

HW 4 - ASTR404

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Q1)

Equations for pressure and energy

```
In[1465]:= eqs = {P -> n k T + 1/3 a T^4, u -> 3/2 n k T + a T^4}; eqs // Column // TraditionalForm
```

```
Out[1465]//TraditionalForm=
```

$$P \rightarrow \frac{a T^4}{3} + k n T$$
$$u \rightarrow a T^4 + \frac{3 k n T}{2}$$

Relation between pressure and energy

```
In[1466]:= eqγ = u == P / (γ - 1)
```

```
Out[1466]= u == \frac{P}{-1 + \gamma}
```

Formula for γ

```
In[1467]:= sγ = Solve[eqγ, γ][[1, 1]] /. eqs // Simplify
```

```
Out[1467]= γ -> \frac{15 k n + 8 a T^3}{9 k n + 6 a T^3}
```

a) Limit of γ as $T \rightarrow \infty$

```
In[1468]:= Limit[sγ[[2]], T -> ∞]
```

```
Out[1468]= \frac{4}{3}
```

b) Limit of γ as $T \rightarrow 0$

```
In[1469]:= Limit[sγ[[2]], T -> 0]
```

```
Out[1469]= \frac{5}{3}
```

Q2)

Solving for i_{out} & i_{In}

Equations

```
In[1470]:= eqs2 = {J == (iout + iIn) / 2, F == Pi (iout - iIn) == σ T^4, J == 3 F / (4 Pi) (τv + 2 / 3)} // Echo;
```

$$\gg \left\{ J = \frac{1}{2} (i_{In} + i_{out}), F = \pi (-i_{In} + i_{out}) = T^4 \sigma, J = \frac{3 F \left(\frac{2}{3} + \tau_v \right)}{4 \pi} \right\}$$

Solution

```
In[1471]:= Quiet@Solve[eqs2, {iIn, iout, J, F}][[1, ;; 2]]
```

$$\text{Out[1471]} = \left\{ i_{In} \rightarrow \frac{3 T^4 \sigma \tau_v}{4 \pi}, i_{out} \rightarrow \frac{T^4 \sigma (4 + 3 \tau_v)}{4 \pi} \right\}$$

Depth when isotropy is 1%

Equations

```
In[1472]:= eqs3 = Append[eqs2, 2 (iout - iIn) / (iout + iIn) == 1 / 100 // Echo];
```

$$\gg \frac{2 (-i_{In} + i_{out})}{i_{In} + i_{out}} = \frac{1}{100}$$

Solution

```
In[1473]:= NSolve[eqs3, {τv, iIn, iout, J, F}][[1, 1]]
```

```
Out[1473]= τv → 132.666666666667
```

Q3)

Solving the radiative transfer equation

```
In[1474]:= sRad = i[λ] -> iλ0 E^-τλ0 - Integrate[S[λ] E^-τ[λ], {τ[λ], τλ0, 0}]
```

```
Out[1474]= i[λ] → e^-τλ0 iλ0 - S[λ] (-1 + Cosh[τλ0] - Sinh[τλ0])
```

a) Limit $\tau_{\lambda 0} \gg 1$

```
In[1475]:= Limit[i[λ] /. sRad, τλ0 → Infinity]
```

```
Out[1475]= S[λ]
```

At thermodynamic equilibrium $S(\lambda) = B(\lambda)$.

b) Limit $\tau_{\lambda 0} \ll 1$

Taylor series

In[1476]:= **Series**[**i**[λ] /. **sRad**, { $\tau\lambda\theta$, 0, 1}]

Out[1476]= $i\lambda\theta + (-i\lambda\theta + S[\lambda]) \tau\lambda\theta + O[\tau\lambda\theta]^2$

Thus if $S(\lambda) > I_{\lambda 0}$, we can see that emission lines will be added to the incident light at specific frequencies and vice versa.

Q4)

Solving the stellar structure equations

Out[1456]=
$$\begin{aligned} M'(r) &= 4\pi\rho r^2 \\ P'(r) &= -\frac{G\rho M(r)}{r^2} \end{aligned}$$

Assuming $\rho = \text{constant}$

In[1477]:= **DSolve**[{**M'**[r] == 4 π $r^2 \rho$, **D**[**P**[r], r] == -**G** **M**[r] ρ / r^2 , **M**[0] == 0, **P**[R] == 0}, {**M**[r], **P**[r]}, r] [[1, 2]] // **FullSimplify**

Out[1477]= $P[r] \rightarrow \frac{2}{3} G \pi (-r^2 + R^2) \rho^2$

Q5)

Solving the Lane-Emden equation

$$\frac{\frac{\partial}{\partial \xi} \left(\xi^2 \frac{\partial \theta(\xi)}{\partial \xi} \right)}{\xi^2} + \theta(\xi)^n = 0$$

Assuming $n = 0$ and solving

In[1478]:= **se** = **DSolve**[{ $\xi^{-2} D[\xi^2 D[\theta[\xi], \xi], \xi] + 1 == 0$, $\theta[0] == 1$, $\theta'[0] == 0$ }, $\theta[\xi]$, ξ] [[1, 1]]

Out[1478]= $\theta[\xi] \rightarrow \frac{1}{6} (6 - \xi^2)$

Location of first zero

In[1479]:= **Solve**[$\theta[\xi] == 0 \&\& \xi \geq 0$ /. **se**, ξ] [[1, 1]]

Out[1479]= $\xi \rightarrow \sqrt{6}$