

# Sec 11.2: Sturm–Liouville Boundary Value Problems

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Instructor: Daesung Kim

## Recall: Eigenvalues

We recall the definition of eigenvalues and eigenfunctions for a general differential operator.

$$y'' + \lambda y = 0 \quad y(0) = y(L) = 0.$$

### Definition

Let  $L[y]$  be a differential operator. (Ex:  $L[y] = -y''$ )

We say that  $\lambda$  is a real (or complex) eigenvalue of  $L$  with **a homogeneous boundary conditions** if  $\lambda \in \mathbb{R}$  (or  $\lambda \in \mathbb{C}$ ) and the differential equation  $L[y] = \lambda y$  with the given homogeneous boundary conditions has **nontrivial solutions**.

The nontrivial solutions are called **eigenfunctions**.

$$\textcircled{1} \quad L[y] = -y'' = \lambda y \Rightarrow y'' + \lambda y = 0.$$

$$\textcircled{2} \quad L[y] = -x^2 y'' + 2y' = \lambda y \Rightarrow x^2 y'' - 2y' + \lambda y = 0.$$

# The Sturm–Liouville boundary problem

We study the Sturm–Liouville boundary problem, which consists of a differential equation of the form

$$L[y] = -(py')' + qy$$

$$\textcircled{\Delta} \quad (p(x)y')' - q(x)y + \lambda r(x)y = 0$$

$$= \lambda r y$$

on  $[0, 1]$  with boundary conditions  $p \cdot y'' + p' \cdot y' - qy + \lambda r y = 0$ .

$$\left\{ \begin{array}{l} \alpha_1 y(0) + \alpha_2 y'(0) = 0, \\ \beta_1 y(1) + \beta_2 y'(1) = 0 \end{array} \right.$$

with  $\alpha_1^2 + \alpha_2^2 > 0$  and  $\beta_1^2 + \beta_2^2 > 0$ .

We further assume that  $p, p', q, r$  are continuous on  $[0, 1]$  and  $p(x), r(x) > 0$  for all  $x \in [0, 1]$ . In this case, the problem is called **regular**.

$$u_x = (p \cdot u_x)_x - q(x) y$$

## Examples of the Sturm–Liouville boundary problem

### Example

If  $p(x) = r(x) = 1$  and  $q(x) = 0$ , then the problem is  $y'' + \lambda y = 0$  with

$$\alpha_1 y(0) + \alpha_2 y'(0) = 0,$$

$$\beta_1 y(1) + \beta_2 y'(1) = 0.$$

$$(p \cdot y')' - q y + \lambda r y = 0$$

$$p=r=1 \quad \& \quad q=0,$$

$$\begin{aligned} (y')' - 0 \cdot y + \lambda y \\ = y'' + \lambda y = 0. \end{aligned}$$

# Examples of the Sturm–Liouville boundary problem

## Example

If  $p(x) = x^k$ ,  $q(x) = 0$ , and  $r(x) = x^{k-2}$ , then

$$\begin{aligned}(p(x)y')' - q(x)y + \lambda r(x)y &= x^k y'' + kx^{k-1}y' + \lambda x^{k-2}y \\&= x^{k-2}(x^2 y'' + kxy' + \lambda y) \\&= 0.\end{aligned}$$

$= 0$

Thus, the problem  $x^2 y'' + kxy' + \lambda y = 0$  with //

$$\alpha_1 y(1) + \alpha_2 y'(1) = 0,$$

$$\beta_1 y(2) + \beta_2 y'(2) = 0.$$

$$\begin{aligned}L[y] &= -x^2 y'' - kxy' \\&= \lambda y\end{aligned}$$

also belongs to the class.

## Lagrange's Identity

$$\underline{(py')' - qy + \lambda ry = 0}$$

Let  $u$  and  $v$  be functions on  $[0, 1]$  with continuous second derivatives and  $L$  a differential operator defined by  $L[y] = \underline{-(p(x)y')' + q(x)y}$ . Then,

$$= \lambda ry$$

$$\int_0^1 (L[u]v - uL[v]) dx = -p(x) [u'(x)v(x) - u(x)v'(x)]_0^1$$

Integration by parts.

$$= p(0) (u'(0)v(0) - u(0)v'(0)) - p(1) (u'(1)v(1) - u(1)v'(1)) = 0$$

Furthermore, if  $u$  and  $v$  satisfies the boundary condition

$$\begin{cases} \alpha_1 y(0) + \alpha_2 y'(0) = 0, \\ \beta_1 y(1) + \beta_2 y'(1) = 0 \end{cases}$$

with  $\alpha_1^2 + \alpha_2^2 > 0$  and  $\beta_1^2 + \beta_2^2 > 0$ , then

$$\int_0^1 (L[u]v - uL[v]) dx = 0.$$

$$\begin{aligned} L[y] &= -(py')' + qy \\ &= \lambda ry \end{aligned}$$

Remark  $(L[u], v) = \int_0^1 L[u] \cdot v \, dx = \int_0^1 L[v] \cdot u \, dx = (u, L[v])$

The Lagrange's identity with the boundary condition can be written as

$$(L[u], v) = (u, L[v])$$

where the inner product on  $[0, 1]$  is defined by

$$(f, g) = \int_0^1 f(x)g(x) \, dx.$$

The Lagrange's identity also holds for complex-valued functions  $f$  and  $g$  with the complex inner product

$$(f, g) = \int_0^1 f(x) \overline{g(x)} \, dx.$$

conjugate bar

$$\overline{(a + bi)} = a - bi$$

## Remark

The Lagrange's identity with the boundary condition can be written as

$$(L[u], v) = (u, L[v])$$

where the inner product on  $[0, 1]$  is defined by

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$$L[u] = \lambda r u$$

$$L[v] = \lambda r v$$

Rmk: If  $u, v$  are solutions to  $S$ - $T$  bdy problem, then  $(\underline{L[u]}, v) = (u, \underline{L[v]})$ .



