

Last class:

$\alpha(s)$: curve in spacetime

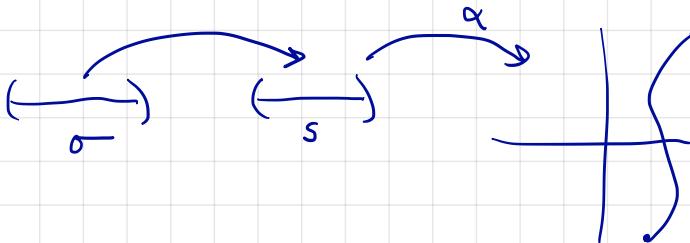
$$\alpha(s) = \begin{bmatrix} ct(s) \\ \vec{x}(s) \end{bmatrix}$$

$\alpha'(s)$ transforms as a vector. A curve is causal if α' always is.

$$\int_a^b \frac{1}{c} \sqrt{g(\alpha', \alpha')} ds = \Delta\tau, \quad \text{change in proper time along curve.}$$

elapsed time for the traveller from $s=a$ to $s=b$.

A curve can be reparameterized:



$$\alpha(s(\sigma))$$

We require $\frac{ds}{d\sigma} > 0$ so that

there is no double backtracking.

between points
along a causal curve

Theorem: Proper time \int is independent of the parametrization.

Pf: Let $\alpha(s)$ be a causal curve and consider the proper time between $\alpha(s_0)$ and $\alpha(s_1)$.

Let

$$\beta(\sigma) = \alpha(s(\sigma)).$$

be a reparameterization with $s_0 = s(\sigma_0)$, $s_1 = s(\sigma_1)$.

Then $\beta'(\sigma) = \alpha'(s(\sigma)) \frac{ds}{d\sigma}$ and

$$\sqrt{g(\beta', \beta')} = \sqrt{g(\alpha', \alpha')} \frac{ds}{d\sigma}, \text{ using } \frac{ds}{d\sigma} > 0.$$

$$\begin{aligned} \text{By change of vars, } \int_{s_0}^{s_1} \frac{1}{c} \sqrt{g(\beta', \beta')} d\sigma &= \int_{\sigma_0}^{\sigma_1} \frac{1}{c} \sqrt{g(\alpha', \alpha')} \frac{ds}{d\sigma} d\sigma \\ &= \int_{s_0}^{s_1} \frac{1}{c} \sqrt{g(\alpha', \alpha')} ds. \end{aligned}$$

Def: A causal curve $\alpha^{(s)}$ is parametrized by proper time

if $g(\alpha', \alpha') = -c^2$ along the curve.

Note, for such a curve,

$$\int_a^b \frac{1}{c} \sqrt{g(\alpha', \alpha')} ds = \int_a^b 1 ds = b-a.$$

How much proper time elapses from $s=a$ to $s=b$? $b-a$.

The parameter encodes this.

If α is a curve in the coord system and $\begin{bmatrix} \vec{x} \\ \vec{\tau} \end{bmatrix} = C^{-1} L C \begin{bmatrix} \vec{x} \\ \vec{\tau} \end{bmatrix} + \vec{z}$,

then $\hat{\alpha} = L\alpha + \vec{z}$. (The C 's are built in).

So $\hat{\alpha}' = L\alpha'$. In particular,

$$g(\hat{\alpha}', \hat{\alpha}') = g(L\alpha', L\alpha') = g(\alpha', \alpha')$$

So if α is parameterized by arclength so is $\hat{\alpha}$.

A geometric analogy:

Def: A curve in the plane (\mathbb{R}^2) or \mathbb{R}^3 parameterized by arc length if $|\alpha'(s)|=1$ for all s .

e.g. $\alpha(s) = [R \cos(s/R), R \sin(s/R)]$ is parameterized by arc length. $\alpha'(s) = [\cos(s/R), \sin(s/R)]$.

How long to go around the circle? $2\pi R$.

How far around the circle? $2\pi R$.

As s goes up by $\frac{1}{r}$, so does arc length

Given any curve in the plane with $\alpha' \neq 0$

we can reparameterize it by arclength.

$$\beta(\sigma) = \alpha(s(\sigma))$$

$$\beta'(\sigma) = \alpha'(s(\sigma)) s'(\sigma)$$

$$1 = |\alpha'(s(\sigma))| \frac{ds}{d\sigma}$$

$$s_0 \quad \frac{ds}{d\sigma} = \frac{1}{|\alpha'(s(\sigma))|} \quad (\text{use } |\alpha'(s)| \neq 0)$$

which is an ODE to solve for s . In fact

$$|\alpha'(s)| ds = d\sigma$$

$$s_0 \quad \sigma = \int_{s_0}^s |\alpha'(t)| dt.$$

Now solve for $s(\sigma)$. \rightarrow possible, in principle: $\frac{ds}{d\sigma} > 0$.

