

co2amp

Mikhail Polyanskiy

September 5, 2014

Contents

1	General notes	2
1.1	Program capabilities	2
1.2	Availability	3
1.3	Acknowledgements	3
2	Getting started manual	4
2.1	Basic concepts	4
2.2	Input pulse and calculation grid	5
2.3	Optical layout and beam propagation	5
2.3.1	Components	5
2.3.2	Layout	6
2.3.3	Beam propagation	7
2.4	Active medium and amplification	8
2.5	Program output	8
3	Models	9
3.1	Molecular dynamics	9
3.1.1	Pumping by electric discharge	9
3.1.2	Pumping and vibrational relaxation dynamics	13
3.2	Amplification	15
3.3	Molecular constants	16
3.4	Propagation	16
3.5	Optical elements	23
3.6	Properties of optical materials	24
	Bibliography	26

Chapter 1

General notes

1.1 Program capabilities

1. Ultrashort pulse amplification in CO₂ active medium
 - Rotational numbers up to J=60
 - Hot- and sequence- bands
 - Isotopic CO₂
2. Molecular dynamics
 - Realistic pumping
 - Collisional relaxation processes
 - Stimulated transitions
 - Independent consideration of active media regions at different elongation from optical axis
3. Diffraction-based beam propagation
 - Beam manipulation with common optical elements
 - Arbitrary optical configuration
4. Linear dispersion and non-linear effects in optical materials
 - Pulse chirping
 - Kerr lensing
 - Self-phase modulation
5. Advanced optics
 - Chirped-pulse amplification
 - Trains of pulses
 - Staging (program output as an input for the next stage stage)
6. User's interface
 - Easy parameter specification
 - Graphical output
 - Save/recall your work

1.2 Availability

The entire code base of the `co2amp` program is published in the GitHub website (<https://github.com/polyanskiy/co2amp>) and is freely available for use, modification and re-distribution under the terms of the General Public License (GPL) (<http://www.gnu.org/copyleft/gpl.html>). The binary package is provided in the form of a Windows installer that contains pre-compiled executables, documentation and the source-code (all found in the installation directory after installation) at <https://github.com/polyanskiy/co2amp/releases>. User interface is built using a cross-platform library QT (<http://qt-project.org/>) and thus should compile under other platforms supported by QT (MacOS, Linux). Third-party components used in the package (`gnuplot`, `7-zip`) are also freely available for multiple platforms (<http://www.gnuplot.info/>, <http://www.7-zip.org/>).

1.3 Acknowledgements

Viktor Platonenko provided the Mathcad code for pulse amplification in the CO₂ active medium that was used as a starting point of the development of the `co2amp` code; he also provided valuable input on the early stage of the work.

Chapter 2

Getting started manual

2.1 Basic concepts

The **co2amp** code allows simulating propagation of an ultrashort pulse through an arbitrary optical system that can include CO₂ amplifiers. Pulse amplification and fast molecular dynamics (stimulated transitions, rotational relaxation) calculations are performed in time domain in the time-frame moving with the pulse. Processes that are much slower than the pulse duration (discharge pumping, vibrational relaxation) are modelled separately in the laboratory time-frame. The program input parameters include initial pulse characteristics, optical configuration, active medium composition and excitation parameters (e.g. discharge profile), and number of nodes in the calculation grids for the time- (pulse time frame) and space- (radial) coordinates. Axial symmetry is assumed at all times. Optical system can include multiple amplifiers if they have same gas composition and pumping dynamics; sequential amplification in different amplifiers can be modeled using the staging option as described later. Temporal pulse shape (averaged across the beam) and beam profile (averaged for the entire pulse duration) at every element of the optical layout are saved and can be accessed in both graphical and tabulated-numerical representation. Complete pulse/beam information (complex field in every node of the time-space calculation grid) at the output of the system can also be saved and used as an input for another system (staging). The user interface program allows saving all inputs and outputs of the calculations as a single compressed file ('.co2' or '.co2x' extension). The difference between the two formats is that the files with extension ending with 'x' ("extended") include the complete output field record and thus are suitable for staged calculations. Fig. 2.1 shows the user interface of the **co2amp** program described in the following sections.

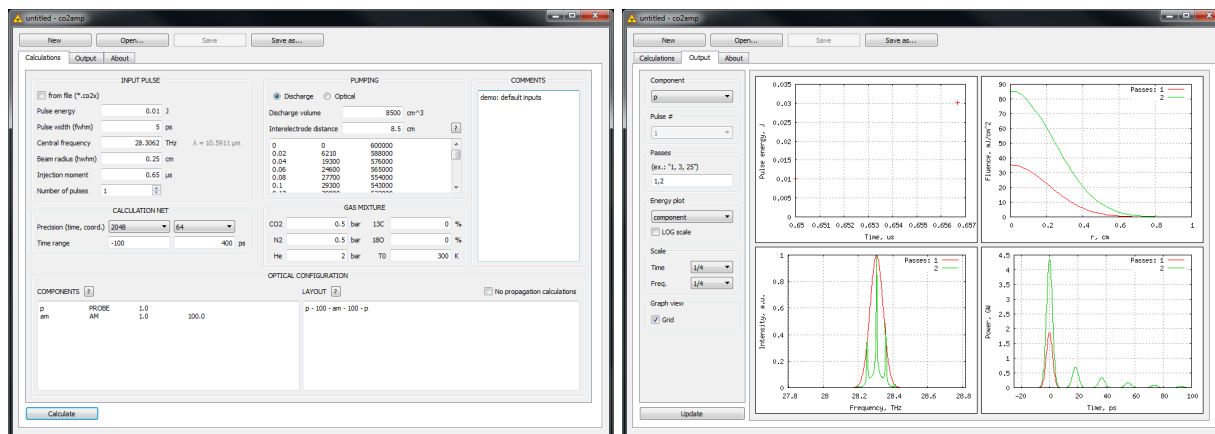


Figure 2.1: **co2amp** user interface: Input- (left) and output- (right) tabs.

2.2 Input pulse and calculation grid

Unless the output of another simulation is used as an input (staged calculations), both temporal shape of the input pulse and the input beam profile are assumed to be Gaussian; and the pulse is assumed to be transform-limited (no initial chirping). Pulse energy, duration, central frequency, and beam radius are entered in the corresponding fields in the input tab. "Injection moment" parameter specifies the time-delay between the beginning of the pumping of the active medium and pulse injection into the optical system. It is also possible to simulate amplification of a train of identical and equidistant pulses; number of pulses in the train and time-delay between them must be specified in this case. Number of nodes in the calculation grid has to be specified for the radial coordinate and for the fast time-frame associated with the pulse (time-frame for the slow processes uses a fixed 1-ns step and does not require a user input). Number of nodes is always a power of two that allows utilization of fast Fourier transform (FFT) algorithms. Calculations with a larger number of nodes are usually more accurate but, on the other hand, take longer and require more computer memory (both calculation time and required memory are roughly proportional to the product of the number of nodes in the time- and space- grids). This is therefore recommended to start from running the simulation with a smaller number of nodes and then repeat it several times, each time with a denser grid. Absence of considerable change in the program output with increase of the number of nodes will indicate that the grid density is sufficient. Limits of the pulse time-frame ("time range") specify the time interval considered in the calculations. Initially, the pulse is centered at $t = 0$ and it tends to shift to the longer delays upon amplification due to the limited medium response time and spectrum modulation. Thus, optimum time range is usually asymmetric with a short negative and a longer positive part. Time-step $\Delta t = (t_{max} - t_{min}) / (n_0 - 1)$, where t_{max} and t_{min} define the time range and n_0 is the number of nodes in the time grid, must be small enough to accurately describe the pulse profile at all stages of the propagation through the optical system. It is also important to keep in mind that time range and number of nodes in the time grid also define the range and step in the frequency domain: $\Delta\nu = 1 / (t_{max} - t_{min})$ and $(\nu_{max} - \nu_{min}) = 1 / \Delta t$. This means that the time range must be long enough to provide sufficient resolution in the frequency domain and at the same time the time step must be sufficiently short to provide a bandwidth that fits the whole spectral region of interest. Maximum radial coordinate is defined separately for every element of the optical system. This is described in the next section. Identifying an appropriate calculation grid is very important for building an accurate model of an optical system. An effort put in this part of the simulation process will pay off by fast and reliable calculations.

2.3 Optical layout and beam propagation

The **co2amp** code uses a concept of optical surfaces similar to that employed in the ray-tracing algorithms. In this concept optical system is represented by a series of thin optical components (surfaces) each of which can alter the wave front. Between surfaces, beam propagates undisturbed. In order to specify the optical system one has to 1) describe all optical components and 2) describe the location of the components in respect to each other. In **co2amp**'s user interface program, this is done by editing text in the two dedicated fields ("Components" and "Layout"). The format of these fields is described below.

