

HW 4 - ASTR540

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Q1

a)

Given equation

$$\text{In[174]:= eqPr} = P \rightarrow -1/3 \text{ Egr} / V$$

$$\text{Out[174]= } P \rightarrow -\frac{\text{Egr}}{3 V}$$

Solving for R given M and ρ

$$\text{In[175]:= sR} = \\ R \rightarrow \text{Assuming}[M / \rho > 0, \text{Abs@ToRadicals@Solve}[M / (4/3 \pi R^3) == \rho, R, \text{Reals}][[1, 1, 2]] // \text{Simplify}]$$

$$\text{Out[175]= } R \rightarrow \frac{\left(\frac{3}{\pi}\right)^{1/3} \left(\frac{M}{\rho}\right)^{1/3}}{2^{2/3}}$$

Substituting for R in the first equation

$$\text{In[193]:= sP} = \text{eqPr} // . \{ \text{Egr} \rightarrow -G M^2 / R, V \rightarrow 4/3 \pi R^3, \text{sR} \} // \text{ToRadicals}$$

$$\text{Out[193]= } P \rightarrow \frac{2^{2/3} G M^2 \left(\frac{\pi}{3}\right)^{1/3}}{3 \left(\frac{M}{\rho}\right)^{4/3}}$$

This is equivalent to the given expression.

b)

Given that the pressures are equal

$$\text{Out[170]= } \frac{a T^4}{3} = \frac{k \rho T}{m}$$

Solving for T

In[177]:= **sT = Solve** $\left[\frac{1}{3} a T^4 == \rho k T / m, T\right]$ [[3, 1]]

Out[177]= $T \rightarrow \frac{3^{1/3} k^{1/3} \rho^{1/3}}{a^{1/3} m^{1/3}}$

Substituting T in the total pressure

In[198]:= **sP2 = P -> 2 × 1 / 3 a T^4 /. sT**

Out[198]= $P \rightarrow \frac{2 \times 3^{1/3} k^{4/3} \rho^{4/3}}{a^{1/3} m^{4/3}}$

This is equal to the required expression.

c)

Equating the two pressures

In[201]:= **eqM = Equal @@ (P /. {{sP}, {sP2}})**

Out[201]=
$$\frac{2^{2/3} G M^2 \left(\frac{\pi}{3}\right)^{1/3}}{3 \left(\frac{M}{\rho}\right)^{4/3}} == \frac{2 \times 3^{1/3} k^{4/3} \rho^{4/3}}{a^{1/3} m^{4/3}}$$

Solving for M

In[255]:= **sM = Solve** $[eqM, M]$ [[1, 1]] // **FullSimplify**

Out[255]= $M \rightarrow 9 \sqrt{\frac{6}{\pi}} \left(\frac{k^{4/3}}{a^{1/3} G m^{4/3} \rho^{2/3}}\right)^{3/2} \rho$

Substituting numerical values

In[209]:= **UnitConvert** $\left[sM[[2]] /. \{k \rightarrow k, m \rightarrow 1/2 m_p, G \rightarrow G, a \rightarrow a, \rho \rightarrow \text{Sun (star)}["Density"]\}, "SolarMass"\right]$

Out[209]= 113.7 M_{\odot}

Q2

a)

Formula for mean mass

In[222]:= **eqμ = 1/μ == mH*(xC/mC + xe/me);**

Substituting numbers

```
In[232]:= eqμ /. {mH -> 1 mp, me -> 1 me, xC -> 1, xe -> 6 me / mC, mC -> 12 mp}
```

$$\text{Out[232]} = \frac{1}{\mu} == \frac{7}{12}$$

This gives the required mean mass of 12/7.

b)

Pressure from the virial theorem

```
In[298]:= eqPv = P == 1 / 3 G M^2 / (R M / ρ)
```

$$\text{Out[298]} = P == \frac{G M \rho}{3 R}$$

Equating it to thermal pressure

```
In[245]:= eqR = eqPv /. {P -> ρ k T / m, ρ -> M / (4 / 3 Pi R^3)}
```

$$\text{Out[245]} = \frac{3 k M T}{4 m \pi R^3} == \frac{G M^2}{4 \pi R^4}$$

Solving for R

```
In[296]:= sR2 = Solve[eqR, R][[1, 1]]
```

$$\text{Out[296]} = R \rightarrow \frac{G m M}{3 k T}$$

Ratio of radius with respect to Sun

```
In[305]:= rS = Divide @@
```

$$\left(\text{sR2}[[2]] /. \{G \rightarrow G, k \rightarrow k\} /. \left\{ \left\{ M \rightarrow 10 \text{ Sun(star)}["\text{Mass}"], m \rightarrow 12/7 m_p, T \rightarrow 6 \times 10^8 \text{ K} \right\}, \right. \right. \\ \left. \left. \left\{ M \rightarrow \text{Sun(star)}["\text{Mass}], m \rightarrow 1/2 m_p, T \rightarrow 1.5 \times 10^7 \text{ K} \right\} \right\} \right)$$

$$\text{Out[305]} = 0.85714285714286$$

c)

