

A fully adaptive, high-order, fast Poisson solver for complex two-dimensional geometries



David Stein

Dan Fortunato
Flatiron Institute

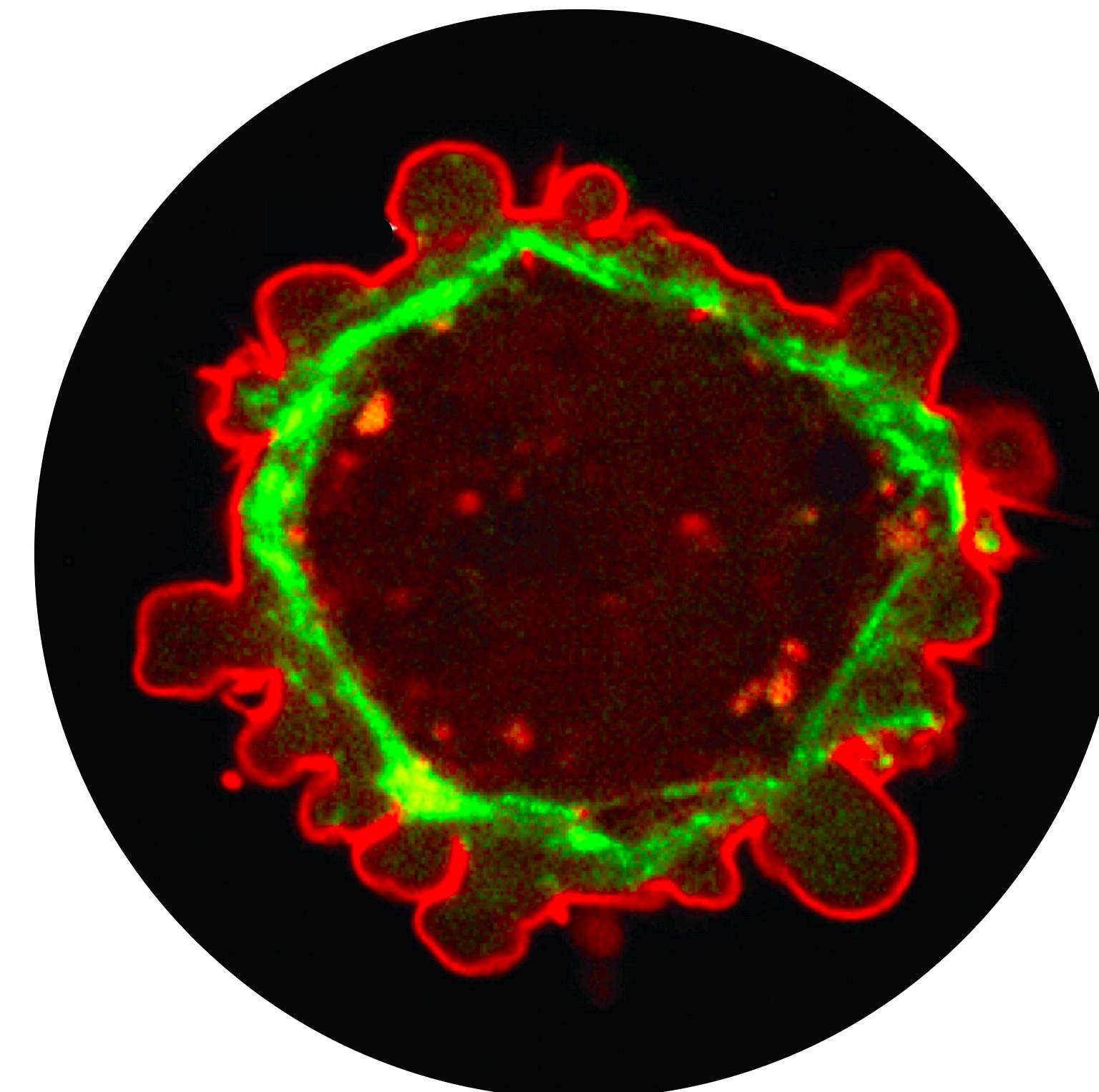


Alex Barnett

Introduction

The challenge of geometry

Geometry can drive the dynamics of physical phenomena, but can pose a numerical challenge.

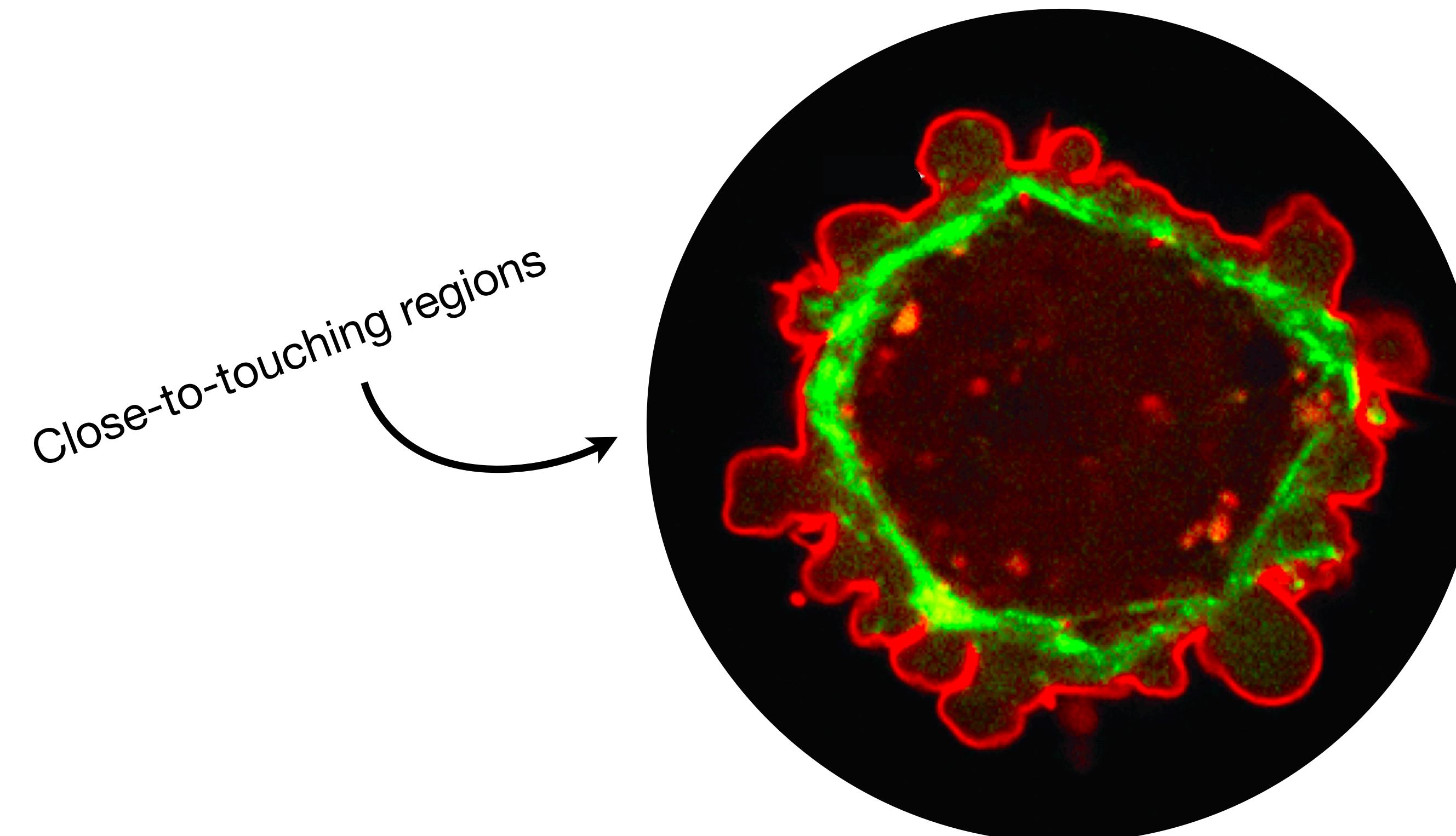


[Charras, 2008]

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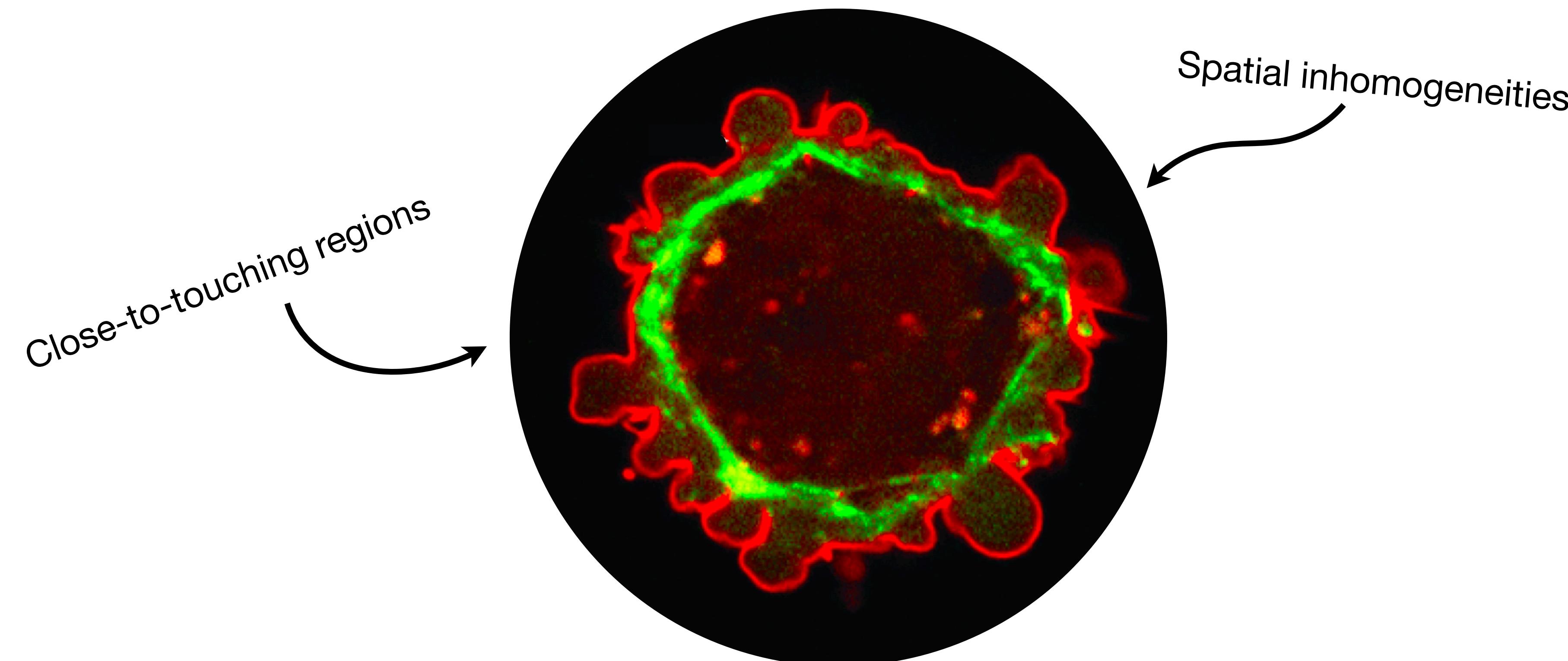


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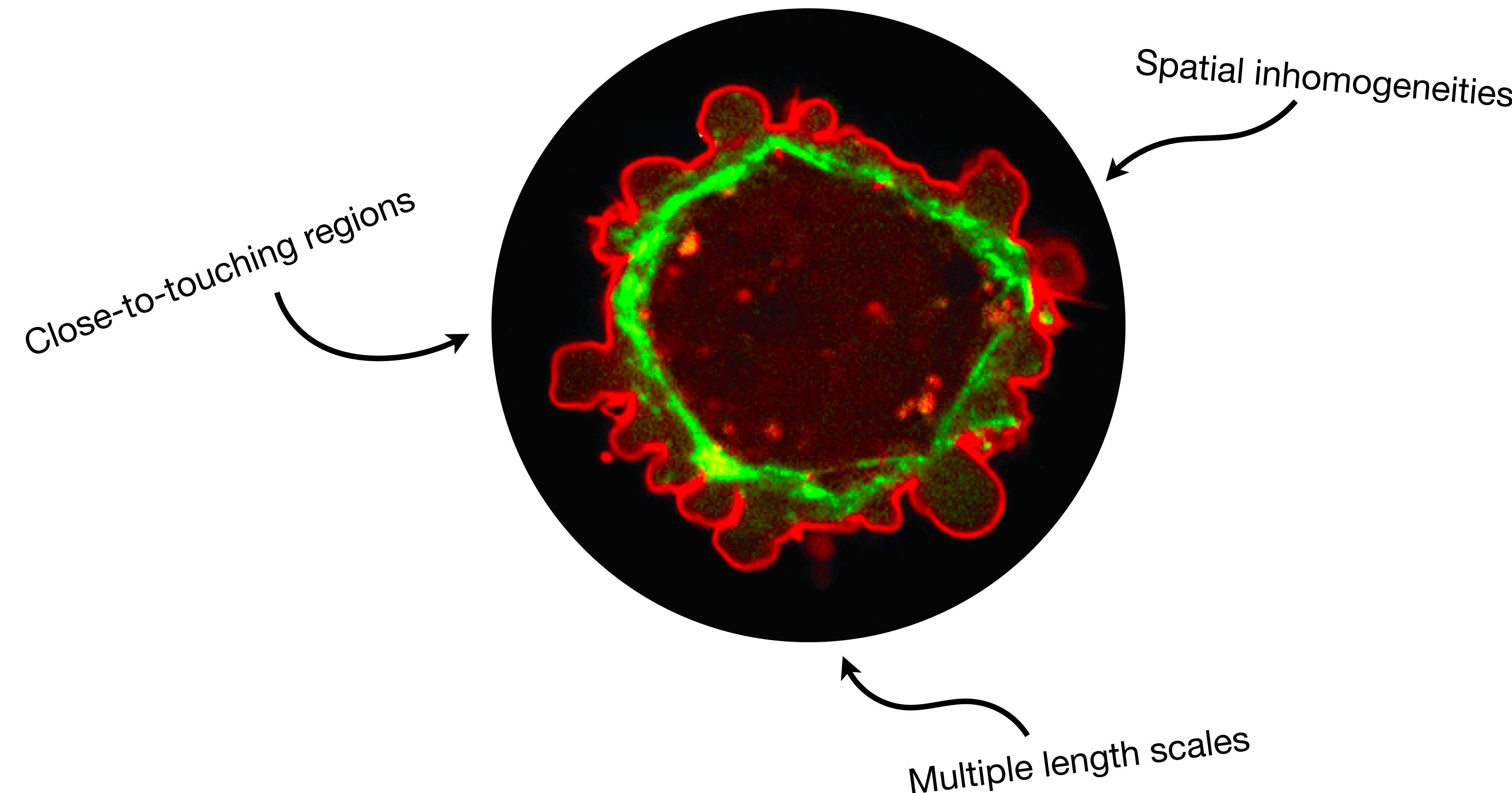


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Inhomogeneous PDEs

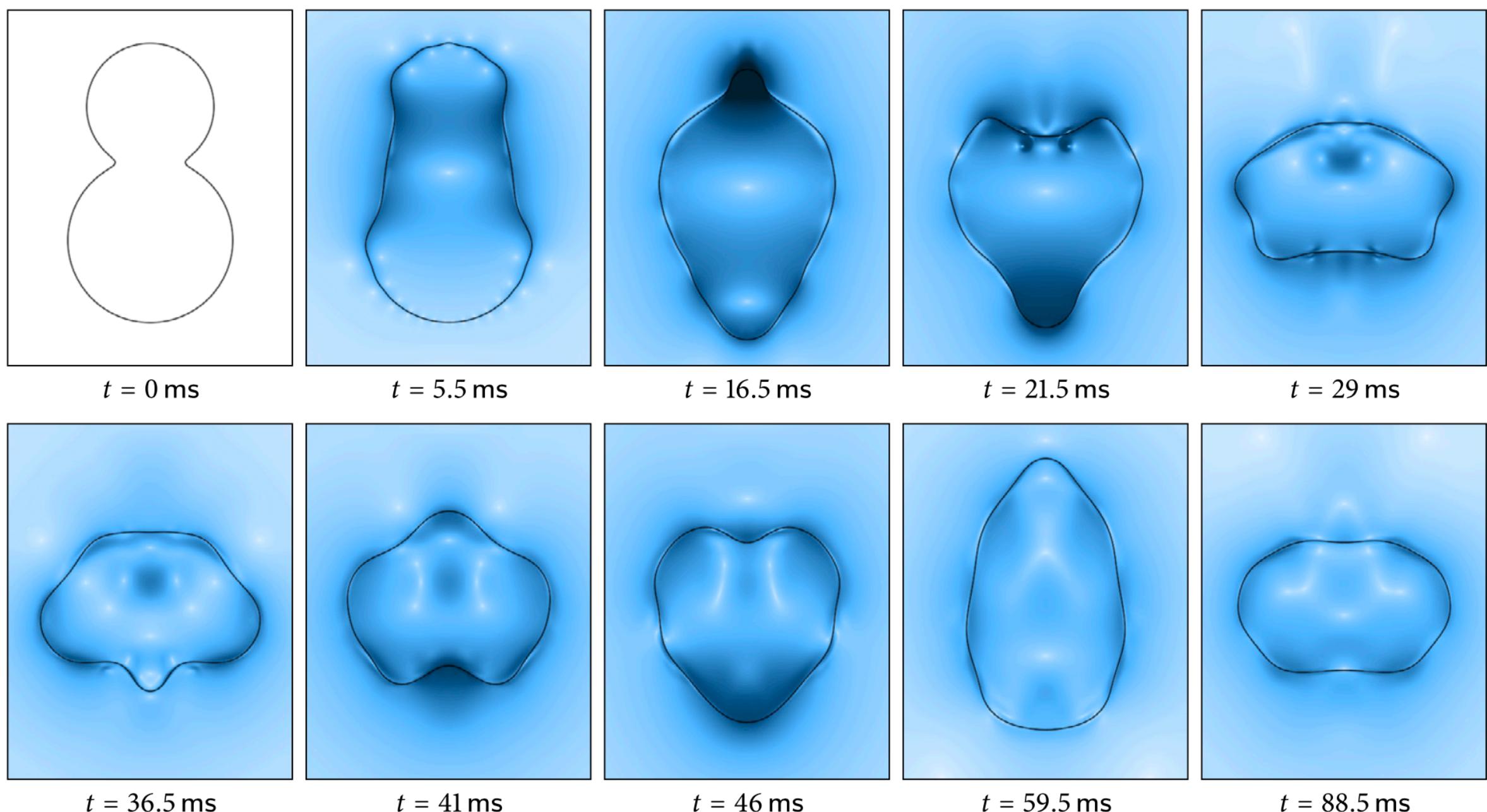
Introduction

Many applications involve solving an inhomogeneous elliptic BVP.

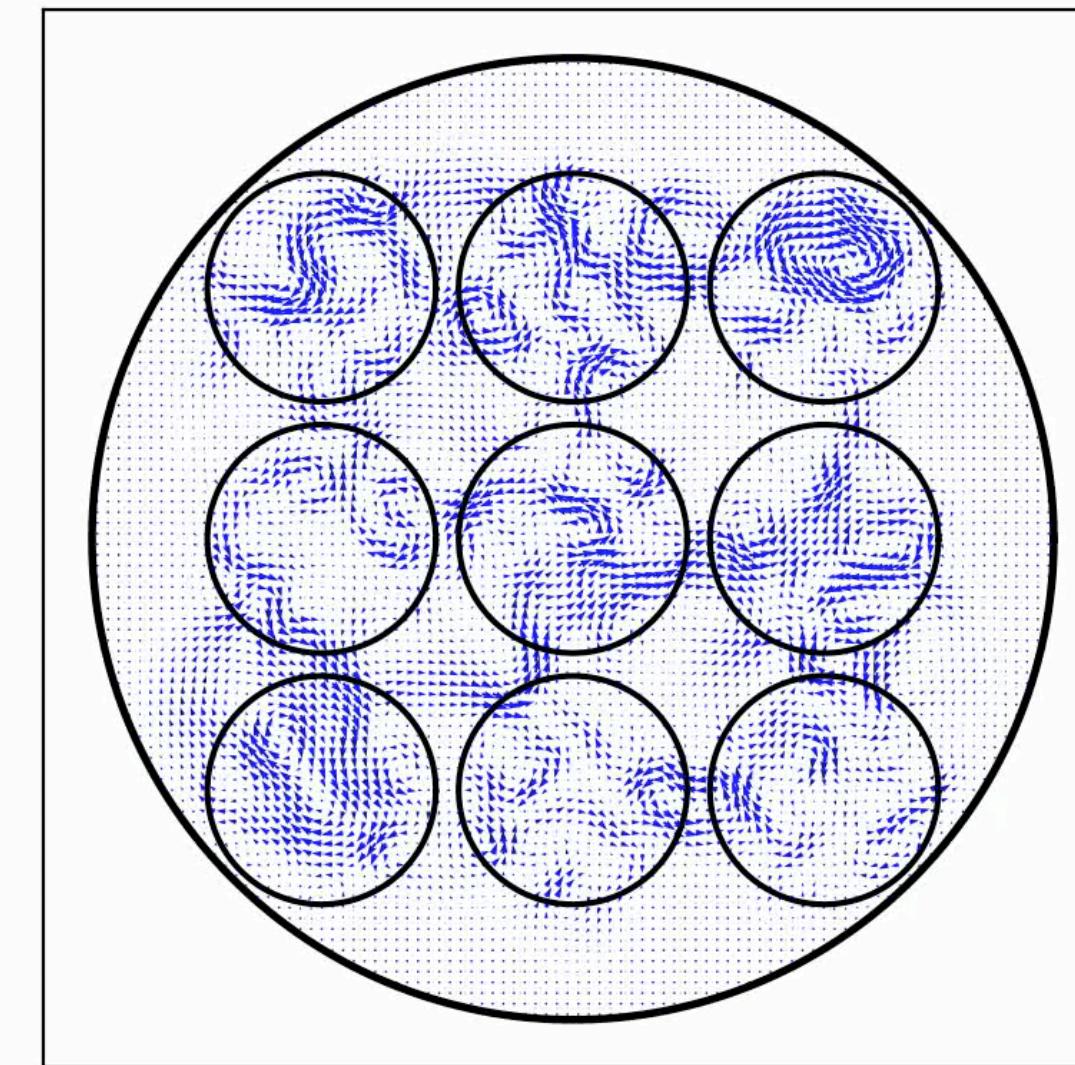
$$Lu = f \quad \text{in } \Omega \quad (\text{e.g. Poisson, Helmholtz, Stokes, ...})$$

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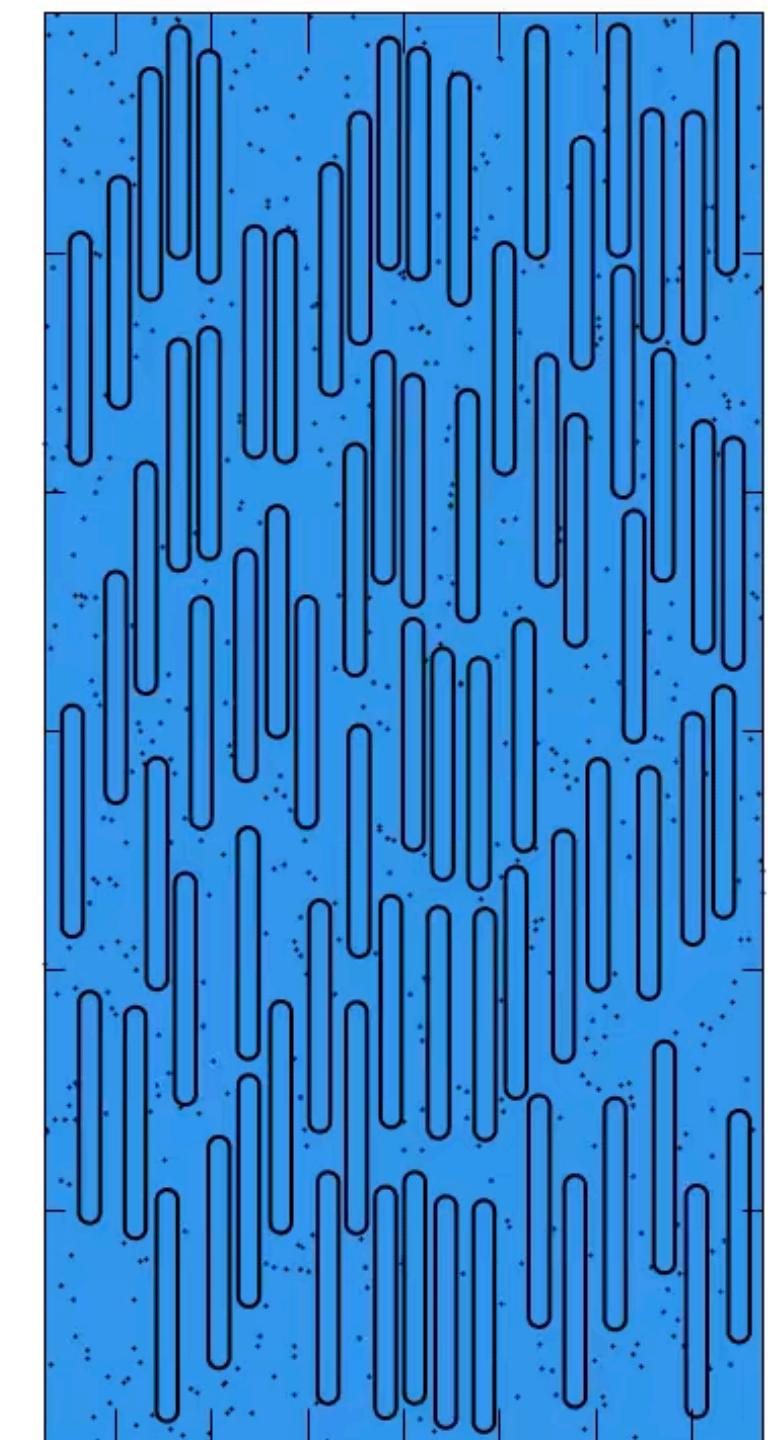
Bubble collision [Saye, 2017]



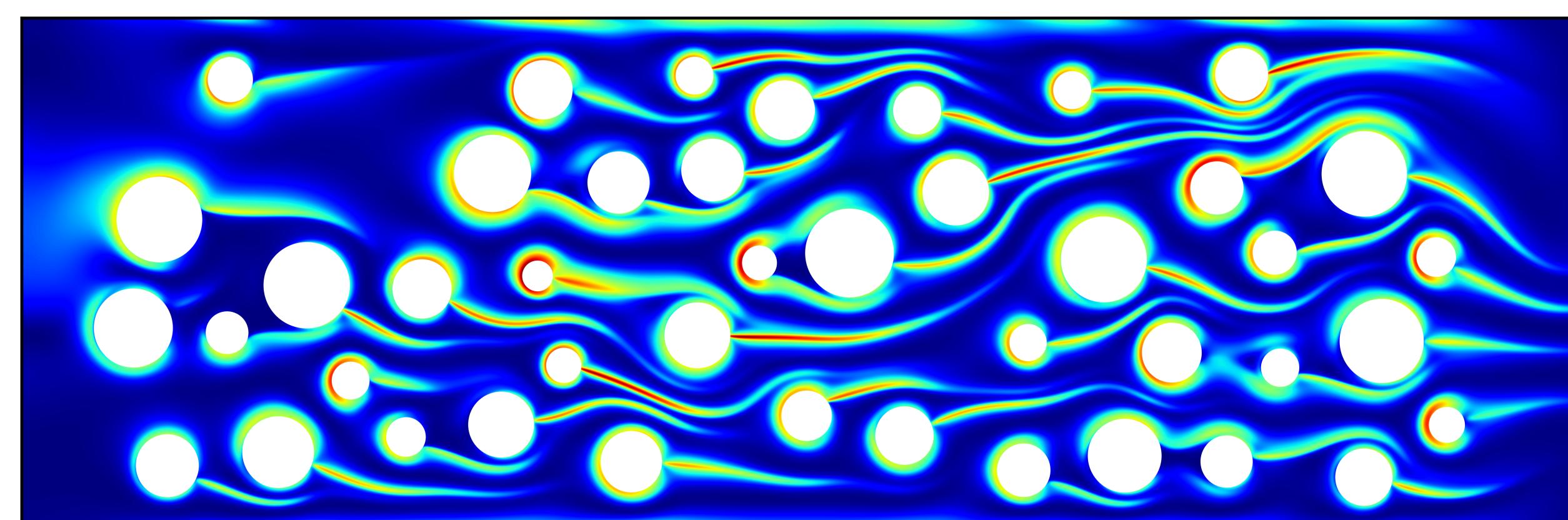
Active droplets [Stein, 2021]



Fluid-structure interaction
[Rycroft et al., 2020]



Non-Newtonian fluids [Stein et al., 2019]