2.3.1 Components

Each optical component is described in a separate line by several tab- or space- separated entries:

- Entry 1 '**ID**': An arbitrary user-defined alpha-numeric string for identifying the component;
- Entry 2 '**Type**': Type of the component; must be one of the types listed in the table;
- Entry 3 '**Field**': Maximum radial coordinate in centimeters to be considered in the simulations;
- Entry 4 '**Parameter 1**' and Entry 5 '**Parameter 2**': Meaning of these fields depend on the type of the component as summarized in the table:

<i>Type</i>	Description	<i>Parameter 1</i>	<i>Parameter 2</i>
AM	Active medium	Length [cm]	-
PROBE	Passive surface, may be used as limiting aperture	-	-
MASK	Opaque circular screen	Radius [cm]	-
ABSORBER	Absorber	Transmittance	-
LENS	Ideal lens (no spherical or chromatic aberrations)	Focal length [cm]	-
WINDOW	Transparent flat window (zero reflection- and absorption- losses)	Window material: KCl, NaCl, ZnSe, GaAs, CdTe, Ge, or Si	Thickness [cm]
STRETCHER	Stretcher or compressor	Pulse chirping [ps/THz] (positive for red-chirp)	-

The following example defines several 2.5-cm diameter components: A probe surface, two 120-cm-long amplifier sections, a 5-cm-thick NaCl window and an absorber for modelling 4% reflection losses:

```

p    PROBE      1.25
am1  AM         1.25 120
am2  AM         1.25 120
win  WINDOW     1.25 NaCl 5
abs  ABSORBER   1.25 0.96

```

2.3.2 Layout

Optical configuration is specified as a sequence of the components identified by their *ID*'s. All components are separated by non-negative distances expressed in centimeters. Components and distances must be separated by spaces, tabs, new-line symbols or dashes. The following is an example of an optical configuration that uses the components defined in the previews example:

```
p-0-abs-0-win-0-abs-100-am1-200-am2-100-abs-0-win-0-abs-0-p
```

This example defines the configuration shown in the Fig. 2.2.

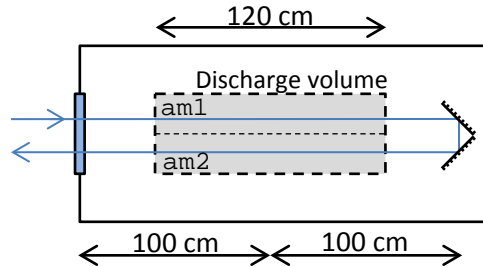


Figure 2.2: Example optical configuration.

In this example beam propagates twice through the same amplifier. However, we consider this configuration as two identical but independent active volumes (**am1** and **am2**) because at each pass the beam interacts with a different zone of the active medium.

2.3.3 Beam propagation

When two components of an optical configuration are separated by a non-zero distance, beam propagation between them is simulated using the Huygens-Fresnel diffraction integral in the assumption of axial symmetry. Propagation is simulated separately for each moment of time from the time calculation grid. Optical surface model assumes that components of the optical layout are infinitely thin. Propagation is calculated between the middle points of the components; for the purpose of simulating pulse interaction with a prolonged optical element, beam is assumed to be collimated inside this element. For instance, the sequence of calculations for the previous example is:

1. Interaction with the window (including reflection losses);
2. 100 cm propagation to the middle of the first pass through the amplifier (no amplification so far);
3. Amplification by a 120-cm amplifier section assuming that the beam is collimated;
4. 200 cm (100×2) propagation to the middle of the second pass through the active volume;
5. Amplification by the second 120-cm amplifier section assuming that the beam is collimated;
6. 100 cm propagation to the window;
7. Interaction with the windows.

Accuracy of the model can be improved if one divides a long amplifier into several shorter sub-sections. For instance, we can modify our example configuration as shown in the Fig. 2.3.

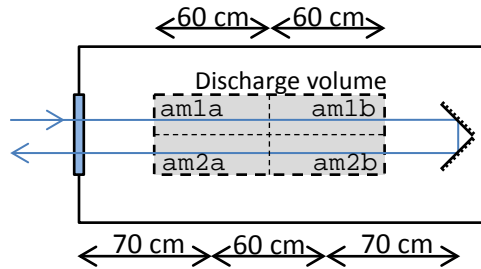


Figure 2.3: Example configuration modified for improved accuracy.

Components and Layout entries for the modified configuration are as follows:

Components:

p	PROBE	1.25		
am1a	AM	1.25	60	
am1b	AM	1.25	60	
am2a	AM	1.25	60	
am2b	AM	1.25	60	
win	WINDOW	1.25	NaCl	5
abs	ABSORBER	1.25	0.96	

Layout:

p-0-abs-0-win-0-abs-70-am1a-60-am1b-140-am2b-60-am2a-70-abs-0-win-0-abs-0-p

Population dynamics in all amplifier sections is modeled separately and thus splitting a long amplifier into shorter sections provides also a more realistic model of the active medium.

2.4 Active medium and amplification

All active medium sections used in a single simulation must have same gas composition and pumping conditions (staged simulations can be used for complex systems with two or more non-identical amplifiers). Composition (including isotopic enrichment of carbon dioxide) and initial temperature of the active medium are specified in the "Gas Mixture" fields of the input tab of the user interface. Pumping by electric discharge is the primary pumping scheme of the **co2amp** code. Rudimentary support for optical pumping is also included but not discussed here. For discharge pumping, geometry of the discharge volume (discharge volume and the distance between electrodes) and discharge profile (tabulated values of discharge current expressed in amperes and voltage in volts at time moments in microseconds) must be provided. Pumping dynamics calculations are done using a Boltzmann equation that takes into account elastic collisions between electrons and molecules and inelastic collisions with molecular rotational-, vibrational- and electronic- excitations and ionization. Tabulated empirical values of corresponding cross-sections as functions of electron energy are used. The Boltzmann equation is solved repeatedly in order to accurately describe the variation of pumping efficiency by the electric field varying during the discharge.

2.5 Program output

The output of the program is given as a temporal pulse structure and its spectrum (both averaged across the beam) and beam profile (averaged on the duration of the pulse) at each element of the optical layout. The user can chose an optical component to display from the list of all the available components. In the case if the selected component is used several times in the optical configuration, it is also possible to specify which passes through the component will be displayed. Also, integral pulse energy can be provided either at every pass through a selected component or at all passes through all components of the system layout. Output for the layout components of type **AM** (active medium) provides additional information that includes gain, discharge profile, population dynamics, and the dynamics of pumping energy distribution (fractions of discharge energy going into excitation of laser levels and excitation of molecular translations and ionization). Right-mouse-click on a figure allows the user to copy either the graphics or the numerical data for future use.

Chapter 3

Models

3.1 Molecular dynamics

Simulations of active medium pumping by electric discharge and vibrational relaxation are done following Karlov and Konev [1].

3.1.1 Pumping by electric discharge

Pumping is described by the Boltzmann equation in the following form (Karlov and Konev cite [2] and [3])

$$\begin{aligned}
 -\frac{1}{3} \left(\frac{E}{N} \right)^2 \frac{d}{du} \left[u \left(\sum_j y_j Q_{mj}(u) \right)^{-1} \frac{df}{du} \right] = \\
 1.09 \times 10^{-3} \frac{d}{du} \left[u^2 f \sum_j \frac{y_j}{M_j} Q_{mj}(u) \right] + \sum_{j=1,2} y_j C_j \frac{d}{du} (uf) + 6B y_2 \frac{d}{du} (uQ(u)f) \\
 + \sum_j y_j \sum_k (u + u_{jk}) Q_{jk}(u + u_{jk}) f(u + u_{jk}) - (uf) \sum_j y_j \sum_k Q_{jk}(u)
 \end{aligned} \tag{3.1}$$

where left part describes energy of electrons in the electric field, the first component of the sum of the right part represents energy transfer via elastic collisions between electrons and molecules, second and third components describe collisions with molecular rotation excitation, and the two last components relate to inelastic collisions with transfer of the energy u_{jk} into vibrational and electronic excitation and ionization.

Electron energy u is expressed in eV;

Ration of the field to the full molecular density E/N is expressed in the units of $10^{-16} \text{ V}\cdot\text{cm}^2$;

y_j are the relative molecule concentrations ($j = 1$ corresponds to CO_2 , $j = 2$ to N_2 and $j = 3$ to He);

$M_1 = 44$, $M_2 = 28$, $M_3 = 4$ are the molar masses,

$C_1 = 8.2 \times 10^{-4} \text{ eV}\cdot\text{\AA}^2$ [4];

$C_2 = 5.06 \times 10^{-4} \text{ eV}\cdot\text{\AA}^2$ [5];

$B = 2.5 \times 10^{-4}$ is the N_2 rotational constant;

Effective crosssections are expressed in \AA ; their numerical values in the nodes are given in the tables below (linear interpolation must be used for determining the values in intermediate points); the data and citations are reproduced from [1].

The following notation for cross-sections is used:

Q_{m1} - Transport cross-section of CO_2 [6];

Q_{m2} - Transport cross-section of N_2 [5];

Table 3.1: Tabulated data on the cross-sections for discharge pumping

u_i	Q_{m1}	u_i	Q_{m2}	u_i	Q_{m3}	u_i	Q
0	140	0	1.4	0	5	0.0015	0
0.04	84	0.001	1.4	0.01	5.4	0.05	0.1
0.1	55	0.002	1.6	0.1	5.8	0.25	0.65
0.3	21	0.008	2	0.2	6.2	0.5	1.15
0.5	10.8	0.01	2.2	1	6.5	0.8	2
0.6	9.4	0.04	4	2	6.1	1	2.65
1	5.7	0.08	6	7	5	1.5	5.6
1.7	5	0.1	6.5	10	4.1	1.8	7.5
2	5.1	0.2	8.8	20	3	1.9	8.2
2.5	6	0.3	9.8			2	8.6
3	7.7	0.4	10			2.15	8.95
4.1	9.4	1	10			2.43	9
5	14.5	1.2	11			2.6	8.9
7.4	10	1.4	12.5			2.75	8.4
10	11.7	1.8	20			2.9	7.65
20	16	2	25			3.25	6.2
27	16.3	2.5	30			3.6	5.1
50	13	3	26			4	4.5
		4	15			4.5	4.16
		5	12			5	3.97
		7	10			5.5	3.93
		10	10			7	4.17
		14	11			9	4.46
		18	12.2			11	4.42
		20	12			15	3.94
		30	10			22	3.15
		100	10			25	3.05

