Expressions for the Limber approximation

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1 Full-sky expressions

The angular power spectrum between two contributions is:

$$C_{\ell}^{ij} = 4\pi \int_0^\infty \frac{dk}{k} \, \mathcal{P}_{\Phi}(k) \Delta_{\ell}^i(k) \Delta_{\ell}^j(k). \tag{1}$$

The expressions for density, RSD, magnification, lensing convergence and CMB lensing are:

$$\Delta_{\ell}^{D}(k) = \int dz p_{z}(z)b(z)T_{\delta}(k,z)j_{\ell}(k\chi(z))$$
(2)

$$\Delta_{\ell}^{RSD}(k) = \int dz \frac{(1+z)p_z(z)}{H(z)} T_{\theta}(k,z) j_{\ell}^{"}(k\chi(z))$$
(3)

$$\Delta_{\ell}^{M}(k) = -\ell(\ell+1) \int \frac{dz}{H(z)} W^{M}(z) T_{\phi+\psi}(k,z) j_{\ell}(k\chi(z)), \tag{4}$$

$$\Delta_{\ell}^{L}(k) = -\frac{\ell(\ell+1)}{2} \int \frac{dz}{H(z)} W^{L}(z) T_{\phi+\psi}(k,z) j_{\ell}(k\chi(z)), \tag{5}$$

$$\Delta_{\ell}^{IA}(k) = -\sqrt{\frac{(\ell+2)!}{(\ell-2)!}} \int dz \, p_z(z) \, b_{IA}(z) \, T_{\delta}(k,z) \, \frac{j_{\ell}(k\chi(z))}{(k\chi(z))^2},\tag{6}$$

$$\Delta_{\ell}^{C}(k) = \frac{\ell(\ell+1)}{2} \int_{0}^{\chi_*} d\chi \frac{\chi_* - \chi}{\chi \chi_*} T_{\phi+\psi}(k, z) j_{\ell}(k\chi), \tag{7}$$

where

$$W^{M}(z) \equiv \int_{z}^{\infty} dz' p_{z}(z') \frac{2 - 5s(z')}{2} \frac{\chi(z') - \chi(z)}{\chi(z')}$$

$$\tag{8}$$

$$W^{L}(z) \equiv \int_{z}^{\infty} dz' p_{z}(z') \frac{\chi(z') - \chi(z)}{\chi(z')}$$
(9)

2 Limber approximation

The Limber approximation is

$$j_{\ell}(x) \simeq \sqrt{\frac{\pi}{2\ell+1}} \,\delta\left(\ell + \frac{1}{2} - x\right).$$
 (10)

Thus for each k and ℓ we can define a radial distance $\chi_{\ell} \equiv (\ell + 1/2)/k$

3 Expressions in the Limber approximation

The expressions above can be written as follows in the Limber approximation. First, the power spectrum can be rewritten as

$$C_{\ell}^{ij} = \frac{2}{2\ell+1} \int_0^\infty dk \, P_{\delta}(k, z=0) \, \tilde{\Delta}_{\ell}^i(k) \tilde{\Delta}_{\ell}^j(k). \tag{11}$$

where

$$\tilde{\Delta}_{\ell}^{D}(k) = p_{z}(\chi_{\ell}) b(\chi_{\ell}) D(\chi_{\ell}) H(\chi_{\ell})$$
(12)

$$\tilde{\Delta}_{\ell}^{RSD}(k) = \frac{1+8\ell}{(2\ell+1)^2} p_z(\chi_{\ell}) f(\chi_{\ell}) D(\chi_{\ell}) H(\chi_{\ell}) - \frac{4}{2\ell+3} \sqrt{\frac{2\ell+1}{2\ell+3}} p_z(\chi_{\ell+1}) f(\chi_{\ell+1}) D(\chi_{\ell+1}) H(\chi_{\ell+1}) \longrightarrow 0$$
(13)

$$\tilde{\Delta}_{\ell}^{ISW}(k) = \frac{3\Omega_{M,0}H_0^2}{k^2}H(\chi_{\ell})g(\chi_{\ell})\left[1 - f(\chi_{\ell})\right]$$
(14)

$$\tilde{\Delta}_{\ell}^{M}(k) = 3\Omega_{M,0}H_{0}^{2}\frac{\ell(\ell+1)}{k^{2}}\frac{D(\chi_{\ell})}{a(\chi_{\ell})\chi_{\ell}}W^{M}(\chi_{\ell}) \longrightarrow 3\Omega_{M,0}H_{0}^{2}\frac{\chi_{\ell}D(\chi_{\ell})}{a(\chi_{\ell})}W^{M}(\chi_{\ell})$$

$$\tag{15}$$

$$\tilde{\Delta}_{\ell}^{L}(k) = \frac{3}{2} \Omega_{M,0} H_0^2 \sqrt{\frac{(\ell+2)!}{(\ell-2)}} \frac{1}{k^2} \frac{D(\chi_{\ell})}{a(\chi_{\ell})\chi_{\ell}} W^M(\chi_{\ell}) \longrightarrow \frac{3}{2} \Omega_{M,0} H_0^2 \frac{\chi_{\ell} D(\chi_{\ell})}{a(\chi_{\ell})} W^M(\chi_{\ell})$$

$$\tag{16}$$

$$\tilde{\Delta}_{\ell}^{IA}(k) = p_z(\chi_{\ell}) \, b_I(\chi_{\ell}) \, D(\chi_{\ell}) \, H(\chi_{\ell}) \, \frac{\sqrt{(\ell+2)(\ell+1)\ell(\ell-1)}}{(\ell+1/2)^2} \tag{17}$$

$$\tilde{\Delta}_{\ell}^{C}(k) = \frac{3}{2} \Omega_{M,0} H_0^2 \frac{\ell(\ell+1)}{k^2} \frac{D(\chi_{\ell})}{a(\chi_{\ell})\chi_{\ell}} \frac{\chi_* - \chi_{\ell}}{\chi_*} \Theta(\chi_{\ell} - \chi_*) \longrightarrow \frac{3}{2} \Omega_{M,0} H_0^2 \frac{\chi_{\ell} D(\chi_{\ell})}{a(\chi_{\ell})} \frac{\chi_* - \chi_{\ell}}{\chi_*} \Theta(\chi_{\ell} - \chi_*)$$

$$\tag{18}$$