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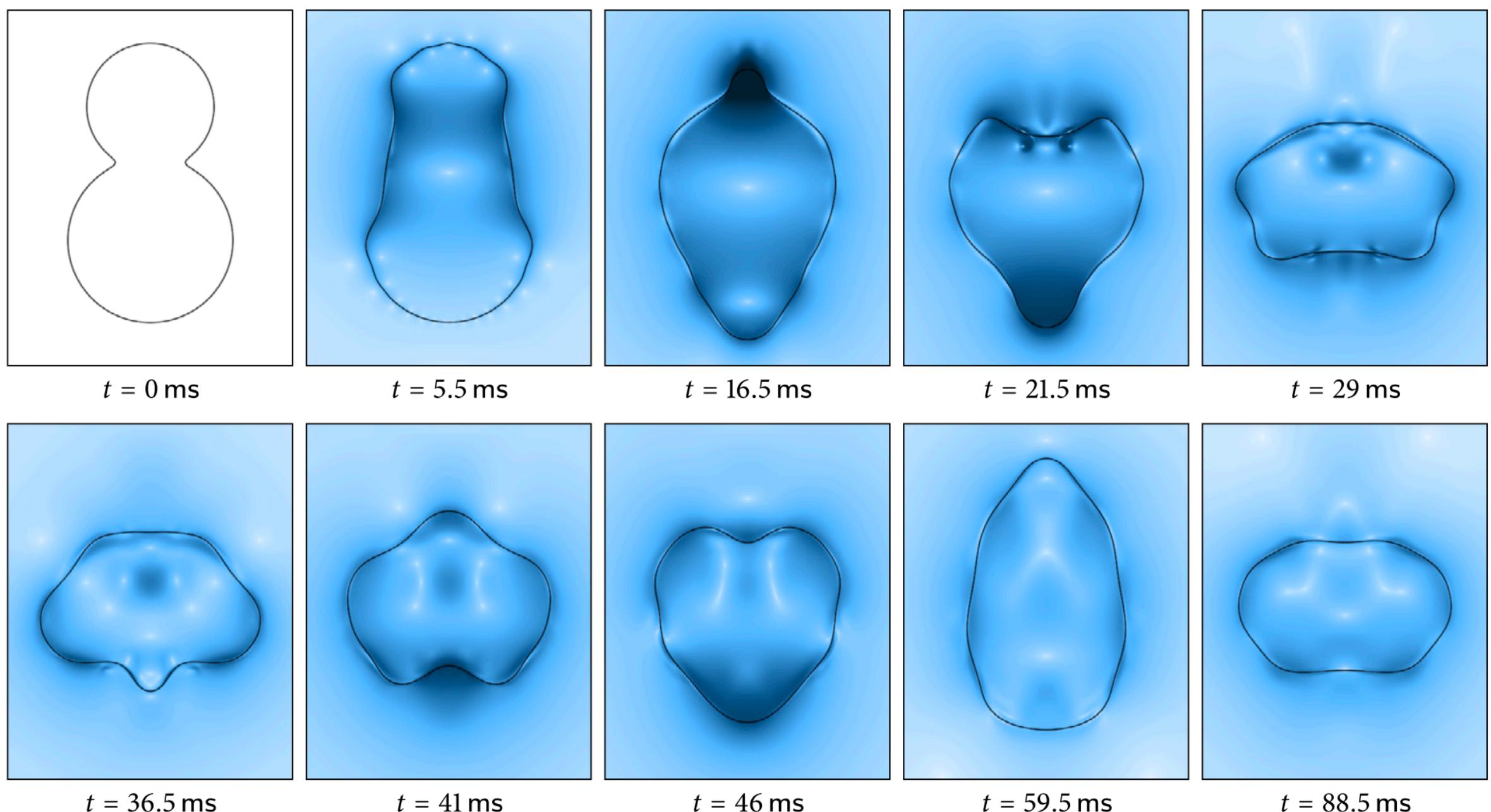
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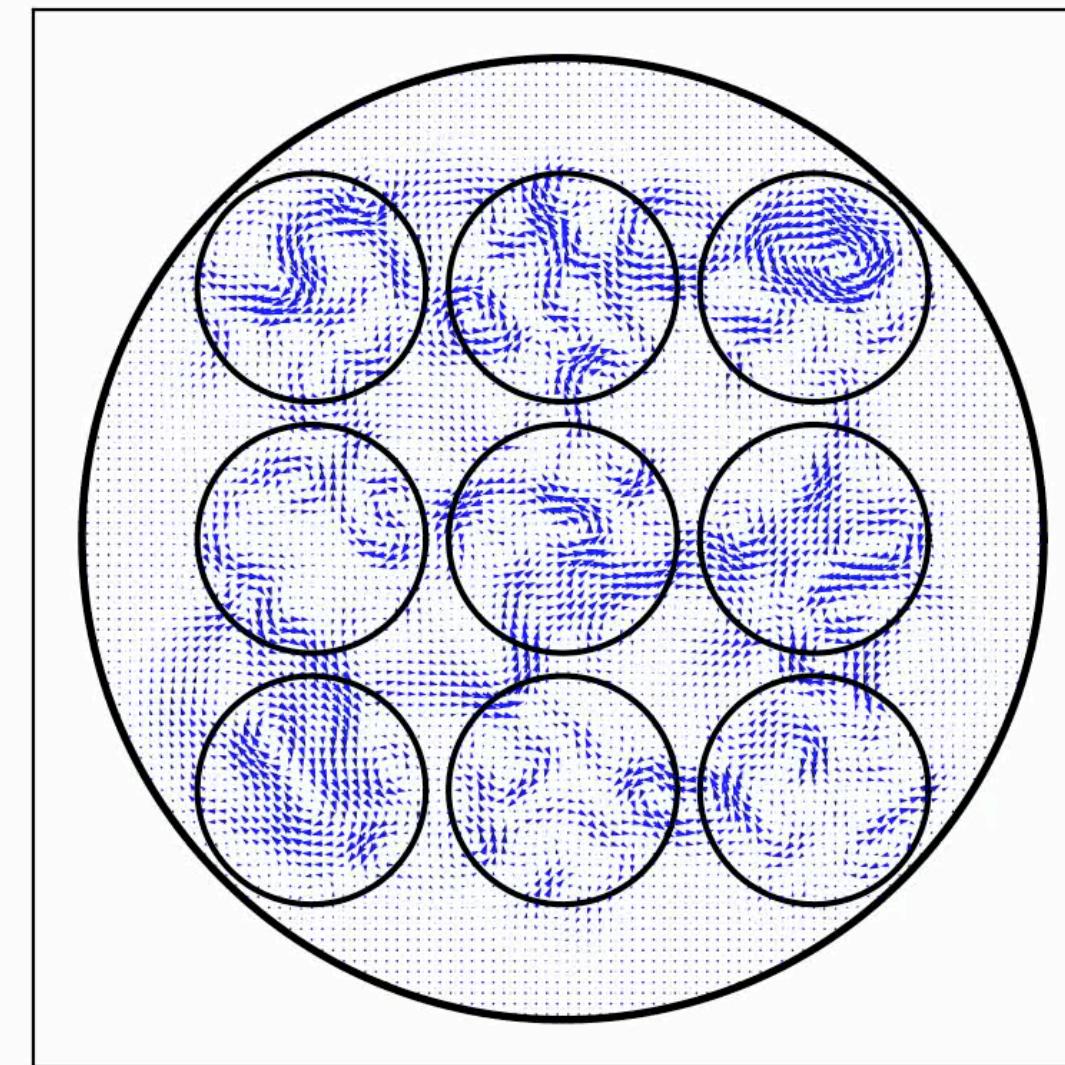
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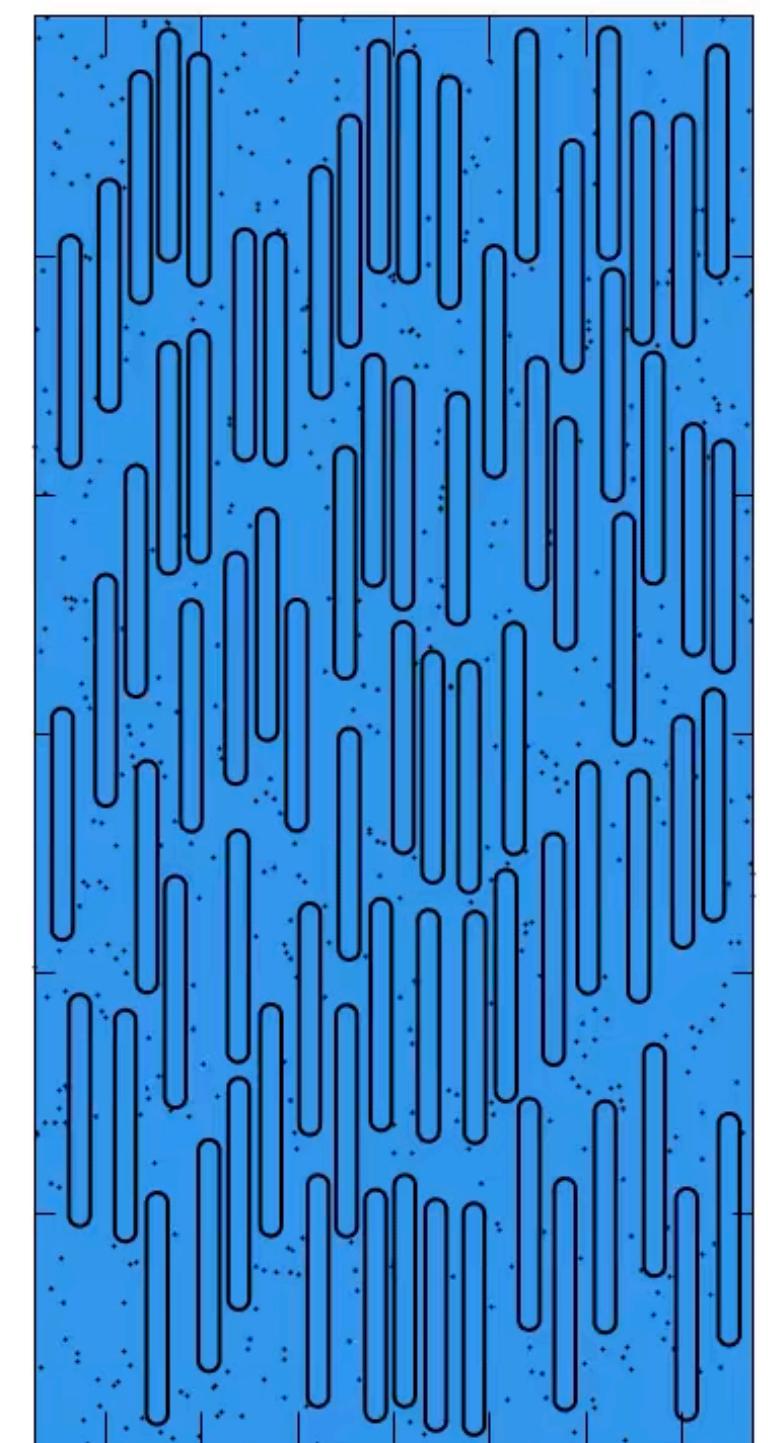
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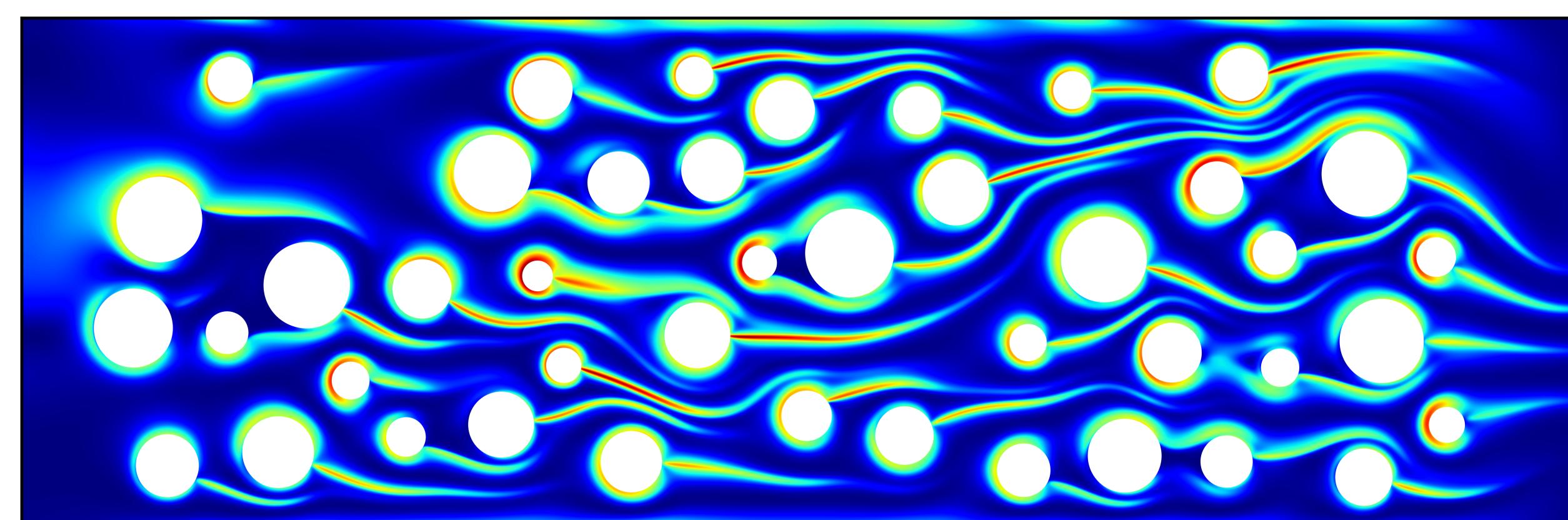
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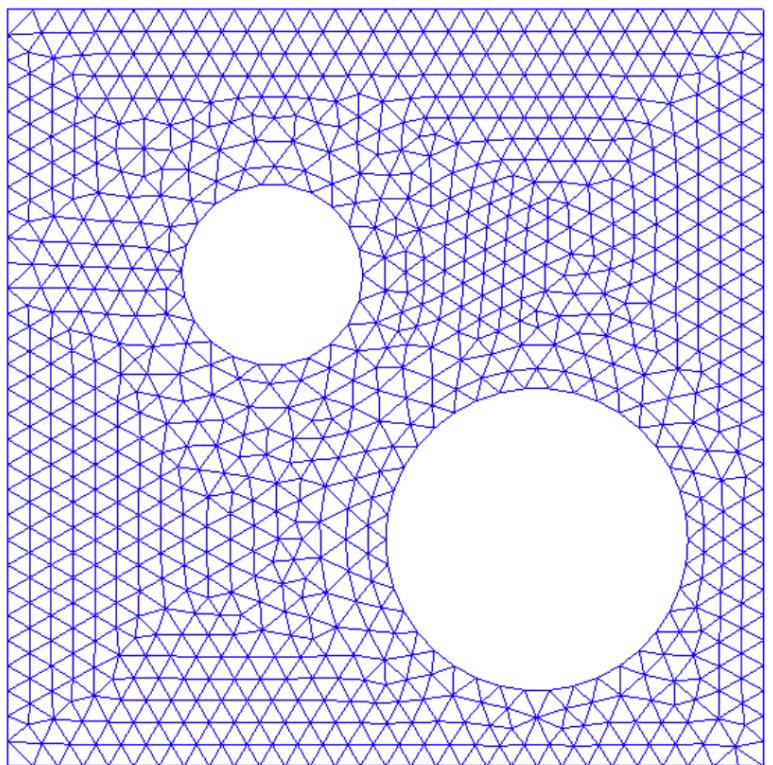
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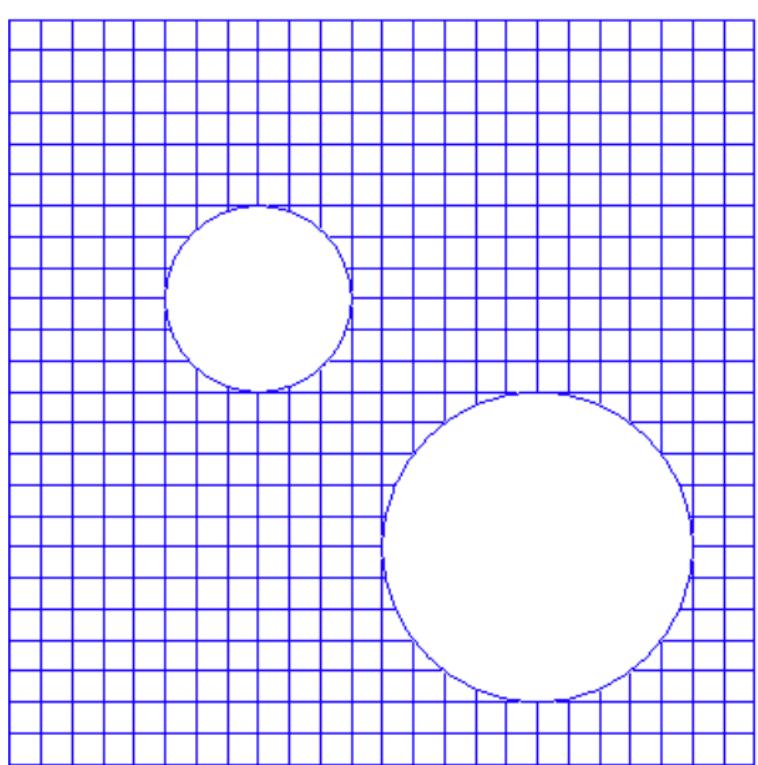
Inhomogeneous PDEs

Many approaches for inhomogeneous BVP

Traditional



Conforming



Cut-cell

- Mesh generation (or cut-cell generation)
- Directly discretize with FEM, FDM, SEM, ...

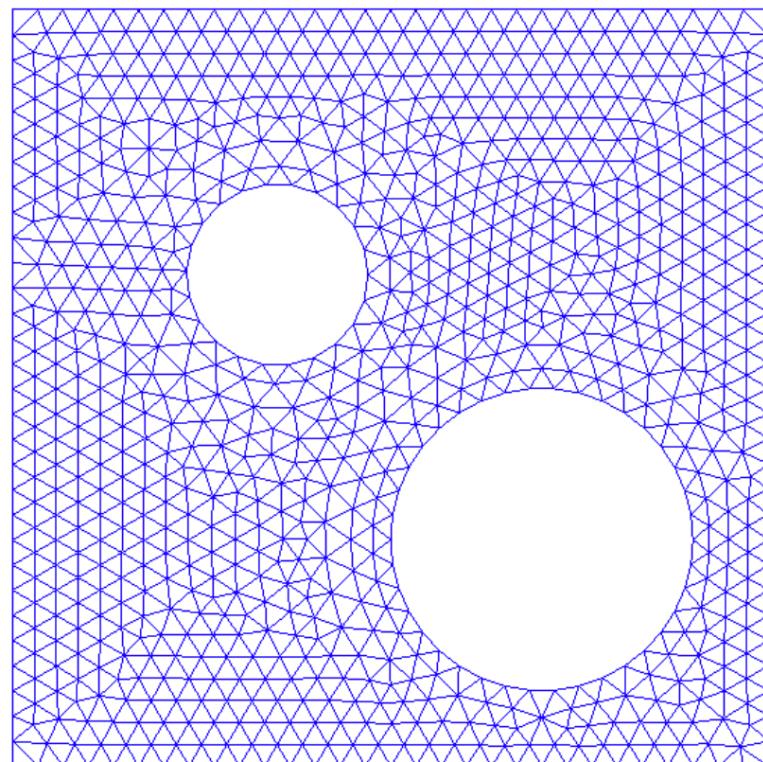
$$Ax = b$$

Discretize on a mesh.
Solve linear system for volume DoFs.

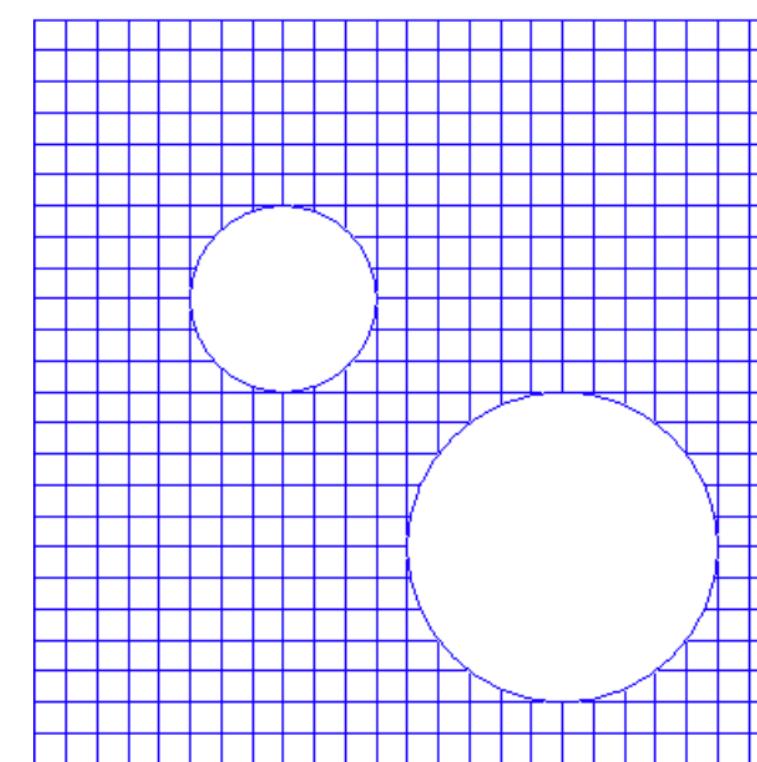
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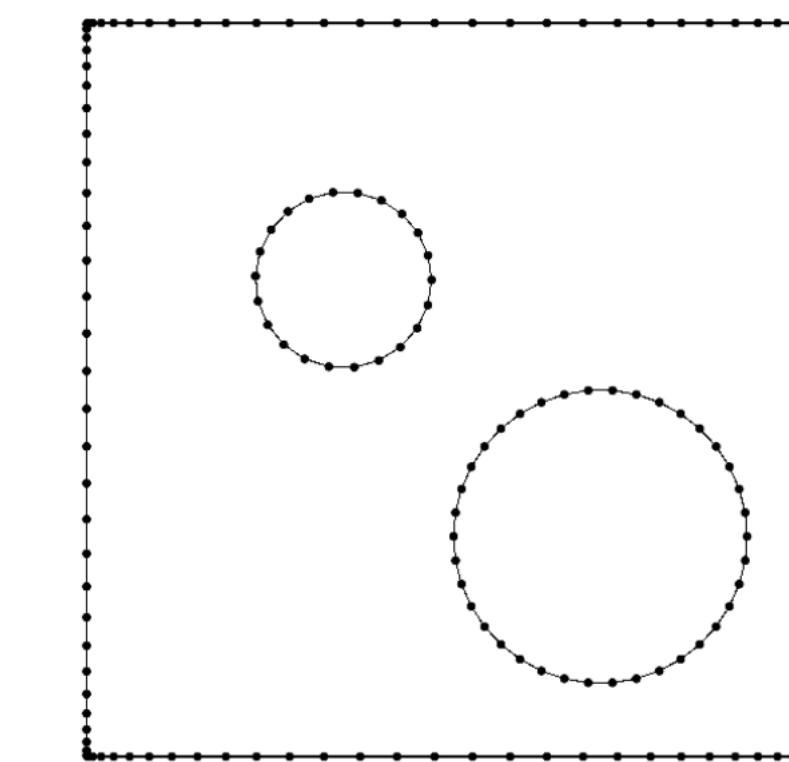
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Analysis-based



- Write solution as $u = v + w$

particular solution *homogeneous solution*

- Find **some** (any!) function V such that

$$Lv = f \quad \text{in } \Omega$$

No solve required. Convolve
with Green's function.

- Compute u_H to satisfy boundary conditions:

$$Lw = 0 \quad \text{in } \Omega$$

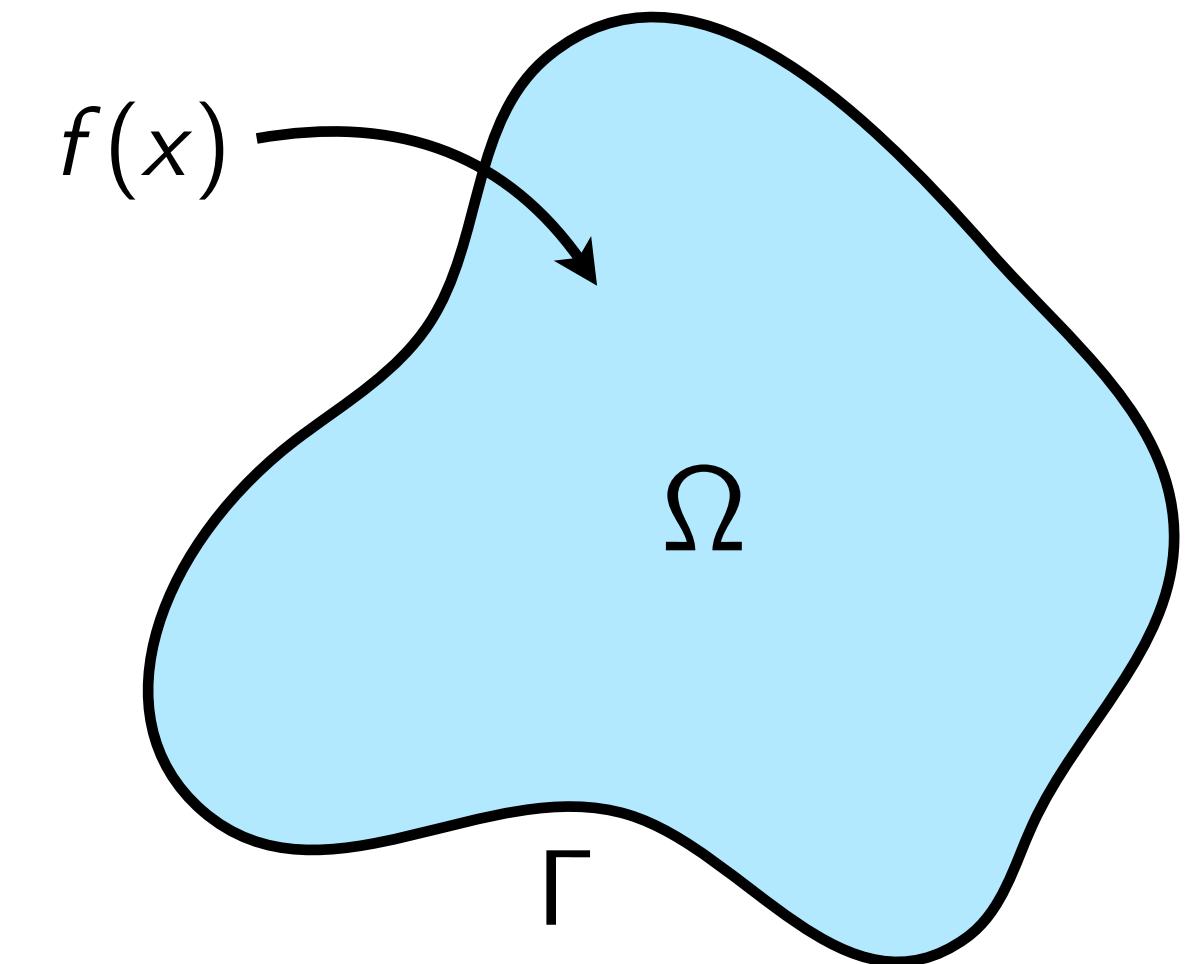
$$Bw = g - Bv \quad \text{on } \Gamma$$

Solve linear system for
boundary DoFs using BIE.

Potential theory

Volume potentials

In many applications $f(x)$ is only given inside Ω . So we have a few options...