But in practice, this is awful.

- 1) Solve an impossible integral
- 2) Solve an impossible algebraic eq.

If given a timelike curve, we can always parameterize it by proper time.

Now, given my curve $\alpha(s)$, we can determine its tangent vector at any point if it were parameterized by proper time. It's just

$$c \frac{\alpha'(s)}{\sqrt{g(\alpha', \alpha')}} \quad \text{which has length } c.$$

This rescaling is the infinitesimal version of reparameterizing.

Def The 4-velocity of a timelike curve $\alpha(s)$ is

$$c \frac{\alpha'(s)}{\sqrt{g(\alpha', \alpha')}}$$

It's just the velocity of a reparameterized curve.

When α' is timelike, I'll write $|\alpha'|$ rather than $\sqrt{-g(\alpha', \alpha')}$

E.g. $\alpha(t) = \begin{bmatrix} c \\ x(t) \end{bmatrix}$, parameterized by coordinate time.

$$\begin{aligned} |\alpha'|^2 &= g(\alpha', \alpha') = c^2 - |\mathbf{x}'|^2 \\ &= c^2 - |\mathbf{v}|^2 \\ &= c^2 (1 - \frac{|\mathbf{v}|^2}{c^2}) \\ &= c^2 \gamma^{-2} \end{aligned}$$

$$\begin{aligned} \frac{c\alpha'}{|\alpha'|} &= \frac{c\alpha'}{c\gamma^{-1}} = \gamma\alpha' = \gamma \begin{bmatrix} c \\ x'(t) \end{bmatrix} \\ &= \gamma \begin{bmatrix} c \\ \vec{v} \end{bmatrix}. \end{aligned}$$

Your text uses notation \mathbf{V} for $\alpha(s)$'s 4-velocity.

The 4 is old-fashioned.

And although reparametrizing is hard,

computing $\frac{d\gamma}{ds}$ is easy.

$$\beta(z) = \alpha(s(z))$$

$$c = |\beta'(z)| = |\alpha'(s(z))| \frac{ds}{dz}$$

$$\boxed{\frac{d\gamma}{ds} = \frac{|\alpha'(s)|}{c}}$$

$$\frac{ds}{dz} = \frac{c}{|\alpha'|}$$

If the curve is parameterized by coordinate time t

$$\alpha' = \begin{bmatrix} c \\ \vec{v} \end{bmatrix} \quad \text{and} \quad |\alpha'|^2 = c^2(1 - \frac{\vec{v}^2}{c^2})$$

$$|\alpha'| = \frac{c}{\gamma(\vec{v})}$$

$$\boxed{\frac{d\gamma}{dt} = \gamma(\vec{v})}$$

$$\frac{dt}{dz} = \gamma(|\vec{v}|)$$

expresses
time dilation.

E.g. Radial motion

$$\vec{v}(t) = \begin{bmatrix} ct \\ R\cos(\omega t) \\ R\sin(\omega t) \\ 0 \end{bmatrix}$$

$$\omega \leq \frac{c}{R}$$

for otherwise you travel $2\pi R$

in time $\frac{2\pi}{\omega}$ with speed $R\omega$.

$$\begin{aligned}\frac{d\vec{r}}{dt} &= \frac{1}{c} \left| \frac{d\vec{r}}{dt} \right| = \frac{1}{c} \sqrt{c^2 - R^2 \omega^2} \\ &= \sqrt{1 - \left(\frac{R\omega}{c} \right)^2}\end{aligned}$$

Or: $|v|^2 = (R\omega)^2$ so $\frac{d\vec{r}}{dt} = \gamma(|v|) =$

How much proper time elapses in a single rotation?

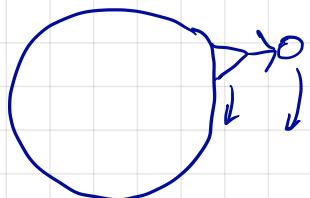
Time for a rotation: $\frac{2\pi}{\omega}$

$$\frac{2\pi}{\omega} \sqrt{1 - \left(\frac{R\omega}{c} \right)^2} \approx \frac{2\pi}{\omega}, \text{ if } |R\omega| \ll c.$$

(Consequence:

Orbiting at light speed, time does not pass

(Consequence



Head traveling faster.

Your head is younger than your feet

Next HW: how much younger?

Curvature of a curve in the plane

α'' mixes up what the curve is doing with how you parameterize it.

