

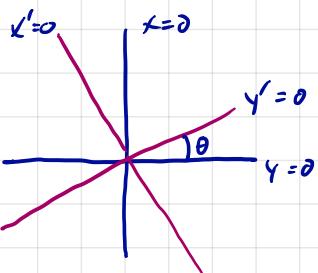
Last class:

2-d Lorentz transforms

O' moving with velocity v relative to O , intersect at

$$t=0 \quad x=0 \quad (ct'=0, x'=0).$$

$$\begin{pmatrix} ct' \\ x' \end{pmatrix} = \gamma \begin{pmatrix} 1 & -\frac{v}{c} \\ \frac{v}{c} & 1 \end{pmatrix} \begin{pmatrix} ct \\ x \end{pmatrix}$$



passive (relabelling)

$$\begin{pmatrix} x' \\ y' \end{pmatrix} = \begin{pmatrix} \cos(-\theta) & -\sin(-\theta) \\ \sin(-\theta) & \cos\theta \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix}$$

(e.g. $y' < 0$ if $x > 0, y = 0$)

(active)

On the other hand, the rotation that takes the xy axes to the x',y' axis is

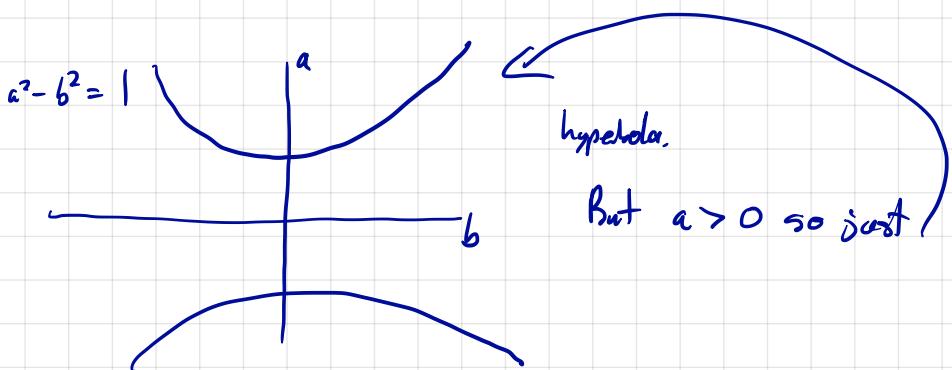
$$\begin{pmatrix} x' \\ y' \end{pmatrix} = \begin{pmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix}$$

I want a similar picture for Lorentz transformations.

$$L_v = \gamma \begin{pmatrix} 1 & v \\ v & 1 \end{pmatrix} = \begin{pmatrix} a & b \\ b & a \end{pmatrix} \quad \left(\begin{array}{l} \text{transformation that} \\ \text{maps } O's \text{ configuration onto } O' \text{'s} \end{array} \right)$$

$$\begin{aligned} a^2 - b^2 &= \gamma^2 \left[1 - \left(\frac{v}{c}\right)^2 \right] \\ &= \frac{1}{1 - \left(\frac{v}{c}\right)^2} \left(1 - \left(\frac{v}{c}\right)^2 \right) = 1 \end{aligned}$$

[those maps are "area" preserving. Spacetime has a notion of area or volume]



Moreover, you are familiar with $\cos^2 \theta + \sin^2 \theta = 1$.

If $a^2 + b^2 = 1$ there is a unique $\theta \in [0, 2\pi)$ with
 $a = \cos \theta$, $b = \sin \theta$.

