

HW 1 - ASTR404

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Q3) Binary star data

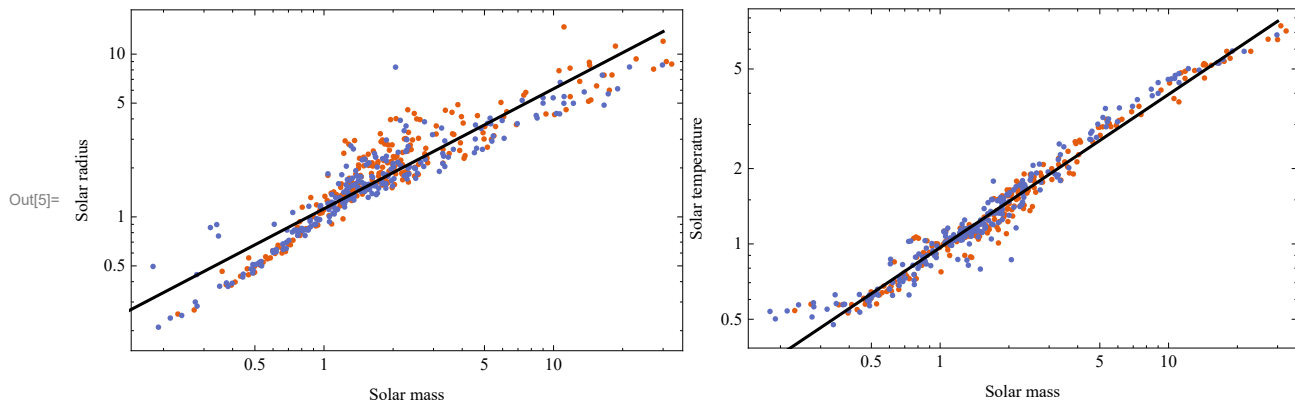
Importing data in solar units

```
In[1]:= tSun = QuantityMagnitude@Sun(star) ["EffectiveTemperature"];
SetDirectory["C:\\Users\\dan7g\\Google Drive\\Acads\\ASTR404\\"];
dataS =
  SemanticImport["eker_2014_simple.dat"] [All, {"T1" -> (#/tSun &), "T2" -> (#/tSun &)}];
```

a) & b) Radius and temperature vs mass

Scatter plots and best fits

```
In[3]:= SetOptions[ListLogLogPlot, PlotStyle -> PointSize[.01],
  PlotTheme -> "Scientific", Frame -> True, PlotRange -> All];
fit[s_] := LinearModelFit[Sort[Join@@ (Log10@Values@Normal@dataS[All, #] & /@
  {"M1", s <> "1"}, {"M2", s <> "2"})], {1, x}, x];
Show[ListLogLogPlot[Values@Normal@dataS[All, #] & /@ {"M1", #1 <> "1"}, {"M2", #1 <> "2"}],
  FrameLabel -> {"Solar mass", #2}, ImageSize -> 320],
  LogLogPlot[10^fit[#1][Log10[M]], {M, 0, 30}, PlotStyle -> Black]] &@@@
{"R", "Solar radius"}, {"T", "Solar temperature"} // Row
```



We can see that both the radius and the temperature increase monotonically with increase in mass of the stars. Moreover, they seem to fit a power law relationship since a linear fit to the log-log plot matches closely.

There is a larger spread in the radius when mass increases whereas temperature appears to converge to the fit with increasing mass.

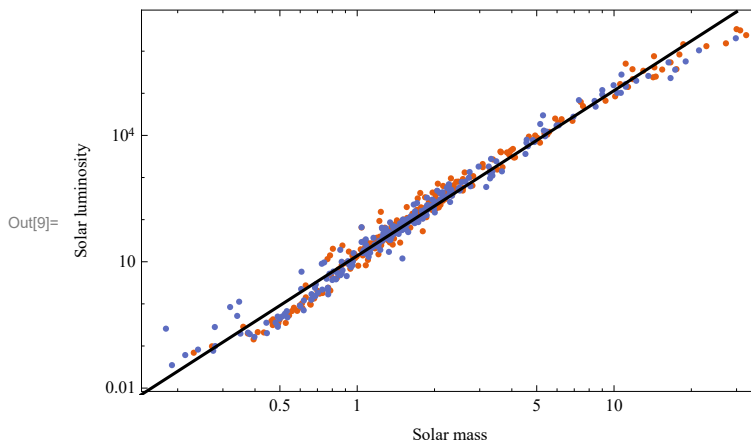
Fitting functions in solar units

```
In[6]:= # == 10^fit[#] [Log10@M] & /@ {"R", "T"} // Column
R == 1.1222945049359 M^0.73689311038015
Out[6]:= T == 0.96705300226301 M^0.6120798178381
```

c) Luminosity vs mass

Scatter plot along with best fit

```
In[7]:= dataL = Transpose[{Normal@dataS[All, "M" <> #],
    Normal@dataS[#^4 &, "T" <> #] * Normal@dataS[4 Pi #^2 &, "R" <> #]}] & /@ {"1", "2"};
fitL = LinearModelFit[Join@@Log10@dataL, {1, x}, x];
Show[ListLogLogPlot[dataL, FrameLabel -> {"Solar mass", "Solar luminosity"}],
    LogLogPlot[10^fitL[Log10@M], {M, 0, 30}, PlotStyle -> Black]
```



We can see that the luminosity also increases monotonically with increase in mass of the stars. Moreover, the luminosity appears to fit a power law relationship with mass since the linear fit to the log-log plot matches closely. The spread of the points appear to be higher for low and high masses.

