = dx\vec{a}_x + dy\vec{a}_y + dz\vec{a}_z$$

- Cylindrical

$$d\vec{L} = d\rho\vec{a}_\rho + \rho d\phi\vec{a}_\phi + dz\vec{a}_z$$

- Spherical

$$d\vec{L} = dr\vec{a}_r + r d\theta\vec{a}_\theta + r \sin(\theta) d\phi\vec{a}_\phi$$

Magnetic Field Intensities

Line Segment

$$\vec{H} = \frac{I}{4\pi\rho} (\cos(\alpha_1) + \cos(\alpha_2))\vec{a}_\phi$$

$$\vec{H} = \frac{I}{4\pi\rho} \frac{2L}{\sqrt{l^2 + 4\rho^2}} \vec{a}_\phi$$

Circular coil

$$\overline{H} = \frac{I\rho^2}{2(\rho^2 + h^2)^{\frac{3}{2}}}a_z$$

Solenoid

$$\overline{H} = nIa_y$$

Amperes Law

$$\oint \overline{H} \cdot d\overline{l} = I_{enclosed} = \int_S \overline{J} \cdot d\overline{S}$$

$$\nabla \times \overline{H} = \overline{J}$$

Applications

Infinite Line Current

$$\overline{H} = \frac{I}{2\pi\rho}a_\phi$$

Infinite Sheet of current

$$\overline{H} = \pm \frac{k}{2} \times a_n$$

Infinite Length Coaxial Transmission Line

Hollow Cylinder

$$\overline{H} = 0 ; \rho < a$$

$$\overline{H} = \frac{I}{2\pi\rho}a_\phi ; \rho > a$$

Solid Cylinder

$$\overline{H} = \frac{I\rho}{2\pi a^2}a_\phi ; \rho < a$$

$$\overline{H} = \frac{I}{2\pi\rho}a_\phi ; \rho > a$$

Coaxial Cable

$$\overline{H} = \frac{I\rho}{2\pi a^2}a_\phi ; \rho < a$$

$$\overline{H} = \frac{I}{2\pi\rho}a_\phi ; a \leq \rho \leq b$$

$$\overline{H} = \frac{I}{2\pi\rho} \left[1 - \frac{\rho^2 - b^2}{r_1^2 - b^2} \right] a_\phi ; b \leq \rho \leq r_1$$

$$\overline{H} = 0 ; \rho > r_1$$

Toroid

$$\overline{H} = 0 ; \rho < \rho_o - a$$

$$\overline{H} = \frac{NI}{2\pi\rho} ; \rho_o - a < \rho < \rho_o + a$$

$$\overline{H} = 0 ; \rho > \rho_o + a$$

Stoke Law

$$\oint \overline{H} \cdot d\overline{l} = \int_S (\nabla \times \overline{H}) \cdot d\overline{S}$$

Magnetic Field Density

$$\overline{B} = \mu \overline{H}$$

Magnetic Flux

$$\Psi = \int_S \overline{B} \cdot d\overline{S}$$

Displacement Current

$$\overline{J}_d = \varepsilon \frac{d\overline{E}}{dt}$$

Maxwells Equations

Name	Integral Form	Point Form
Ampere's Law	$\oint \overline{H} \cdot d\overline{l} = \int_S \overline{J} \cdot d\overline{S}$	$\nabla \times \overline{H} = \overline{J} + \varepsilon \frac{d\overline{E}}{dt}$
Gauss' Law	$\oint \overline{D} \cdot d\overline{S} = \int_V \rho_v dV$	$\nabla \cdot \overline{E} = \frac{\rho_v}{\varepsilon}$
None Existence of magnetic dipoles	$\oint_S \overline{B} \cdot d\overline{S} = 0$	$\nabla \cdot \overline{B} = 0$

Name	Integral Form	Point Form
Faraday's Law	$\int \vec{E} \cdot d\vec{l} = 0$	$\nabla \times \vec{E} = -\frac{d\vec{B}}{dt}$