Table 3.2: Tabulated data on the cross-sections for discharge pumping - continued 1

u_i	Q_{11}	u_i	Q_{12}	u_i	Q_{13}	u_i	Q_{14}	u_i	Q_{15}
0.083	0	0.167	0	0.252	0	2.37	0	2.37	0
0.085	0.36	0.2	0.54	2.7	0.25	3	0.26	3	0.17
0.09	1.04	0.25	0.82	3	0.4	3.5	0.52	3.65	0.33
0.1	1.6	0.3	0.82	3.3	0.6	4	0.5	3.8	0.31
0.12	1.84	0.5	0.68	3.6	0.65	4.5	0.22	4	0.21
0.14	2.12	0.7	0.56	4.5	0.23	4.6	0.1	4.3	0.1
0.16	2.16	1	0.47	4.6	0.1	5	0	5	0
0.2	2.08	1.4	0.45	5	0				
0.3	1.76	2	0.55						
0.4	1.52	3	1.15						
0.5	1.28	3.9	1.83						
0.6	1.08	4.5	1.4						
0.8	0.8	5	0.4						
1	0.58	6	0.28						
1.2	0.48	10	0.2						
1.6	0.34	20	0.1						
1.8	0.35								
2	0.4								
2.5	0.64								
3	1.04								
3.7	1.4								
4	1.36								
4.2	1.2								
4.5	0.92								
5	0.53								
6	0.4								
8	0.36								
9	0.28								
10	0.16								
10.1	0								
$u_{11} = 0.083 \text{ eV}$		$u_{12} = 0.167 \text{ eV}$		$u_{13} = 0.252 \text{ eV}$		$u_{14} = 0.339 \text{ eV}$		$u_{15} = 0.422 \text{ eV}$	
u_i	Q_{16}	u_i	Q_{17}	u_i	Q_{18}	u_i	Q_{19}	u_i	$Q_{1,10}$
2.5	0	0.29	0	7	0	10.5	0	13.8	0
3	0.19	0.3	0.44	8	0.5	11.5	0.56	15	0.1
3.6	0.245	0.35	0.65	8.4	0.6	14	0.8	16	0.13
4	0.21	0.4	0.73	9	0.46	20	1.2	17	0.17
5.07	0	0.5	0.84	10	0.175	30	2	30	1.55
		0.8	1	10.5	0	50	4	40	2.1
		1	1						
		2	0.78						
		6	0.37						
		10	0.25						
		50	0						
$u_{16} = 2.5 \text{ eV}$		$u_{17} = 0.29 \text{ eV}$		$u_{18} = 7 \text{ eV}$		$u_{19} = 10.5 \text{ eV}$		$u_{1,10} = 13.8 \text{ eV}$	

Table 3.3: Tabulated data on the cross-sections for discharge pumping - continued 2

u_i	Q_{21}	u_i	Q_{22}	u_i	Q_{23}	u_i	Q_{24}	u_i	Q_{25}
0.29	0	1.83	0	1.9	0	2.05	0	2.1	0
0.5	0.0052	1.9	0.208	2	0.416	2.1	0.416	2.15	0.208
0.8	0.0083	2	1.46	2.1	1.33	2.2	1.16	2.2	0.541
1	0.0104	2.05	2.29	2.2	1.87	2.26	1.58	2.3	0.915
1.2	0.0166	2.1	1.66	2.3	1.25	2.55	0	2.46	1.12
1.3	0.0728	2.2	0.79	2.36	0.208	2.75	0.832	2.5	1.12
1.4	0.135	2.35	0.208	2.42	0	2.77	0	2.6	0.208
1.6	0.25	2.45	1.98	2.5	0.499	3	0.208	2.62	0
1.8	0.52	2.5	1.78	2.61	0.915	3.05	0.208	2.68	0
1.9	0.832	2.62	0.208	2.7	0.624	3.25	0	2.8	0.416
2	3.02	2.75	1.04	2.75	0.208			2.9	0.75
2.05	3.12	2.95	1.66	2.8	0			3	0
2.1	2.08	3.05	0.624	2.92	0.416			3.2	0.25
2.15	1.25	3.2	0.208	3	0.208			3.3	0.125
2.2	0.832	3.4	0.208	3.25	0.208			3.35	0
2.3	2.9	4	0	3.31	0				
2.45	1.04								
2.53	1.25								
2.6	1.75								
2.62	2.08								
2.68	1.73								
2.73	0.416								
2.85	0.32								
2.92	0.416								
3.12	0.728								
3.3	0.52								
4	0								
$u_{21} = 0.29$ eV		$u_{22} = 0.58$ eV		$u_{23} = 0.87$ eV		$u_{24} = 1.16$ eV		$u_{25} = 1.45$ eV	
u_i	Q_{26}	u_i	Q_{27}	u_i	Q_{28}	u_i	Q_{29}	u_i	$Q_{2,10}$
2.3	0	2.4	0	2.6	0	5	0	6.8	0
2.4	0.75	2.5	0.208	2.7	0.208	5.9	0.41	7.1	0.57
2.5	1.04	2.75	0.75	2.9	0.29	6.1	0.41	8.1	0.57
2.55	1.12	3	0	3	0.208	7	0.07	8.6	0.25
2.6	1.04	3.2	0.166	3.1	0	9	0	9.5	0.12
2.65	0.624	3.3	0.146	3.2	0			20.7	0
2.7	0.416	3.4	0	3.3	1.04				
2.8	0.208			3.4	0				
2.9	0.125								
3	2.5								
3.1	0.166								
3.2	0								
$u_{26} = 1.74$ eV		$u_{27} = 2.03$ eV		$u_{28} = 2.32$ eV		$u_{29} = 5$ eV		$u_{2,10} = 6.8$ eV	
u_i	$Q_{2,11}$	u_i	$Q_{2,12}$	u_i	$Q_{2,13}$	u_i	$Q_{2,14}$	u_i	$Q_{2,15}$
8.4	0	11.25	0	12.5	0	14	0	15.6	0
8.7	0.42	13.8	0.41	13	0.4	14.3	1.7	18	0.1
9.1	0.42	14	1	13.6	0.4	14.8	1.7	20	0.21
10	0.3	14.7	1	14	0.16	15.6	0.2	50	2.52
20.7	0	15	0.25	20.7	0	20.6	0.2	100	2.52
		65	0			25.4	2.8		
						100	2.8		
$u_{2,11} = 8.4$ eV		$u_{2,12} = 11.25$ eV		$u_{2,13} = 12.5$ eV		$u_{2,14} = 14$ eV		$u_{2,15} = 15.6$ eV	

Q_{m3} - Transport cross-section of He [6];
 Q - Cross-section of resonant excitation of N_2 rotation [7, 8];
 Q_{11} - Cross-section of the process $(000) \rightarrow (01^10)$ [6];
 Q_{12} - Cross-section of the process $(000) \rightarrow (100 + 020)$ [6];
 $Q_{13} \dots Q_{16}$ - Cross-sections of resonant processes around 3.8 eV [6];
 Q_{17} - Cross-section of the process $(000) \rightarrow (001)$ [6];
 $Q_{18} \dots Q_{1,10}$ - Cross-sections of electronic excitation and ionization of CO_2 [4];
 $Q_{21} \dots Q_{28}$ - Cross-sections of the process $N_2(v=0) \rightarrow N_2(v=1 \dots 8)$ [9–11];
 $Q_{29} \dots Q_{2,15}$ - Cross-sections of electronic excitation and ionization of N_2 [11].

Equation 3.1 is solved numerically using the tridiagonal matrix algorithm. Rate constant ω_{jk} and electron drift speeds v_d are defined as:

$$\omega_{jk} \left[\frac{\text{cm}^3}{\text{s}} \right] = 5.93 \times 10^{-9} \int_0^\infty u Q_{jk}(u) f(u) du \quad (3.2)$$

$$v_d \left[\frac{\text{cm}}{\text{s}} \right] = -5.93 \times 10^7 \left(\frac{1}{3} \frac{E}{N} \right) \frac{df}{du} \int_0^\infty u \left(\sum_j y_j Q_{mj}(u) \right)^{-1} du \quad (3.3)$$

Fraction of electron energy transmitted via inelastic processes is defined as:

$$z_{jk} = 10^6 \frac{y_j u_{jk} \omega_{jk}}{\left(\frac{E}{N} \right) v_d} \quad (3.4)$$

Fraction of electron energy transmitted to translations and rotations are:

$$z_t = 5.93 \times 10^7 \frac{\int_0^\infty u^2 \left(\sum_j \frac{y_j}{M_j} Q_{mj}(u) \right) f(u) du}{\left(\frac{E}{N} \right) v_d} \quad (3.5)$$

$$z_r = 5.93 \times 10^7 \frac{\sum_{j=1,2} y_j C_j \int_0^\infty u f(u) du + 6 y_2 B \int_0^\infty u Q(u) f(u) du}{\left(\frac{E}{N} \right) v_d} \quad (3.6)$$

Then, for 4-temperature model we get:

$q_2 = \sum_{k=1}^6 z_{1k}$ - fraction of energy transferred to CO_2 symmetric stretch (ν_1) and bending (ν_2) modes;
 $q_3 = z_{17}$ - fraction of energy transferred to CO_2 asymmetric stretch mode (ν_3);
 $q_4 = \sum_{k=1}^8 z_{2k}$ - fraction of energy transferred to N_2 vibrations;
 $q_T = z_t + z_r$ - fraction of energy transferred to translation and rotation;
 $q_{ei} = \sum_{k=9}^{15} z_{2k} + \sum_{k=8}^{10} z_{1k}$ - fraction of energy spent on electronic levels excitation and ionization.