Potential theory

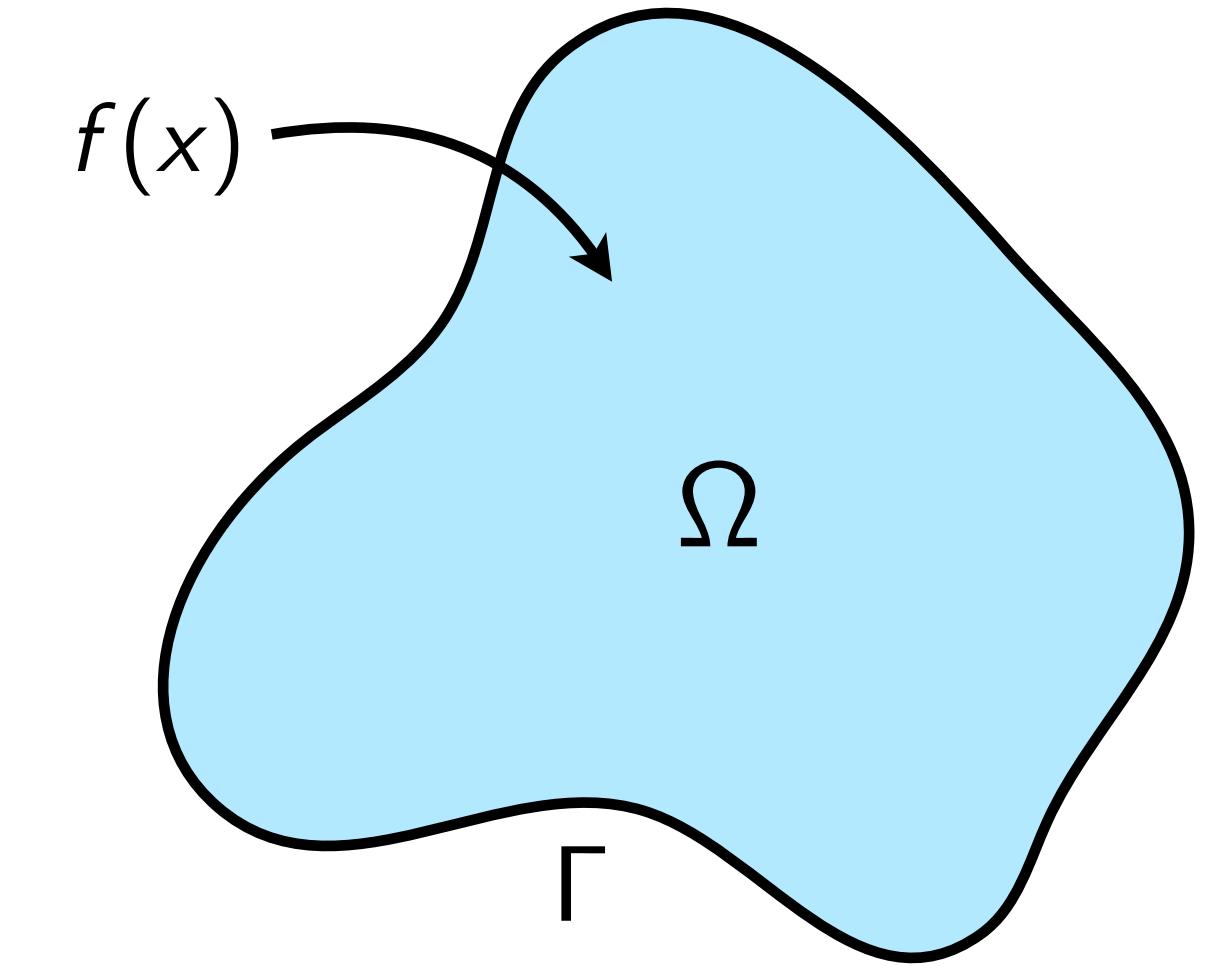
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- Build accurate quadrature scheme over Ω (e.g. adaptive boxes with cut cells near boundary) and compute

$$v(x) = \int_{\Omega} \Phi(x, y) f(y) dy \quad \left(\text{Poisson: } \Phi(x, y) = \frac{1}{2\pi} \log |x - y| \right)$$

free space fundamental solution



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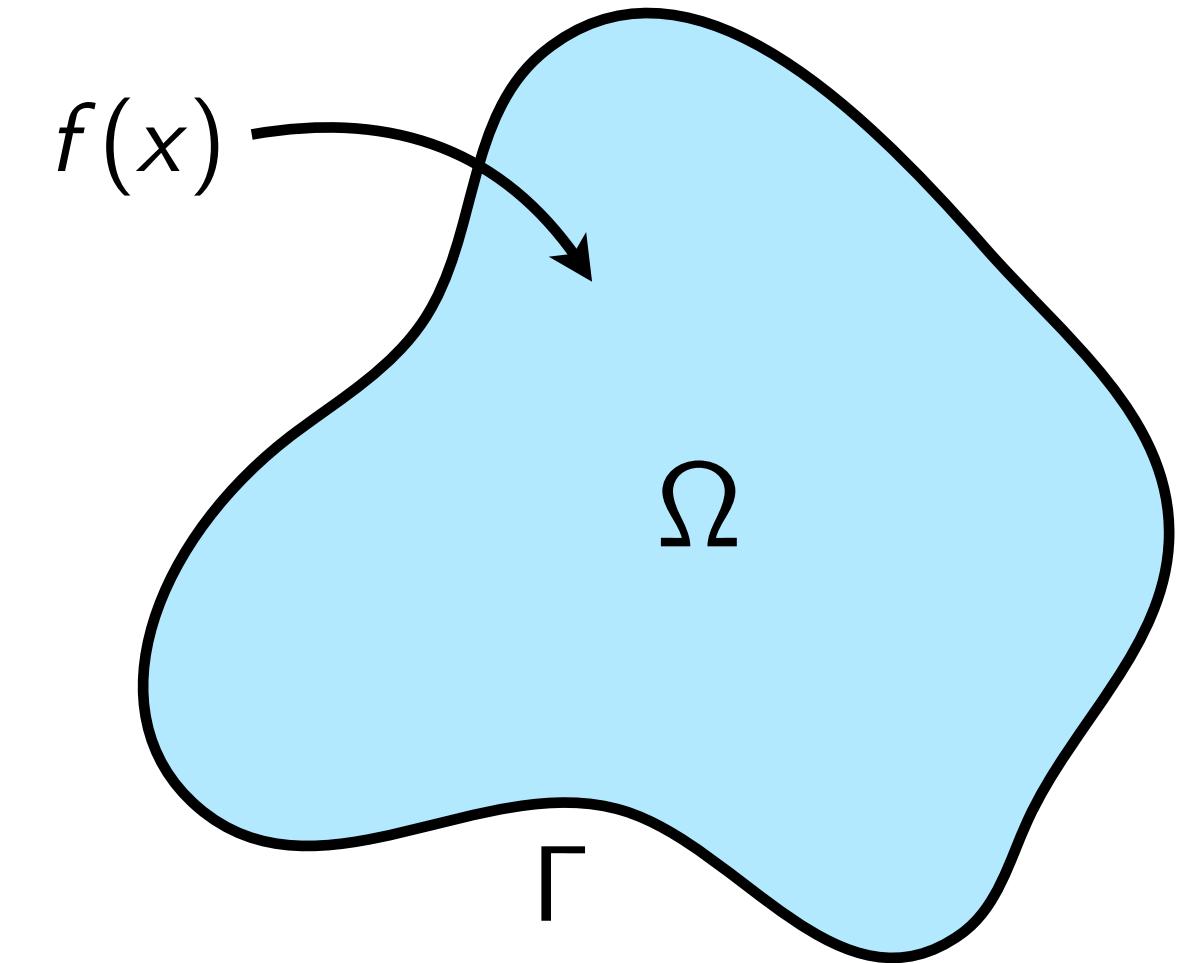
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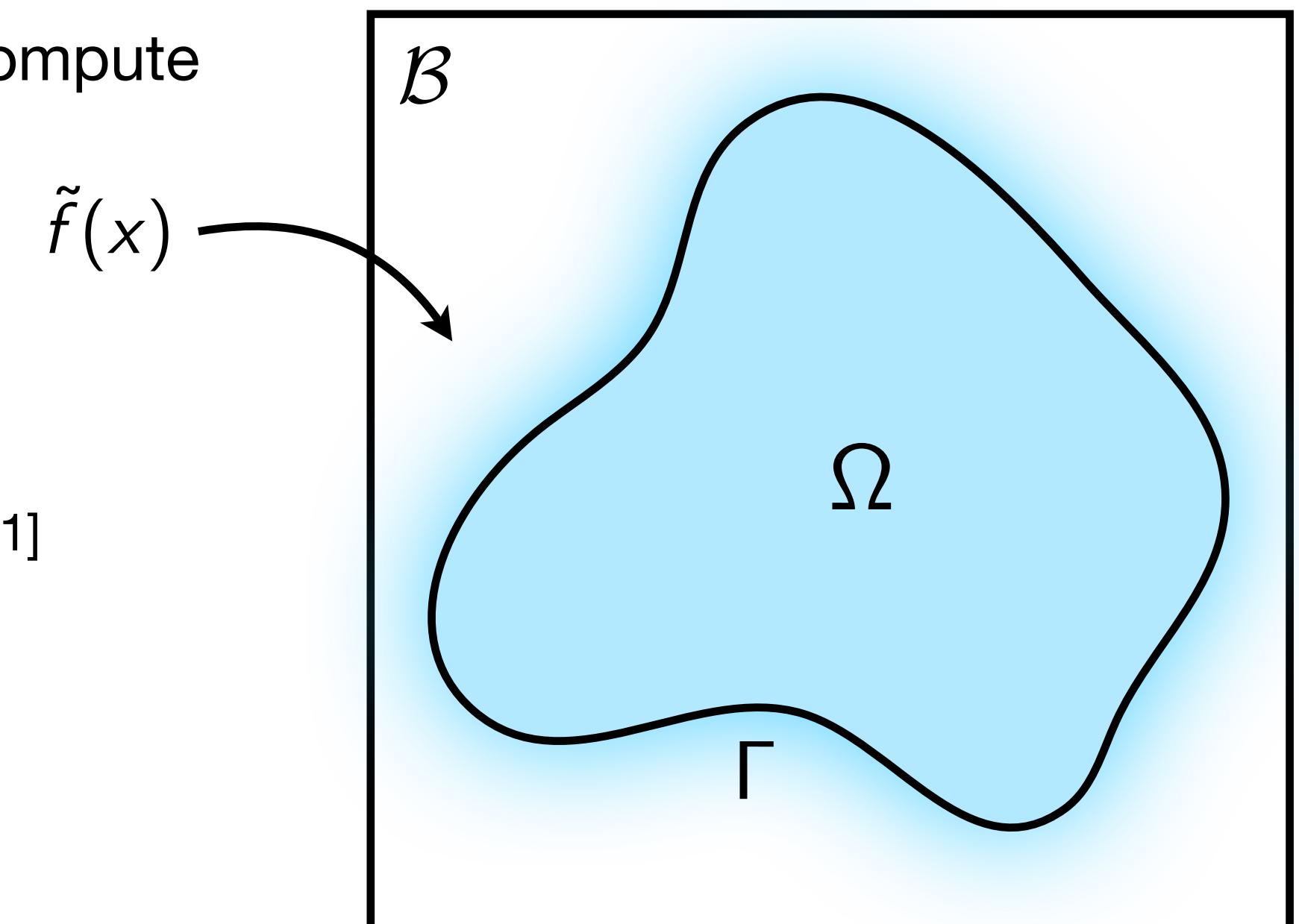
free space fundamental solution



- Extend f to \tilde{f} outside Ω (“function extension”). Adaptively resolve \tilde{f} and compute

$$v(x) = \int_{\mathcal{B}} \Phi(x, y) \tilde{f}(y) dy$$

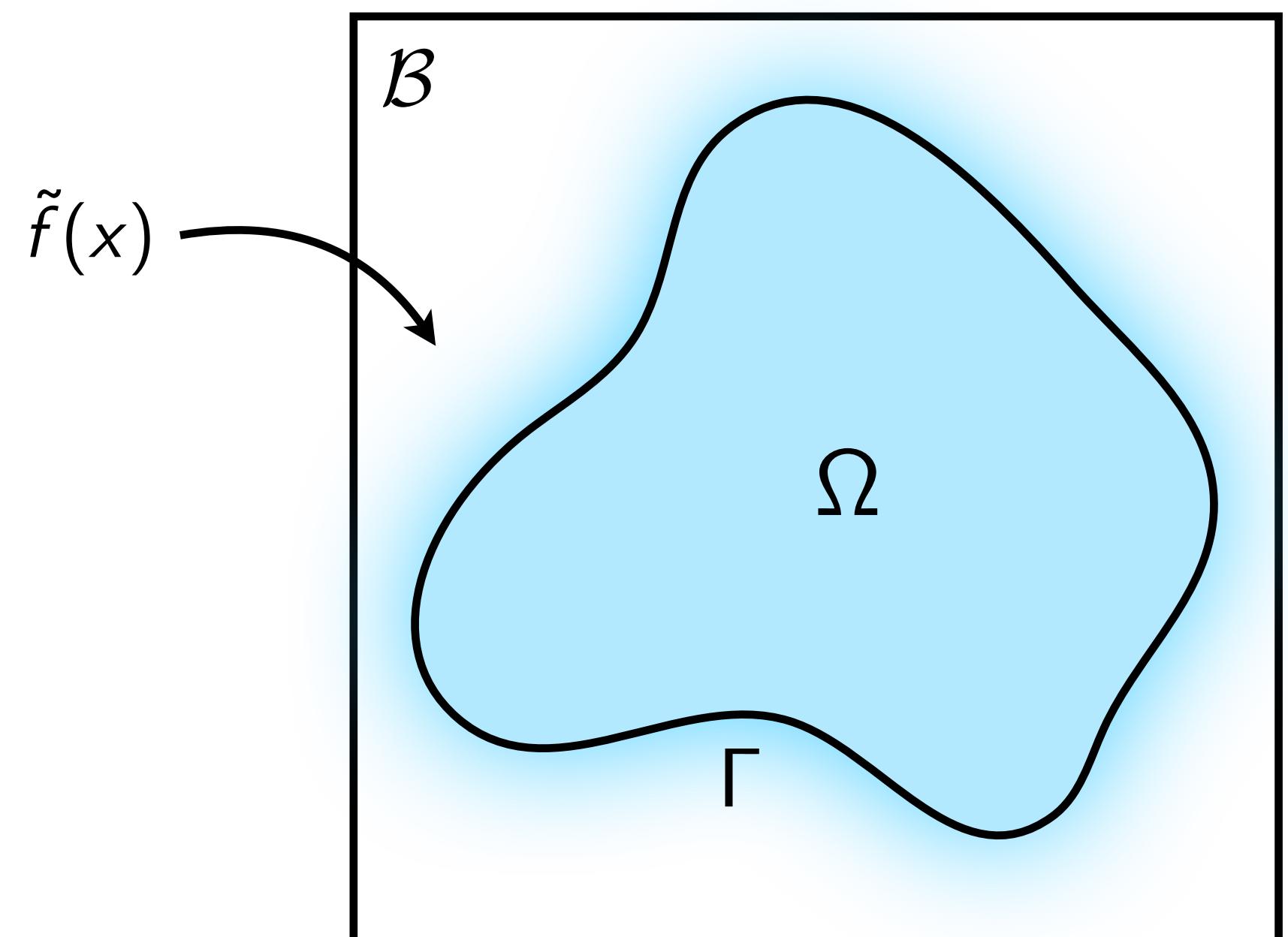
- Quadrature tables for boxes can be **precomputed**.
- Apply FMM to adaptive quadtree of \tilde{f} (“box code”). [Ethridge & Greengard, 2001]
- Want v as smooth as U so no effort is wasted.



Function extension

Prior work

- Finite difference extension, second-order accurate [Mayo, 1984]
- Fourier continuation [Bruno & Lyon, 2010], [Bruno & Paul, 2020]
- Immersed boundary smooth extension [Stein, Guy, & Thomases, 2015]
- C^k polyharmonic extension + box code, fourth-order accurate [Askham & Cerfon, 2017]
- Partition of unity extension [Fryklund, Lehto, & Tornberg, 2018], [Fryklund & Greengard, 2022]

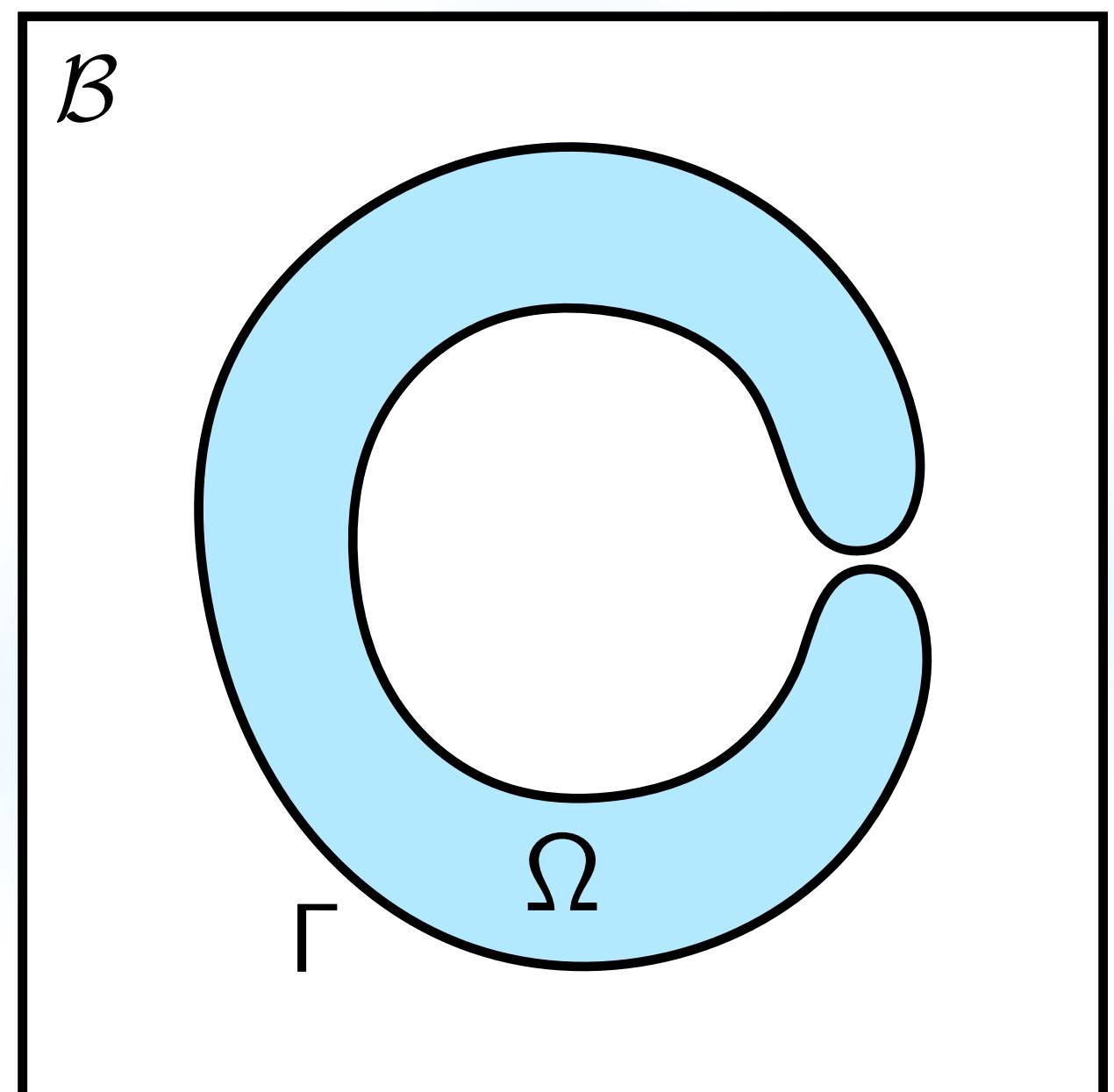


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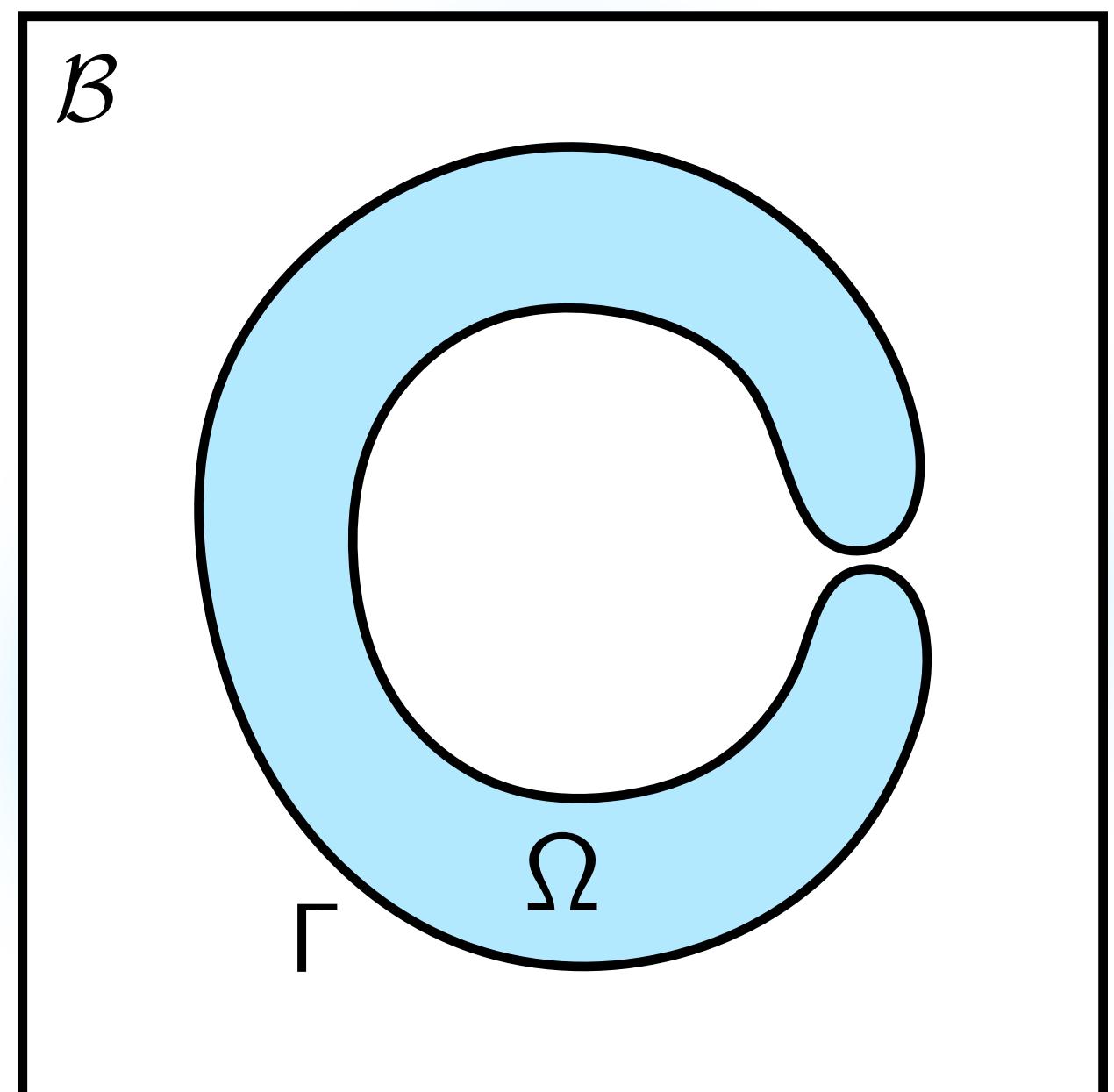
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- multiscale geometry
- multiscale data
- close-to-touching regions

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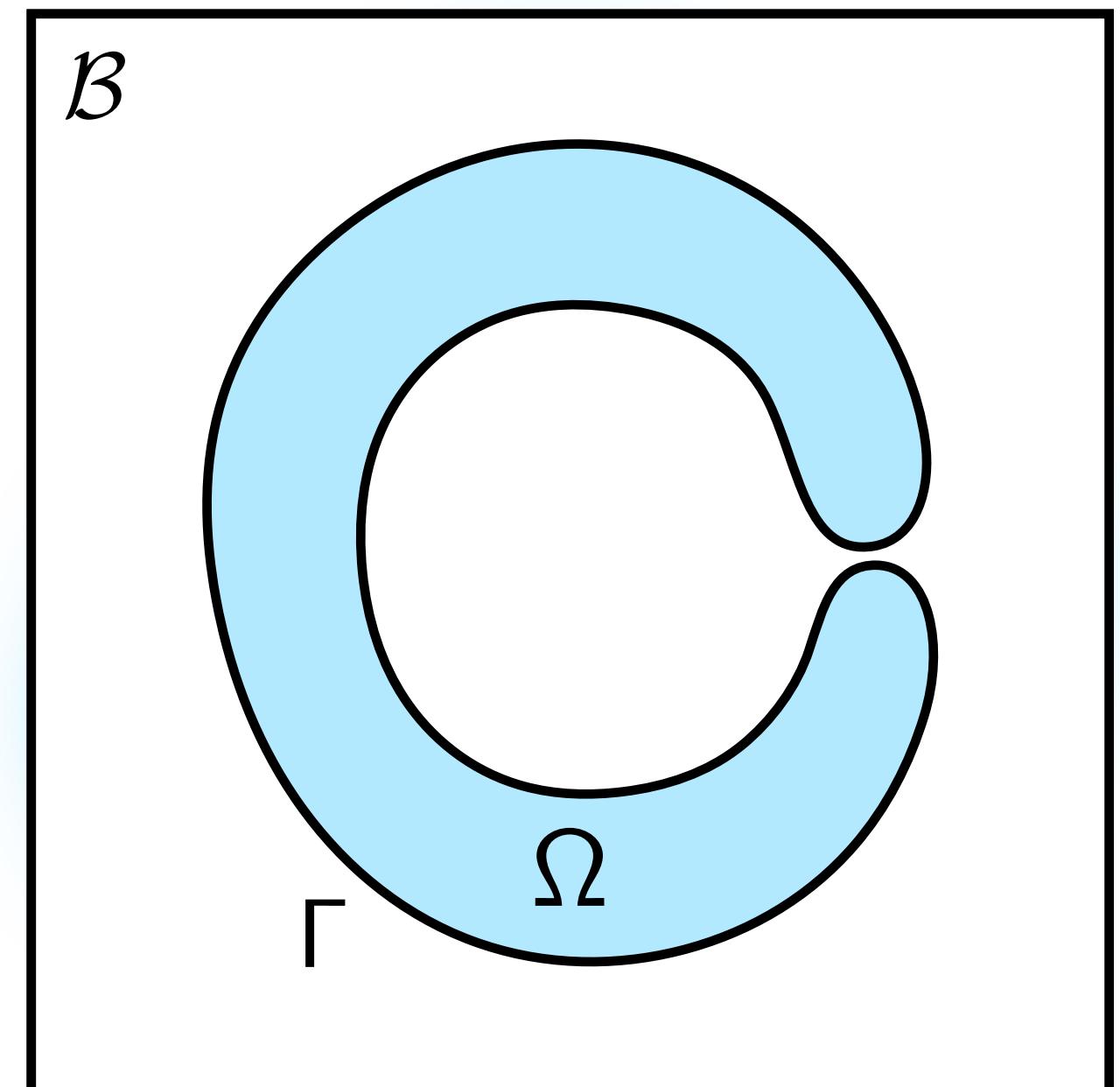
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Q: What if we truncate instead of extend?