Recall the hyperbolic trig functions:

$$\cosh z = \frac{e^z + e^{-z}}{2} \quad \sinh z = \frac{e^z - e^{-z}}{2}$$

$$\text{Exercise: } \cosh^2 z - \sinh^2 z = 1$$

In fact, if $a^2 - b^2 = 1$ there is a unique $\gamma \in \mathbb{R}$

$$a = \cosh \gamma$$

$$b = \sinh \gamma$$

$$\text{In fact: } \frac{b}{a} = \tanh(\gamma), \quad \gamma = \operatorname{arctanh}(b/a)$$

$$\text{So } Y = \cosh(\gamma) \quad \gamma \text{ rapidly.}$$

Recall the angle sum formula

$$e^{i\theta} = \cos \theta + i \sin \theta$$

$$e^{i(\theta_1 + \theta_2)} = \cos \theta_1 + \theta_2 + i \sin (\theta_1 + \theta_2)$$

↓

$$(\cos \theta_1 + i \sin \theta_1)(\cos \theta_2 + i \sin \theta_2)$$

$$= [\cos \theta_1 \cos \theta_2 - \sin \theta_1 \sin \theta_2]$$

$$+ i [\sin \theta_1 \cos \theta_2 + \sin \theta_2 \cos \theta_1]$$

Exercise:

$$\cosh(z_1 + z_2) = \cosh(z_1) \cosh(z_2) + \sinh(z_1) \sinh(z_2)$$

$$\sinh(z_1 + z_2) = \sinh(z_1) \cosh(z_2) + \sinh(z_2) \cosh(z_1)$$

(Do this from the definition).

$$\mathcal{L}_\phi = \begin{pmatrix} \cosh \phi & \sinh \phi \\ \sinh \phi & \cosh \phi \end{pmatrix}$$

$$\mathcal{L}_{\phi_2} \mathcal{L}_{\phi_1} = \begin{pmatrix} C_2 & S_2 \\ S_2 & C_2 \end{pmatrix} \begin{pmatrix} C_1 & S_1 \\ S_1 & C_1 \end{pmatrix}$$

$$= \begin{pmatrix} C_1 C_2 + S_1 S_2 & C_2 S_1 + S_2 C_1 \\ C_2 S_1 + S_2 C_1 & C_1 C_2 + S_1 S_2 \end{pmatrix}$$

$$= \begin{pmatrix} \cosh(\phi_1 + \phi_2) & \sinh(\phi_1 + \phi_2) \\ \sinh(\phi_1 + \phi_2) & \cosh(\phi_1 + \phi_2) \end{pmatrix}$$

$$= \mathcal{L}_{\phi_1 + \phi_2}$$

The Lorentz matrices are closed under mult.

They contain the id, \mathcal{L}_0 . $\mathcal{L}_{-\phi} = \mathcal{L}_\phi^{-1}$ by above.

So they form a group.

$SO(1,1)$ (proper, orthochronous Lorentz transformations)
 we'll see why shortly.

Suppose O'' is traveling with velocity v relative to O' ,
 and O' is traveling with velocity 0
 and all agree on a common origin.

What is the transformation from O to O''
 coordinates?

$$\varphi_1 = \operatorname{arctanh}\left(-\frac{v}{c}\right)$$

$$\varphi_2 = \operatorname{arctanh}\left(-\frac{w}{c}\right)$$

$$\begin{pmatrix} t'' \\ x'' \end{pmatrix} = \begin{pmatrix} c_2 & s_2 \\ s_2 & c_2 \end{pmatrix} \begin{pmatrix} c_1 & s_1 \\ s_1 & c_1 \end{pmatrix} \begin{pmatrix} t \\ x \end{pmatrix}$$

$$= \begin{pmatrix} c_3 & s_3 \\ s_3 & c_3 \end{pmatrix} \begin{pmatrix} t \\ x \end{pmatrix}$$

$$c_3 = \cosh(\varphi_1 + \varphi_2)$$

$$s_3 = \sinh(\varphi_1 + \varphi_2)$$

So O'' is moving w.r.t. O with velocity z

$$\tanh(\gamma_1 + \gamma_2) = \left(-\frac{z}{c}\right)$$

$$\begin{aligned}-\frac{z}{c} &= \tanh(\gamma_1 + \gamma_2) = \frac{\sinh(\gamma_1 + \gamma_2)}{\cosh(\gamma_1 + \gamma_2)} \\&= \frac{s_1 c_2 + s_2 c_1}{c_1 c_2 + s_1 s_2} \\&= \frac{s_1 c_1 + s_2 c_2}{1 + \frac{s_1}{c_1} \frac{s_2}{c_2}}\end{aligned}$$