3.1.2 Pumping and vibrational relaxation dynamics

3-temperature model is used for describing vibrational dynamics of the laser system

T_2 - vibrational temperature of ν_1 and ν_2 vibrations of CO_2 ;
 T_3 - vibrational temperature of the ν_3 vibration of CO_2 ;
 T_4 - vibrational temperature of N_2 .

Vibrational temperatures are related to average number of quanta e_x in the corresponding vibrations as:

$$\begin{aligned}
e_2 &= \frac{2}{\exp(960/T_2) - 1} \\
e_3 &= \frac{1}{\exp(3380/T_3) - 1} \\
e_4 &= \frac{1}{\exp(3350/T_4) - 1}
\end{aligned} \tag{3.7}$$

("2" in the first equation is due to 2-fold degeneracy of the bend vibration levels).
Pumping/relaxation dynamics is described by the following equations [1]

$$\begin{aligned}
\frac{de_4}{dt} &= p_{e4} - r_a(e_4 - e_3) \\
\frac{de_3}{dt} &= p_{e3} + r_c(e_4 - e_3) - r_3 f_3 \\
\frac{de_2}{dt} &= f_2 (p_{e2} + 3r_3 f_3 - r_2(e_2 - e_{2T}))
\end{aligned} \tag{3.8}$$

where

$$\begin{aligned}
p_{e4} &= 0.8 \times 10^{-3} \frac{q_4}{ny_2} W(t); \quad p_{e3} = 0.8 \times 10^{-3} \frac{q_3}{ny_1} W(t); \quad p_{e2} = 2.8 \times 10^{-3} \frac{q_2}{ny_1} W(t); \\
f_2 &= \frac{2(1 + e_2)^2}{2 + 6e_2 + 3e_2^2}; \quad f_3 = e_3(1 + e_2/2)^3 - (1 + e_3)(e_2/2)^3 \exp(-500/T); \\
r_a &= kny_1; \quad r_c = kny_2; \quad r_2 = k_2 n; \quad r_3 = k_3 n; \\
k_2 &= \sum_{i=1}^3 y_i k_{2i}; \quad k_3 = \sum_{i=1}^3 y_i k_{3i}; \\
n &= 273 \frac{p[\text{bar}]}{T_0[\text{K}]},
\end{aligned}$$

where $W(t)$ is measured in kW/cm^3 , p_e - in μs , and constants k are calculated using the following expressions [12, 13]:

$$\begin{aligned}
k &= 240/T^{1/2}; \\
k_{31} &= A(t) \exp(4.138 + 7.945x - 631.24x^2 + 2239x^3); \\
k_{32} &= A(t) \exp(-1.863 + 213.3x - 2796.2x^2 + 9001.9x^3); \\
k_{33} &= A(t) \exp(-3.276 + 291.4x - 3831.8x^2 + 12688x^3); \\
k_{21} &= 1.16 \times 10^3 \exp(-59.3x); \\
k_{22} &= 8.55 \times 10^2 \exp(-69x); \\
k_{22} &= 1.3 \times 10^3 \exp(-40.6x),
\end{aligned}$$

where $x = T^{-1/3}$, $A(t) = (T/273)(1 + e_{2T}/2)^{-3}$, and temperature is expressed in K.

Finally, the gas temperature dynamics is described by the Equation 3.9

$$\frac{dT}{dt} = \frac{y_1}{C_V} (500r_3 f_3 + 960r_2(e_2 - e_{2T})) + 2.7 \frac{W(t)q_T}{nC_V}, \tag{3.9}$$

where $C_V = 2.5(y_1 + y_2) + 1.5y_3$.

3.2 Amplification

Amplification is simulated using the equations from Volkin's paper [14]

$$\begin{aligned}\frac{\partial \varepsilon}{\partial z} &= - \sum_J \rho_J, \\ \frac{\partial \rho_J}{\partial t} + \left(i(\omega_c - \omega_{0J}) + \frac{1}{T_2} \right) \rho_J &= - \frac{\sigma_J n_J \varepsilon}{2T_2}, \\ \frac{\partial n_J}{\partial t} + \frac{n_J - n_J^0}{\tau_R} &= 4(\rho_J \varepsilon^* + c.c.),\end{aligned}\tag{3.10}$$

where $\varepsilon(z, t)$ is the complex field envelope, ω_c the carrier frequency, ω_{0J} the frequency in the line center and σ_J the stimulated emission (or absorption) cross-section in the line center:

$$\sigma_J[\text{m}^2] = \frac{(\lambda_J[\text{m}])^2 A_J[\text{s}^{-1}]}{4} \times \frac{T_2[\text{s}]}{\pi},$$

where the first term defines the integral cross-section σ_0 of the rotational line [15] and the second term is the maximum of the normalized Lorentzian profile of a line with linewidth $\Delta\omega_{FWHM} = 2\pi\Delta\nu = 2/T_2$,

$$n_J^0 = z(J)(N^U - N^L),$$

where $z(J)$ is the Boltzmann distribution:

$$\begin{aligned}z(J) &= 2 \frac{hB}{kT} (2J+1) \exp\left(-\frac{hB}{kT} J(J+1)\right) & 626, 636, 828, 838 \\ z(J) &= \frac{hB}{kT} (2J+1) \exp\left(-\frac{hB}{kT} J(J+1)\right) & 628, 638\end{aligned}$$

$$h = 6.62606957 \times 10^{-34} \text{ J}\cdot\text{s}$$

$$k = 1.3806488 \times 10^{-23} \text{ J/K}$$

$$B: [\text{s}^{-1}]$$

$$T: [\text{K}]$$

Field amplitude is related to the intensity as follows:

$$I[\text{W/m}^2] = 2h[\text{J}\cdot\text{s}]\nu_c[\text{s}^{-1}]\varepsilon^2\tag{3.11}$$

Additional equations:

$$\begin{aligned}\frac{d}{dt} N^U &= 2 \sum_J (\rho_J \varepsilon_J^* + c.c.), \\ \frac{d}{dt} N^{L10} &= -2 \sum_{J10} (\rho_J \varepsilon_J^* + c.c.), \\ \frac{d}{dt} N^{L9} &= -2 \sum_{J9} (\rho_J \varepsilon_J^* + c.c.)\end{aligned}\tag{3.12}$$

Relaxation times are defined from the following equations:

$$T_2[\text{s}] = \frac{10^{-6}}{\pi \times 7.61 \times 750 \times (P_{CO2} + 0.733P_{N2} + 0.64P_{He})}\tag{3.13}$$

$$\tau_R[\text{s}] = \frac{10^{-7}}{750 \times (P_{CO2} + 1.2P_{N2} + 0.6P_{He})} \quad (3.14)$$

P is measured in Bar.

$$\Delta\nu_{FWHM} = \frac{1}{\pi T_2}$$

Useful relations for Lorentzian distribution:

$$\Delta\nu_{HWHM} = \frac{1}{2\pi T_2}; \quad \Delta\omega_{HWHM} = \frac{1}{T_2}$$

3.3 Molecular constants

bla-bla

3.4 Propagation

Beam propagation between layout elements is calculated using Huygens-Fresnel integration at each time-step of the pulse time-frame. In the general case, the field in a given position (x', y') of the output plane is calculated by integrating contributions from each point (x, y) of the input plane taken into account the phase delay.

$$E(x', y') = \iint_x \int_y E(x, y) \frac{e^{ikR}}{i\lambda R} dx dy,$$

where k is the wave number, λ the wavelength, and R is the distance between points (x, y) in and (x', y') in the output plane: $R = \sqrt{(x - x')^2 + (y - y')^2 + \Delta z^2}$. Δz is the distance between the centers of the input and the output planes.

In the case of radially symmetric beam used in our model, calculations can be considerably accelerated by taken into account that the field is the same in all points at equal distance from the beam axis. In this case integration can be done by summing the contributions of concentric rings in the input plane as shown in the Fig. 3.1.

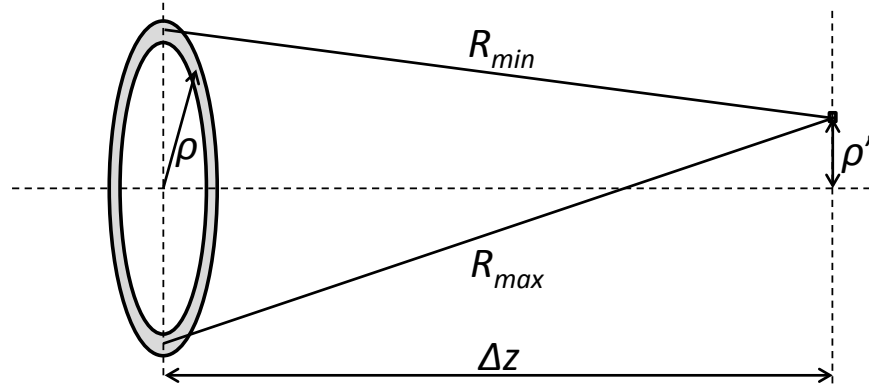


Figure 3.1: Huygens-Fresnel integration in the axial-symmetric system: Contribution of a field from a concentric ring in the input plane to a field in a point in the output plane.