$\tilde{f}(x) \dots ?$



Potential theory

Truncating a volume potential

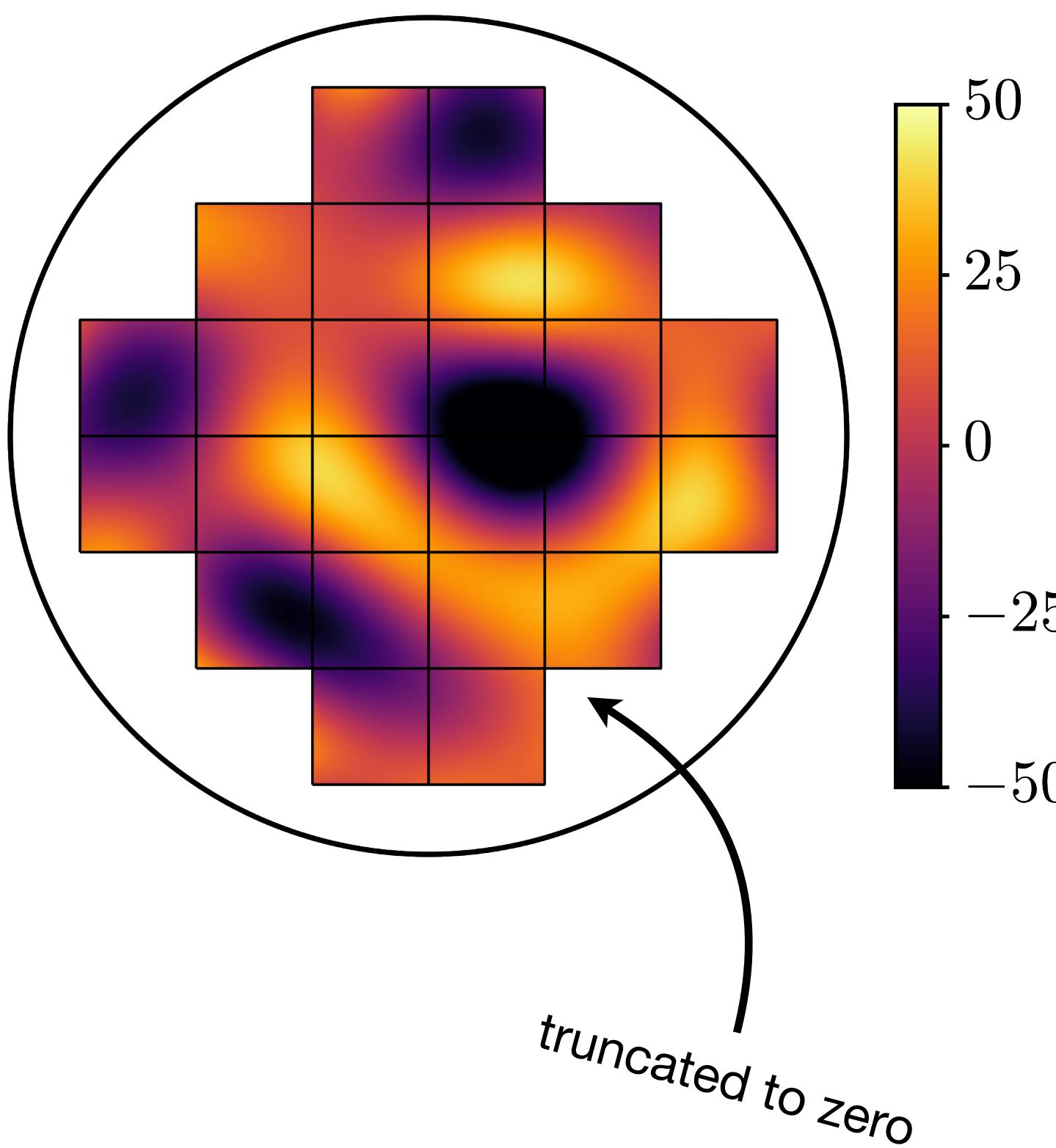
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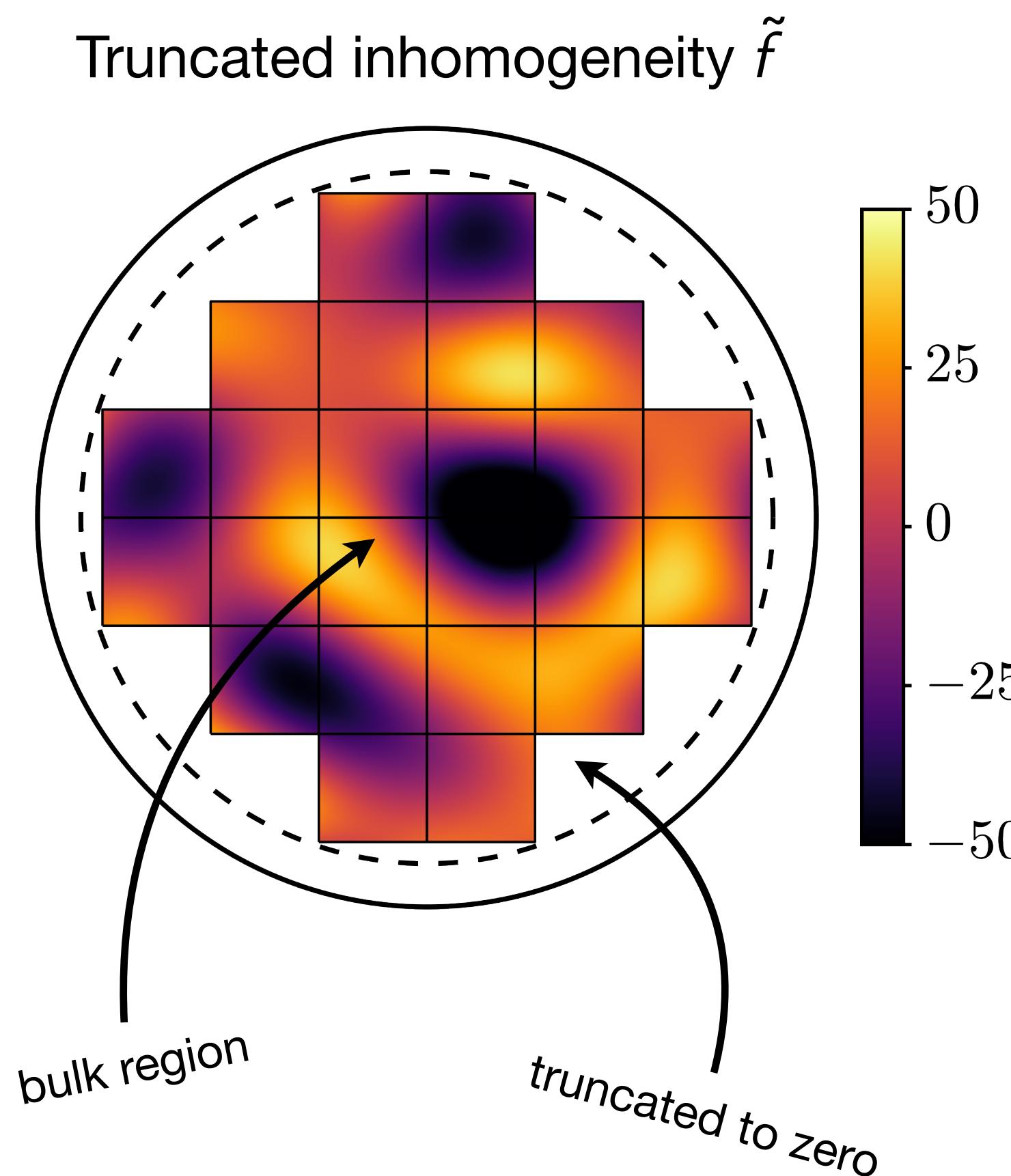
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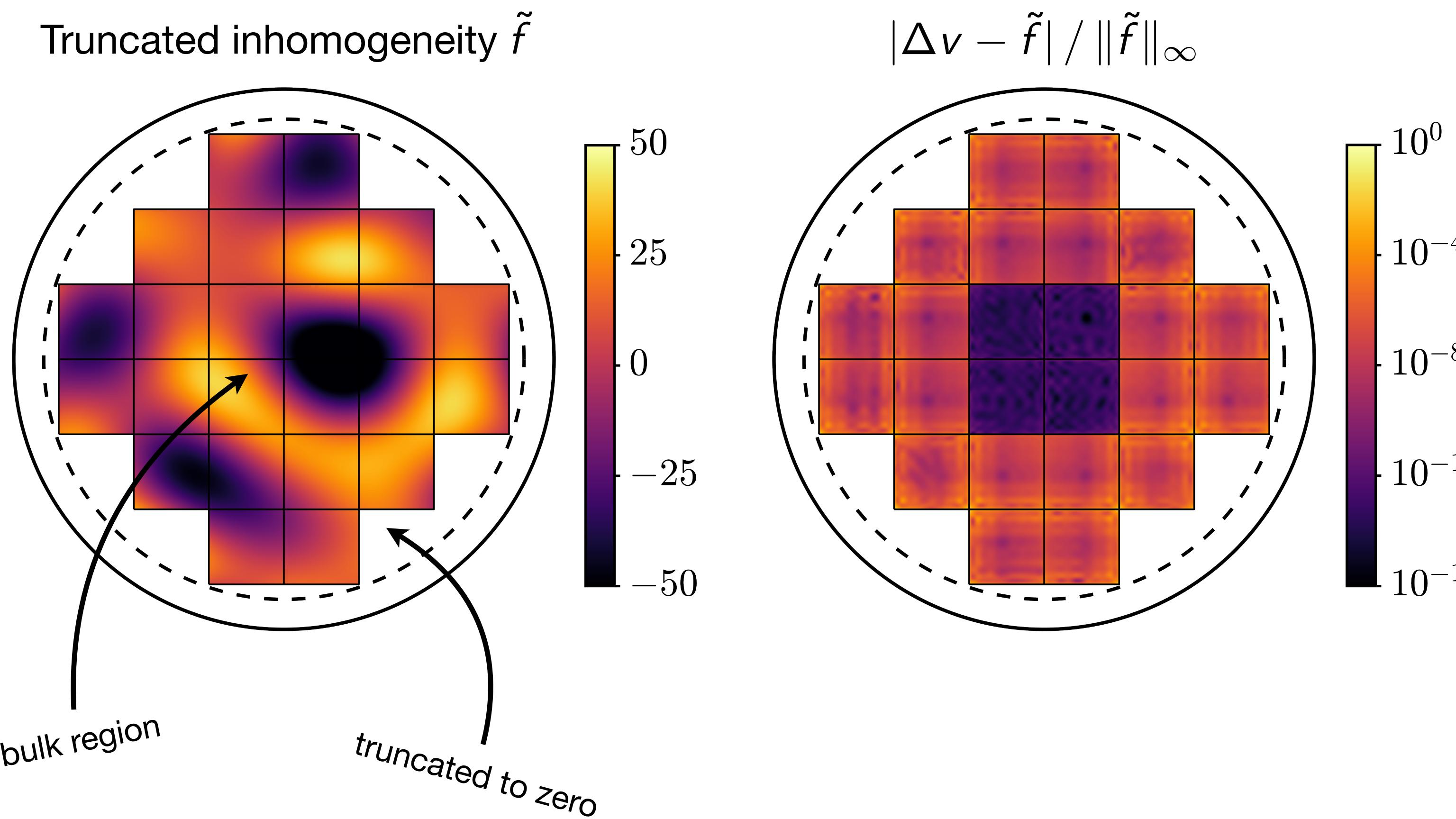
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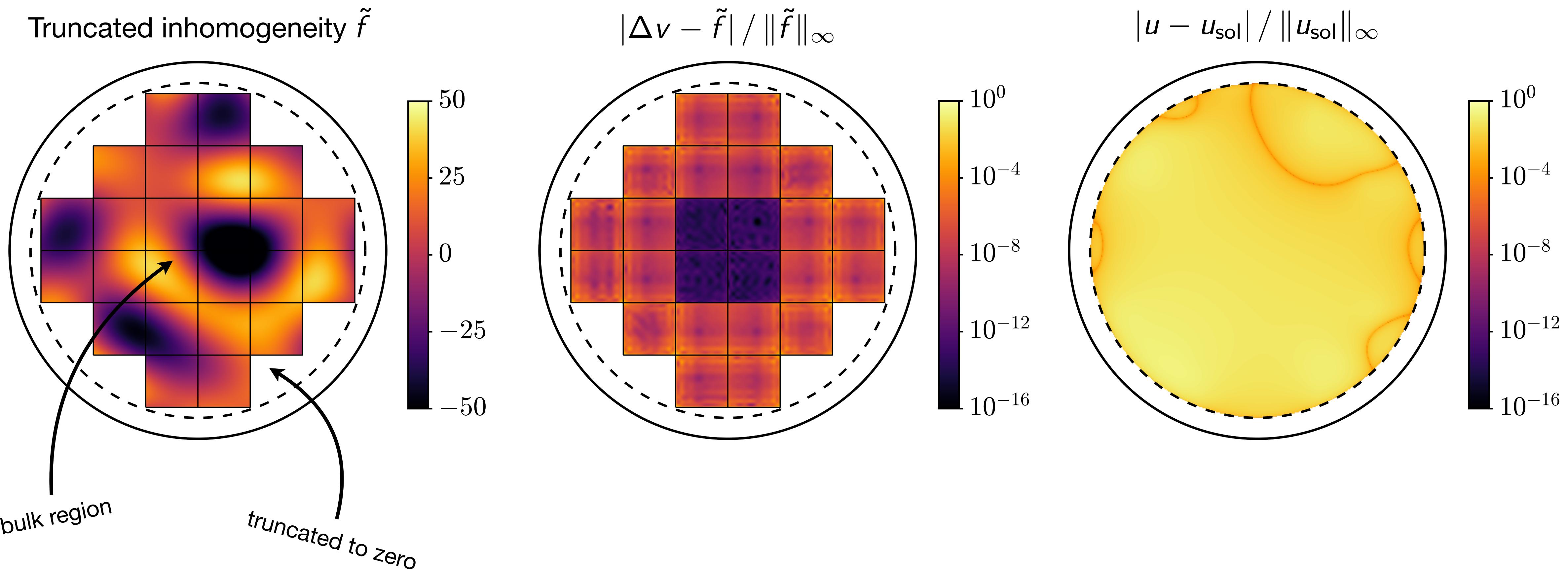
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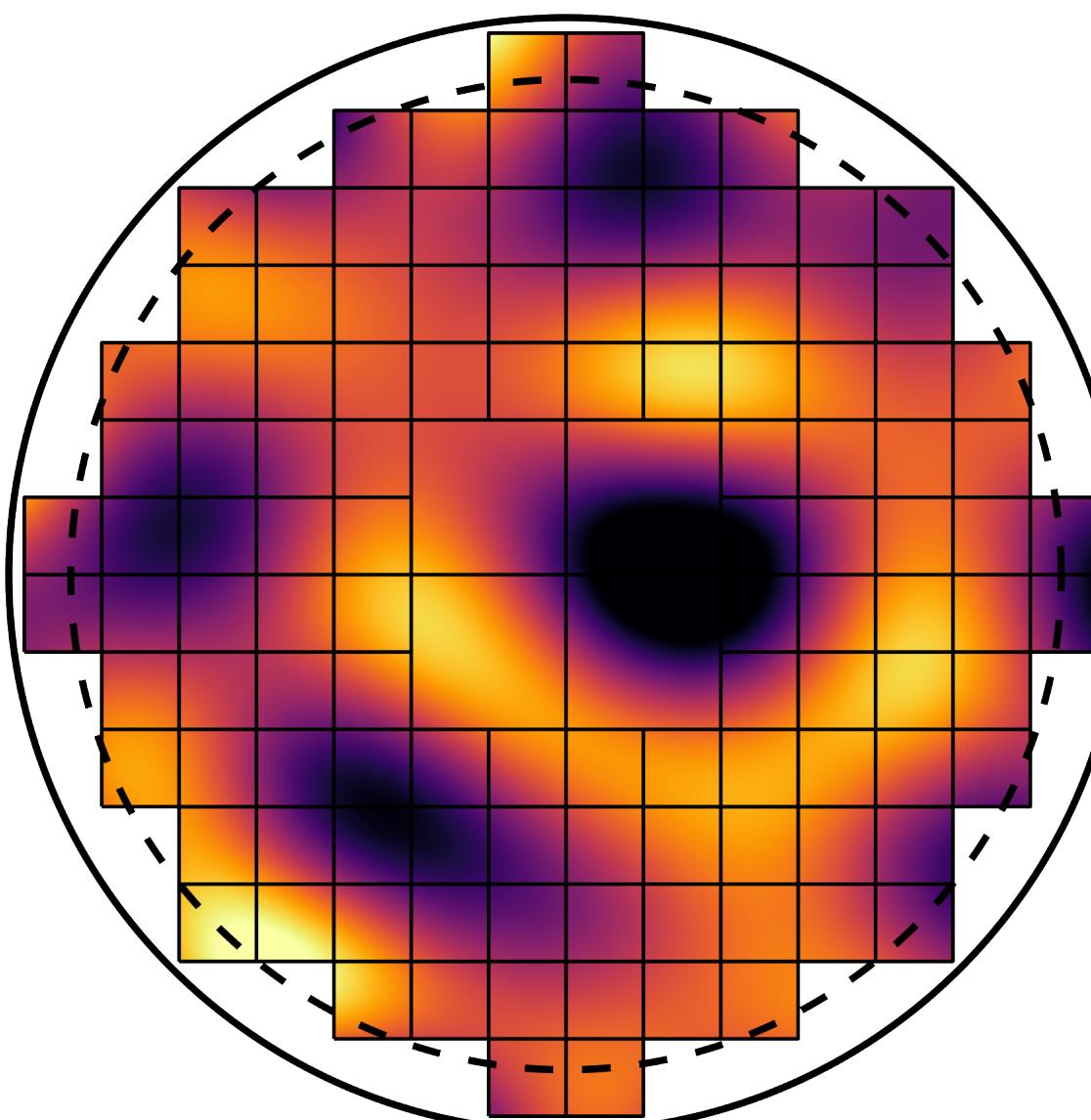


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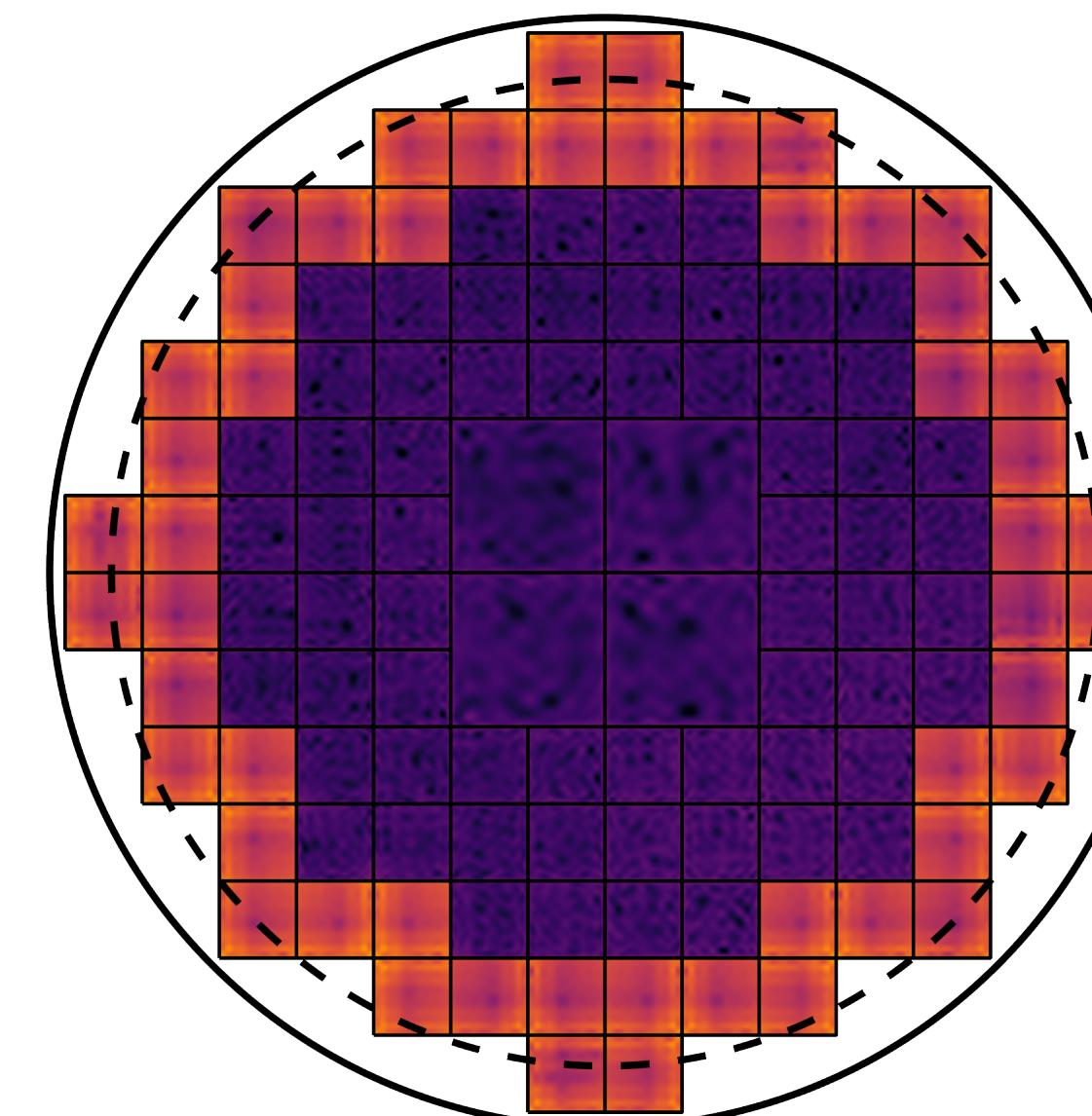
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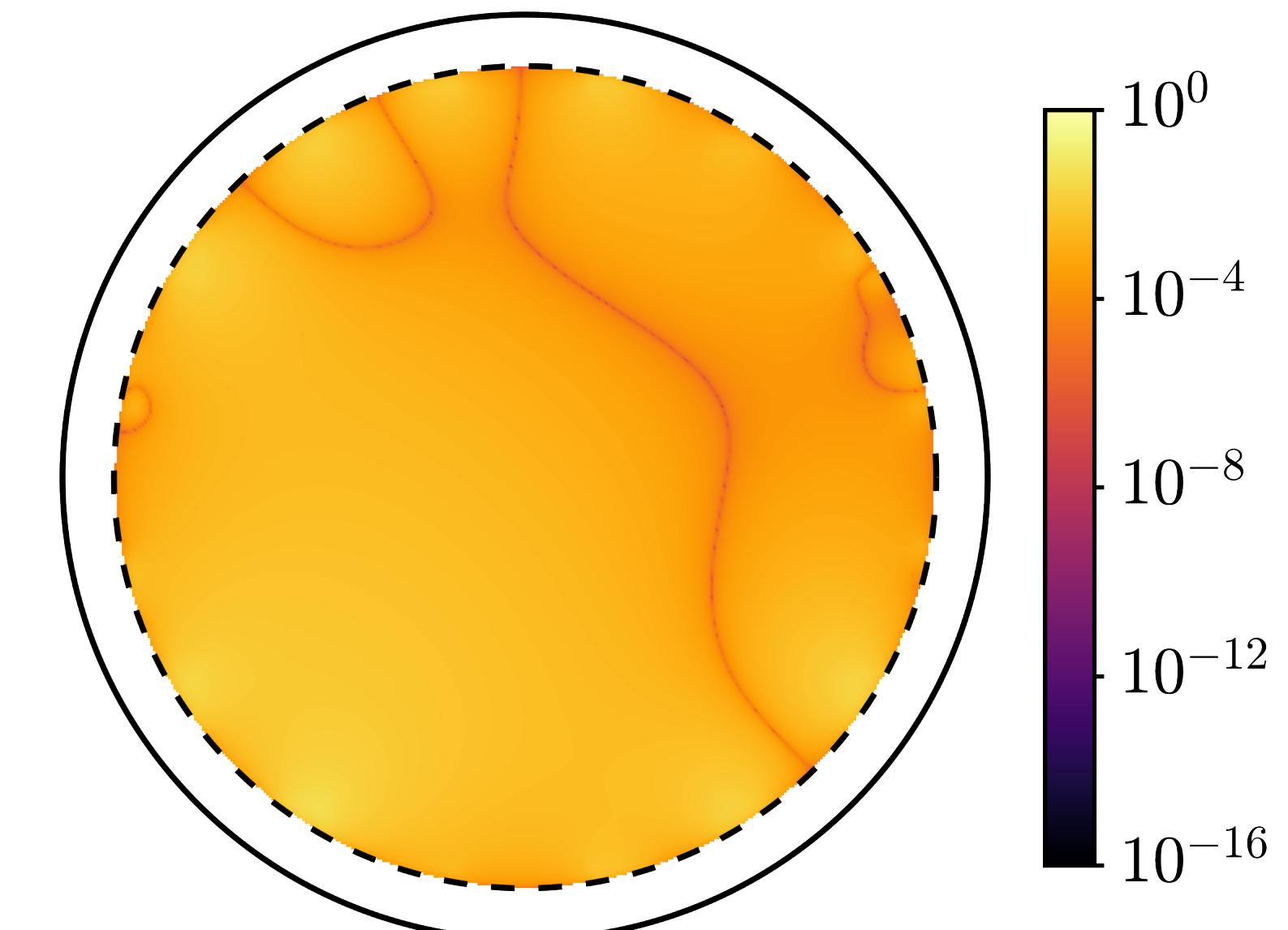
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$|\Delta v - \tilde{f}| / \|\tilde{f}\|_\infty$



$|u - u_{\text{sol}}| / \|u_{\text{sol}}\|_\infty$

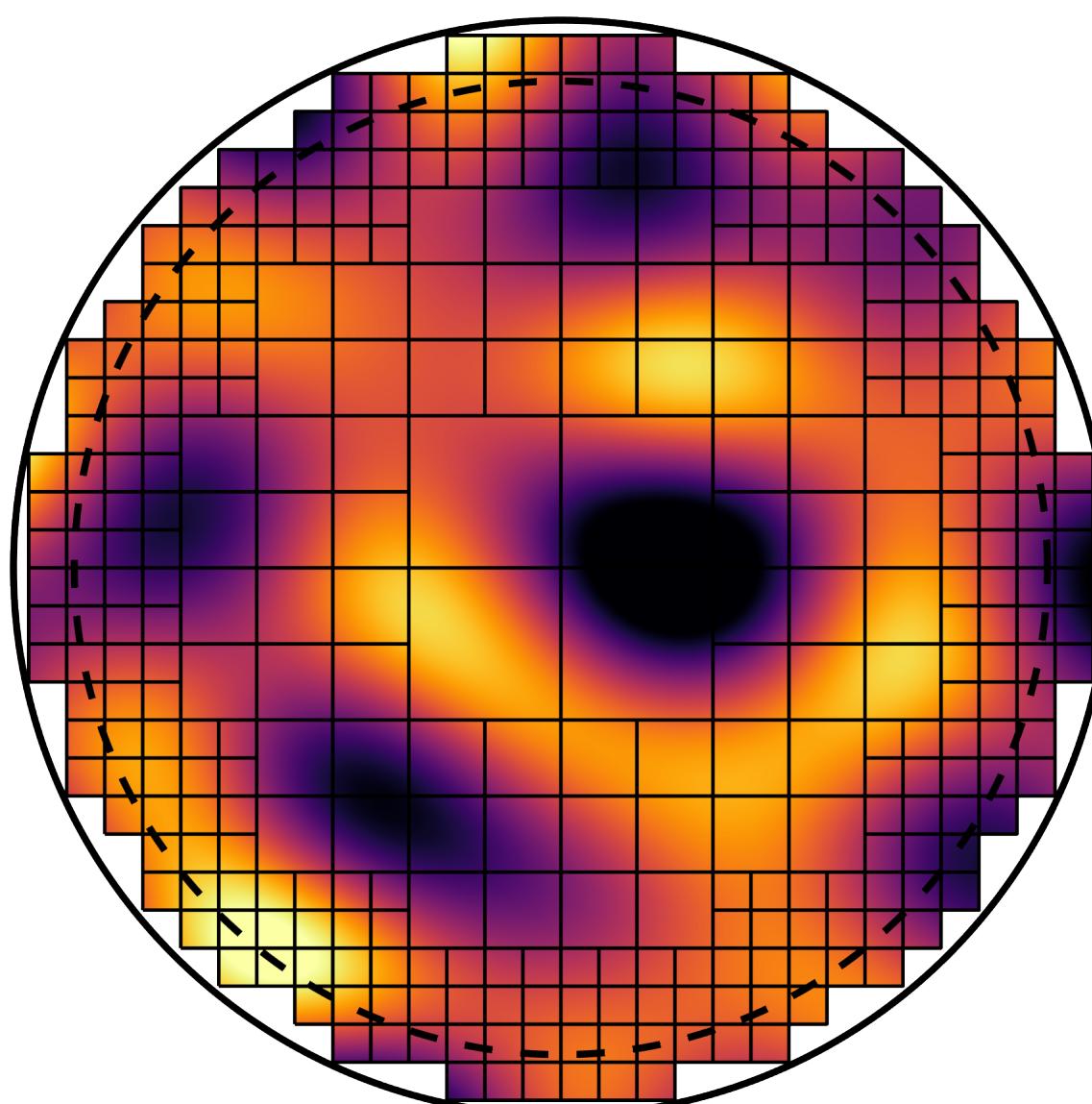


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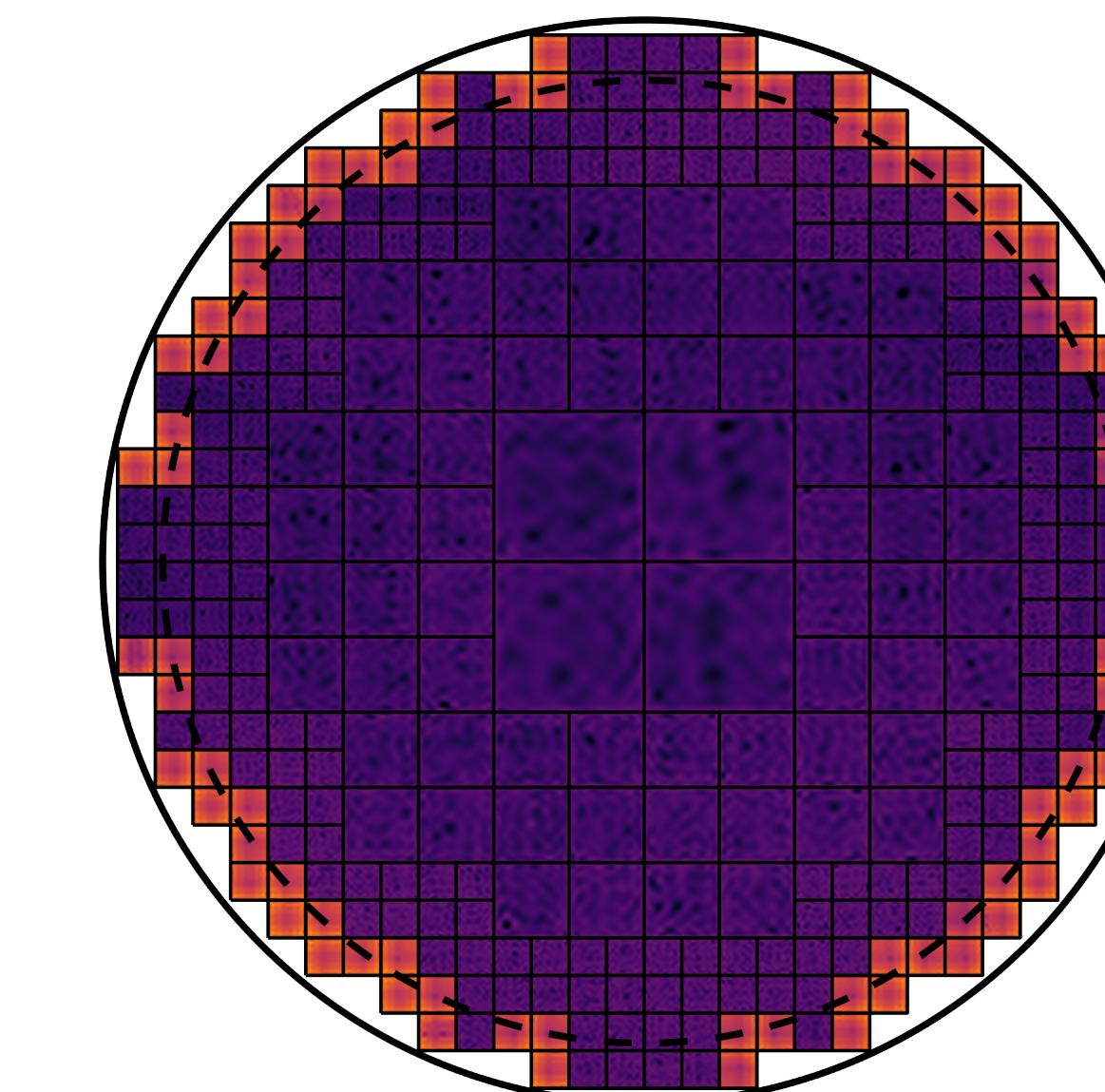
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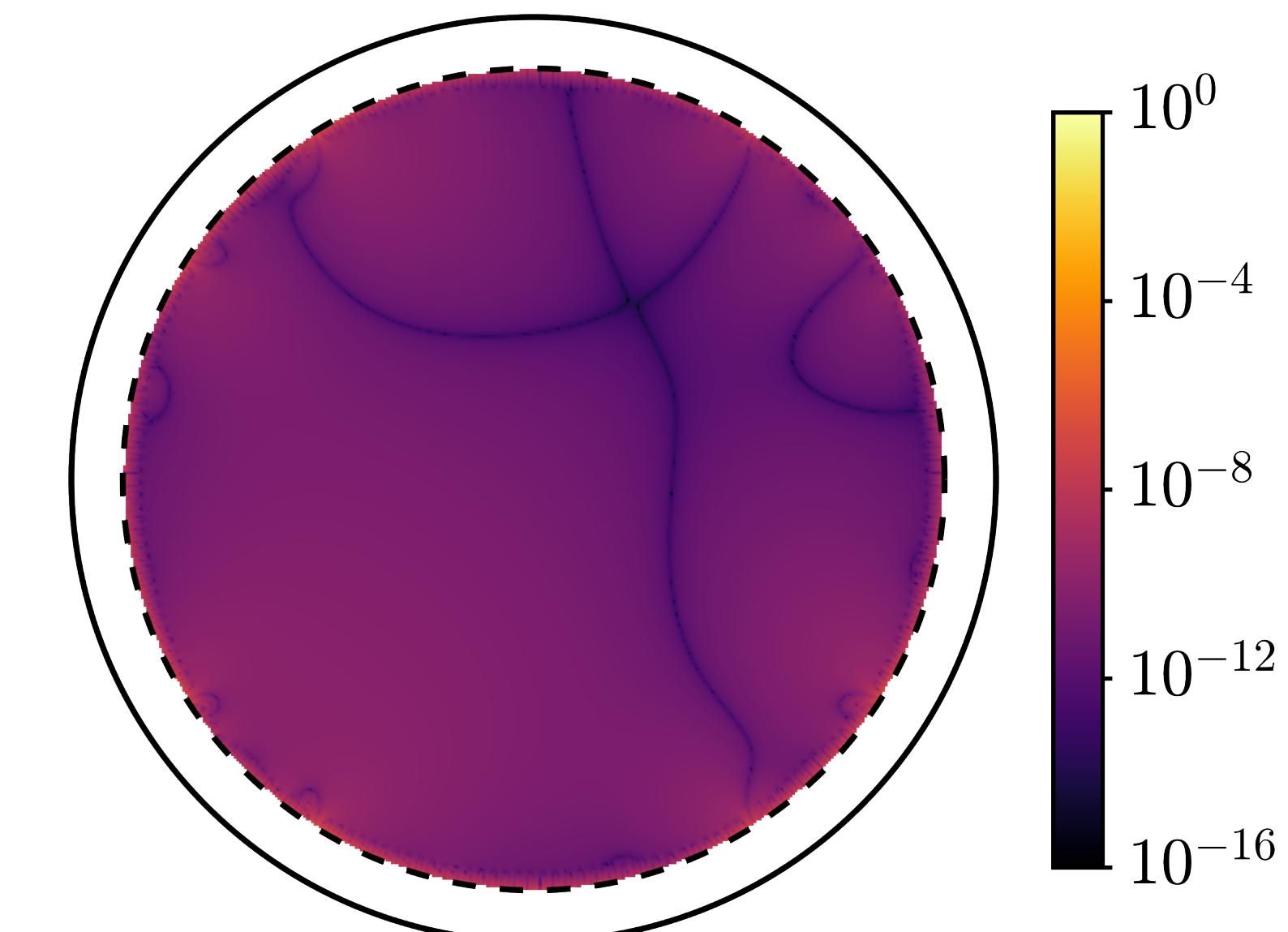
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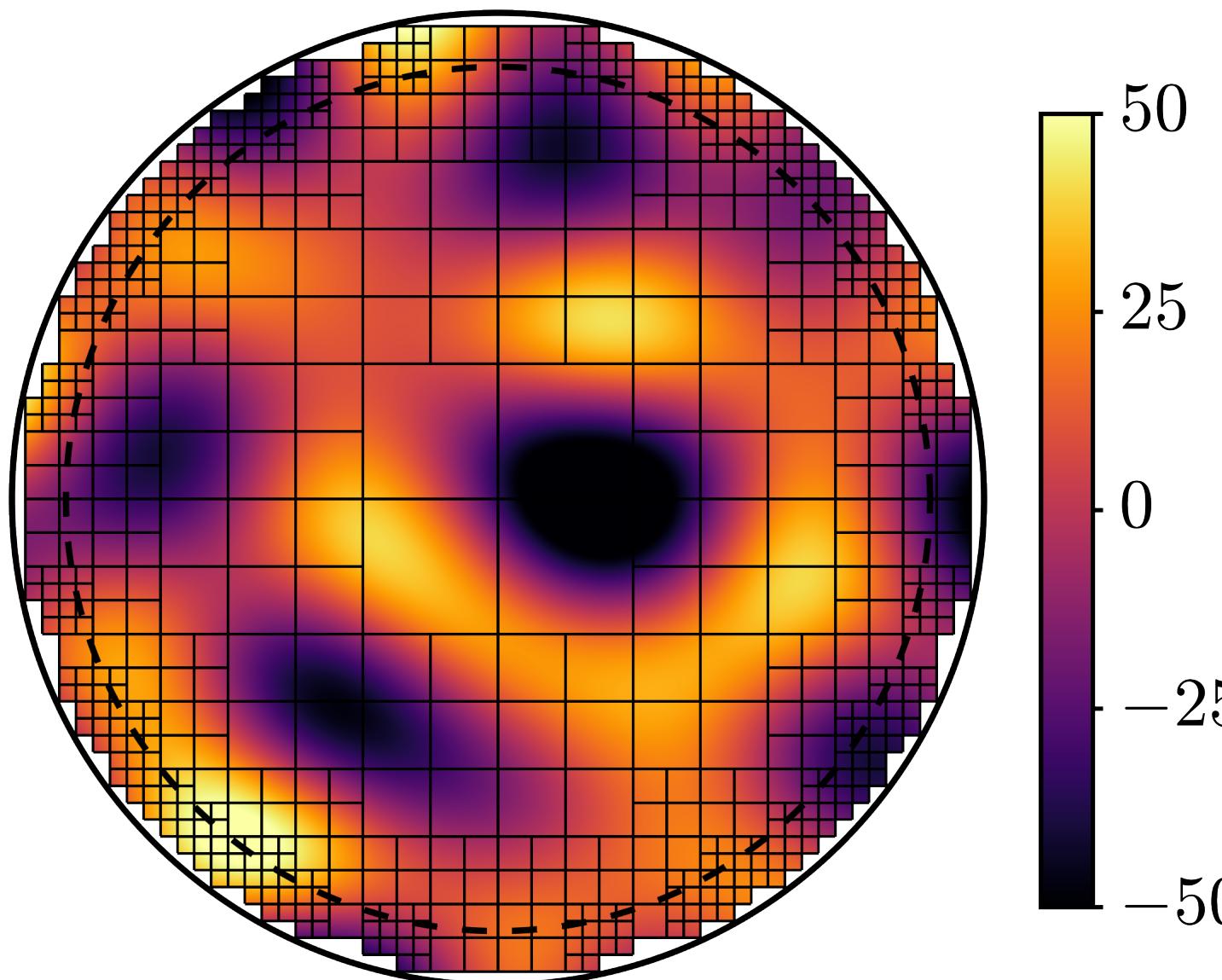