Finding temperature given luminosity

```
In[314]:= Solve[10^7 L⊙ == 4 Pi (rS 1 R⊙)^2 σ T^4, T, Reals][[2, 1]]
```

$$\text{Out[314]} = T \rightarrow 350617.28036314 \text{ K}$$

d)

Fraction of energy released

In[321]:= $ef = (24 - 23.985) / 24$

Out[321]= 0.00062500000000002

e)

Time taken to deplete 10% of carbon

Energy output divided by luminosity:

In[325]:= $\text{UnitConvert}\left[10 M_{\odot} * .1 * ef \, c^2 / 1 \times 10^7 L_{\odot}, \text{"Years"}\right]$

Out[325]= 920.43081892875 yr

Q3

a)

Defining f(E)

In[180]:= $f[En_] = \text{Exp}\left[-En / (k \, T) - \text{Sqrt}\left[EG / En\right]\right]$ Out[180]= $e^{-\sqrt{\frac{EG}{En}} - \frac{En}{k \, T}}$ Finding maximum E_0 In[181]:= $E0 = \text{Solve}\left[D[f[En], En] == 0, En\right][[1, 1, 2]]$ Out[181]= $\frac{EG^{1/3} k^{2/3} T^{2/3}}{2^{2/3}}$

b)

Taylor expansion around E_0 to 2nd orderIn[182]:= $fA = \text{Series}\left[\text{Log}@f[En], \{En, E0, 2\}\right] // \text{Normal} // \text{PowerExpand} // \text{Exp}$ Out[182]= $e^{-\frac{3 \left(En - \frac{EG^{1/3} k^{2/3} T^{2/3}}{2^{2/3}}\right)^2}{2 \times 2^{1/3} EG^{4/3} k^{5/3} T^{5/3}} - \frac{EG^{1/3}}{2^{2/3} k^{1/3} T^{1/3}} - \frac{2^{1/3} EG^{1/3}}{k^{1/3} T^{1/3}}}$

Finding standard deviation

```
In[183]:= sσ = Solve[-1 / (2 σ^2) == Last@CoefficientList[Log@fA // PowerExpand, En], σ][[2, 1]]
Out[183]= σ →  $\frac{2^{1/6} EG^{1/6} k^{5/6} T^{5/6}}{\sqrt{3}}$ 
```

c)

Integrating the approximation for f(E)

```
In[184]:= int = Integrate[fA, {En, -∞, ∞}, Assumptions → {k > 0, T > 0, EG > 0}]
Out[184]=  $\frac{1}{\sqrt{3}} 2^{2/3} e^{-\frac{3 \left(\frac{EG}{kT}\right)^{1/3}}{2^{2/3}}} \sqrt{\pi} (EG k^5 T^5)^{1/6}$ 
```

Checking whether this is equal to RHS

```
In[185]:= int == Sqrt[2 Pi] f[E0] σ /. sσ // PowerExpand // Simplify
Out[185]= True
```

Therefore the integral is equivalent to the given expression.

Q4

Nuclear power density

```
In[186]:= ε[T_] =  $\left( 2^{5/3} \sqrt{2} e^{-\frac{3 \left(\frac{EG}{kT}\right)^{1/3}}{2^{2/3}}} Q_D EG^{1/6} S_0 xA xB \right) / \left( \sqrt{3} (kT)^{2/3} \sqrt{\mu} aA aB mH^2 \right);$ 
```

$\partial \log(\epsilon) / \partial \log(T)$

```
In[265]:= sβ = β -> D[Log@ε[T], T] / D[Log[T], T] /. T -> T0 // PowerExpand // Simplify
Out[265]= β →  $-\frac{2}{3} + \frac{EG^{1/3}}{2^{2/3} k^{1/3} T_0^{1/3}}$ 
```

This is the required value of β .