The field in a given point with radial coordinate ρ' in the output plane is then calculated as

Table 3.4: Einstein coefficients A of laser transitions of $'626'$ CO₂, s⁻¹

J	Regular			Sequence			Hot-e			Hot-f		
	10P	10R	9P	10P	10R	9P	10P	10R	9P	10P	10R	9P
0	-	0.134	-	0.824	0.331	0.812	-	-	-	-	-	-
1	-	-	-	0.267	0.172	0.296	-	0.108	-	0.179	0.137	0.167
2	0.267	-	0.296	-	-	0.192	-	-	-	-	-	-
3	-	-	-	0.493	0.369	0.486	0.191	0.150	0.230	0.191	0.158	0.231
4	0.229	0.183	0.252	-	-	0.205	-	-	-	0.189	0.166	0.228
5	-	-	-	0.456	0.383	0.449	0.190	0.162	0.229	-	-	-
6	0.218	0.188	0.240	-	-	0.211	-	-	-	-	-	-
7	-	-	-	0.440	0.392	0.434	0.187	0.168	0.226	0.186	0.170	0.224
8	0.212	0.191	0.234	-	-	0.216	-	-	-	-	-	-
9	-	-	-	0.432	0.396	0.425	0.185	0.171	0.223	0.184	0.173	0.222
10	0.209	0.193	0.230	-	-	0.219	-	-	-	0.183	0.175	0.219
11	-	-	-	0.425	0.400	0.420	0.183	0.174	0.221	-	-	-
12	0.206	0.194	0.227	-	-	0.221	-	-	-	0.181	0.176	0.218
13	-	-	-	0.421	0.402	0.416	0.182	0.175	0.219	0.178	0.179	0.213
14	0.204	0.195	0.225	-	-	0.224	-	-	-	-	-	-
15	-	-	-	0.417	0.404	0.413	0.181	0.177	0.217	0.180	0.177	0.216
16	0.203	0.196	0.223	-	-	0.226	-	-	-	0.179	0.178	0.218
17	-	-	-	0.413	0.405	0.411	0.179	0.178	0.215	-	-	-
18	0.201	0.197	0.221	-	-	0.228	-	-	-	0.177	0.180	0.221
19	-	-	-	0.410	0.406	0.409	0.178	0.179	0.213	-	-	-
20	0.200	0.197	0.220	-	-	0.230	-	-	-	0.176	0.181	0.222
21	-	-	-	0.407	0.407	0.408	0.177	0.180	0.212	0.175	0.181	0.220
22	0.198	0.197	0.219	-	-	0.231	-	-	-	0.174	0.182	0.224
23	-	-	-	0.404	0.407	0.406	0.176	0.180	0.211	-	-	-
24	0.197	0.197	0.218	-	-	0.233	-	-	-	0.173	0.182	0.225
25	-	-	-	0.400	0.407	0.405	0.175	0.181	0.209	0.172	0.183	0.205
26	0.196	0.197	0.217	-	-	0.235	-	-	-	0.171	0.183	0.203
27	-	-	-	0.397	0.407	0.404	0.174	0.182	0.208	0.170	0.183	0.202
28	0.195	0.197	0.216	-	-	0.237	-	-	-	-	-	-
29	-	-	-	0.395	0.406	0.403	0.174	0.182	0.207	0.169	0.184	0.201
30	0.193	0.197	0.215	-	-	0.238	-	-	-	0.168	0.184	0.200
31	-	-	-	0.392	0.406	0.402	0.173	0.183	0.205	0.168	0.184	0.200
32	0.192	0.196	0.215	-	-	0.240	-	-	-	0.168	0.184	0.200
33	-	-	-	0.388	0.406	0.402	0.172	0.183	0.204	0.168	0.184	0.200
34	0.191	0.196	0.214	-	-	0.242	-	-	-	0.168	0.184	0.200
35	-	-	-	0.385	0.405	0.401	0.171	0.183	0.203	0.168	0.184	0.200
36	0.189	0.196	0.213	-	-	0.244	-	-	-	0.168	0.184	0.200
37	-	-	-	0.383	0.403	0.401	0.170	0.184	0.202	0.168	0.184	0.200
38	0.188	0.195	0.213	-	-	0.246	-	-	-	0.168	0.184	0.200
39	-	-	-	0.380	0.402	0.400	0.169	0.184	0.201	0.168	0.184	0.200
40	0.187	0.194	0.212	-	-	0.248	-	-	-	0.168	0.184	0.200
41	-	-	-	0.376	0.401	0.399	0.168	0.185	0.199	0.168	0.184	0.200
42	0.185	0.194	0.212	-	-	0.249	-	-	-	0.168	0.184	0.200
43	-	-	-	0.373	0.400	0.400	0.167	0.185	0.198	0.168	0.184	0.200
44	0.184	0.193	0.211	-	-	0.251	-	-	-	0.168	0.184	0.200
45	-	-	-	0.369	0.398	0.400	0.167	0.185	0.197	0.168	0.184	0.200
46	0.182	0.192	0.211	-	-	0.253	-	-	-	0.168	0.184	0.200
47	-	-	-	0.366	0.397	0.400	0.166	0.185	0.196	0.168	0.184	0.200
48	0.181	0.191	0.210	-	-	0.255	-	-	-	0.168	0.184	0.200
49	-	-	-	0.362	0.394	0.400	0.165	0.186	0.194	0.168	0.184	0.200
50	0.179	0.191	0.210	-	-	0.257	-	-	-	0.168	0.184	0.200
51	-	-	-	0.359	0.393	0.399	0.164	0.186	0.193	0.168	0.184	0.200
52	0.178	0.190	0.209	-	-	0.259	-	-	-	0.168	0.184	0.200
53	-	-	-	0.354	0.391	0.401	0.163	0.186	0.192	0.168	0.184	0.200
54	0.176	0.189	0.209	-	-	0.262	-	-	-	0.168	0.184	0.200
55	-	-	-	0.351	0.388	0.401	0.162	0.186	0.191	0.168	0.184	0.200
56	0.174	0.187	0.209	-	-	0.264	-	-	-	0.168	0.184	0.200
57	-	-	-	0.347	0.385	0.400	0.161	0.186	0.189	0.168	0.184	0.200
58	0.173	0.186	0.209	-	-	0.266	-	-	-	0.168	0.184	0.200
59	-	-	-	0.344	-	-	0.160	0.186	0.188	0.168	0.184	0.200
60	0.171	-	0.208	-	-	-	-	-	-	0.159	-	-

Table 3.5: Einstein coefficients A of laser transitions of $^{628}\text{CO}_2$, s^{-1}

J	Regular			Sequence			Hot-e			Hot-f		
	10P	10R	9P	9R	10P	10R	9P	9R	10P	10R	9P	9R
0	-	0.101	-	0.190	- no data -			- no data -			- no data -	
1	0.303	0.122	0.569	0.229								
2	0.202	0.131	0.379	0.246								
3	0.181	0.136	0.340	0.256								
4	0.172	0.140	0.323	0.262								
5	0.167	0.142	0.314	0.266								
6	0.163	0.144	0.307	0.270								
7	0.161	0.145	0.303	0.273								
8	0.159	0.147	0.299	0.275								
9	0.157	0.148	0.296	0.277								
10	0.156	0.149	0.294	0.279								
11	0.155	0.150	0.292	0.280								
12	0.154	0.150	0.290	0.282								
13	0.153	0.151	0.288	0.283								
14	0.152	0.152	0.287	0.284								
15	0.151	0.152	0.285	0.285								
16	0.151	0.153	0.284	0.286								
17	0.150	0.153	0.283	0.287								
18	0.149	0.154	0.282	0.288								
19	0.149	0.154	0.281	0.289								
20	0.148	0.155	0.280	0.290								
21	0.148	0.155	0.279	0.290								
22	0.147	0.156	0.278	0.291								
23	0.146	0.156	0.277	0.292								
24	0.146	0.157	0.276	0.293								
25	0.145	0.157	0.275	0.293								
26	0.145	0.157	0.274	0.294								
27	0.144	0.158	0.273	0.295								
28	0.144	0.158	0.272	0.295								
29	0.143	0.158	0.271	0.296								
30	0.143	0.159	0.271	0.297								
31	0.142	0.159	0.270	0.297								
32	0.142	0.159	0.269	0.298								
33	0.141	0.160	0.268	0.298								
34	0.141	0.160	0.267	0.299								
35	0.140	0.160	0.266	0.300								
36	0.140	0.161	0.266	0.300								
37	0.139	0.161	0.265	0.301								
38	0.139	0.161	0.264	0.301								
39	0.138	0.161	0.263	0.302								
40	0.138	0.162	0.262	0.302								
41	0.137	0.162	0.262	0.303								
42	0.137	0.162	0.261	0.303								
43	0.136	0.163	0.260	0.304								
44	0.136	0.163	0.259	0.304								
45	0.135	0.163	0.258	0.304								
46	0.135	0.163	0.258	0.305								
47	0.134	0.164	0.257	0.305								
48	0.134	0.164	0.256	0.306								
49	0.133	0.164	0.255	0.306								
50	0.133	0.164	0.254	0.307								
51	0.133	0.165	0.254	0.307								
52	0.132	0.165	0.253	0.308								
53	0.132	0.165	0.252	0.308								
54	0.132	0.165	0.251	0.309								
55	0.132	0.165	0.250	0.309								
56	0.132	0.165	0.250	0.309								
57	0.132	0.165	0.249	0.310								
58	0.132	0.165	0.248	0.310								
59	0.132	0.165	0.248	0.310								
60	0.132	-	0.248	-								

Table 3.6: Einstein coefficients A of laser transitions of $^{828}\text{CO}_2$, s^{-1}