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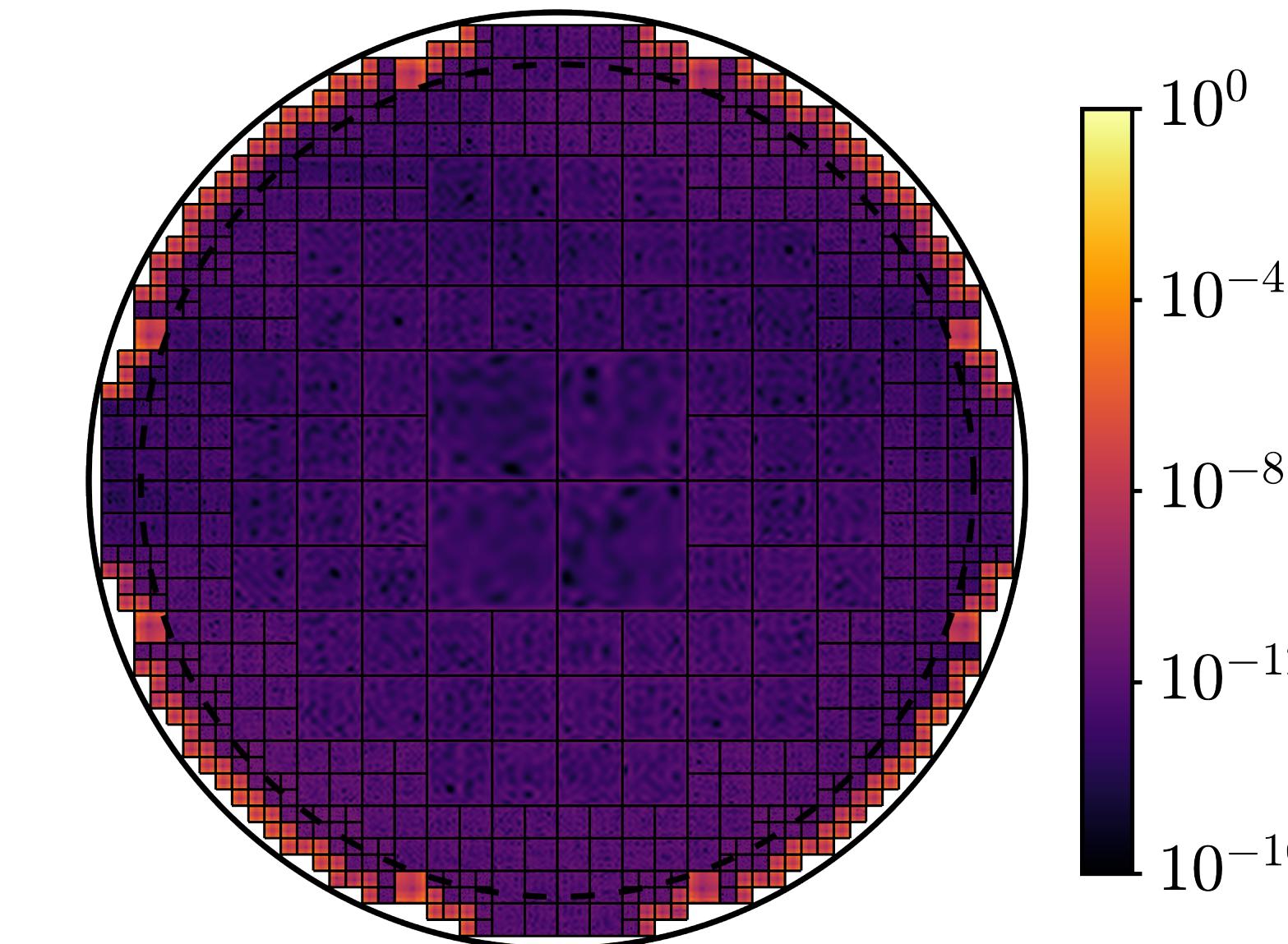
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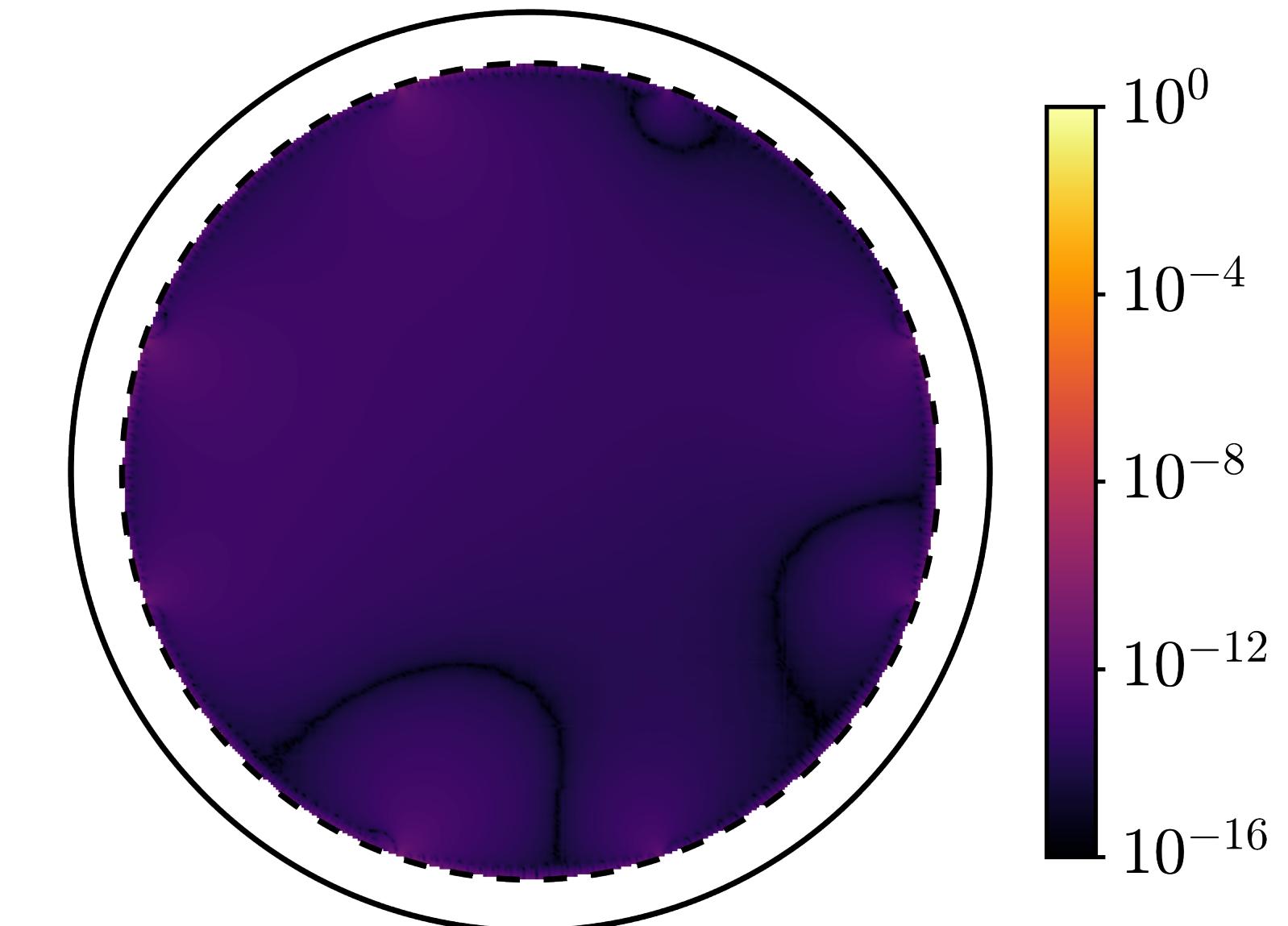
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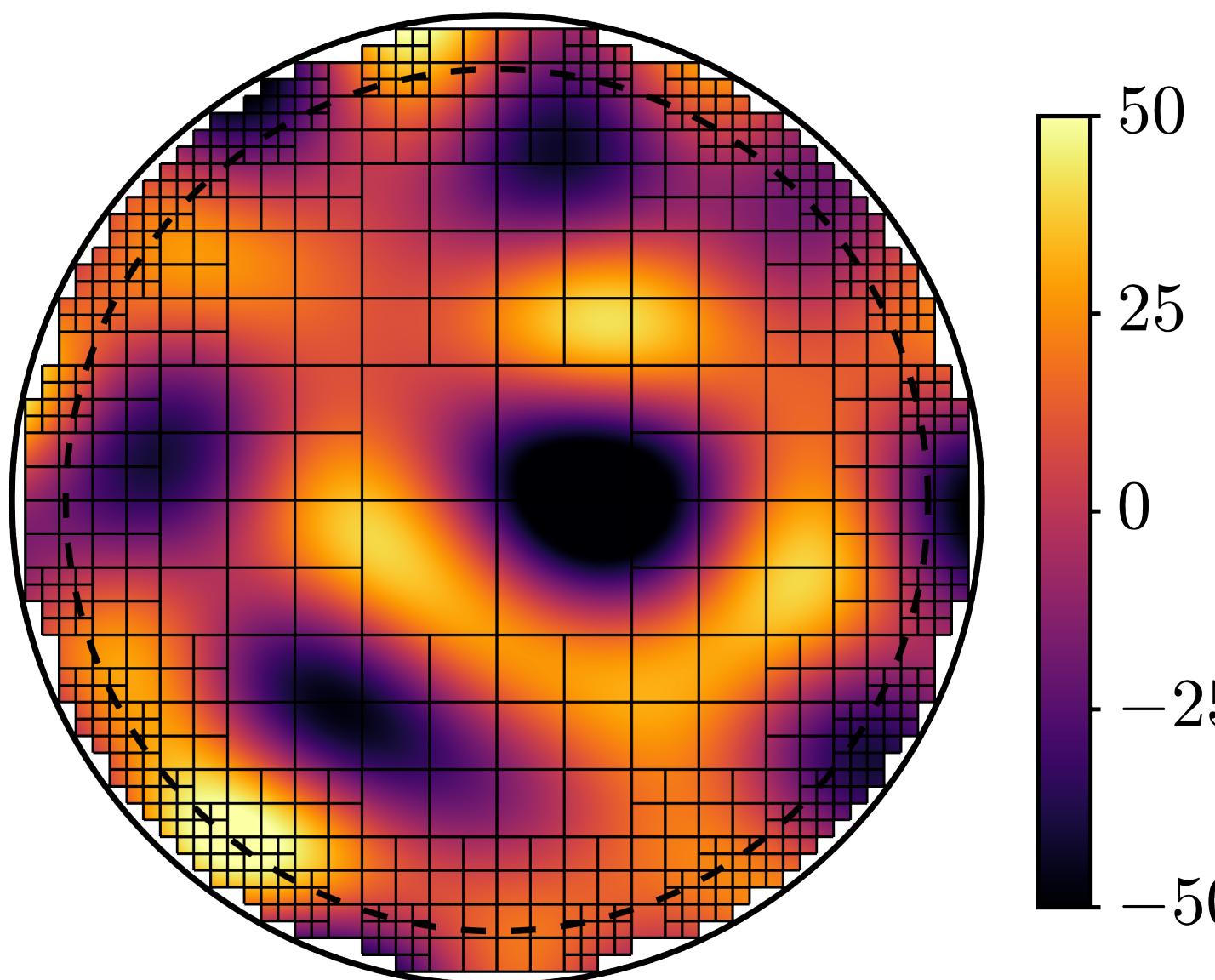


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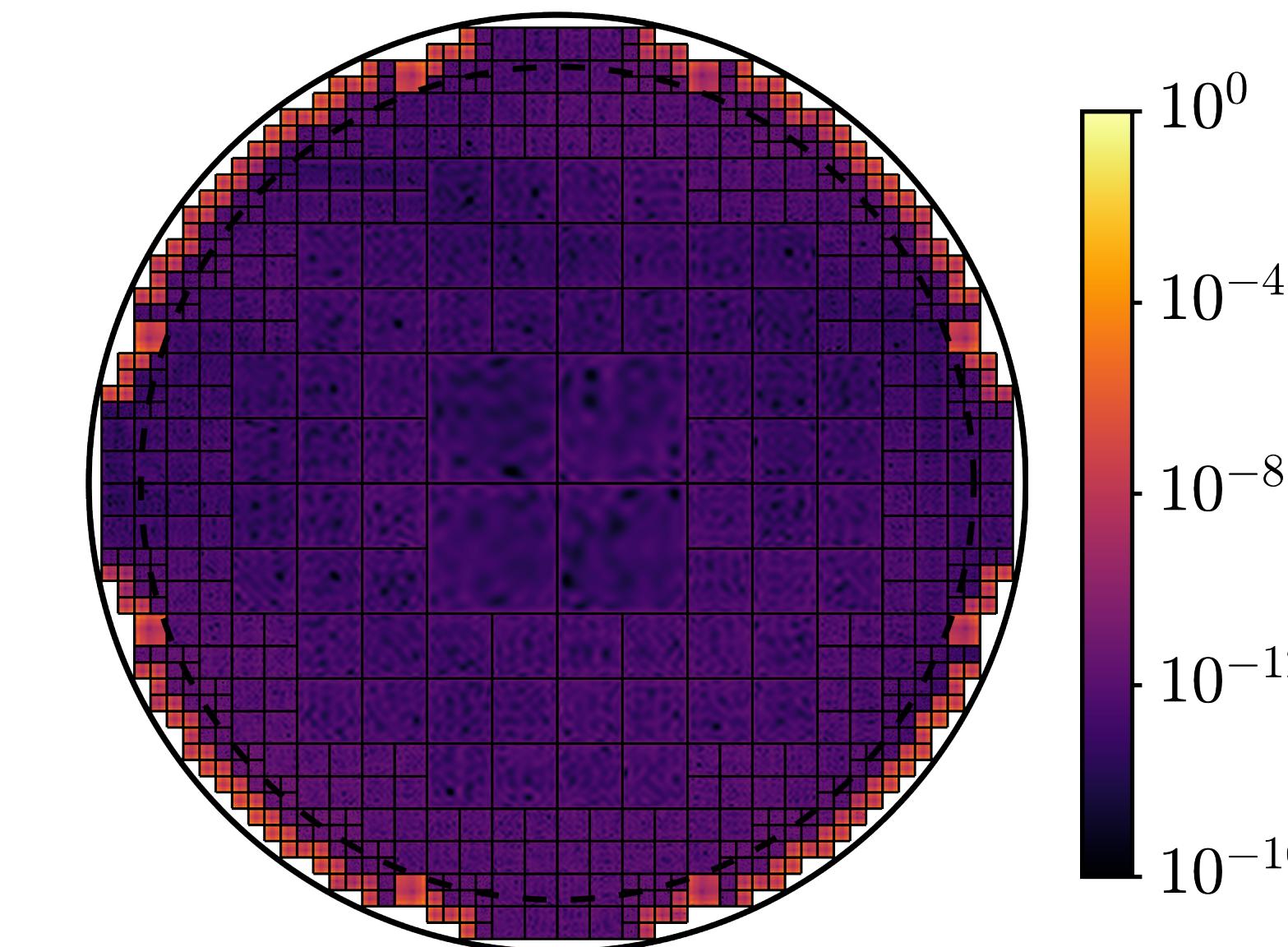
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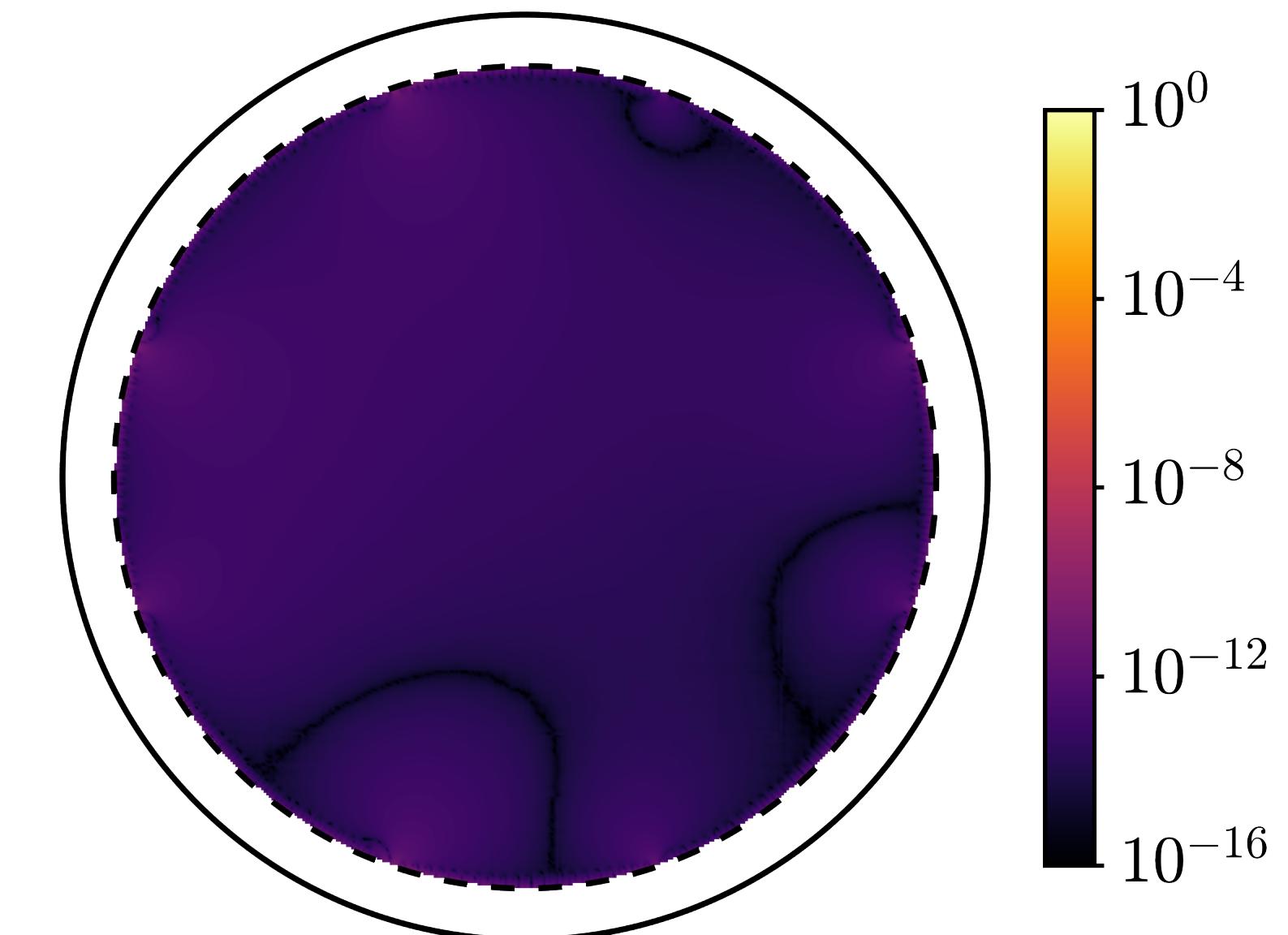
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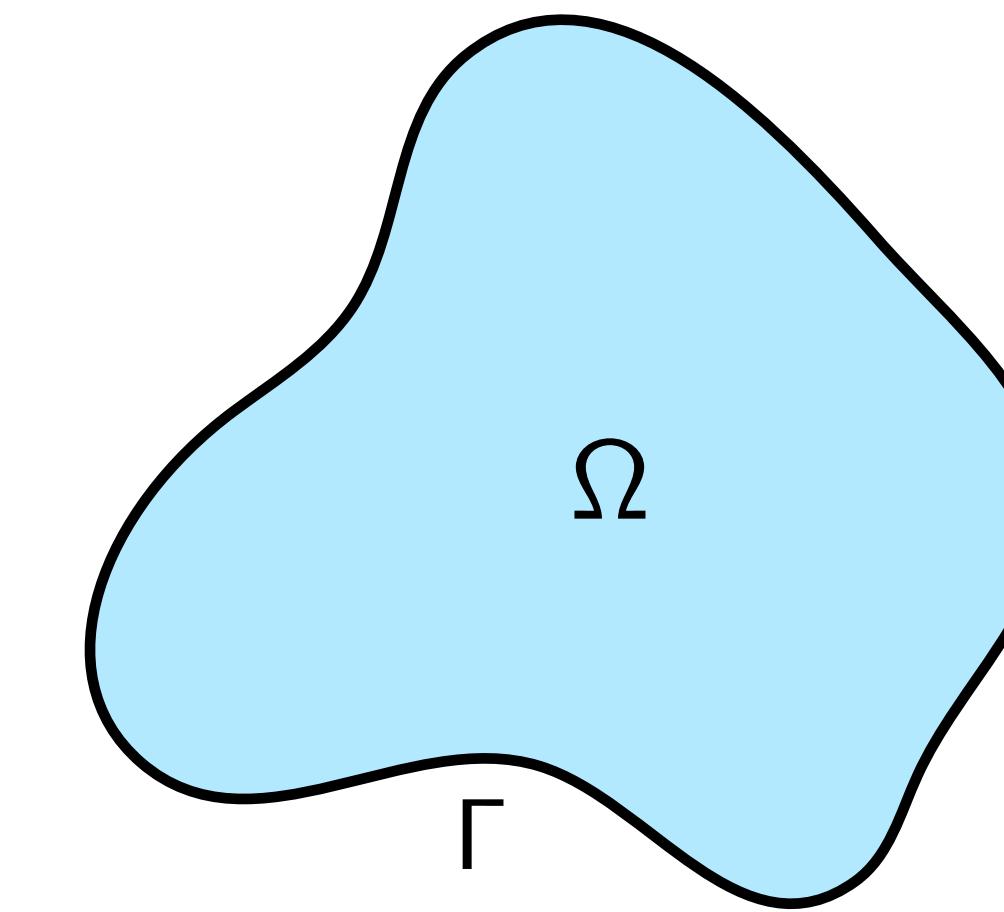


Theorem (F., Stein, & Barnett)

Let the truncated region $\text{supp}(f - \tilde{f})$ be well separated from each bulk box. Then the volume potential v induced by \tilde{f} is as smooth on each bulk box as that induced by the original f .