Numerical value for p + p

```
In[269]:= sβ /. {T0 -> 1.57 × 10^7 K, k -> k, EG -> 500 keV}
Out[269]= β → 3.854104659328
```

Numerical value for $p + C$

```
In[270]:= sβ /. {T0 → 1.57 × 107 K, k → k, EG → 35.5 MeV }
```

```
Out[270]= β → 18.053023481094
```

Q5

a)

Defining $E(v)$ and $f(v)$

```
In[188]:= En[v_] := -ψ[r] - 1/2 v^2
f[v_] := If[En[v] > 0, F En[v]^(n - 3/2), 0]
```

Integrating $f(v)$ over all v

```
In[190]:= sρ = ρ[r] -> Integrate[4 Pi f[v] v^2, {v, 0, Sqrt[-2 ψ[r]]}, Assumptions → ψ[r] < 0 && n > 1/2]
```

```
Out[190]= ρ[r] → (2 √2 F π3/2 Gamma[-1/2 + n] (-ψ[r])n) / Gamma[1 + n]
```

b)

Substituting ψ in the Poisson equation

```
In[191]:= eqP = Laplacian[ψ[r], {r, θ, φ}, "Spherical"] == 4 Pi G ρ[r] /. sρ;
eqP // TraditionalForm
```

```
Out[191]//TraditionalForm=
```

$$\psi''(r) + \frac{2\psi'(r)}{r} = \frac{8\sqrt{2}\pi^{5/2}FG\Gamma(n - \frac{1}{2})(-\psi(r))^n}{\Gamma(n+1)}$$

Lane-Emden equation

```
In[192]:= eqLE = 1/ξ^2 D[ξ^2 D[θ[ξ], ξ], ξ] + θ[ξ]^n == 0 // Simplify; eqLE // TraditionalForm
```

```
Out[192]//TraditionalForm=
```

$$\theta''(\xi) + \frac{2\theta'(\xi)}{\xi} + \theta(\xi)^n = 0$$

Thus we can see that this differential equation is the same as the Lane-Emden equation with the variables scaled appropriately as follows:

$$\psi \rightarrow c \frac{1}{n} \theta$$

```
Out[264]= r → √(c 1/n) ξ
```

$$c \rightarrow \frac{(-1)^{-n} \Gamma(1+n)}{8\sqrt{2}FG\pi^{5/2}\Gamma(-\frac{1}{2}+n)}$$

c)

Substituting $A / \sqrt{\frac{r^2}{R0^2} + 1}$ for $\psi(r)$

$$\begin{aligned} \text{In[259]:= } \text{eqA} &= \text{eqP} /. n \rightarrow 5 /. \psi \rightarrow (A / \text{Sqrt}[1 + (r/R0)^2] \&) \\ \text{Out[259]= } &\frac{3 A r^2}{\left(1 + \frac{r^2}{R0^2}\right)^{5/2} R0^4} - \frac{3 A}{\left(1 + \frac{r^2}{R0^2}\right)^{3/2} R0^2} == - \frac{7 A^5 F G \pi^3}{8 \sqrt{2} \left(1 + \frac{r^2}{R0^2}\right)^{5/2}} \end{aligned}$$

Checking whether a solution exists

$$\begin{aligned} \text{In[275]:= } \text{SA} &= \text{Solve}[\text{eqA}, A][[-1]][[1]] // \text{Simplify} \\ \text{Out[275]= } A &\rightarrow \frac{\left(\frac{3}{7}\right)^{1/4} 2^{7/8} \left(1 + \frac{r^2}{R0^2}\right)^{1/4}}{\pi^{3/4} (F G (r^2 + R0^2))^{1/4}} \end{aligned}$$

Therefore for the above value of A, a solution of this form exists.

Computing density

$$\begin{aligned} \text{In[282]:= } \text{SP} &= \\ &\rho[r] \rightarrow \text{Laplacian}[\psi[r], \{r, \theta, \phi\}, \text{"Spherical"}] / (4 \pi G) /. \psi \rightarrow (A / \text{Sqrt}[1 + (r/R0)^2] \&) /. \text{SA} // \\ &\text{FullSimplify} \\ \text{Out[282]= } \rho[r] &\rightarrow - \left(\left(3 \left(\frac{3}{7} \right)^{1/4} F^2 G R0^2 \right) / \left(2 \times 2^{1/8} \pi^{7/4} \left(1 + \frac{r^2}{R0^2} \right)^{1/4} (F G (r^2 + R0^2))^{9/4} \right) \right) \end{aligned}$$

Asymptotic limit

Using Taylor expansion for the reciprocal of r:

$$\begin{aligned} \text{In[326]:= } \text{Normal@Series}[\rho[r] /. \text{SP} /. r \rightarrow 1/u, \{u, 0, 6\}] /. u \rightarrow 1/r // \text{Simplify} \\ \text{Out[326]= } & - \frac{3 \left(\frac{3}{7}\right)^{1/4} F \left(\frac{1}{R0^2}\right)^{3/4} R0^4}{2 \times 2^{1/8} (F G)^{5/4} \pi^{7/4} r^5} \end{aligned}$$

Thus we can see that the density scales as r^{-5} .