J	Regular			Sequence			Hot-e			Hot-f		
	10P	10R	9P	9R	10P	10R	9P	9R	10P	10R	9P	9R
0	-	0.038	-	0.121	-	-	-	-	-	-	-	-
1	-	-	-	-	-	-	-	-	-	-	-	-
2	0.075	0.048	0.240	0.156	-	-	-	-	-	-	-	-
3	-	-	-	-	-	-	-	-	-	-	-	-
4	0.064	0.051	0.204	0.166	-	-	-	-	-	-	-	-
5	-	-	-	-	-	-	-	-	-	-	-	-
6	0.061	0.053	0.195	0.171	-	-	-	-	-	-	-	-
7	-	-	-	-	-	-	-	-	-	-	-	-
8	0.059	0.054	0.190	0.175	-	-	-	-	-	-	-	-
9	-	-	-	-	-	-	-	-	-	-	-	-
10	0.059	0.054	0.187	0.178	-	-	-	-	-	-	-	-
11	-	-	-	-	-	-	-	-	-	-	-	-
12	0.058	0.054	0.184	0.179	-	-	-	-	-	-	-	-
13	-	-	-	-	-	-	-	-	-	-	-	-
14	0.057	0.055	0.183	0.182	-	-	-	-	-	-	-	-
15	-	-	-	-	-	-	-	-	-	-	-	-
16	0.057	0.055	0.181	0.183	-	-	-	-	-	-	-	-
17	-	-	-	-	-	-	-	-	-	-	-	-
18	0.056	0.055	0.179	0.185	-	-	-	-	-	-	-	-
19	-	-	-	-	-	-	-	-	-	-	-	-
20	0.056	0.055	0.178	0.187	-	-	-	-	-	-	-	-
21	-	-	-	-	-	-	-	-	-	-	-	-
22	0.056	0.055	0.178	0.187	-	-	-	-	-	-	-	-
23	-	-	-	-	-	-	-	-	-	-	-	-
24	0.055	0.055	0.177	0.189	-	-	-	-	-	-	-	-
25	-	-	-	-	-	-	-	-	-	-	-	-
26	0.055	0.055	0.176	0.191	-	-	-	-	-	-	-	-
27	-	-	-	-	-	-	-	-	-	-	-	-
28	0.055	0.055	0.175	0.192	-	-	-	-	-	-	-	-
29	-	-	-	-	-	-	-	-	-	-	-	-
30	0.054	0.055	0.174	0.193	-	-	-	-	-	-	-	-
31	-	-	-	-	-	-	-	-	-	-	-	-
32	0.054	0.055	0.174	0.195	-	-	-	-	-	-	-	-
33	-	-	-	-	-	-	-	-	-	-	-	-
34	0.054	0.055	0.174	0.196	-	-	-	-	-	-	-	-
35	-	-	-	-	-	-	-	-	-	-	-	-
36	0.053	0.055	0.173	0.198	-	-	-	-	-	-	-	-
37	-	-	-	-	-	-	-	-	-	-	-	-
38	0.053	0.055	0.173	0.200	-	-	-	-	-	-	-	-
39	-	-	-	-	-	-	-	-	-	-	-	-
40	0.052	0.054	0.172	0.201	-	-	-	-	-	-	-	-
41	-	-	-	-	-	-	-	-	-	-	-	-
42	0.052	0.054	0.172	0.202	-	-	-	-	-	-	-	-
43	-	-	-	-	-	-	-	-	-	-	-	-
44	0.052	0.054	0.171	0.204	-	-	-	-	-	-	-	-
45	-	-	-	-	-	-	-	-	-	-	-	-
46	0.051	0.054	0.171	0.205	-	-	-	-	-	-	-	-
47	-	-	-	-	-	-	-	-	-	-	-	-
48	0.051	0.054	0.170	0.207	-	-	-	-	-	-	-	-
49	-	-	-	-	-	-	-	-	-	-	-	-
50	0.050	0.054	0.170	0.208	-	-	-	-	-	-	-	-
51	-	-	-	-	-	-	-	-	-	-	-	-
52	0.050	0.053	0.170	0.210	-	-	-	-	-	-	-	-
53	-	-	-	-	-	-	-	-	-	-	-	-
54	0.049	0.053	0.170	0.213	-	-	-	-	-	-	-	-
55	-	-	-	-	-	-	-	-	-	-	-	-
56	0.049	0.052	0.170	0.214	-	-	-	-	-	-	-	-
57	-	-	-	-	-	-	-	-	-	-	-	-
58	0.049	0.052	0.170	0.216	-	-	-	-	-	-	-	-
59	-	-	-	-	-	-	-	-	-	-	-	-
60	0.048	-	-	0.169	-	-	-	-	-	-	-	-

Table 3.7: Einstein coefficients A of laser transitions of $'636'$ CO₂, s⁻¹

J	Regular			Sequence			Hot-e			Hot-f		
	10P	10R	9P	9R	10P	10R	9P	9R	10P	10R	9P	9R
0	-	0.147	-	0.077	-	-	-	-	-	-	-	-
1	-	-	-	-	-	-	-	-	-	-	-	-
2	0.293	0.189	0.153	0.099	-	-	-	-	-	-	-	-
3	-	-	-	-	-	-	-	-	-	-	-	-
4	0.251	0.201	0.131	0.106	-	-	-	-	-	-	-	-
5	-	-	-	-	-	-	-	-	-	-	-	-
6	0.239	0.207	0.124	0.109	-	-	-	-	-	-	-	-
7	-	-	-	-	-	-	-	-	-	-	-	-
8	0.233	0.210	0.121	0.111	-	-	-	-	-	-	-	-
9	-	-	-	-	-	-	-	-	-	-	-	-
10	0.229	0.212	0.119	0.113	-	-	-	-	-	-	-	-
11	-	-	-	-	-	-	-	-	-	-	-	-
12	0.226	0.214	0.118	0.114	-	-	-	-	-	-	-	-
13	-	-	-	-	-	-	-	-	-	-	-	-
14	0.224	0.215	0.117	0.115	-	-	-	-	-	-	-	-
15	-	-	-	-	-	-	-	-	-	-	-	-
16	0.222	0.216	0.116	0.116	-	-	-	-	-	-	-	-
17	-	-	-	-	-	-	-	-	-	-	-	-
18	0.220	0.216	0.115	0.117	-	-	-	-	-	-	-	-
19	-	-	-	-	-	-	-	-	-	-	-	-
20	0.219	0.217	0.114	0.118	-	-	-	-	-	-	-	-
21	-	-	-	-	-	-	-	-	-	-	-	-
22	0.217	0.217	0.114	0.119	-	-	-	-	-	-	-	-
23	-	-	-	-	-	-	-	-	-	-	-	-
24	0.216	0.217	0.113	0.120	-	-	-	-	-	-	-	-
25	-	-	-	-	-	-	-	-	-	-	-	-
26	0.215	0.217	0.113	0.121	-	-	-	-	-	-	-	-
27	-	-	-	-	-	-	-	-	-	-	-	-
28	0.213	0.217	0.112	0.122	-	-	-	-	-	-	-	-
29	-	-	-	-	-	-	-	-	-	-	-	-
30	0.212	0.217	0.112	0.122	-	-	-	-	-	-	-	-
31	-	-	-	-	-	-	-	-	-	-	-	-
32	0.210	0.217	0.112	0.123	-	-	-	-	-	-	-	-
33	-	-	-	-	-	-	-	-	-	-	-	-
34	0.209	0.216	0.111	0.124	-	-	-	-	-	-	-	-
35	-	-	-	-	-	-	-	-	-	-	-	-
36	0.207	0.216	0.111	0.125	-	-	-	-	-	-	-	-
37	-	-	-	-	-	-	-	-	-	-	-	-
38	0.206	0.215	0.111	0.126	-	-	-	-	-	-	-	-
39	-	-	-	-	-	-	-	-	-	-	-	-
40	0.204	0.215	0.111	0.127	-	-	-	-	-	-	-	-
41	-	-	-	-	-	-	-	-	-	-	-	-
42	0.203	0.214	0.110	0.128	-	-	-	-	-	-	-	-
43	-	-	-	-	-	-	-	-	-	-	-	-
44	0.201	0.214	0.110	0.129	-	-	-	-	-	-	-	-
45	-	-	-	-	-	-	-	-	-	-	-	-
46	0.199	0.213	0.110	0.130	-	-	-	-	-	-	-	-
47	-	-	-	-	-	-	-	-	-	-	-	-
48	0.198	0.212	0.110	0.131	-	-	-	-	-	-	-	-
49	-	-	-	-	-	-	-	-	-	-	-	-
50	0.196	0.211	0.110	0.132	-	-	-	-	-	-	-	-
51	-	-	-	-	-	-	-	-	-	-	-	-
52	0.194	0.210	0.110	0.133	-	-	-	-	-	-	-	-
53	-	-	-	-	-	-	-	-	-	-	-	-
54	0.192	0.209	0.109	0.133	-	-	-	-	-	-	-	-
55	-	-	-	-	-	-	-	-	-	-	-	-
56	0.191	0.208	0.109	0.133	-	-	-	-	-	-	-	-
57	-	-	-	-	-	-	-	-	-	-	-	-
58	0.191	0.208	0.109	0.133	-	-	-	-	-	-	-	-
59	-	-	-	-	-	-	-	-	-	-	-	-
60	0.191	-	-	0.109	-	-	-	-	-	-	-	-

Table 3.8: Einstein coefficients A of laser transitions of $^{13}\text{C}^{18}\text{O}_2$, s^{-1}