$$\frac{s_1}{c_1} = \tanh(\gamma_1) = -\frac{v}{c}$$

$$\frac{s_2}{c_2} = \tanh(\gamma_2) = -\frac{w}{c}$$

$$\frac{z}{c} = \frac{\frac{v}{c} + \frac{w}{c}}{1 + \left(\frac{v}{c}\right)\left(\frac{w}{c}\right)}$$

← addition of velocities
↓
↓ last term pair

Exercise: If $a, b \in (-1, 1)$

$a+b < 1$ true

$$1 + \frac{b}{a} \leq \frac{1}{a} + b \quad (\Rightarrow 0)$$

$$\frac{b-1}{a} \leq b-1 \quad \checkmark$$

$$\frac{a+b}{1+ab} \in (-1, 1) \text{ also.}$$

Similarly, suppose O' is moving with velocity v w.r.t. O

And it observes a particle traveling with velocity w .

What velocity does O observe?

$$\begin{pmatrix} ct' \\ x' \end{pmatrix} = \begin{pmatrix} ct' \\ vt' + x'_0 \end{pmatrix} \quad \text{assuming at } x'_0 \text{ when } t'=0$$

$$= \begin{pmatrix} ct' \\ \left(\frac{v}{c}\right)ct' + x'_0 \end{pmatrix}$$

$$\begin{pmatrix} ct \\ x \end{pmatrix} = \begin{pmatrix} c & s \\ s & c \end{pmatrix} \begin{pmatrix} ct' \\ \left(\frac{v}{c}\right)ct' + x'_0 \end{pmatrix}$$

What matters: $\begin{pmatrix} c & s \\ s & c \end{pmatrix} \begin{pmatrix} ct' \\ \frac{v}{c}ct' \end{pmatrix} = \begin{bmatrix} c + s\left(\frac{v}{c}\right) \\ s + c\left(\frac{v}{c}\right) \end{bmatrix} ct'$

$$\begin{pmatrix} c & s \\ s & c \end{pmatrix} L_w \begin{pmatrix} 1 \\ 0 \end{pmatrix} ct' = L_v L_w \begin{pmatrix} 1 \\ 0 \end{pmatrix} ct'$$

$$\frac{dx}{dt} = \frac{dx}{dt'} \cdot \frac{dt'}{dt} = \frac{[s + c(\omega_0)]c}{[c + s(\omega_0)]} = \frac{s/c + (\omega_0)c}{1 + \frac{s}{c}\omega_0}$$

$$= \frac{\frac{v}{c} + \frac{\omega}{c} \cdot c}{1 + \frac{v}{c}\omega}$$

$$= \frac{v + \omega c}{1 + \frac{v}{c}\omega}$$

(Velocity addition formula again)

$$F = (E_2, x_2)$$

$$E = (E_1, x_1)$$

Interval: $c^2(t_2 - t_1)^2 - (x_2 - x_1)^2 = I(E, F)$

$$X = \begin{pmatrix} dt_2 - dt_1 \\ x_2 - x_1 \end{pmatrix} = C(F - E)$$

Interval: $X^T \underbrace{\begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}}_G X$