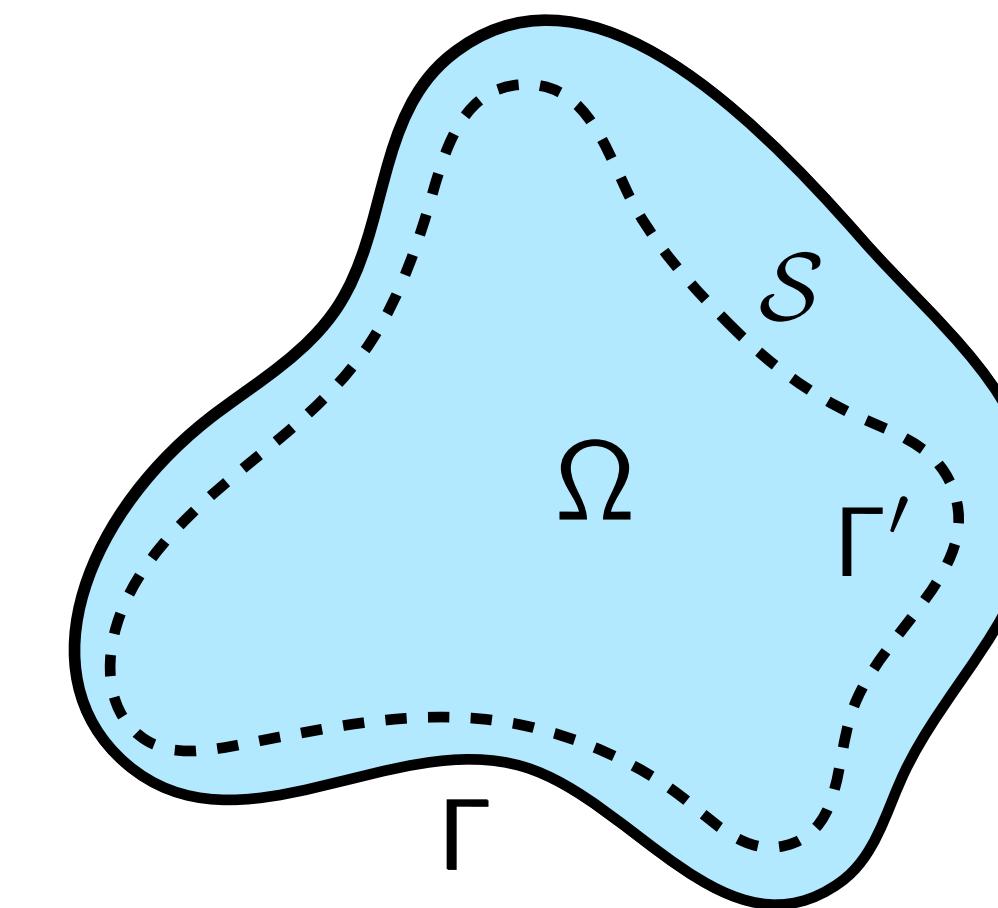
Inhomogeneous PDEs

Our approach



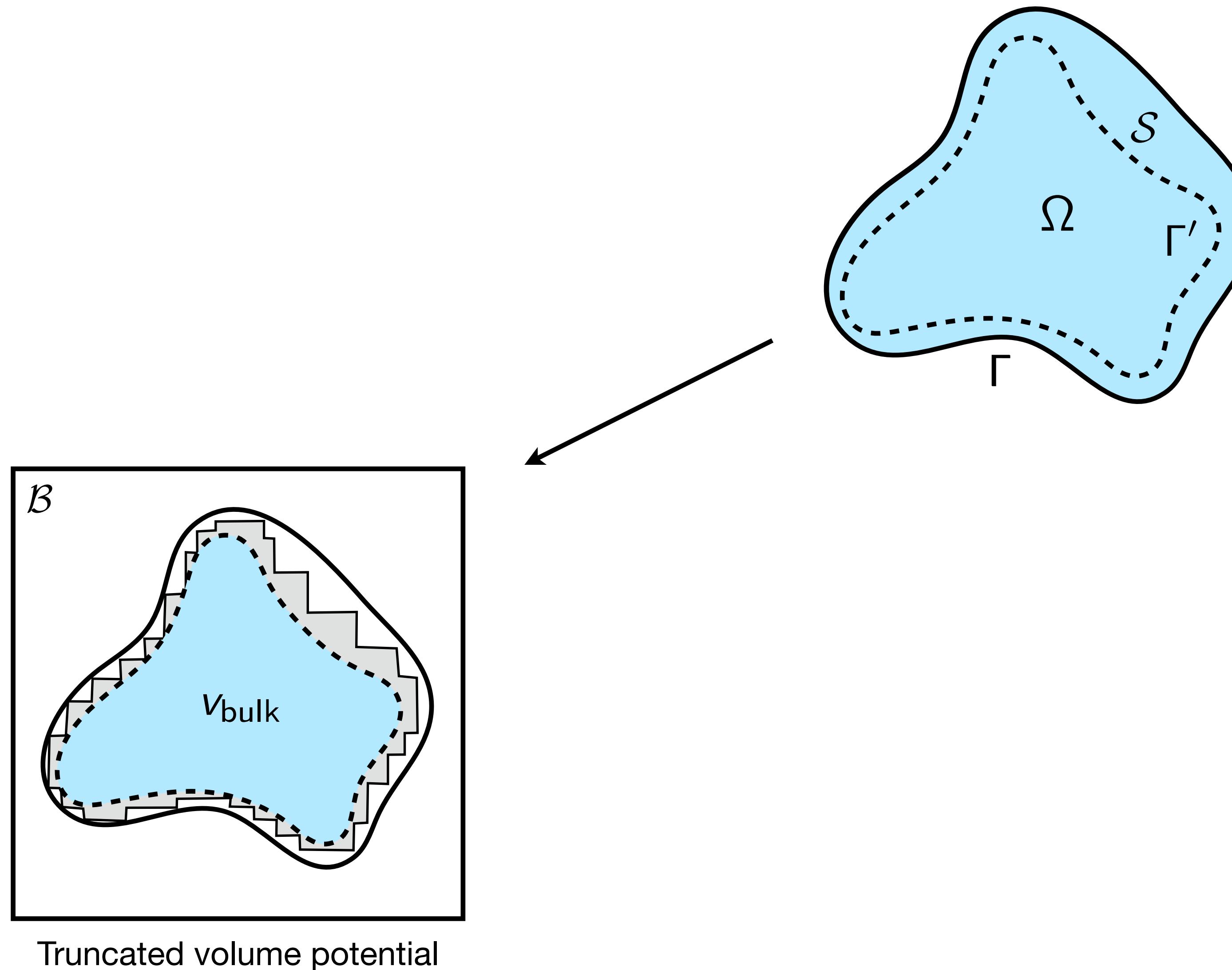
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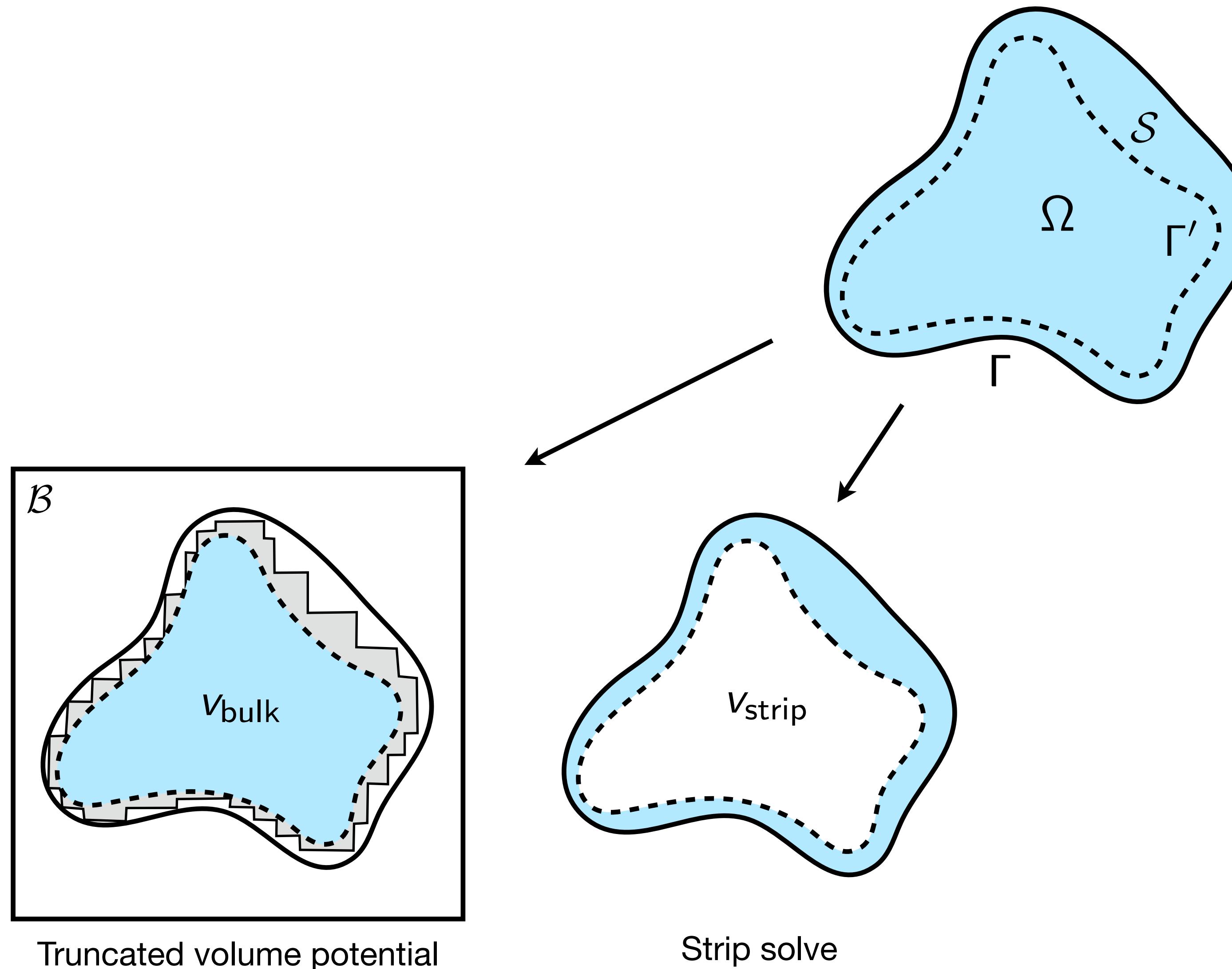
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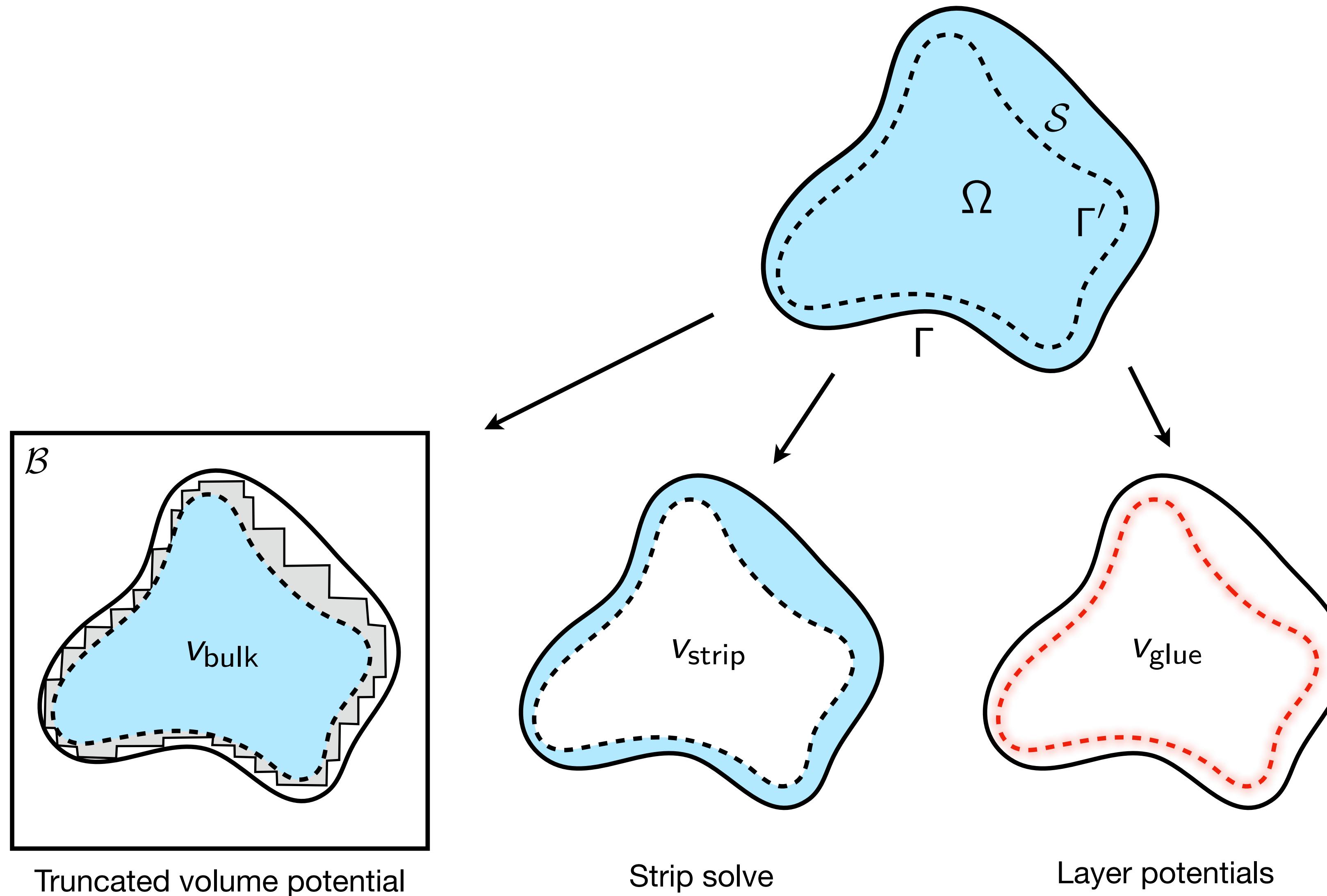
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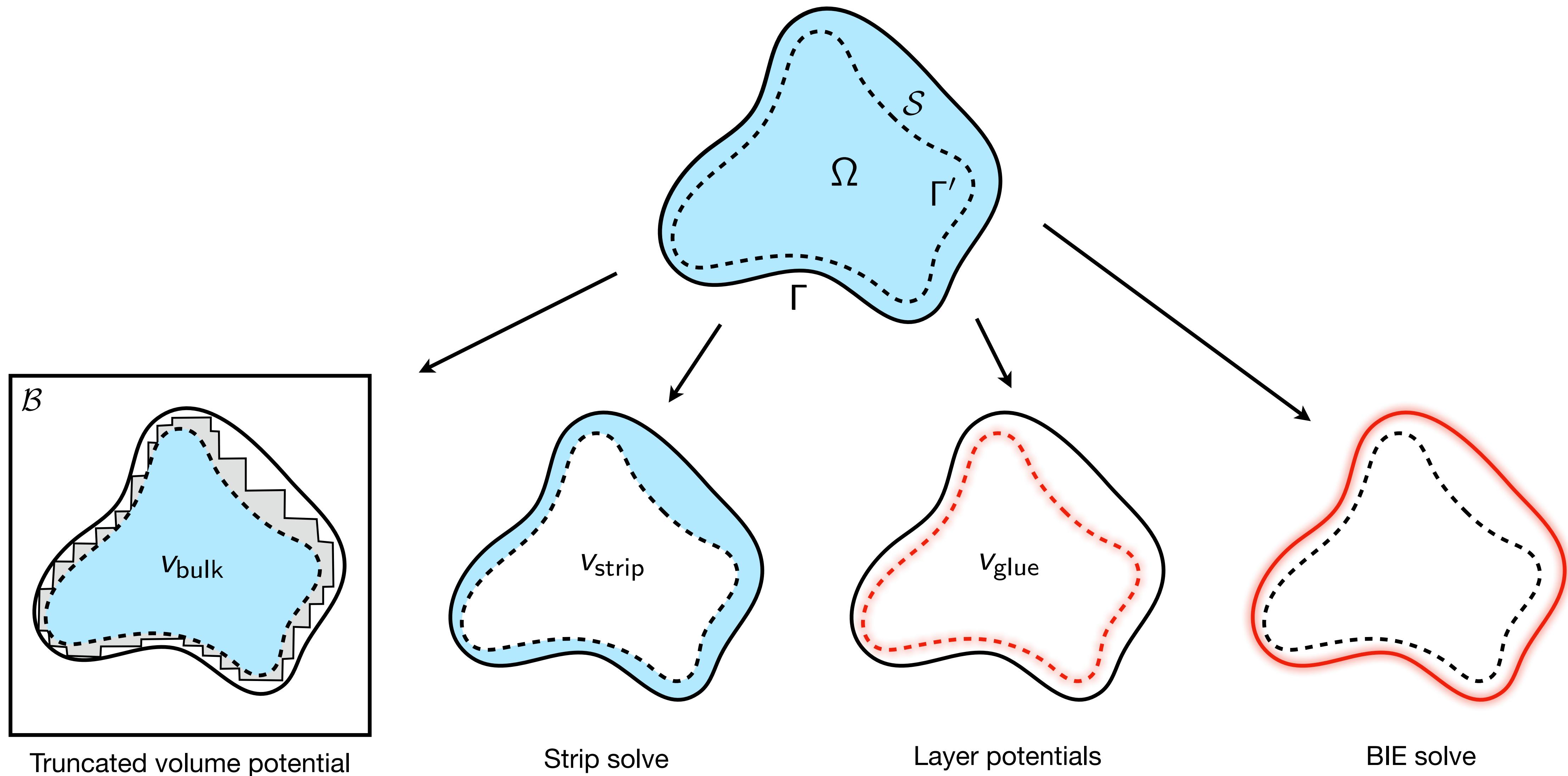
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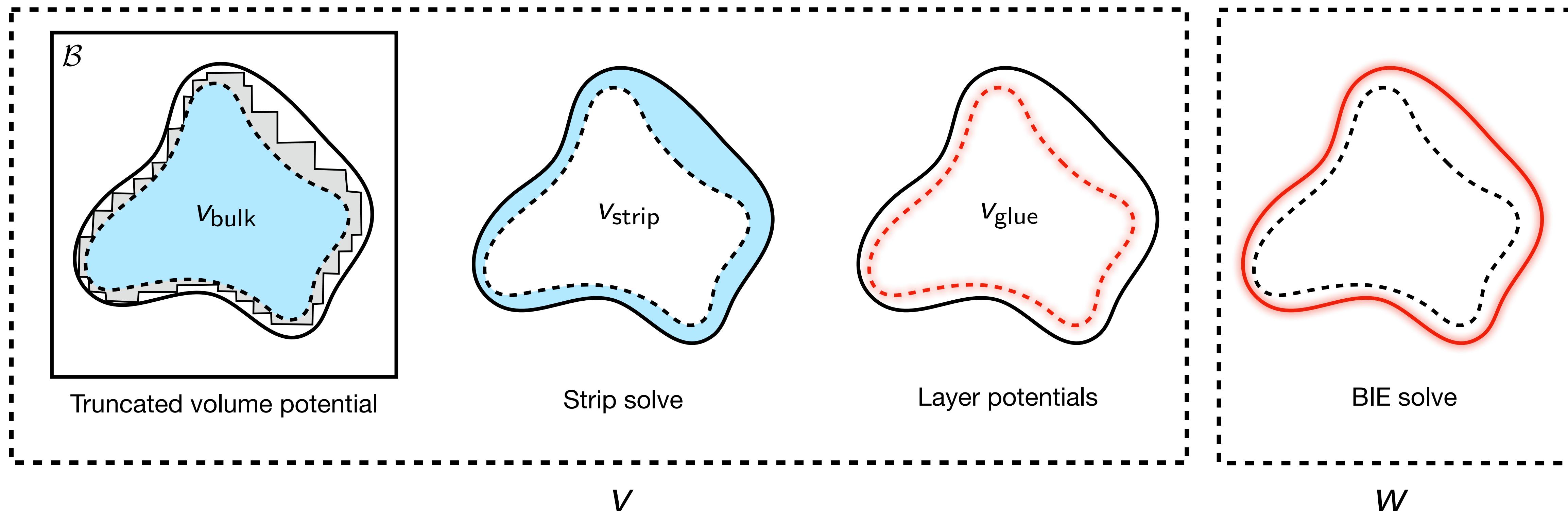
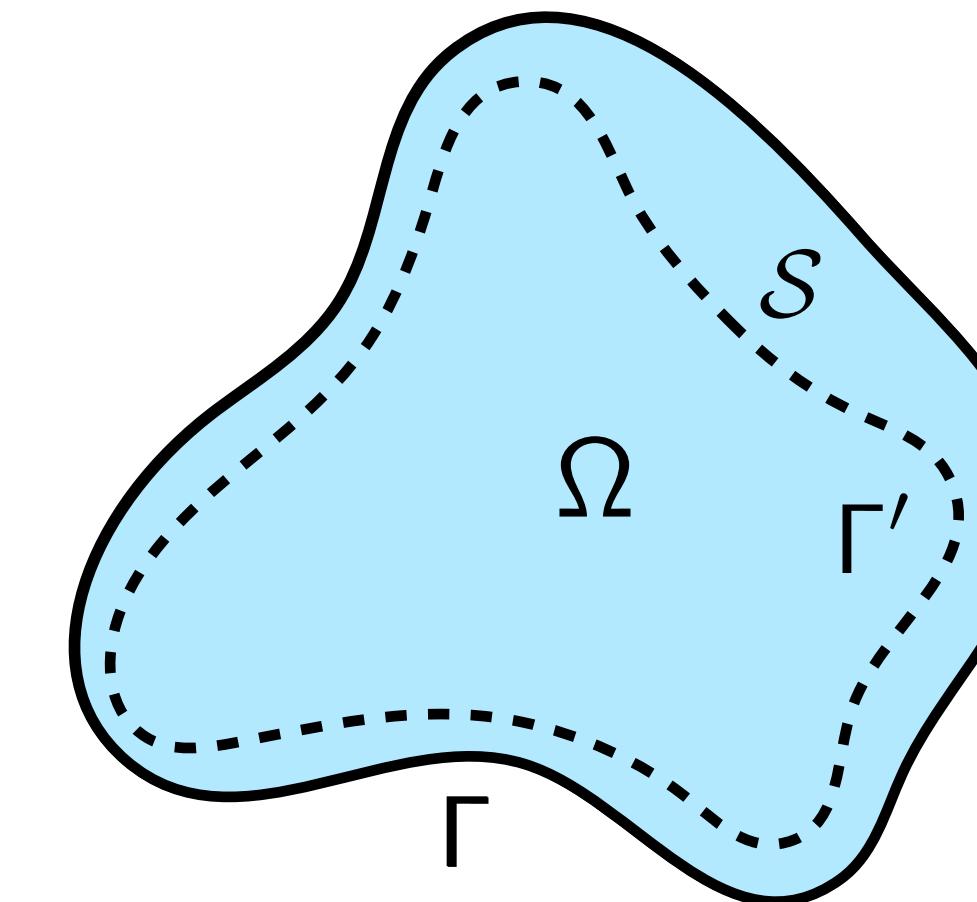
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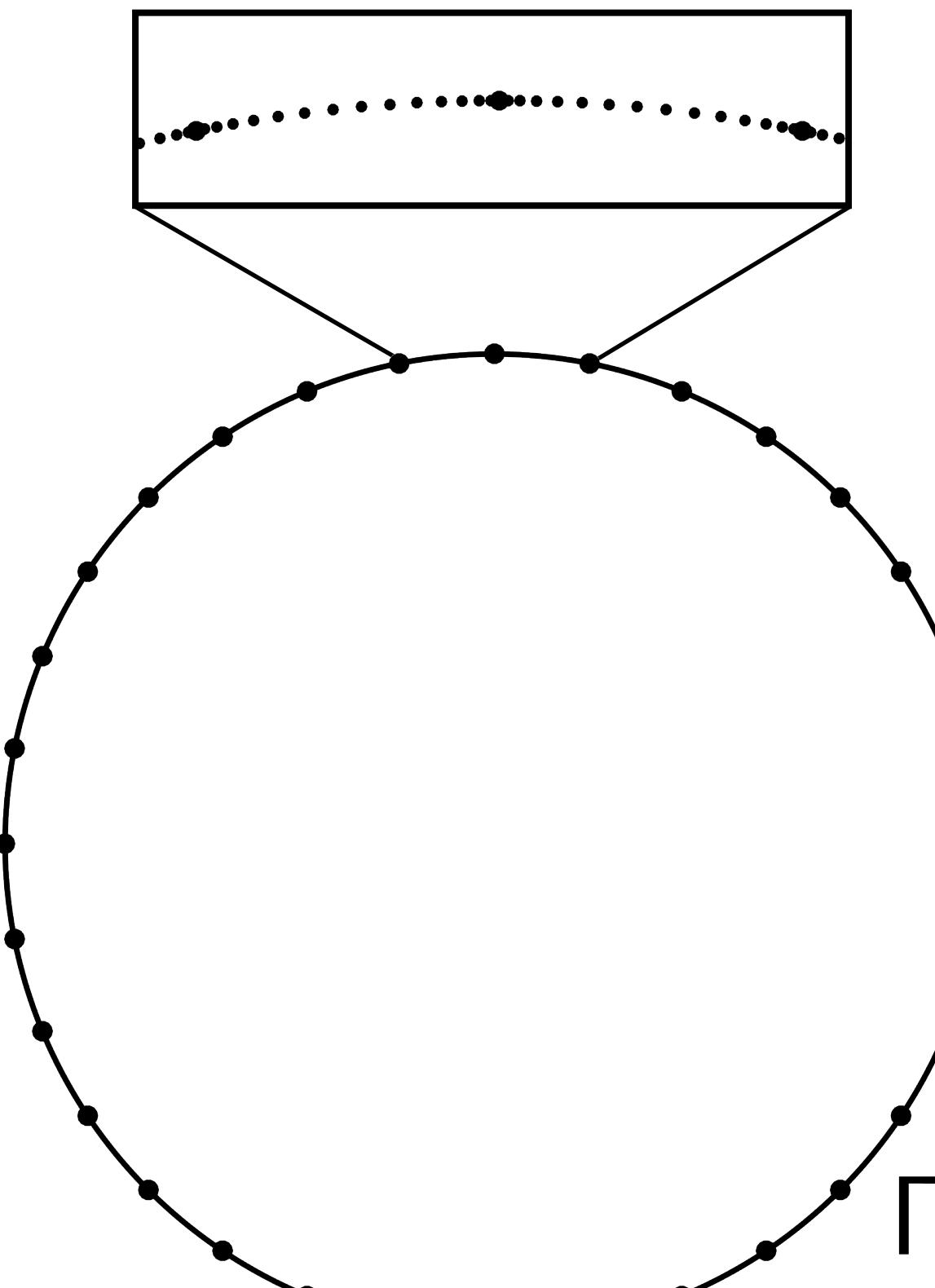
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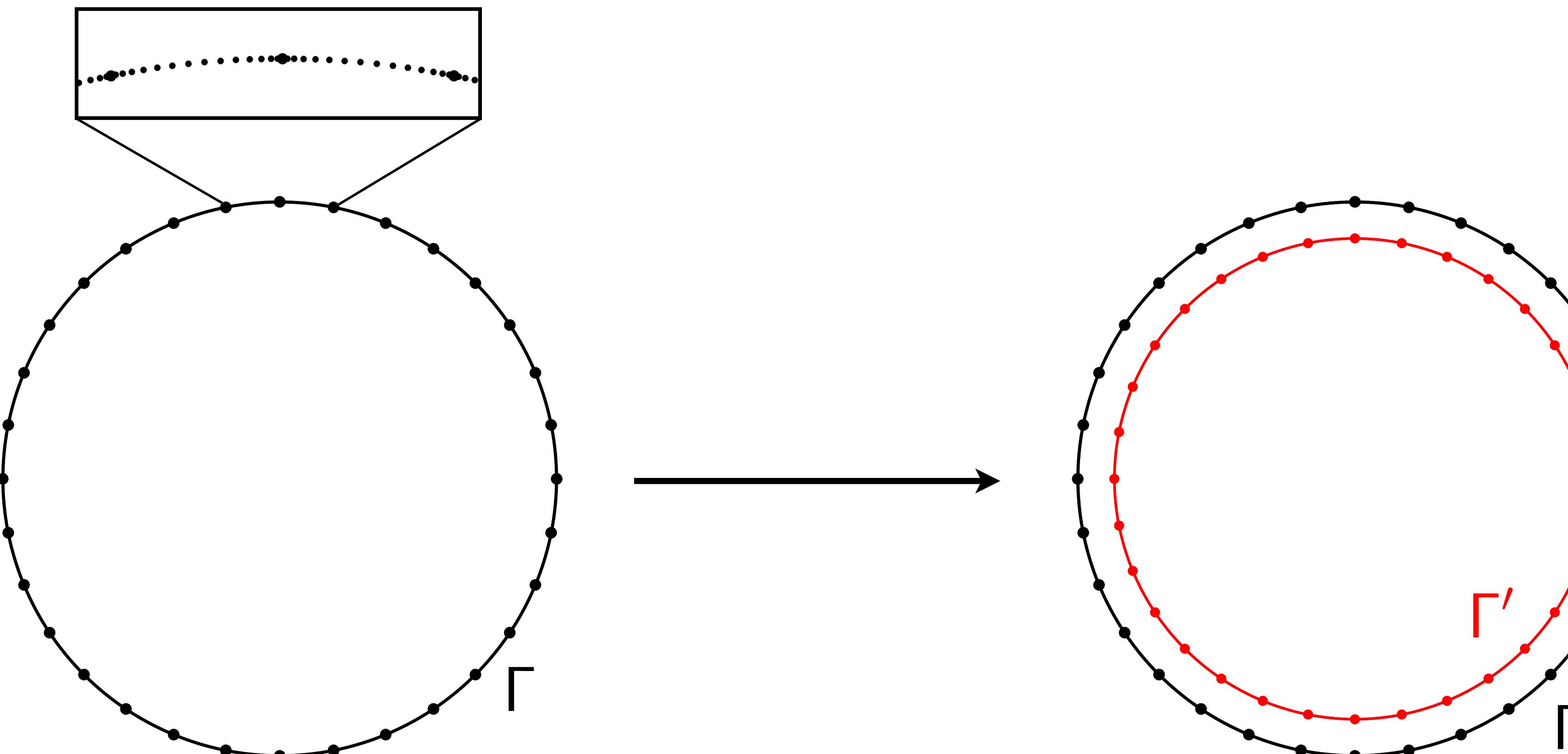
Defining the truncation region

What happens on an adaptive geometry?