J	Regular			Sequence			Hot-e			Hot-f		
	10P	10R	9P	9R	10P	10R	9P	9R	10P	10R	9P	9R
0	-	0.382	-	0.330	- no data -			- no data -			- no data -	
1	0.126	0.256	0.109	0.221								
2	0.151	0.231	0.131	0.199								
3	0.161	0.221	0.140	0.190								
4	0.167	0.216	0.144	0.186								
5	0.170	0.213	0.147	0.183								
6	0.171	0.211	0.149	0.181								
7	0.173	0.209	0.150	0.180								
8	0.173	0.209	0.151	0.179								
9	0.174	0.208	0.152	0.179								
10	0.174	0.208	0.152	0.178								
11	0.174	0.208	0.152	0.178								
12	0.174	0.208	0.153	0.178								
13	0.173	0.208	0.153	0.179								
14	0.173	0.208	0.153	0.179								
15	0.173	0.208	0.153	0.179								
16	0.172	0.208	0.153	0.179								
17	0.171	0.209	0.153	0.180								
18	0.171	0.209	0.152	0.180								
19	0.170	0.209	0.152	0.181								
20	0.169	0.210	0.152	0.181								
21	0.169	0.210	0.152	0.182								
22	0.168	0.210	0.152	0.182								
23	0.167	0.211	0.152	0.182								
24	0.167	0.211	0.152	0.183								
25	0.166	0.211	0.152	0.184								
26	0.165	0.212	0.151	0.184								
27	0.164	0.212	0.151	0.185								
28	0.163	0.212	0.151	0.186								
29	0.162	0.213	0.150	0.186								
30	0.162	0.213	0.150	0.187								
31	0.161	0.213	0.150	0.188								
32	0.160	0.214	0.150	0.188								
33	0.159	0.214	0.150	0.189								
34	0.158	0.214	0.150	0.190								
35	0.157	0.215	0.150	0.191								
36	0.156	0.215	0.149	0.192								
37	0.155	0.215	0.149	0.192								
38	0.154	0.216	0.149	0.193								
39	0.153	0.216	0.149	0.194								
40	0.152	0.216	0.149	0.195								
41	0.151	0.217	0.148	0.195								
42	0.151	0.217	0.148	0.196								
43	0.149	0.217	0.148	0.197								
44	0.148	0.217	0.148	0.198								
45	0.148	0.217	0.148	0.199								
46	0.146	0.218	0.147	0.200								
47	0.145	0.218	0.147	0.201								
48	0.144	0.218	0.147	0.202								
49	0.143	0.218	0.147	0.203								
50	0.142	0.218	0.147	0.204								
51	0.141	0.218	0.147	0.205								
52	0.140	0.219	0.147	0.206								
53	0.139	0.219	0.146	0.207								
54	0.138	0.219	0.146	0.208								
55	0.137	0.219	0.146	0.209								
56	0.136	0.219	0.146	0.210								
57	0.135	0.219	0.146	0.211								
58	0.134	0.219	0.146	0.212								
59	0.133	0.219	0.146	0.213								
60	0.132	-	0.145	-								

Table 3.9: Einstein coefficients A of laser transitions of $^{838}\text{CO}_2$, s^{-1}

J	Regular			Sequence			Hot-e			Hot-f		
	10P	10R	9P	9R	10P	10R	9P	9R	10P	10R	9P	9R
0	-	0.053	-	0.070	-	-	-	-	-	-	-	-
1	-	-	-	-	-	-	-	-	-	-	-	-
2	0.105	0.068	0.138	0.090	-	-	-	-	-	-	-	-
3	-	-	-	-	-	-	-	-	-	-	-	-
4	0.090	0.072	0.118	0.096	-	-	-	-	-	-	-	-
5	-	-	-	-	-	-	-	-	-	-	-	-
6	0.086	0.074	0.112	0.098	-	-	-	-	-	-	-	-
7	-	-	-	-	-	-	-	-	-	-	-	-
8	0.083	0.075	0.109	0.101	-	-	-	-	-	-	-	-
9	-	-	-	-	-	-	-	-	-	-	-	-
10	0.082	0.076	0.107	0.102	-	-	-	-	-	-	-	-
11	-	-	-	-	-	-	-	-	-	-	-	-
12	0.081	0.076	0.106	0.103	-	-	-	-	-	-	-	-
13	-	-	-	-	-	-	-	-	-	-	-	-
14	0.080	0.077	0.105	0.105	-	-	-	-	-	-	-	-
15	-	-	-	-	-	-	-	-	-	-	-	-
16	0.080	0.077	0.104	0.105	-	-	-	-	-	-	-	-
17	-	-	-	-	-	-	-	-	-	-	-	-
18	0.079	0.077	0.103	0.106	-	-	-	-	-	-	-	-
19	-	-	-	-	-	-	-	-	-	-	-	-
20	0.079	0.077	0.103	0.107	-	-	-	-	-	-	-	-
21	-	-	-	-	-	-	-	-	-	-	-	-
22	0.078	0.077	0.102	0.108	-	-	-	-	-	-	-	-
23	-	-	-	-	-	-	-	-	-	-	-	-
24	0.077	0.077	0.102	0.109	-	-	-	-	-	-	-	-
25	-	-	-	-	-	-	-	-	-	-	-	-
26	0.077	0.077	0.101	0.110	-	-	-	-	-	-	-	-
27	-	-	-	-	-	-	-	-	-	-	-	-
28	0.077	0.077	0.101	0.111	-	-	-	-	-	-	-	-
29	-	-	-	-	-	-	-	-	-	-	-	-
30	0.076	0.077	0.100	0.111	-	-	-	-	-	-	-	-
31	-	-	-	-	-	-	-	-	-	-	-	-
32	0.075	0.077	0.100	0.112	-	-	-	-	-	-	-	-
33	-	-	-	-	-	-	-	-	-	-	-	-
34	0.075	0.077	0.100	0.113	-	-	-	-	-	-	-	-
35	-	-	-	-	-	-	-	-	-	-	-	-
36	0.074	0.077	0.099	0.114	-	-	-	-	-	-	-	-
37	-	-	-	-	-	-	-	-	-	-	-	-
38	0.074	0.077	0.099	0.115	-	-	-	-	-	-	-	-
39	-	-	-	-	-	-	-	-	-	-	-	-
40	0.073	0.076	0.099	0.116	-	-	-	-	-	-	-	-
41	-	-	-	-	-	-	-	-	-	-	-	-
42	0.073	0.076	0.099	0.116	-	-	-	-	-	-	-	-
43	-	-	-	-	-	-	-	-	-	-	-	-
44	0.072	0.076	0.098	0.117	-	-	-	-	-	-	-	-
45	-	-	-	-	-	-	-	-	-	-	-	-
46	0.071	0.075	0.098	0.118	-	-	-	-	-	-	-	-
47	-	-	-	-	-	-	-	-	-	-	-	-
48	0.071	0.075	0.098	0.119	-	-	-	-	-	-	-	-
49	-	-	-	-	-	-	-	-	-	-	-	-
50	0.070	0.075	0.098	0.120	-	-	-	-	-	-	-	-
51	-	-	-	-	-	-	-	-	-	-	-	-
52	0.070	0.075	0.098	0.121	-	-	-	-	-	-	-	-
53	-	-	-	-	-	-	-	-	-	-	-	-
54	0.069	0.074	0.098	0.122	-	-	-	-	-	-	-	-
55	-	-	-	-	-	-	-	-	-	-	-	-
56	0.068	0.073	0.098	0.123	-	-	-	-	-	-	-	-
57	-	-	-	-	-	-	-	-	-	-	-	-
58	0.068	0.073	0.098	0.124	-	-	-	-	-	-	-	-
59	-	-	-	-	-	-	-	-	-	-	-	-
60	0.067	-	-	0.097	-	-	-	-	-	-	-	-

$$E(\rho') = \int_{\rho} E(\rho) \frac{e^{-ikR}}{i\lambda R} J_0(k\Delta R/2) dS \quad (3.15)$$

where $dS = \pi(\rho + d\rho/2)^2 - \pi(\rho - d\rho/2)^2$ is the area of the ring, $R = (R_{max} + R_{min})/2$, $\Delta R = R_{max} - R_{min}$, $R_{max} = \sqrt{(\rho + \rho')^2 + \Delta z^2}$, $R_{min} = \sqrt{(\rho - \rho')^2 + \Delta z^2}$, and J_0 is the Bessel function.

3.5 Optical elements

Optical elements change the electric field as described by the following formulas. In these formulas $E(t, \rho)$ is the field in the input of the element and $E'(t, \rho)$ - after passing through the element.

Probe

(Passive surface, may be used as limiting aperture)

$$E'(t, \rho) = E(t, \rho),$$

Mask

(Opaque circular screen)

$$\begin{aligned} E'(t, \rho) &= 0, & \rho < R \\ E'(t, \rho) &= E(t, \rho), & \rho \geq R \end{aligned}$$

where R is the radius of the mask.

Absorber

(Ideal ND filter)

$$E'(t, \rho) = E(t, \rho) \sqrt{T},$$

where T is the transmittance.

Lens

(Ideal lens without aberrations)

$$E'(t, \rho) = E(t, \rho) \exp \left(\pi i \frac{\rho^2}{\lambda_c F} \right),$$

where F is the focal length of the lens and λ_c is the central wavelength.

Window

(Flat transparent window with zero reflection- and absorption- losses)

Linear dispersion

$$\begin{aligned} \hat{E}(\nu, \rho) &= \mathcal{F}(E(t, \rho)), \\ \hat{E}'(\nu, \rho) &= \hat{E}(\nu, \rho) \exp \left(-2\pi i \nu \frac{\Theta}{c} \left(n_0(\nu) - n_0(\nu_c) - \nu_c \frac{dn_0}{d\nu} \right) \right), \\ E'(t, \rho) &= \mathcal{F}^{-1}(\hat{E}'(\nu, \rho)), \end{aligned}$$

where $\hat{E}(\nu, \rho)$ is the field in the frequency domain, \mathcal{F} and \mathcal{F}^{-1} the Fourier transform and inverse Fourier transform functions, ν the frequency, ν_c the central frequency, n_0 is the linear component of the index of refraction and Θ is the thickness of the window.