Coordinates for O' $\begin{pmatrix} c\theta \\ s\theta \\ x_i' \end{pmatrix} = L_{-v} \begin{pmatrix} c\theta \\ s\theta \\ x \end{pmatrix}$

$$X' = L_{-v} X$$

$$(X')^T G X = X^T \underbrace{L_{-v}^T G L_{-v}}_{} X$$

$$\begin{aligned} &\xrightarrow{\text{Simplifying}} \begin{pmatrix} c & -s \\ -s & c \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} c & -s \\ -s & c \end{pmatrix} = \begin{pmatrix} c & -s \\ -s & c \end{pmatrix} \begin{pmatrix} c & -s \\ -s & c \end{pmatrix} \\ &= \begin{pmatrix} c^2 - s^2 & 0 \\ 0 & -c^2 + s^2 \end{pmatrix} \\ &= \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \end{aligned}$$

$$= X^T G X.$$

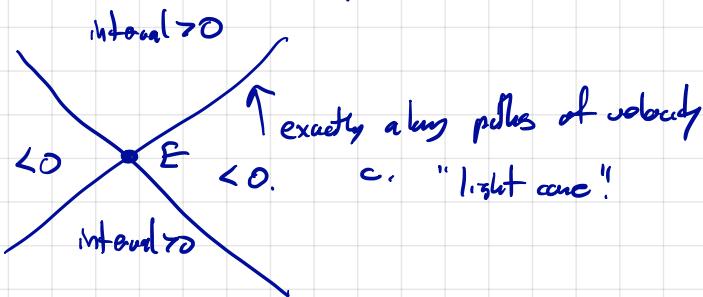
The interval between two points, in any of these coordinate systems, is

$$(c\Delta t)^2 - (\Delta x)^2$$

This is the fundamental quantity of 2-d spacetime that replaces the notion of distance in Euclidean geometry.

Let us suppose E is origin.

$$I(O, F) = 0 \text{ when } |\Delta x| = c|\Delta t|$$



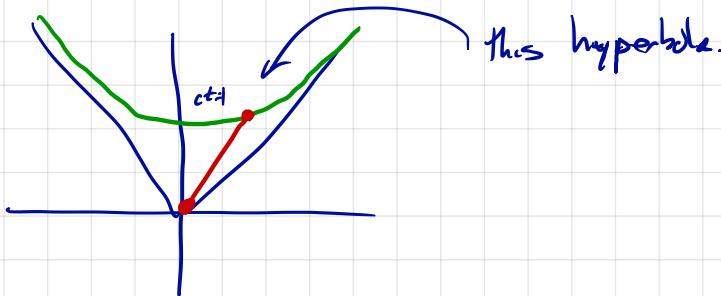
Where are all points with interval = 1?

Here we some. Start with $\begin{pmatrix} 1 \\ 0 \end{pmatrix}$. Now send through φ .

$$\begin{bmatrix} c & s \\ s & c \end{bmatrix} \begin{bmatrix} 1 \\ 0 \end{bmatrix} = \begin{bmatrix} c \\ s \end{bmatrix}$$

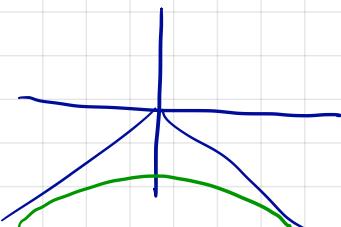
Recall: $c^2 - s^2 = 1$

$$(ct)^2 - x^2 = 1$$



We missed some:

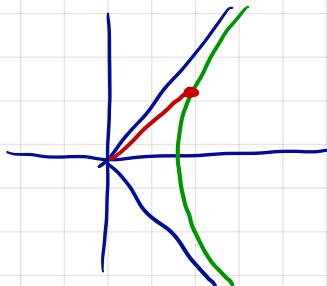
$$\begin{pmatrix} ct \\ s \end{pmatrix} = \begin{pmatrix} -1 \\ 0 \end{pmatrix} \text{ also has interval 1!}$$



The transformations we are working with only take the upper branch to the upper branch, etc.

What about -1 ? Start with $\begin{bmatrix} 0 \\ 1 \end{bmatrix}$

Now send through \mathcal{T}_f to get $\begin{bmatrix} s \\ c \end{bmatrix}$ $s^2 - c^2 = -1 \checkmark$



$$c > 0$$

so always
stays on right
side.

Similarly for left-hand branch.