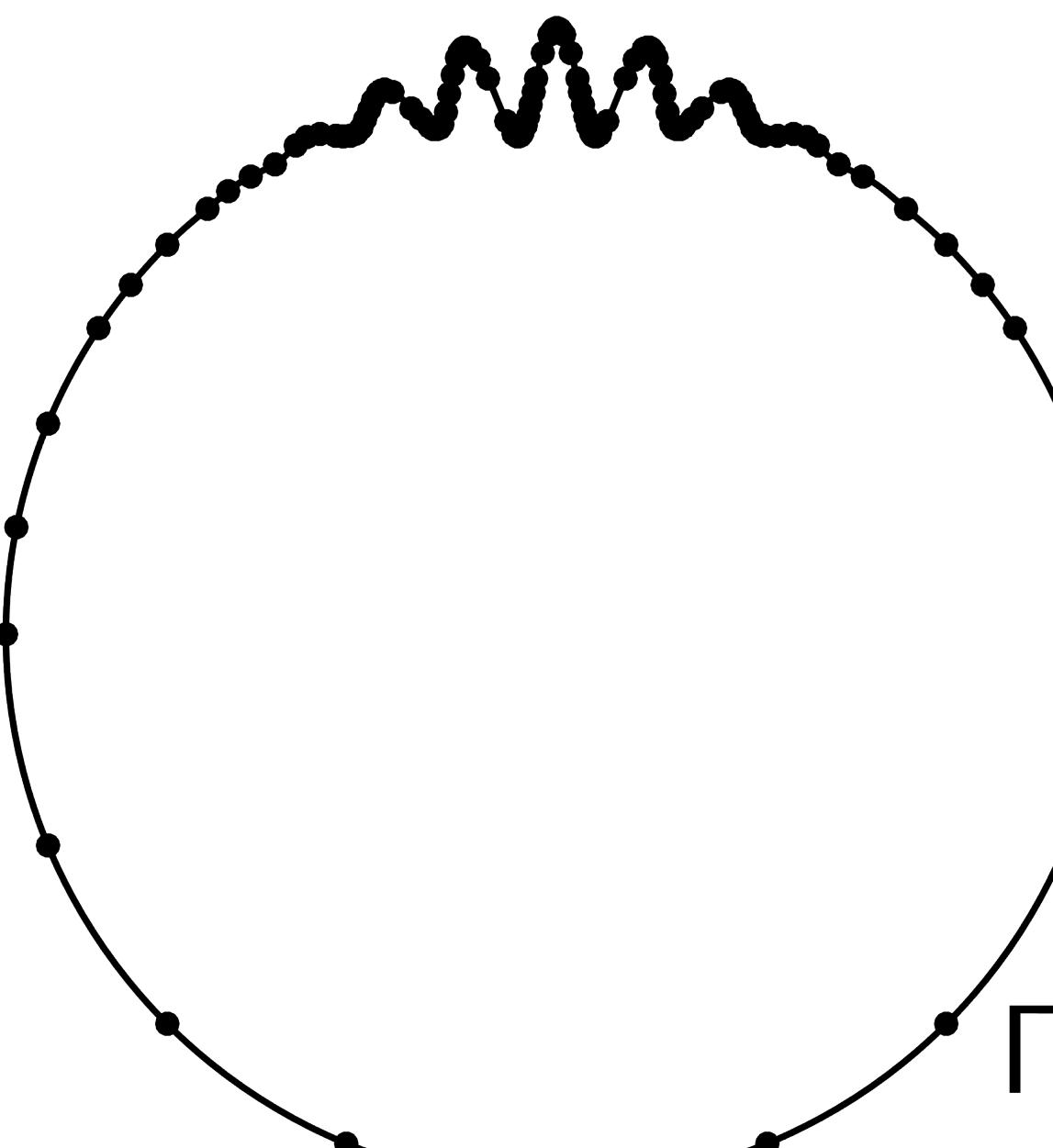
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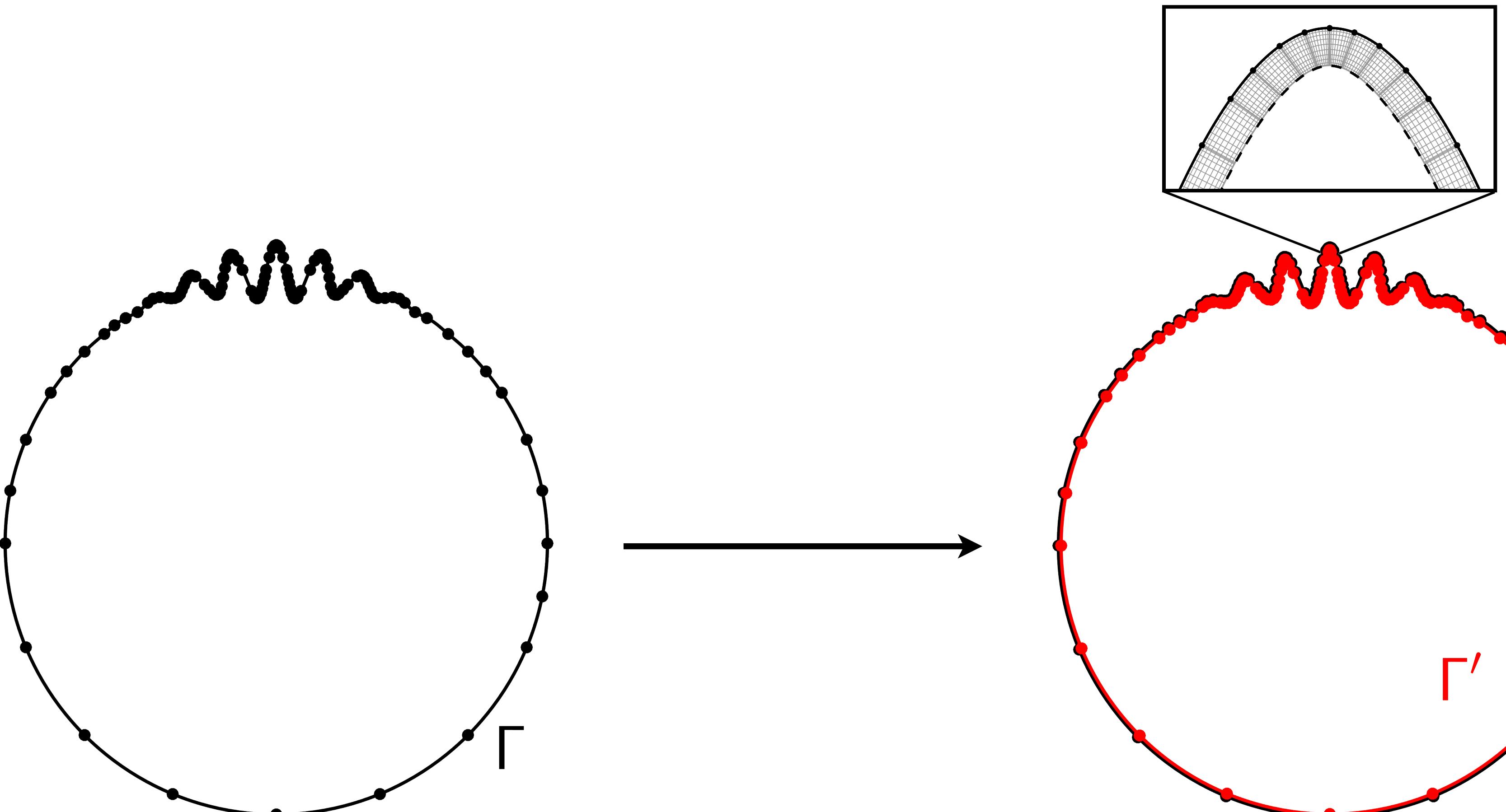
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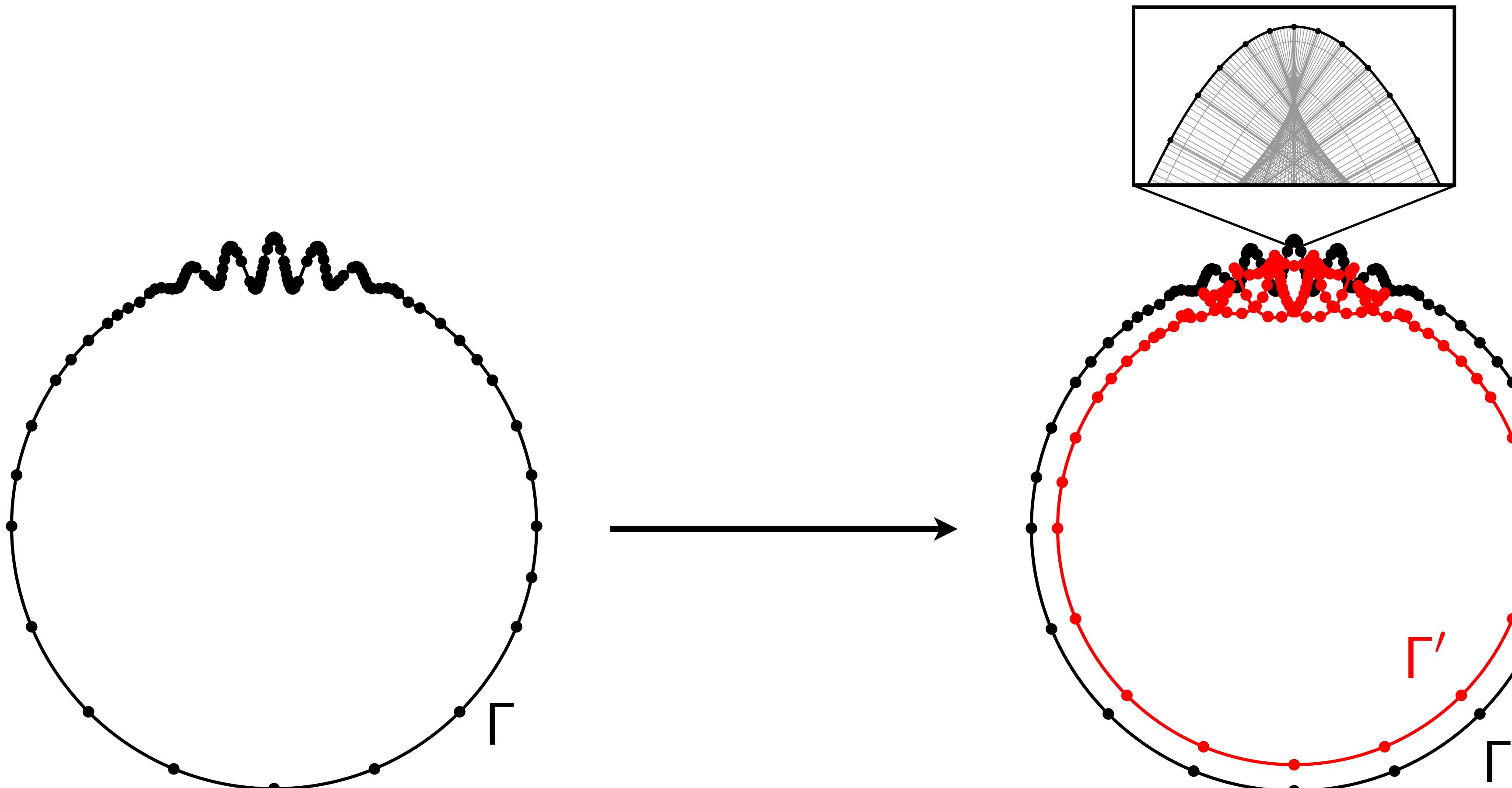
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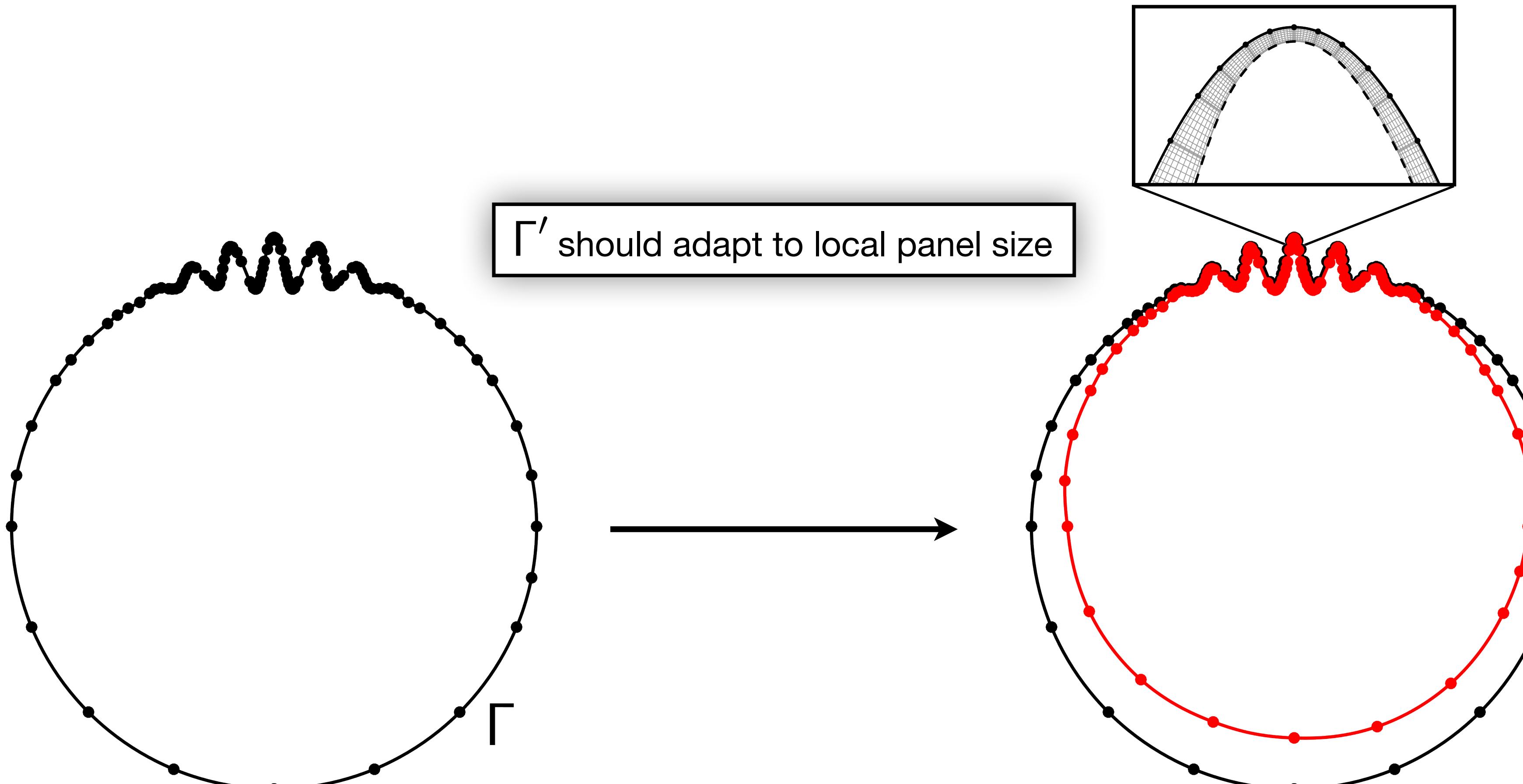
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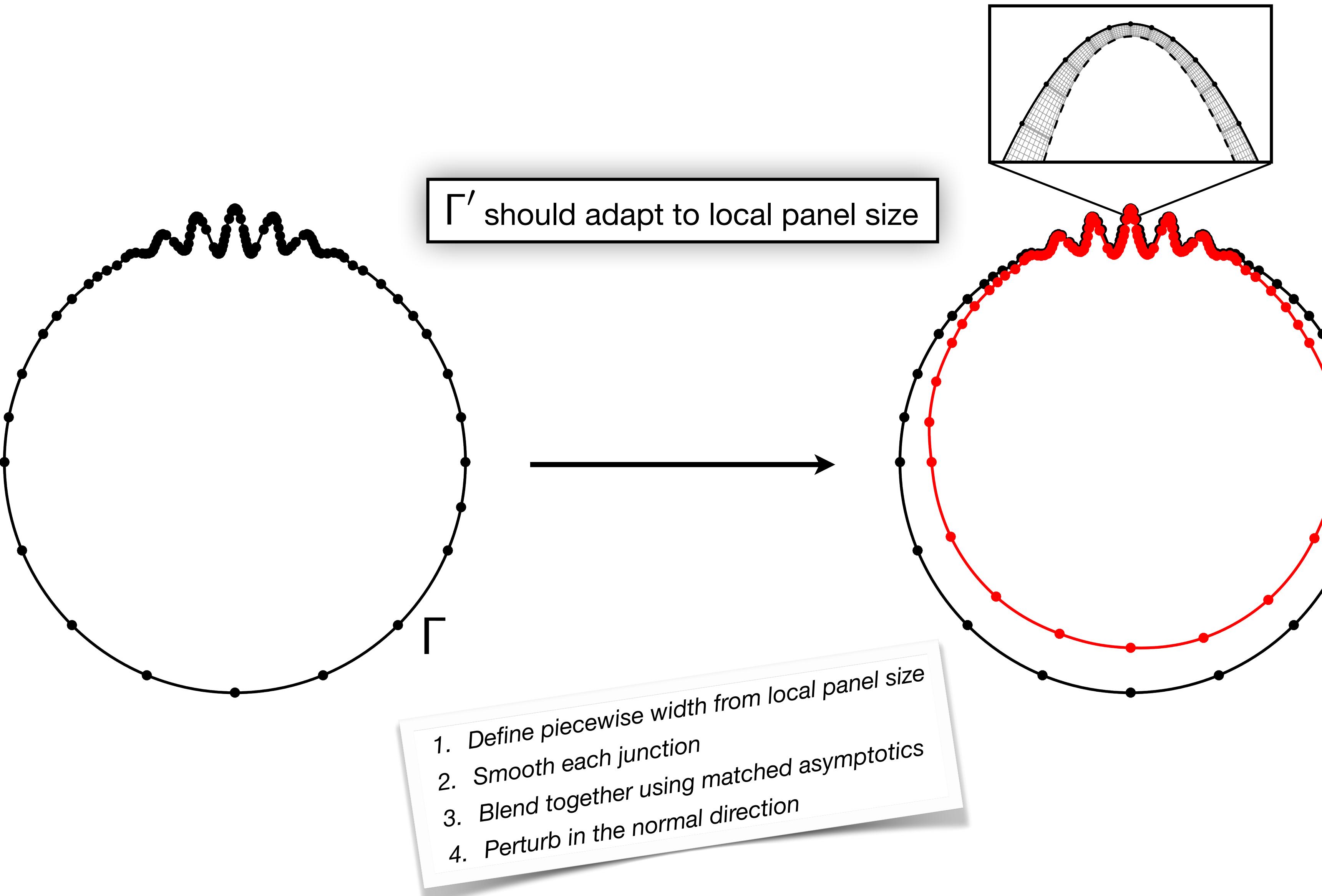
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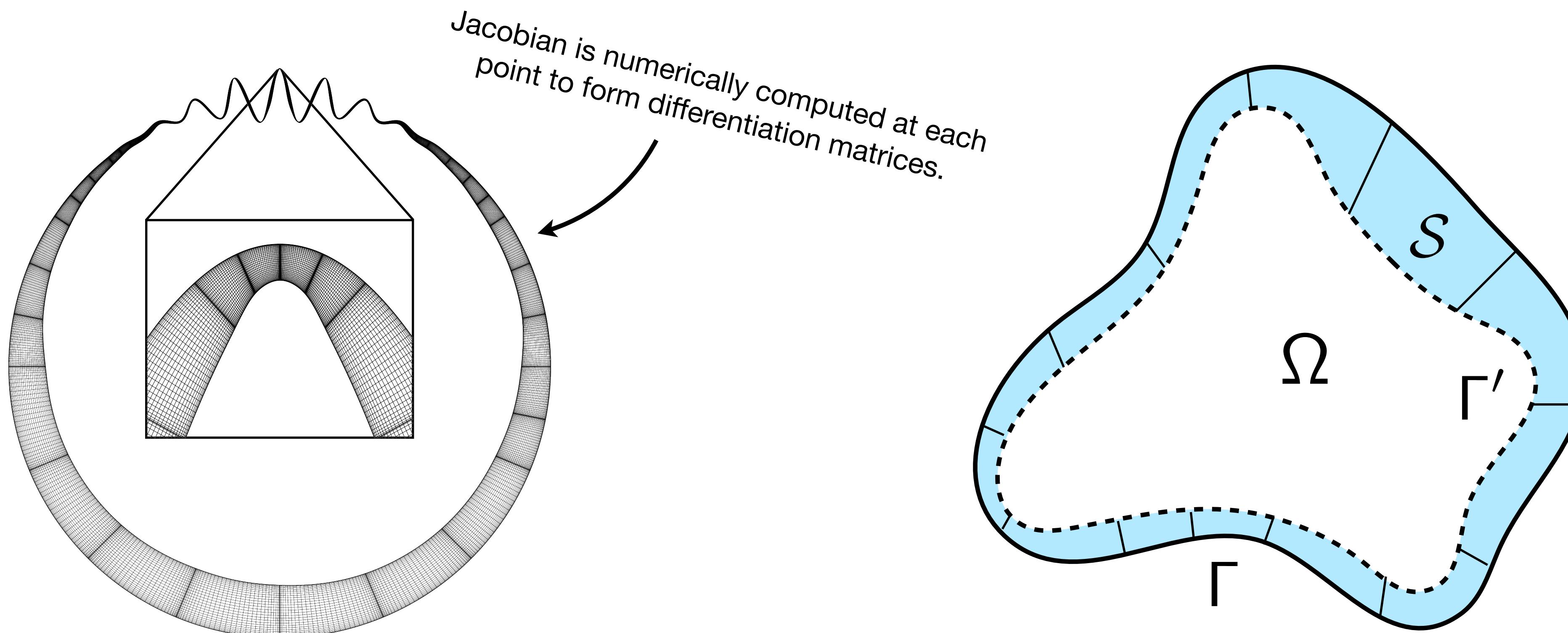


Solving the strip problem

Spectral collocation

$$\begin{aligned}\Delta v_{\text{strip}}(\mathbf{x}) &= f(\mathbf{x}), & \mathbf{x} \in \mathcal{S}, \\ v_{\text{strip}}(\mathbf{x}) &= 0, & \mathbf{x} \in \Gamma \cup \Gamma'.\end{aligned}$$

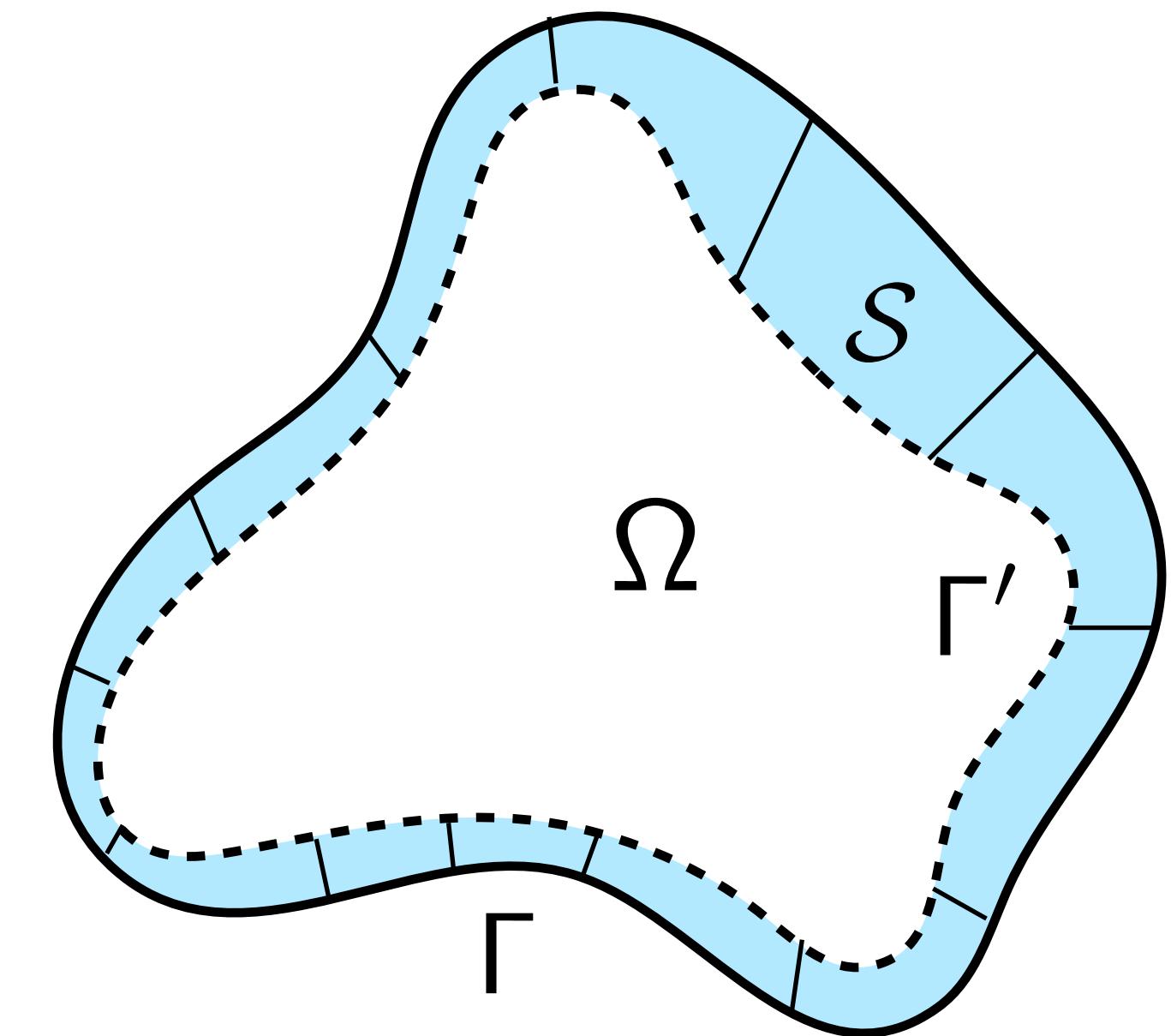
We use a spectral collocation method in \mathcal{S} and solve using the “1D” HPS method.



Solving the strip problem

A fast direct solver for the strip

We use the hierarchical Poincaré–Steklov scheme to build a **fast direct solver** in \mathcal{S} .



[Martinsson, 2013]

[Gillman & Martinsson, 2014]

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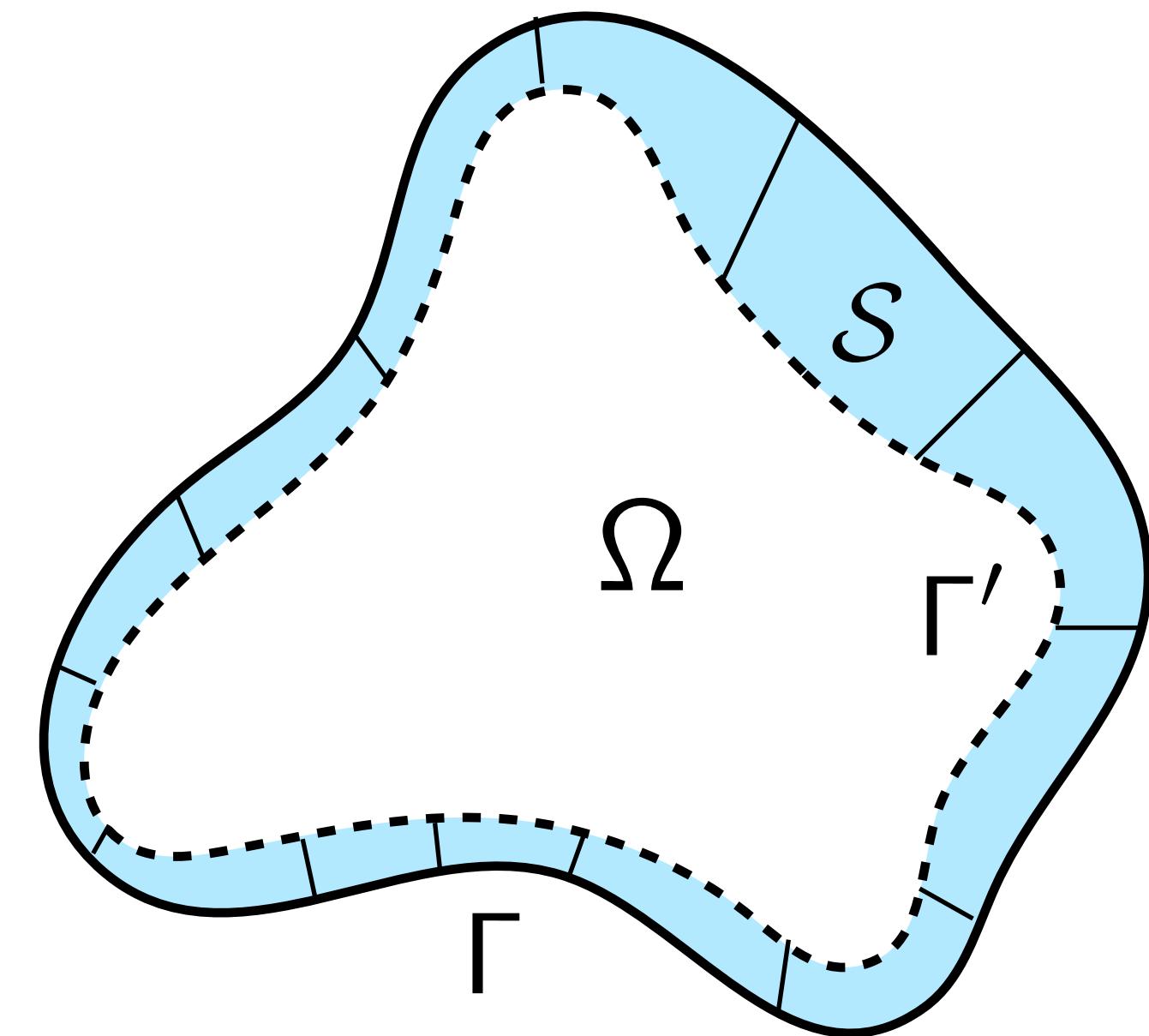
We use the hierarchical Poincaré–Steklov scheme to build a **fast direct solver** in \mathcal{S} .

