

# Optimization using NLOPT

Yoon Kyung Eunnie Lee

Created: November 15, 2014

Last updated: May 1, 2015

## 1 Definition of the BEM problem

### BEM Matrix $\underline{\underline{\mathbf{M}}}$

The first meeting reviewed the definition and size of the matrices (and vectors) that constitute the BEM problem. The surface current in response to an input field is obtained from:

$$\underline{\underline{\mathbf{M}}} \underline{\underline{\mathbf{C}}} = \underline{\underline{\mathbf{F}}}, \quad (1a)$$

$$\underline{\underline{\mathbf{C}}} = \underline{\underline{\mathbf{M}}}^{-1} \underline{\underline{\mathbf{F}}}. \quad (1b)$$

where the surface current  $\underline{\underline{\mathbf{C}}}$  and input field  $\underline{\underline{\mathbf{F}}}$  are vectors of length  $2N$ , where  $N$  is the number of edges in the mesh.  $\underline{\underline{\mathbf{M}}} = \underline{\underline{\mathbf{M}}}^T$  is the symmetric BEM matrix. (computed using SCUFF) *Note: Avoid using matrix inversion whenever possible. For example, avoid  $\underline{\underline{\mathbf{M}}}^{-1}$ .*

### Power/Force/Torque matrix $\underline{\underline{\mathbf{Q}}}$

For power, force, and torque, BEM computes them via the PFT matrix  $\underline{\underline{\mathbf{Q}}}$ , which obeys

$$\underline{\underline{\mathbf{Q}}} = \underline{\underline{\mathbf{Q}}}^\dagger,$$

where  $\underline{\underline{\mathbf{Q}}}^\dagger = (\overline{\underline{\underline{\mathbf{Q}}}})^T$  denotes complex conjugate.  $\underline{\underline{\mathbf{M}}}$  and  $\underline{\underline{\mathbf{Q}}}_{OPFT}$  are independent of the incident field.

## 2 FOM

### Definition

The second meeting reviewed how to obtain the scalar value *Figure of Merit* (FOM) and its gradient. The FOM is obtained from  $\underline{\underline{\mathbf{Q}}}$  as:

$$FOM = \underline{\underline{\mathbf{C}}}^\dagger \underline{\underline{\mathbf{Q}}} \underline{\underline{\mathbf{C}}}, \quad (2)$$

where the  $\underline{\underline{\mathbf{C}}}$  is obtained from eq. 1a.

## Differential of FOM

The differential ( $\delta$ , also called 'grad') of FOM can be obtained from applying the gradient to eq. 1a:

$$\delta(\underline{\underline{\mathbf{M}}} \underline{\underline{\mathbf{C}}} = \underline{\underline{\mathbf{F}}}) \quad (3a)$$

$$\underline{\underline{\mathbf{M}}} \delta \underline{\underline{\mathbf{C}}} = \delta \underline{\underline{\mathbf{F}}} \quad (3b)$$

Using  $\delta(\underline{\underline{\mathbf{C}}}^\dagger) = (\delta \underline{\underline{\mathbf{C}}})^\dagger$ ,

$$\delta(FOM) = \delta(\underline{\underline{\mathbf{C}}}^\dagger \underline{\underline{\mathbf{Q}}} \underline{\underline{\mathbf{C}}})$$

$$= \delta \underline{\underline{\mathbf{C}}}^\dagger \underline{\underline{\mathbf{Q}}} \underline{\underline{\mathbf{C}}} + \underline{\underline{\mathbf{C}}}^\dagger \underline{\underline{\mathbf{Q}}} \delta(\underline{\underline{\mathbf{C}}})$$

$$= \delta \underline{\underline{\mathbf{F}}}^\dagger (\underline{\underline{\mathbf{M}}}^{-1})^\dagger \underline{\underline{\mathbf{Q}}} \underline{\underline{\mathbf{C}}} + \underline{\underline{\mathbf{C}}}^\dagger \underline{\underline{\mathbf{Q}}} \underline{\underline{\mathbf{M}}}^{-1} \delta \underline{\underline{\mathbf{F}}}$$

$$\therefore \delta(FOM) = \delta \underline{\underline{\mathbf{F}}}^\dagger [(\underline{\underline{\mathbf{M}}}^{-1})^\dagger \underline{\underline{\mathbf{Q}}} \underline{\underline{\mathbf{C}}}] + [\underline{\underline{\mathbf{C}}}^\dagger \underline{\underline{\mathbf{Q}}} \underline{\underline{\mathbf{M}}}^{-1}] \delta \underline{\underline{\mathbf{F}}}. \quad (3c)$$

