Programming Module Session-8

Solver 1:

Implementation of the transport equation model into cavitatingFoam: "cavitatingTransportFoam"

Solver1: cavitatingTransportFoam

About cavitatingFoam

Barotropic equation of state:
$$\frac{D\rho_m}{Dt} = \Psi \frac{DP}{Dt}$$

Continuity equation:
$$\frac{\partial \rho_{m}}{\partial t} + \nabla \bullet (\rho_{m}U) = 0$$

Momentum equation:
$$\frac{\partial \rho_{\scriptscriptstyle m} U}{\partial t} + \nabla \bullet (\rho_{\scriptscriptstyle m} U U) = -\nabla P + \nabla \bullet \Big[(\mu_{\rm \it eff} (\nabla U + (\nabla U)^{\scriptscriptstyle T}) \Big]$$

 ρ_m , $\rho_{l,sat}$ and $\rho_{v,sat}$ are density of the mixture, liquid and vapor densities at saturation pressure

Solver1: cavitatingTransportFoam

About interPhaseChangeFoam

Continuity equation:
$$\frac{\partial \rho_m}{\partial t} + \nabla \bullet (\rho_m U) = 0$$

$$\text{Momentum equation:} \quad \frac{\partial \rho_{\scriptscriptstyle m} U}{\partial t} + \nabla \bullet (\rho_{\scriptscriptstyle m} U U) = -\nabla P + \nabla \bullet \Big[(\mu_{\it eff} (\nabla U + (\nabla U)^{\scriptscriptstyle T}) \Big]$$

Transport equation:

$$\frac{\partial(\alpha\rho_l)}{\partial t} + \nabla \cdot (\alpha\rho_l \mathbf{U}) + \nabla \cdot [\alpha \mathbf{U}_c(1-\alpha)] = R_c - R_e$$

The mixture density and viscosity:

$$\rho_m = (1 - \alpha)\rho_v + \alpha\rho_l$$

$$\mu_m = (1 - \alpha)\mu_v + \alpha\mu_l$$

Solver1: cavitatingTransportFoam

Kunz Model

Transport equation:

$$\frac{\partial(\alpha\rho_l)}{\partial t} + \nabla \cdot (\alpha\rho_l \mathbf{U}) + \nabla \cdot [\alpha \mathbf{U}_c(1-\alpha)] = R_c - R_e$$

Evaporation source term

$$R_e = Cv \frac{\alpha \rho_v}{t_{\infty}(0.5\rho_l \mathbf{U_{\infty}}^2)} \min[0, P - P_v]$$

Condensation source term

$$R_c = Cc \frac{(1-\alpha)\alpha^2 \rho_v}{t_\infty}$$

U_∞: mean stream velocity

P_v: vapor pressure

 C_c : empirical condensation constants

 C_v : empirical vaporization constants

Solver 2: Implementation of ODE solvers

Solver1: ODE Solvers

ODETest

$$\frac{dy_1}{dx} = -y_2, \qquad \frac{dy_2}{dx} = y_1 - \frac{y_2}{x},$$

$$\frac{dy_3}{dx} = y_2 - \frac{2y_3}{x}, \qquad \frac{dy_4}{dx} = y_3 - \frac{3y_4}{x}$$

Then we calculate the jacobian for x and y:

$$\frac{-dy_2}{dx} = 0, \qquad \frac{d(y_1 - \frac{y_2}{x})}{dx} = \frac{y_2}{x^2}, \qquad \frac{d(y_2 - \frac{2y_3}{x})}{dx} = \frac{2y_3}{x^2}, \qquad \frac{d(y_3 - \frac{3y_4}{x})}{dx} = \frac{3y_4}{x^2}$$

$$\frac{-dy_2}{dy_1} = 0, \qquad \frac{-dy_2}{dy_2} = -1, \qquad \frac{-dy_2}{dy_3} = 0, \qquad \frac{-dy_2}{dy_4} = 0$$

$$\frac{d(y_1 - \frac{y_2}{x})}{dy_1} = 1, \qquad \frac{d(y_1 - \frac{y_2}{x})}{dy_2} = -\frac{1}{x}, \qquad \frac{d(y_1 - \frac{y_2}{x})}{dy_3} = 0, \qquad \frac{d(y_1 - \frac{y_2}{x})}{dy_4} = 0$$

$$\frac{d(y_2 - \frac{2y_3}{x})}{dy_1} = 0, \qquad \frac{d(y_2 - \frac{2y_3}{x})}{dy_2} = 1, \qquad \frac{d(y_2 - \frac{2y_3}{x})}{dy_3} = -\frac{2}{x}, \qquad \frac{d(y_2 - \frac{2y_3}{x})}{dy_4} = 0$$

$$\frac{d(y_3 - \frac{3y_4}{x})}{dy_1} = 0, \qquad \frac{d(y_3 - \frac{3y_4}{x})}{dy_2} = 0, \qquad \frac{d(y_3 - \frac{3y_4}{x})}{dy_3} = 0, \qquad \frac{d(y_3 - \frac{3y_4}{x})}{dy_4} = -\frac{3}{x}$$