## Theorem

*Every eigenvalue of the Sturm–Liouville problem is real.*

$$L[y] = -(p y')' + q y = \lambda r y$$

with Bdry conditions.

Proof Suppose  $\lambda$  is an eigenvalue and  $\phi$  is the corresponding eigenfunction.

$$\textcircled{1} \quad L[\phi] = \lambda r \phi$$

$$\textcircled{2} \quad (L[\phi], \phi) = (\phi, L[\phi])$$

## Theorem

Every eigenvalue of the Sturm–Liouville problem is real.

$$\begin{aligned} (L[\phi], \phi) &= (\phi, L[\phi]) \\ \int_0^1 \underbrace{L[\phi]}_{\lambda r \phi} \cdot \overline{\phi} \, dx &= \int_0^1 \phi \cdot \overline{\underbrace{L[\phi]}_{\lambda r \phi}} \, dx \end{aligned}$$

$$\int_0^1 \lambda r \phi \cdot \overline{\phi} \, dx = \int_0^1 \phi \cdot \overline{\lambda r \phi} \, dx$$

$$r(x) > 0 \Rightarrow r(x) : \text{real} \Rightarrow \overline{r(x)} = r(x)$$

$$\lambda \int_0^1 (r \phi \cdot \overline{\phi}) \, dx = \overline{\lambda} \int_0^1 (r \phi \cdot \overline{\phi}) \, dx$$

## Theorem

Every eigenvalue of the Sturm–Liouville problem is real.

$$(\lambda - \bar{\lambda}) \underbrace{\int_0^1 \underbrace{r(x)}_{>0} \underbrace{\phi(x) \overline{\phi(x)}}_{>0} dx}_{>0 \text{ (Not trivial)}} = 0$$

$$\Rightarrow (\lambda - \bar{\lambda}) = 0$$

$$\Rightarrow \lambda = \bar{\lambda}$$

$$\Rightarrow \lambda : \text{real.}$$



## Theorem

If  $\phi_n$  and  $\phi_m$  are two eigenfunctions of the Sturm–Liouville problem corresponding to distinct eigenvalues  $\lambda_n$  and  $\lambda_m$ , then

$$(r\phi_n, \phi_m) = \int_0^1 \phi_n(x)\phi_m(x)r(x)dx = 0.$$

We call  $\phi_n$  and  $\phi_m$  are orthogonal on  $[0, 1]$  with the weight  $r(x)$ .

Proof ①  $L[\phi_n] = \lambda_n r \phi_n$

$$L[\phi_m] = \lambda_m r \phi_m$$

②  $(L[\phi_n], \phi_m) = (\phi_n, L[\phi_m])$

$$(\lambda_n r \phi_n, \phi_m) = (\phi_n, \lambda_m r \phi_m)$$

$$0 \neq (\lambda_n - \lambda_m) \cdot (r \phi_n, \phi_m) = 0$$



## Theorem

*Every eigenvalue of the Sturm–Liouville problem is simple; that is, if  $\phi_1$  and  $\phi_2$  are the corresponding eigenfunctions of the same eigenvalue  $\lambda$ , then they are linearly dependent. Furthermore, the eigenvalues form an infinite sequence  $\lambda_1 < \lambda_2 < \cdots$  and  $\lim_{n \rightarrow \infty} \lambda_n = \infty$ .*

## Definition

An eigenfunction  $\phi$  is normalized with respect to the weight  $r$  if

$$(r\phi, \phi) = \int_0^1 \phi^2 r(x) dx = 1.$$

## Remark

Let  $\lambda_1 < \lambda_2 < \cdots$  be the eigenvalues and  $\phi_n$  the corresponding normalized eigenfunctions. Then,

$$(r\phi_n, \phi_m) = \int_0^1 \phi_n \phi_m r(x) dx = \begin{cases} 1, & m = n \\ 0 & m \neq n. \end{cases}$$

In this case, we call the set  $\{\phi_n : n = 1, 2, \dots\}$  is orthonormal.

Suppose that we are given a function  $f(x)$  and want to represent it in terms of the normalized eigenfunctions  $\phi_n$  as we did in the heat conduction equation. To be specific, the question is to find  $C_n$  such that

$$f(x) = \sum_{n=1}^{\infty} C_n \phi_n(x).$$

$$= \frac{f(x) + f(x-)}{2}$$

Using the orthogonality and the normalization, it is obvious to guess that

$$C_n = \int_0^1 f(x) \phi_n(x) r(x) dx.$$

$$y'' + \lambda y = 0 \Rightarrow \text{eigenfunctions are } \rightarrow f(x) = \sum \text{sines} \\ \text{cosines} \\ \underline{\text{sine}} \text{ or } \underline{\text{cosine}}$$

## Theorem

Let  $\phi_1, \phi_2, \dots$  be the normalized eigenfunctions of the Sturm–Liouville problem. If  $f(x)$  and  $f'(x)$  are piecewise continuous on  $[0, 1]$ , then

$$\underbrace{\sum_{n=1}^{\infty} C_n \phi_n(x)} = \frac{1}{2}(f(x+) + f(x-))$$

with

$$C_n = \int_0^1 f(x) \phi_n(x) r(x) dx.$$