## 3 Adjoint Current

From eq. 3c, The terms inside the brackets can be grouped together and computed in terms of the adjoint current  $\underline{\underline{\mathbf{C}}}_A$  to avoid matrix inversion. Let us write:

$$\underline{\underline{\mathbf{C}}}_A^T = [\underline{\underline{\mathbf{C}}}^\dagger \underline{\underline{\mathbf{Q}}} \underline{\underline{\mathbf{M}}}^{-1}] \quad (4a)$$

$$\underline{\underline{\mathbf{C}}}_A = [\underline{\underline{\mathbf{C}}}^\dagger \underline{\underline{\mathbf{Q}}} \underline{\underline{\mathbf{M}}}^{-1}]^T = \underline{\underline{\mathbf{M}}}^{-1} \overline{\underline{\underline{\mathbf{Q}}}} \overline{\underline{\underline{\mathbf{C}}}} \quad (4b)$$

$$\overline{\underline{\underline{\mathbf{C}}}_A} = (\underline{\underline{\mathbf{M}}}^{-1}) \underline{\underline{\mathbf{Q}}} \underline{\underline{\mathbf{C}}} = (\underline{\underline{\mathbf{M}}}^{-1})^\dagger \underline{\underline{\mathbf{Q}}} \underline{\underline{\mathbf{C}}} \quad (4c)$$

$$\therefore \boxed{\underline{\underline{\mathbf{M}}} \underline{\underline{\mathbf{C}}}_A = \overline{\underline{\underline{\mathbf{Q}}}} \overline{\underline{\underline{\mathbf{C}}}}} \quad (4d)$$

Notice the similarity between eq. 4d and  $\underline{\underline{\mathbf{M}}} \underline{\underline{\mathbf{C}}} = \underline{\underline{\mathbf{F}}}$ . Both can be solved with LU factorization. eq. 3c now becomes:

$$\delta(FOM) = \delta \underline{\underline{\mathbf{F}}}^\dagger [\overline{\underline{\underline{\mathbf{C}}}_A}] + [\underline{\underline{\mathbf{C}}}_A^T] \delta \underline{\underline{\mathbf{F}}} \quad (5a)$$

$$\therefore \boxed{\delta(FOM) = 2 \cdot \text{Re} \left[ \underline{\mathbf{C}}_A^T \delta \underline{\mathbf{F}} \right]}. \quad (5b)$$

Using the complex vectors  $\underline{\mathbf{F}}$  and  $\underline{\mathbf{F}}^\dagger$  separately, we get two differential equations for FOM.

$$\left( \frac{\delta(FOM)}{\delta(\underline{\mathbf{F}})} = \underline{\mathbf{C}}_A^T \right) \quad (6a)$$

$$\left( \frac{\delta(FOM)}{\delta(\underline{\mathbf{F}}^\dagger)} = \underline{\mathbf{C}}_A \right) \quad (6b)$$

## 4 Input Field Decomposition

### Construction of $\underline{\mathbf{F}}$

The total incident field vector  $\underline{\mathbf{F}}$  is computed inside BEM in the following way:

$$F_i = \langle \phi^{inc}, \beta_i \rangle. \quad (7a)$$

In integral form, this is equivalent to:

$$\boxed{F_i = \int_{S_i} \phi^{inc} \cdot \beta_i \, dA, \quad (i = 1..2N)} \quad (7b)$$

$$\boxed{\underline{\mathbf{F}} = \int_S \phi^{inc} \cdot \underline{\beta} \, dA,} \quad (7c)$$

where  $\phi^{inc}$  is the total incident field in  $(x, y, z)$ , and  $\beta_i$  is the unit normal for surface current in  $(x, y, z)$ , per edge  $i = 1..2N$ . Keep in mind that  $\beta_i$  and  $\phi^{inc}$  are vectors in 3D, while  $\underline{\mathbf{F}}$  is a 2N-dimensional vector, and  $\underline{\beta}$  is a  $(2N \times 3)$  matrix.