Magnetic Boundary Conditions

$$\mu_1 \vec{H}_{1n} = \mu_2 \vec{H}_{2n} \leftarrow \text{or} \rightarrow \vec{B}_{1n} = \vec{B}_{2n}$$

$$H_{1t} - H_{2t} = k$$

Refraction

$$\frac{\tan(\theta_1)}{\tan(\theta_2)} = \frac{\mu_1}{\mu_2}$$

Note

θ wrt to normal to plane

Force due to Magnetic Fields

$$\vec{F} = I \vec{L} \times \vec{B}$$

$$\vec{F} = Q \vec{u} \times \vec{B}$$

Force on Current Element

$$d\vec{F} = I d\vec{l} \times \vec{B}$$

$$d\vec{F} = \vec{K} d\vec{S} \times \vec{B}$$

$$d\vec{F} = \vec{J} dV \times \vec{B}$$

Force on charged particle

- Lorentz Force Equation

$$\vec{F} = Q(\vec{E} + \vec{u} \times \vec{B})$$

Force between Current Elements

$$\vec{F}_1 = -\vec{F}_2 = \frac{\mu_o I_1 I_2}{4\pi} \int_{l_1} \int_{l_2} \frac{d\vec{l}_1 \times (d\vec{l}_2 \times \vec{a}_{R_{21}})}{|\vec{R}_{21}|^2}$$

Faradays Law

$$V_{emf} = \frac{-Nd\Psi}{dt} = \oint \overline{E} d\overline{l}$$

Stationary loop Varying magnetic Field

$$V_{emf} = \oint \overline{E} \cdot d\overline{l} = - \int_S \frac{d\overline{B}}{dt} \cdot d\overline{S}$$

- $V_{emf} \rightarrow$ Transformer EMF

Moving loop Static magnetic Field

$$V_{emf} = \oint \overline{E} \cdot d\overline{l} = \oint_L (\overline{u} \times \overline{B}) \cdot d\overline{l}$$

- Using Stokes Theorem

$$\nabla \times \overline{E} = \nabla \times (\overline{u} \times \overline{B})$$

Moving loop varying Magnetic Field

$$V_{emf} = \oint \overline{E} \cdot d\overline{l} = - \int_S \frac{d\overline{B}}{dt} \cdot d\overline{S} + \oint_L (\overline{u} \times \overline{B}) \cdot d\overline{l}$$

Part 2

Phasors

$$z = a + ib \rightarrow r\angle\theta = re^{j\theta}$$

$$X = \cos(\omega t - \beta x) = \text{Re}\{\overline{X}_o e^{j\omega t}\}$$

- We can define the phasor X_s such that

$$\overline{X}(x, y, z, t) = \text{Re}\{\overline{X}_s(x, y, z) e^{j\omega t}\}$$

Maxwell equations

Name	Phasor Form
Ampere's Law	$\nabla \times \overline{H}_s = \overline{J}_s + j\omega \overline{D}_s$
Gauss' Law	$\nabla \cdot \overline{E}_s = \frac{\rho_v}{\epsilon}$

Name	Phasor Form
None Existence of magnetic dipoles	$\nabla \cdot \bar{B}_s = 0$
Faraday's Law	$\nabla \times \bar{E}_s = -j\omega \bar{B}_s$

Wave Propagation #todo

- General Wave Equation:

$$\frac{d^2 X}{dt^2} - u^2 \nabla^2 X = 0$$

- Or

$$\frac{d^2 E}{dt^2} - \frac{1}{\mu\epsilon} \nabla^2 E = 0$$

- Solution to Wave Equation:

$$X(x, y, z, t) = A_o e^{-\alpha x} \cos(\omega t - \beta x) = \text{Re}\{\bar{X}_s(x, y, z) e^{j\omega t}\}$$

$$\beta = \frac{\omega}{u} = \frac{2\pi}{\lambda}$$

Lossy Dielectrics

- Wave Equation:

$$\nabla^2 \bar{E}_s - \gamma^2 \bar{E}_s = 0$$

γ - Propagation constant

$$\gamma = \sqrt{j\omega\mu(\sigma + j\omega\epsilon)}$$

- $\gamma = \alpha + j\beta$

α - Attenuation constant

$$\alpha = \omega \sqrt{\frac{\mu\epsilon}{2} \left[\sqrt{1 + \left(\frac{\sigma}{\omega\epsilon}\right)^2} - 1 \right]}$$