How curvy is it? a) Reparam by arc length

b) $|\alpha''|$ tells you curvature.

e.g. $\alpha(s) = (R \cos(\omega s), R \sin(\omega s))$



$$\alpha'' = -\omega^2 (R \cos(\omega s), R \sin(\omega s))$$

$$|\alpha''| = \omega^2 R$$

↑ ↑
part for the curve's *shape*
part for the parameterization.

We can eliminate the parameterization dependence by looking at curves parameterized by arc length.

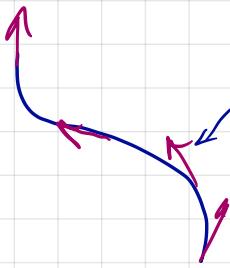
$$\omega = \frac{1}{R}$$

Now $\alpha'' = -\frac{1}{R^2} \alpha$ and $|\alpha''| = \frac{1}{R}$

Tiny circle, huge $1/R$. Big circle, tiny $1/R$.

We call this quantity $|\alpha''|$ when α is formed by arc length the *curvature* of the curve.

Units: $1/L$.



always unit length.

How can α' change?

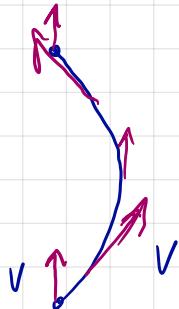
$$\alpha' \cdot \alpha' = 1$$

$$\text{so } \frac{d}{ds} \alpha' \cdot \alpha' = 0$$

$$\Leftrightarrow = 2\alpha' \cdot \alpha''$$

So α'' is always perp to α' . All α' can do is rotate.

Space like picture



$$|V| = c$$

Analogously, 4-velocity: $\frac{d\alpha}{d\tau} = V$

4-acceleration: $\frac{d^2\alpha}{d\tau^2}$

$$V(t) = c \frac{\alpha'}{|\alpha'|} \quad \frac{d\tau}{dt} = \frac{1}{c} |\alpha'|$$

$$\frac{dV}{d\tau} = \frac{dV}{dt} \cdot \underbrace{\frac{dt}{d\tau}}_{\gamma(|v|)} = \frac{dV}{dt} \cdot \frac{c}{|\alpha'|}$$

Typically a mess

$$\alpha = \begin{bmatrix} c \\ R\cos(\omega t) \\ R\sin(\omega t) \end{bmatrix}$$

$$|\alpha'| = \sqrt{c^2 - R\omega^2}$$

$$V = \frac{1}{\sqrt{1 - (\frac{R\omega}{c})^2}} \begin{bmatrix} c \\ -R\omega \sin \\ R\omega \cos \end{bmatrix}$$

$$\frac{dV}{d\tau} = \frac{1}{\sqrt{1 - (\frac{R\omega}{c})^2}} \begin{bmatrix} 0 \\ -R\omega^2 \cos(\omega t) \\ -R\omega^2 \sin(\omega t) \end{bmatrix} \frac{1}{\sqrt{1 - (\frac{R\omega}{c})^2}}$$

$$= \frac{1}{(1 - (\frac{R\omega}{c})^2)} \begin{bmatrix} 0 \\ -\omega^2 R \cos(\omega t) \\ -\omega^2 R \sin(\omega t) \end{bmatrix}$$

↑ ↑
relativistic correction factor usual acceleration
Newtonian

Is infinite if $R\omega = c$!

e.g. $\alpha(t) = \begin{bmatrix} ct \\ x(t) \end{bmatrix}$

$$\frac{dx}{dt} = \gamma(v) \begin{bmatrix} c \\ v \end{bmatrix}$$

$$\frac{d\gamma}{dt} = \frac{1}{c} \sqrt{c^2 - v^2}$$

$$\frac{d^2\alpha}{dt^2} = \frac{dt}{dz} \frac{d}{dt} \left[\gamma \begin{bmatrix} c \\ v \end{bmatrix} \right]$$

$$dt = \gamma dz$$

$$= \gamma \left[\gamma' \begin{bmatrix} c \\ v \end{bmatrix} + \gamma \begin{bmatrix} 0 \\ v' \end{bmatrix} \right]$$

$$= (\gamma\gamma') \begin{bmatrix} c \\ v \end{bmatrix} + \gamma^2 \begin{bmatrix} 0 \\ v' \end{bmatrix}$$

$$\frac{d}{dt} \frac{1}{\sqrt{1-\frac{v^2}{c^2}}} = -\frac{1}{2} \left(\frac{1}{1-\frac{v^2}{c^2}}\right)^{3/2} (-2) \frac{v}{c} \frac{v'}{c}$$

$$= \gamma^3 \frac{v}{c^2} \frac{dv}{dt}$$

$$\gamma \gamma' = \gamma^4 \frac{|v|}{c^2} \frac{d}{dt} |v|$$