Nonlinear interaction

$$E'(t, \rho) = E(t, \rho) \exp \left(-2\pi i \nu_c \frac{\Theta}{c} n_2 I(t, \rho) \right),$$

where n_2 is the nonlinear refractive index and $I(t, \rho)$ is the field intensity.

Stretcher

(*Stretcher or compressor*)

$$\begin{aligned}\hat{E}(\nu, \rho) &= \mathcal{F}(E(t, \rho)), \\ \hat{E}'(\nu, \rho) &= \hat{E}(\nu, \rho) \exp(-\pi i (\nu - \nu_c)^2 S), \\ E'(t, \rho) &= \mathcal{F}^{-1}(\hat{E}'(\nu, \rho)),\end{aligned}$$

where S is the stretching factor (in s/Hz, positive for red chirp).

3.6 Properties of optical materials

The following expressions and values for linear- (n_0) and nonlinear- (n_2) refractive indexes are used in the program (wavelength λ in the dispersion formulas must be expressed in μm):

CdTe

$$\begin{aligned}n_0 &= \sqrt{1 + \frac{6.1977889\lambda^2}{\lambda^2 - 0.1005326} + \frac{3.2243821\lambda^2}{\lambda^2 - 5279.518}} \quad [16] \\ n_2 &= -2.95 \times 10^{-13} \text{ cm}^2/\text{W at } 1.06 \mu\text{m} \quad [17]\end{aligned}$$

GaAs

$$\begin{aligned}n_0 &= \sqrt{5.372514 + \frac{5.466742\lambda^2}{\lambda^2 - 0.4431307^2} + \frac{0.02429960\lambda^2}{\lambda^2 - 0.8746453^2} + \frac{1.957522\lambda^2}{\lambda^2 - 36.9166^2}} \quad [18] \\ n_2 &= -3.26 \times 10^{-13} \text{ m}^2/\text{W at } 1.06 \mu\text{m} \quad [17]\end{aligned}$$

Ge

$$\begin{aligned}n_0 &= \sqrt{9.28156 + \frac{6.72880\lambda^2}{\lambda^2 - 0.44105} + \frac{0.21307\lambda^2}{\lambda^2 - 3870.1}} \quad [19] \\ n_2 &= 2.83 \times 10^{-13} \text{ cm}^2/\text{W at } 10.6 \mu\text{m} \quad [17]\end{aligned}$$

KCl

$$\begin{aligned}n_0 &= \sqrt{1.26486 + \frac{0.30523\lambda^2}{\lambda^2 - 0.100^2} + \frac{0.41620\lambda^2}{\lambda^2 - 0.131^2} + \frac{0.18870\lambda^2}{\lambda^2 - 0.162^2} + \frac{2.6200\lambda^2}{\lambda^2 - 70.42^2}} \quad [20] \\ n_2 &= 5.62 \times 10^{-16} \text{ cm}^2/\text{W at } 1.06 \mu\text{m} \quad [17]\end{aligned}$$

NaCl

$$\begin{aligned}n_0 &= \sqrt{1.00055 + \frac{0.19800\lambda^2}{\lambda^2 - 0.050^2} + \frac{0.48398\lambda^2}{\lambda^2 - 0.100^2} + \frac{0.38696\lambda^2}{\lambda^2 - 0.128^2} + \frac{0.25998\lambda^2}{\lambda^2 - 0.158^2} + \frac{0.08796\lambda^2}{\lambda^2 - 40.50^2} + \frac{3.17064\lambda^2}{\lambda^2 - 60.98^2} + \frac{0.30038\lambda^2}{\lambda^2 - 120.34^2}} \quad [20] \\ n_2 &= 4.38 \times 10^{-16} \text{ cm}^2/\text{W at } 1.06 \mu\text{m} \quad [17]\end{aligned}$$

Si

$$n_0 = 3.41983 + \frac{0.159906}{\lambda^2 - 0.028} - 0.123109 \left(\frac{1}{\lambda^2 - 0.028} \right)^2 + 1.26878 \times 10^{-6} \lambda^2 - 1.95104 \times 10^{-9} \lambda^4 \quad [21]$$
$$n_2 = 1.0 \times 10^{-13} \text{ cm}^2/\text{W at } 2.2 \text{ } \mu\text{m} \quad [22]$$

ZnSe

$$n_0 = \sqrt{1 + \frac{4.45813734\lambda^2}{\lambda^2 - 0.200859853^2} + \frac{0.467216334\lambda^2}{\lambda^2 - 0.391371166^2} + \frac{2.89566290\lambda^2}{\lambda^2 - 47.1362108^2}} \quad [23]$$
$$n_2 = 2.87 \times 10^{-14} \text{ cm}^2/\text{W at } 1.06 \text{ } \mu\text{m} \quad [17]$$

Bibliography

- [1] N. V. Karlov and Y. B. Konev. High pressure pulsed CO₂ lasers. In A. M. Prokhorov, editor, *Handbook on lasers*. Sovetskoe Radio, Moscow, 1978 (in Russian).
- [2] T. Holstein. Energy distribution of electrons in high frequency gas discharges. *Phys. Rev.*, 70:367–384, 1946.
- [3] W. L. Nighan. Electron energy distributions and collision rates in electrically excited N₂, CO, and CO₂. *Phys. Rev. A*, 2:1989–2000, 1970.
- [4] R. D. Hake and A. V. Phelps. Momentum-transfer and inelastic-collision cross sections for electrons in O₂, CO, and CO₂. *Phys. Rev.*, 158:70–84, 1967.
- [5] L. S. Frost and A. V. Phelps. Rotational excitation and momentum transfer cross sections for electrons in H₂ and N₂ from transport coefficients. *Phys. Rev.*, 127:1621–1633, 1962.
- [6] J. J. Lowke, A. V. Phelps, and B. W. Irwin. Predicted electron transport coefficients and operating characteristics of CO₂-N₂-He laser mixtures. *J. Appl. Phys.*, 44:4664–4671, 1973.
- [7] Y. D. Oksyuk. Excitation of the rotational levels of diatomic molecules by electron impact in the adiabatic approximation. *Sov. J. Exp. Theor. Phys.*, 22:873, 1966.
- [8] N. Chandra and P. G. Burke. Rotational excitation cross sections for e[−]-N₂ scattering. *J. Phys. B: At. Mol. Phys.*, 6:2355–2357, 1973.
- [9] A. V. Phelps. Rotational and vibrational excitation of molecules by low-energy electrons. *Rev. Mod. Phys.*, 40:399–410, 1968.
- [10] G. J. Schulz. Vibrational excitation of nitrogen by electron impact. *Phys. Rev.*, 125:229–232, 1962.
- [11] A. G. Engelhardt, A. V. Phelps, and C. G. Risk. Excitation of the rotational levels of diatomic molecules by electron impact in the adiabatic approximation. *Phys. Rev.*, 135:A1566–A1574, 1964.
- [12] A. S. Biryukov, V. K. Konyukhov, A. I. Lukovnikov, and R. I. Serikov. Relaxation of the vibrational energy of the (00⁰1) level of the CO₂ molecule. *Sov. J. Exp. Theor. Phys.*, 39:610, 1974.
- [13] R. L. Taylor and S. Bitterman. Survey of vibrational relaxation data for processes important in the CO₂-N₂ laser system. *Rev. Mod. Phys.*, 41:26–47, 1969.
- [14] H. C. Volkin. Calculation of short-pulse propagation in a large CO₂-laser amplifier. *J. Appl. Phys.*, 50:1179–1188, 1979.
- [15] R. D. Hake and A. V. Phelps. Einstein coefficients, cross sections, f values, dipole moments, and all that. *arXiv:physics/0202029*, 2002.
- [16] A. G. DeBell, E. L. Dereniak, J. Harvey, J. Palmer J. Nissley, A. Selvarajan, and W. L. Wolfe. Cryogenic refractive indices and temperature coefficients of cadmium telluride from 6 μm to 22 μm. *Appl. Opt.*, 18:3114–3115, 1979.

- [17] M. Sheik-Bahae. Dispersion of bound electron nonlinear refraction in solids. *IEEE J. Quant. Electron.*, 27:1296–1309, 1991.
- [18] T. Skauli, P. S. Kuo, K. L. Vodopyanov, T. J. Pinguet, O. Levi, L. A. Eyres, J. S. Harris, M. M. Fejer, L. Becouarn B. Gerard, and E. Lallier. Improved dispersion relations for gaas and applications to nonlinear optics. *J. Appl. Opt.*, 94:6447–6455, 2003.
- [19] N. P. Barnes and M. S. Piltch. Temperature-dependent sellmeier coefficients and nonlinear optics average power limit for germanium. *J. Opt. Soc. Am.*, 69:178–180, 1979.
- [20] H. H. Li. Refractive index of alkali halides and its wavelength and temperature derivatives. *J. Phys. Chem. Ref. Data*, 5:329–528, 1976.
- [21] D. F. Edwards and E. Ochoa. Infrared refractive index of silicon. *Appl. Opt.*, 19:4130–4131, 1980.
- [22] A. D. Bristow, N. Rotenberg, and H. M. van Driel. Two-photon absorption and kerr coefficients of silicon for 850–2200 nm. *Appl. Phys. Lett.*, 90:191104, 2007.
- [23] B. Tatian. Fitting refractive-index data with the sellmeier dispersion formula. *Appl. Opt.*, 23:4477–4485, 1984.