Fitting function in solar units

```
In[10]:= "L" == 10^fitL[Log10[M]]
Out[10]:= L == 13.842846029206 M^3.9221050549238
```

Q2) Properties of Deneb

Defining parameters of Deneb

```
In[11]:= Rs = Quantity[7.72 × 10^10, "Meter"];
Ts = Quantity[8525., "Kelvin"];
ds = Quantity[440., "Parsecs"];
```

a) Luminosity

$$L = 4 \pi R^2 \sigma T^4$$

SI units

```
In[14]:= Ls = UnitConvert[4 Pi Rs^2 σ Ts^4, "SI"]
Out[14]= 2.2430201949436 × 10^31 W
```

Solar units

```
In[15]:= UnitConvert[Ls, "SolarLuminosity"]
Out[15]= 58 610.873655353 L⊙
```

b) Parallax and distance modulus

Parallax

$$\text{Parallax} = \text{Diameter (in AU)} / \text{Distance (in Parsecs)}$$

```
In[16]:= Quantity[QuantityMagnitude[2 Rs, "AU"] / QuantityMagnitude[ds, "Parsecs"], "ArcSeconds"]
Out[16]= 0.0023456823901778"
```

Distance modulus

$$\mu = 5 \log_{10}(d \text{ in parsecs}) - 5$$

```
In[17]:= 5 Log10[QuantityMagnitude[ds, "Parsecs"]] - 5
Out[17]= 8.2172633824309
```

c) Radiant flux at star's surface

Flux is luminosity divided by surface area of star:

```
In[18]:= UnitConvert[Ls / (4 Pi Rs^2), "SI"]
```

```
Out[18]= 2.9949440880982 × 108 N / (m s)
```

d) Radiant flux at earth's surface

Luminosity divided by area of sphere with radius equal to distance to earth:

```
In[19]:= UnitConvert[Ls / (4 Pi ds^2), "SI"]
```

```
Out[19]= 9.6831548346327 × 10-9 N / (m s)
```

e) Peak wavelength

Looking up Wien's displacement law

```
In[20]:= eqW = FormulaData[{"WienDisplacementLaw", "Wavelength"}]
```

```
Out[20]= λmax ==  $\frac{1 b}{T}$ 
```

Substituting value of temperature

```
In[21]:= eqW[[1]] == UnitConvert[eqW[[2]] /. T -> Ts, "meters"]
```

```
Out[21]= λmax == 3.3991470967742 × 10-7 m
```

Q1) Planck's law limits

a) Rayleigh-Jeans law

Defining Planck function

```
In[22]:= Bv = 2 h c^2 / (Exp[h c / k / T / λ] - 1) / λ^5
```

```
Out[22]= 
$$\frac{2 c^2 h}{\left(-1 + e^{\frac{c h}{k T \lambda}}\right) \lambda^5}$$

```

Taylor expansion for λ^{-1}

```
In[23]:= TE1 = Normal@Series[Bv /. λ -> 1 / x, {x, 0, 10}] /. x -> 1 / λ
```

```
Out[23]= 
$$\frac{c^2 h^6}{15 120 k^5 T^5 \lambda^{10}} - \frac{c^5 h^4}{360 k^3 T^3 \lambda^8} + \frac{c^3 h^2}{6 k T \lambda^6} - \frac{c^2 h}{\lambda^5} + \frac{2 c k T}{\lambda^4}$$

```

Substituting $\lambda = n \cdot h \cdot c / (k \cdot T)$

In[24]:= TE1 /. $\lambda \rightarrow n \cdot h \cdot c / (k \cdot T)$

$$\text{Out[24]} = \frac{k^5 T^5}{15 120 c^3 h^4 n^{10}} - \frac{k^5 T^5}{360 c^3 h^4 n^8} + \frac{k^5 T^5}{6 c^3 h^4 n^6} - \frac{k^5 T^5}{c^3 h^4 n^5} + \frac{2 k^5 T^5}{c^3 h^4 n^4}$$

Therefore for $n \gg 1$, we can approximate the equation using just the last term which corresponds to the following term in the Taylor series:

$$\frac{2 c k T}{\lambda^4}$$

b) Wien's function

Substituting $\lambda = n \cdot h \cdot c / (k \cdot T)$

In[25]:= Bv /. $\lambda \rightarrow n \cdot h \cdot c / (k \cdot T)$

$$\text{Out[25]} = \frac{2 k^5 T^5}{c^3 \left(-1 + e^{\frac{1}{n}} \right) h^4 n^5}$$

Taking limit

In the limit $n \ll 1$, $e^{1/n} - 1 \sim e^{1/n}$. Therefore we can ignore the 1 in the denominator. This gives:

$$\frac{2 c^2 e^{-\frac{c h}{k T \lambda}} h}{\lambda^5}$$

Therefore the constants are:

$$a = 2 c^2 h$$

$$b = \frac{c h}{k}$$