β - Phase shift constant

$$\beta = \omega \sqrt{\frac{\mu\epsilon}{2} \left[\sqrt{1 + \left(\frac{\sigma}{\omega\epsilon}\right)^2} + 1 \right]}$$

$$\beta = \frac{\omega\mu\sigma}{2\alpha}$$

η - Intrinsic Impedance

$$\eta = \sqrt{\frac{J\omega\mu}{\sigma + J\omega\epsilon}}$$

$$\eta = |\eta|e^{J\theta_\eta}$$

$$|\eta| = \frac{\sqrt{\frac{\mu}{\epsilon}}}{\left[1 + \left(\frac{\sigma}{\omega\epsilon}\right)^2\right]^{\frac{1}{4}}}$$

θ_η - Phase

$$\theta_\eta = \frac{1}{2}\tan^{-1}\left(\frac{\sigma}{\omega\epsilon}\right)$$

Note

- \overline{E} leads \overline{H} by angle θ_η
- \overline{H} lags \overline{E} by angle θ_η

Plane Waves in Free Space

- We know that $\gamma^2 = J\omega\mu(\sigma + J\omega\epsilon)$
- For free space $\sigma = 0$, therefore:

$$\gamma = J\omega\sqrt{\mu_o\epsilon_o}$$

- $\alpha = 0$, therefore **no attenuation**

$$\beta = \omega\sqrt{\mu_o\epsilon_o} = \frac{\omega}{c}$$

- $\eta_o = \sqrt{\frac{\mu_o}{\epsilon_o}} = 120\pi$
- \overline{E} and \overline{H} are in phase

Lossless Dielectrics

- For free space $\sigma = 0$, therefore:

$$\gamma = J\omega\sqrt{\mu\varepsilon}$$

- And $\eta = \sqrt{\frac{\mu}{\varepsilon}}$
- And $\beta = \omega\sqrt{\mu\varepsilon}$, therefore:

$$c = \frac{1}{\sqrt{\mu_o\mu_r\varepsilon_o\varepsilon_r}}$$

Good Conductors

Plane Waves

Skin Effect

- Skin depth / penetration is the distance travelled by EM wave getting attenuated by a factor e

$$\delta = \frac{1}{\alpha}$$

- For good conductors,

$$\delta = \frac{1}{\sqrt{\pi\rho\mu\sigma}}$$

- δ is very small for good conductors \therefore waves dies down at high rate
- Very high frequency (RF range) comes from good conductors, δ is small at high frequency \therefore **Skin effect**

Wave Polarization

Linear

$$E_x = E_{ox}\cos(\omega t - \beta z)$$

$$E_y = E_{oy}\cos(\omega t - \beta z)$$

$$z = 0$$

$$\Delta\phi = \phi_y - \phi_x = 0$$

$$E_{ox} = E_{oy} \rightarrow \theta = 45^\circ$$

$$E_{ox} \neq E_{oy} \rightarrow \theta = \tan^{-1}\left(\frac{E_{oy}}{E_{ox}}\right)$$

Circular

$$E_x = E_{ox}\cos(\omega t - \beta z)$$

$$E_y = E_{oy} \cos(\omega t - \beta z - \frac{\pi}{2})$$

$$E_{ox} = E_{oy}$$

$$\Delta\phi = \phi_y - \phi_x = \frac{\pi}{2}$$

Elliptical

$$E_x = E_{ox} \cos(\omega t - \beta z)$$

$$E_y = E_{oy} \sin(\omega t - \beta z)$$

$$E_{ox} \neq E_{oy}$$

$$\Delta\phi = \phi_y - \phi_x = \frac{\pi}{2}$$

$$\left(\frac{E_x}{E_{ox}}\right)^2 + \left(\frac{E_y}{E_{oy}}\right)^2 = 1$$

Poynting Theorem

$$\overline{P} = \overline{E} \times \overline{H}$$

- For perfect dielectric:

$$E_x = E_{ox} \cos(\omega t - \beta z)$$

$$H_y = \frac{E_{ox}}{\eta} \cos(\omega t - \beta z)$$

- Power density:

$$P_z = \frac{E_{ox}^2}{\eta} \cos^2(\omega t - \beta z)$$

- For lossy dielectric, power density in phasor form

$$\langle P_z \rangle = \frac{1}{2} \frac{E_{ox}^2}{|\eta|} e^{-2az} \cos\theta_\eta$$