### Decomposition of $\phi^{inc}$

The incident field  $\phi^{inc}$  can be expanded into any orthonormal basis  $\mathbf{Y}^{lm}$ .

$$\phi^{inc} = \sum c_{lm} \mathbf{Y}^{lm} \quad (8)$$

Both  $\phi$  and  $\mathbf{Y}$  are spatial vectors in 3D.

### Decomposition of $\underline{\mathbf{F}}$

The 2N-dimensional incident field vector  $\underline{\mathbf{F}}$  is a superposition of the  $\underline{\mathbf{F}}^{lm}$  vectors.

$$\underline{\mathbf{F}} = \sum c_{lm} \underline{\mathbf{F}}^{lm}, \quad (9)$$

Each  $\underline{\mathbf{F}}^{lm}$  is separately computed using the same integration method of eq. 7b:

$$F_i^{lm} = \langle \mathbf{Y}^{lm}, \beta_i \rangle. \quad (10a)$$

In integral form, this is equivalent to:

$$\boxed{F_i^{lm} = \int \mathbf{Y}^{lm} \cdot \beta_i \, dA, \quad (i = 1..2N)} \quad (10b)$$

$$\boxed{\underline{\mathbf{F}}^{lm} = \int \mathbf{Y}^{lm} \cdot \underline{\beta} \, dA.} \quad (10c)$$

### Decomposition of FOM

$\delta(FOM)$  can now be computed with respect to the change in each orthonormal mode  $\mathbf{Y}^{lm}$ .

$$\begin{aligned} \delta(FOM) &= 2 \cdot \text{Re} \left[ \underline{\mathbf{C}}_A^T \delta \underline{\mathbf{F}} \right] \\ &= 2 \cdot \text{Re} \left[ \underline{\mathbf{C}}_A^T \sum \delta \left( c_{lm} \underline{\mathbf{F}}^{lm} \right) \right] \\ &= 2 \cdot \text{Re} \left[ \delta c_{lm} \left( \sum \underline{\mathbf{C}}_A^T \underline{\mathbf{F}}^{lm} \right) \right] \\ &= 2 \cdot \text{Re} \left[ \delta c_{lm} \left( \sum \underline{\mathbf{C}}_A^T \int \mathbf{Y}^{lm} \cdot \underline{\beta} \, dA \right) \right] \end{aligned}$$

Therefore,

$$\frac{\delta(FOM)}{\delta c_{lm}} = 2 \cdot \text{Re} \left[ \underline{\mathbf{C}}_A^T \underline{\mathbf{F}}^{lm} \right], \quad (11a)$$

$$\frac{\delta(FOM)}{\delta c_{lm}} = 2 \cdot \text{Re} \left[ \underline{\mathbf{C}}_A^T \int \mathbf{Y}^{lm} \cdot \underline{\beta} \, dA \right]. \quad (11b)$$

Recall that the surface current vector  $\underline{\mathbf{C}}$  and the adjoint current vector  $\underline{\mathbf{C}}_A$  depend on the incident field. For the computation of  $FOM$ , we use  $\underline{\mathbf{C}}$  and  $\underline{\mathbf{C}}_A$  computed using the total incident field vector  $\underline{\mathbf{F}}$ .

## 5 Input to NLOPT

Steven suggested on 2015.03.16 that the following function be constructed:

```
function objective (int n, double *c)
    compute RHS
    sol=A\rhs
    return FOM
```

where \*c is an array of coefficients  $c_{lp}$ . (Steven said, why don't you start with 25 coeff.s, 5 by 5? You would need real and imaginary parts. There is no meaning in trying to make sense of different force/torque from modes. go straight to optimizing. )

### 5.1 Implementation details