Solver1: ODE Solvers

myODETest

$$\frac{d^2y}{dx^2} = y, y(0) = 0, y'(0) = 1$$

Then change our ODE to first order ODEs

$$\frac{dy}{dx} = z, \frac{dz}{dx} = y, y(0) = 0, z(0) = 1$$

Jacobian for x:

$$\frac{dz}{dx} = 0, \frac{dy}{dx} = 0$$

Jacobian for y:

$$\frac{dz}{dy} = 0, \frac{dz}{dz} = 1, \frac{dy}{dy} = 1, \frac{dy}{dz} = 0,$$

Solver 3: Implementation of Electro-Magnetic Solver

Solver 3: Electro-Magnetic Solver

Governing equations

Maxwell's equations: $\nabla \times E = 0$

where E is the electric field strength. $\nabla \cdot B = 0$

where B is the magnetic flux density. $\nabla \times H = J$

where H is the magnetic field strength and J is current density

Charge continuity: $\nabla \cdot J = 0$

Ohm's law: $J = \sigma E$ where σ is the electric conductivity

Constitutive law: $B = \mu_0 H$

where μ_0 is the magnetic permeability of vacuum

Solver 3: Electro-Magnetic Solver

Governing equations

Combining above all equations and Coulomb gauge condition ($\nabla \cdot A = 0$) leads to

Poissons's equation for the magneto static fields: $\nabla^2 A = \mu_0 \sigma(\nabla \phi)$

Laplace's equation for the electric potential: $\nabla \cdot [\sigma(\nabla \phi)] = 0$

Solver 4: Implementation of chtMultiRegionFoam for incompressible fluid

Solver 5:

Modification of stressedFoam for the implementation of material non-linearity

Momentum balance

$$\frac{\partial^2(\rho \mathbf{u})}{\partial t^2} - \nabla \cdot \boldsymbol{\sigma} = \mathbf{0}$$

Constitutive relation (linear elasticity) $\sigma = 2\mu\varepsilon + \lambda tr(\varepsilon)\mathbf{I}$

Stain-displacement relation
$$\varepsilon = \frac{1}{2} [\nabla \mathbf{u} + (\nabla \mathbf{u})^T]$$

u is the solid displacement vector

 ρ is the density

 σ is the stress tensor

I is the unit tensor

 μ and λ are material properties called Lamé parameters

Combining all above equations, the governing equation for linear elastic solid body is written with the only unknown variable u.

$$\frac{\partial^2(\rho \mathbf{u})}{\partial t^2} - \nabla \cdot \left[\mu \nabla \mathbf{u} + \underline{\mu} (\nabla \mathbf{u})^T + \lambda \mathbf{I} tr(\nabla \mathbf{u}) \right] = \mathbf{0}$$

Further rearrangement to improve the convergence

$$\nabla \cdot \boldsymbol{\sigma} = \underbrace{\nabla \cdot (\mu \nabla \mathbf{u})}_{implicit} + \underbrace{\nabla \cdot [\mu (\nabla \mathbf{u})^T + \lambda \mathbf{I} tr(\nabla \mathbf{u})]}_{explicit}$$

$$= \underbrace{\nabla \cdot [(2\mu + \lambda) \nabla \mathbf{u}]}_{implicit} + \underbrace{\nabla \cdot [\mu (\nabla \mathbf{u})^T + \lambda \mathbf{I} tr(\nabla \mathbf{u}) - (\mu + \lambda) \nabla \mathbf{u}]}_{explicit}$$

The traction Displacement boundary

Fixed traction force (T), which can be expressed as: $\mathbf{T} = \boldsymbol{\sigma} \cdot \mathbf{n}$

Where **n** is the surface normal on the boundary.

Since the displacement \mathbf{u} is the computing variable, the traction force is converted into displacement on the boundary.

$$\mathbf{T} = \boldsymbol{\sigma} \cdot \mathbf{n} = \left[\mu \nabla \mathbf{u} + \mu (\nabla \mathbf{u})^T + \lambda \mathbf{I} tr(\nabla \mathbf{u}) \right] \cdot \mathbf{n}$$

$$= \underbrace{\left[(2\mu + \lambda) \nabla \mathbf{u} + \mu (\nabla \mathbf{u})^T + \lambda \mathbf{I} tr(\nabla \mathbf{u}) - (\mu + \lambda) \nabla \mathbf{u} \right] \cdot \mathbf{n}}_{explicit}$$

$$(\nabla \mathbf{u}) \cdot \mathbf{n} = \frac{\mathbf{T} - [\mu(\nabla \mathbf{u})^T + \lambda \mathbf{I} tr(\nabla \mathbf{u}) - (\mu + \lambda)\nabla \mathbf{u}] \cdot \mathbf{n}}{2\mu + \lambda}$$

Implementation of material non-linearity like a plasticity in a solid body

Add plastic feature $\left[2\mu(d\varepsilon_p) + \lambda \mathbf{I}tr\left(d\varepsilon_p\right)\right]$

 $d\varepsilon_p$ stands for the incremental plastic strain tensor

$$\nabla \cdot \left\{ \mu \nabla (d\mathbf{u}) + \mu [\nabla (d\mathbf{u})]^T + \lambda \mathbf{I} tr[\nabla (d\mathbf{u})] - \underbrace{[2\mu (d\boldsymbol{\varepsilon}_p) + \lambda \mathbf{I} tr(d\boldsymbol{\varepsilon}_p)]}_{plasticity} \right\} = \mathbf{0}$$

$$\nabla \cdot [(2\mu + \lambda)\nabla(d\mathbf{u})] =$$

$$-\nabla \cdot \{\mu[\nabla(d\mathbf{u})]^T + \lambda \mathbf{I}tr[\nabla(d\mathbf{u})] - (\mu + \lambda)\nabla(d\mathbf{u})\} +$$

$$\nabla \cdot [2\mu(d\varepsilon_p) + \lambda \mathbf{I}tr(d\varepsilon_p)]$$