Given an inhomogeneity f :

① On each element, compute:

- Solution operator:
- Dirichlet-to-Neumann map:

$$S \in \mathbb{R}^{n^2 \times 4n}$$
$$DtN \in \mathbb{R}^{4n \times 4n}$$



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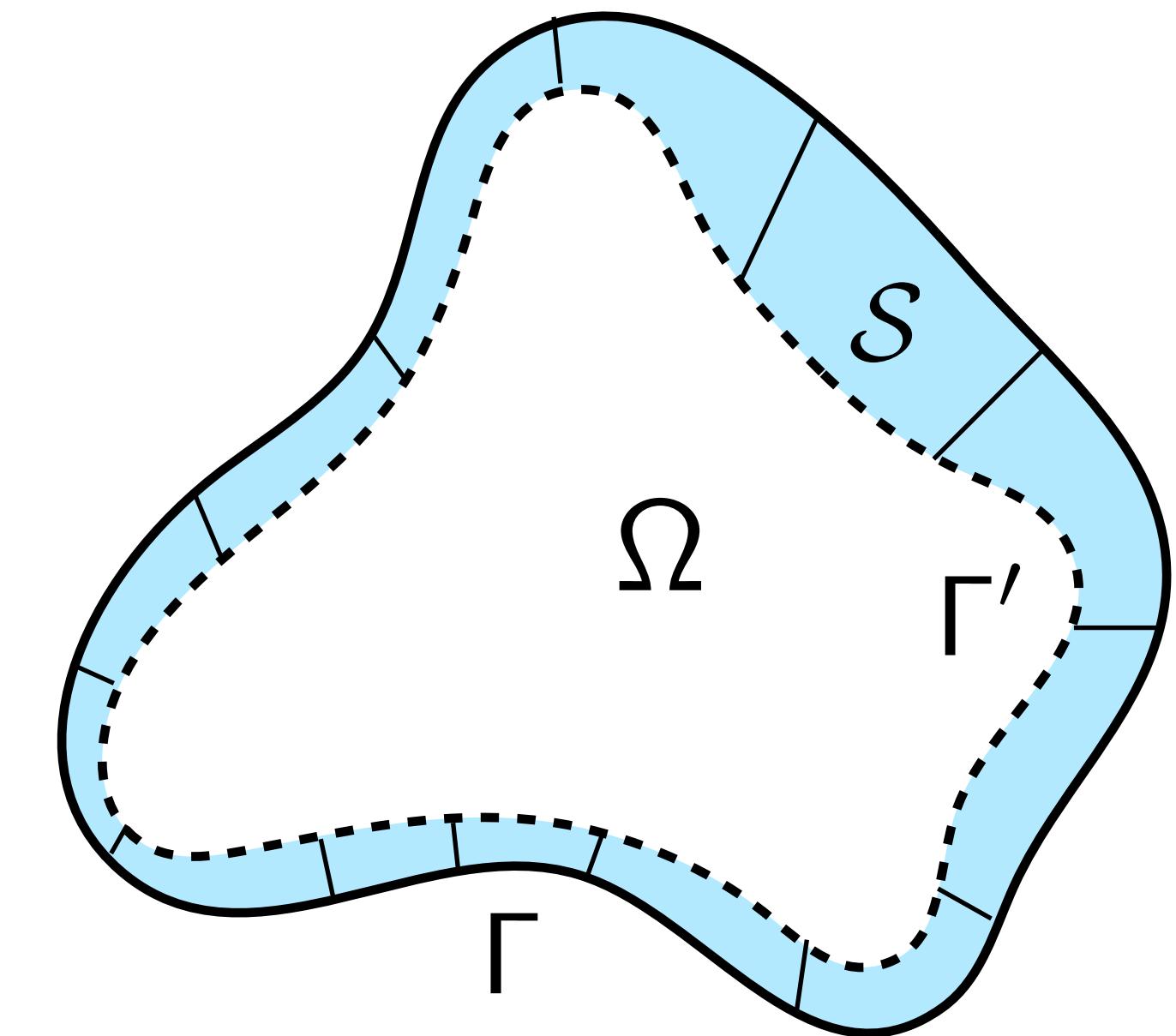
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② Merge adjacent elements pairwise

- Compute S and DtN on parent via Schur complement



[Martinsson, 2013]

[Gillman & Martinsson, 2014]

Solving the strip problem

A fast direct solver for the strip

We use the hierarchical Poincaré–Steklov scheme to build a **fast direct solver** in \mathcal{S} .

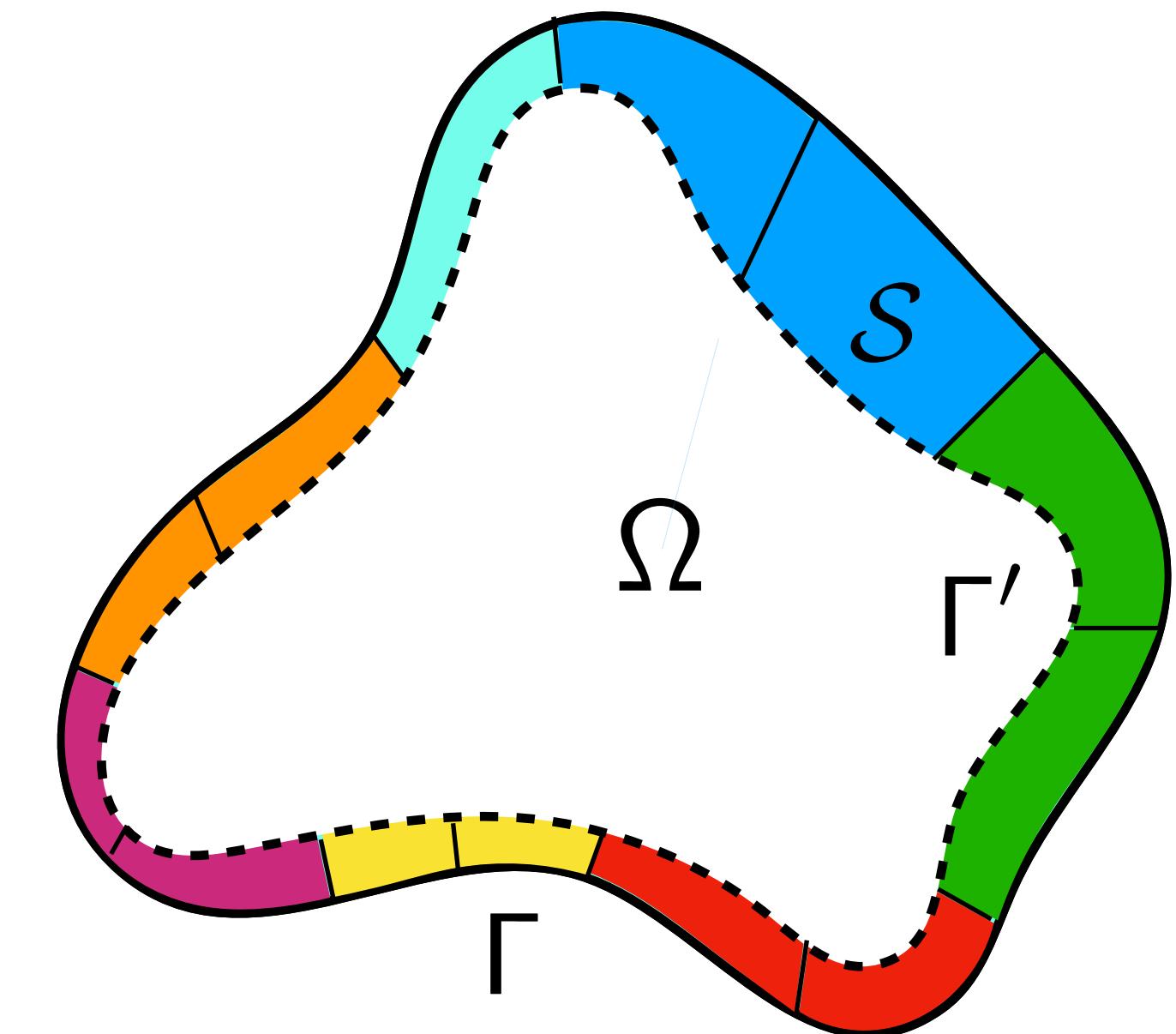
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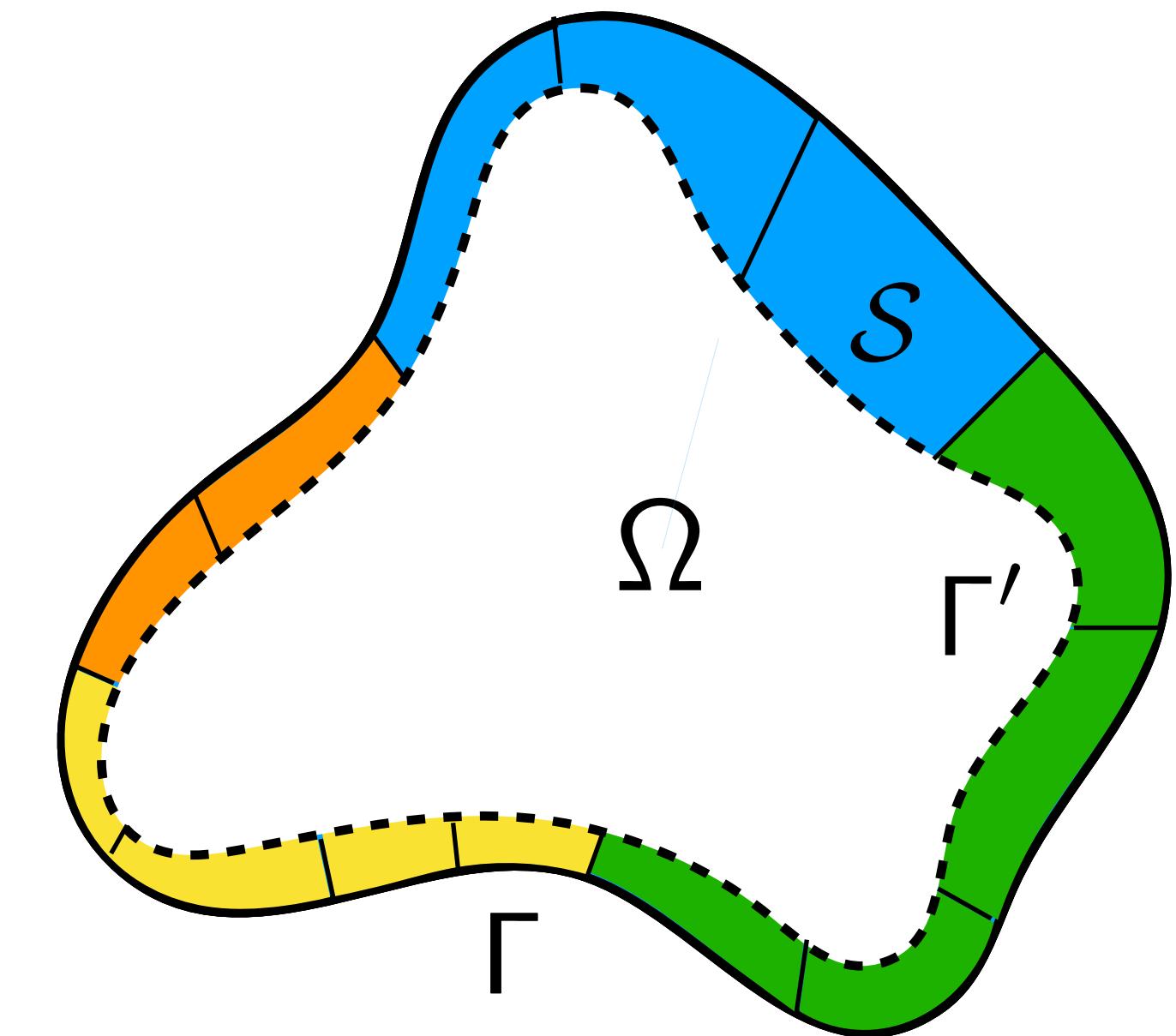
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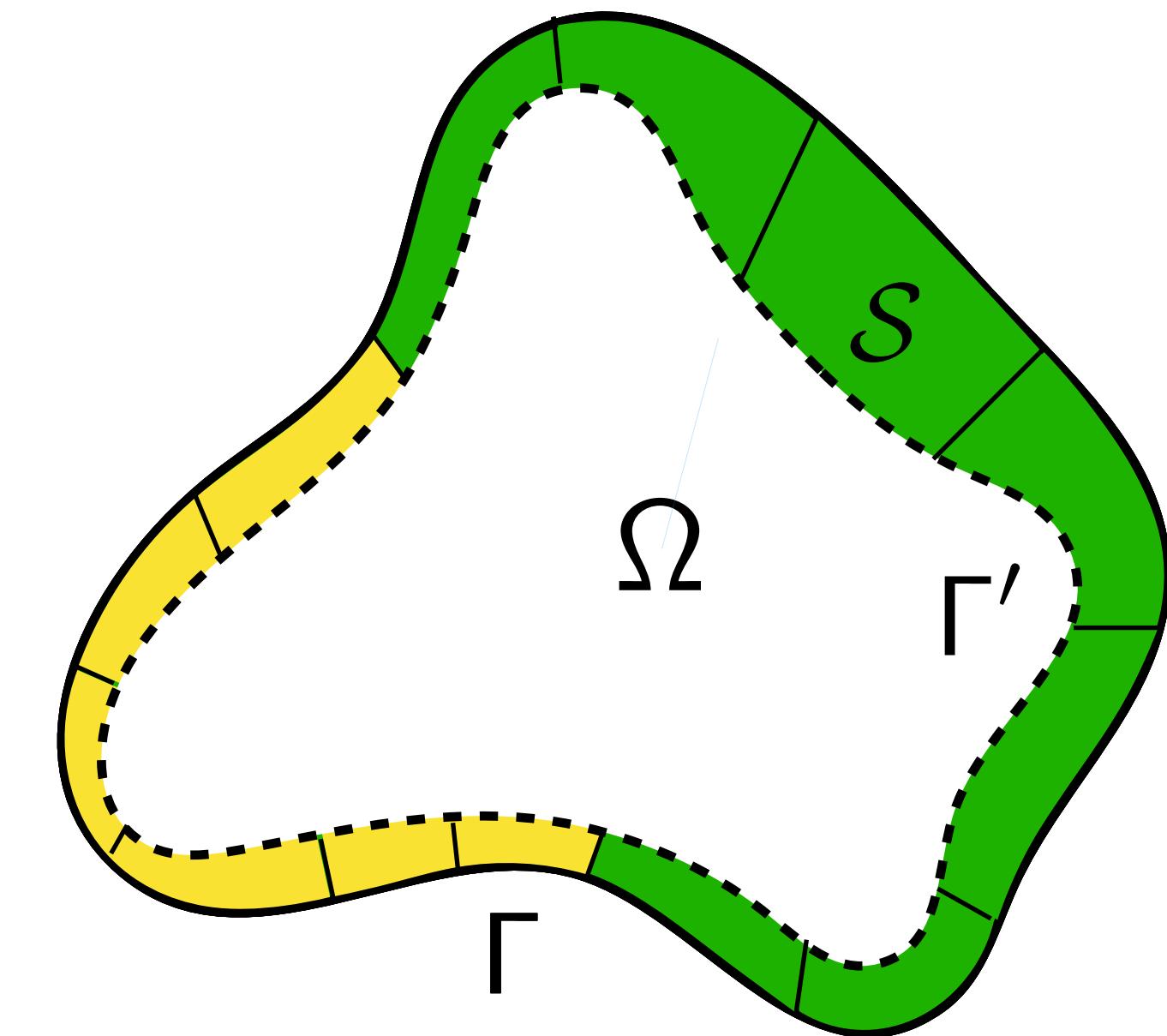
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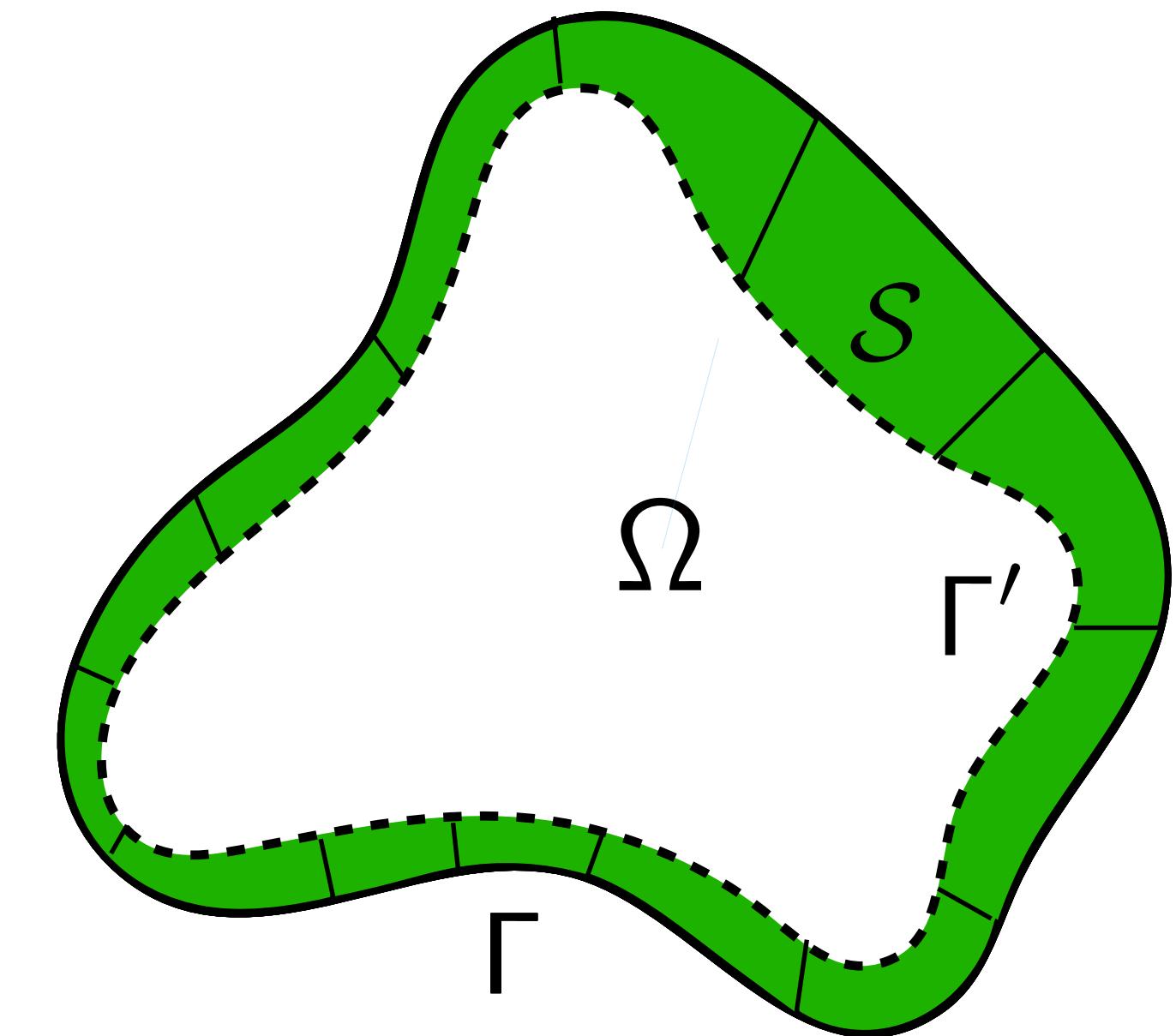
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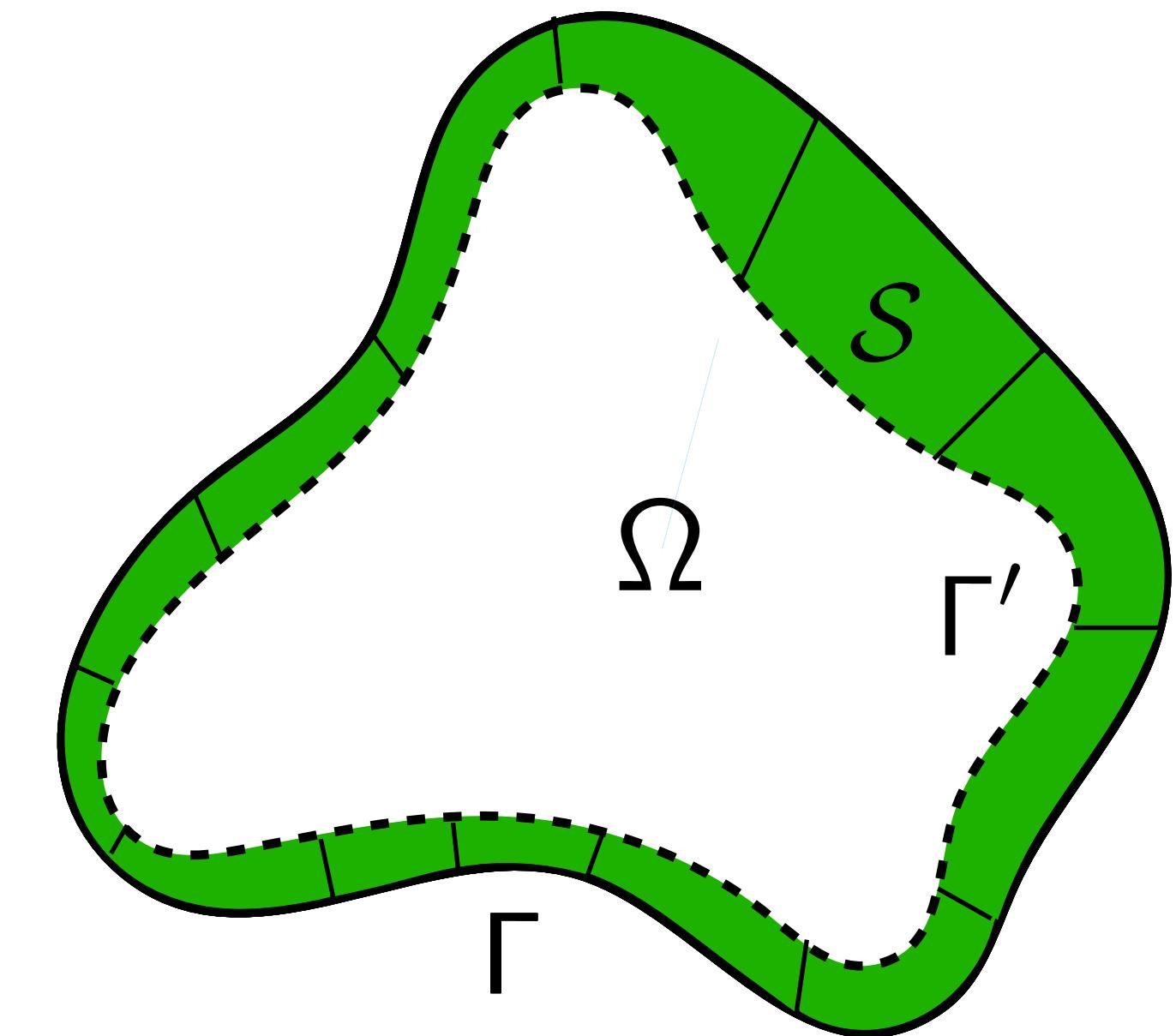
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$$\text{Cost: } \mathcal{O}(p^6 n_{\text{panel}}) + \mathcal{O}(p^3 n_{\text{panel}}) + \mathcal{O}(p^2 n_{\text{panel}}) = \mathcal{O}(n_{\text{panel}})$$

①

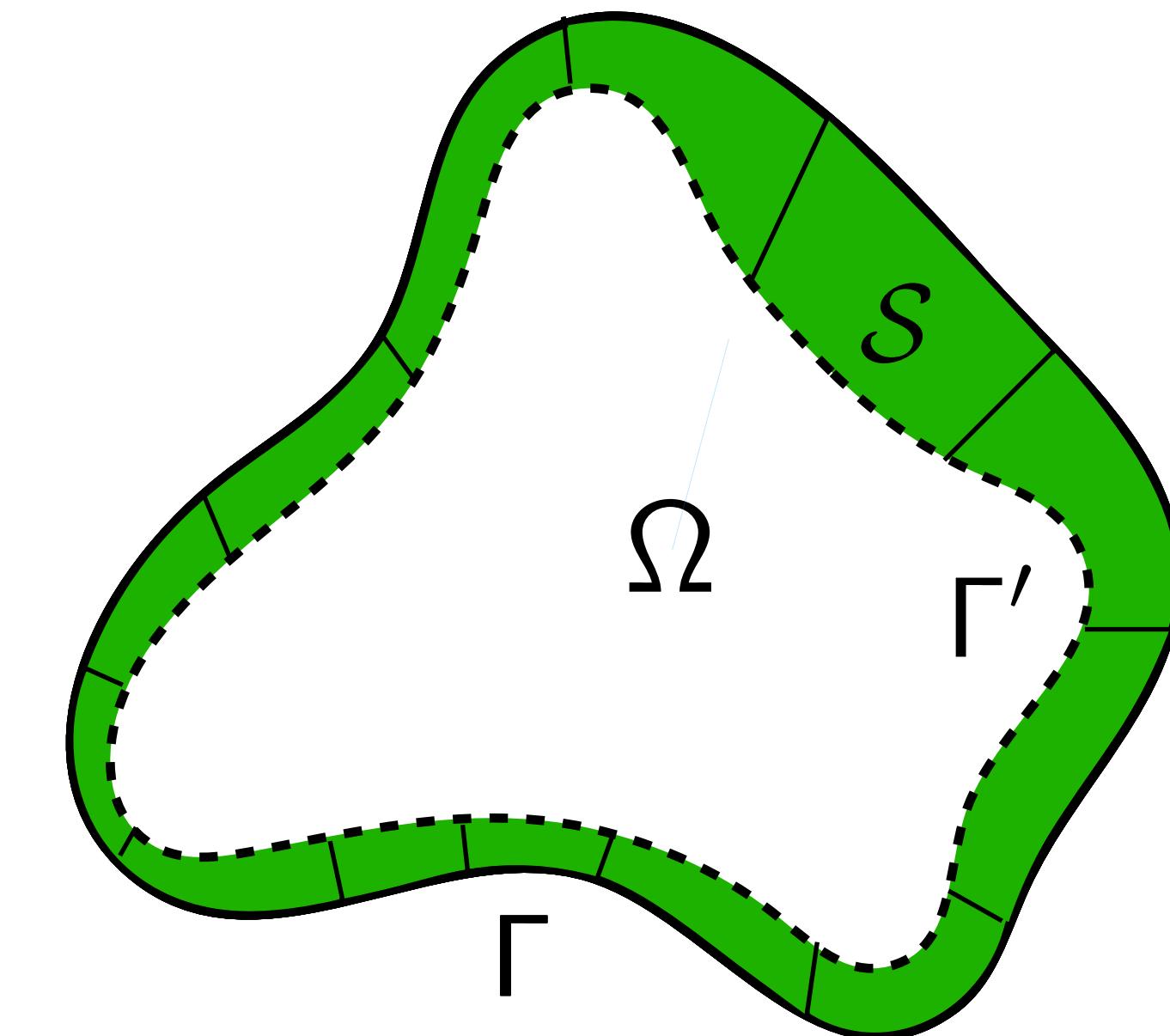
②

③

We typically use $p = 16$ on each panel and upsample the SEM grid to $2p = 32$.

[Martinsson, 2013]

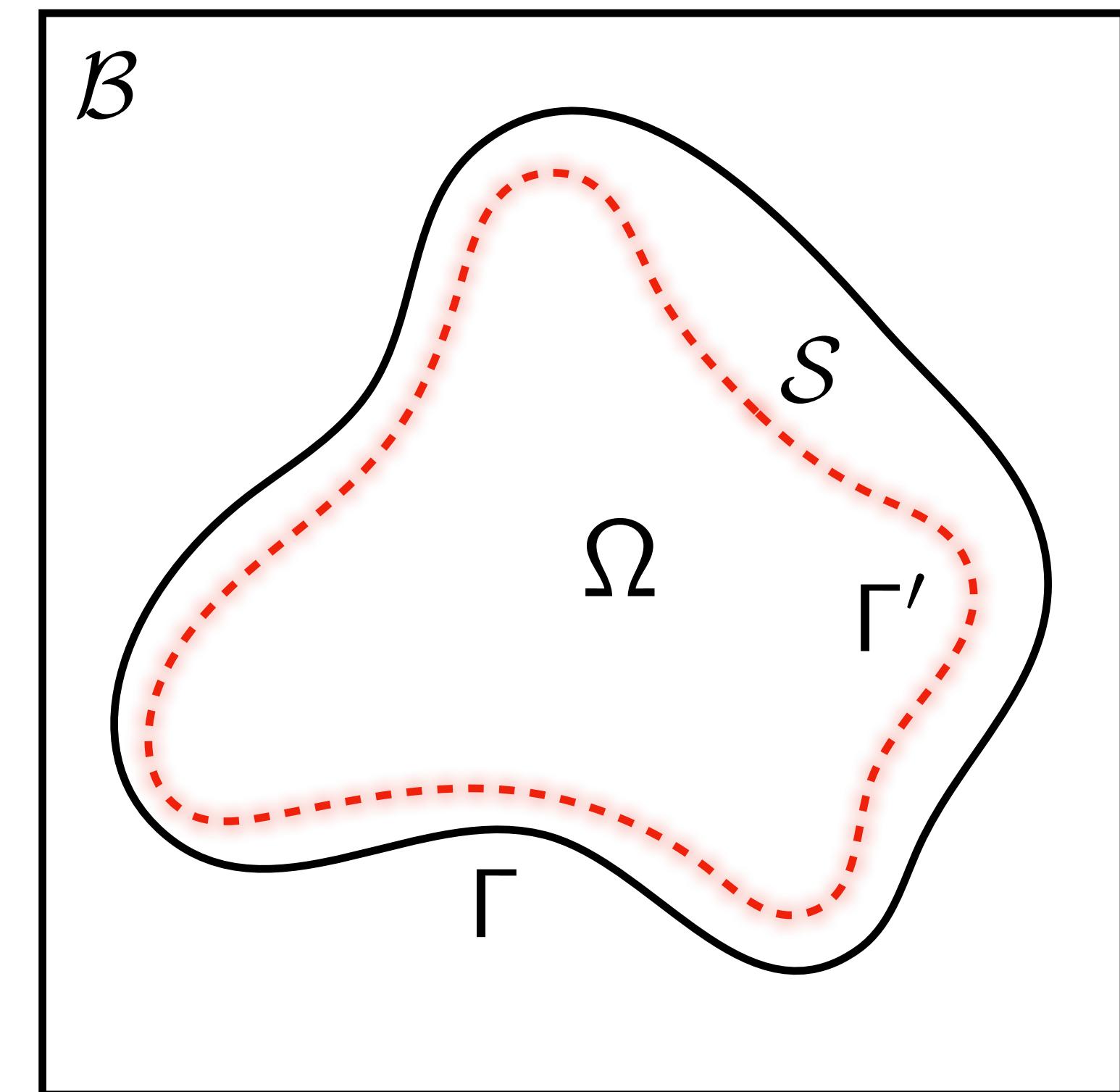
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Patching the solutions

Layer potentials

The solutions v_{bulk} and v_{strip} may differ in value and normal derivative along Γ' .



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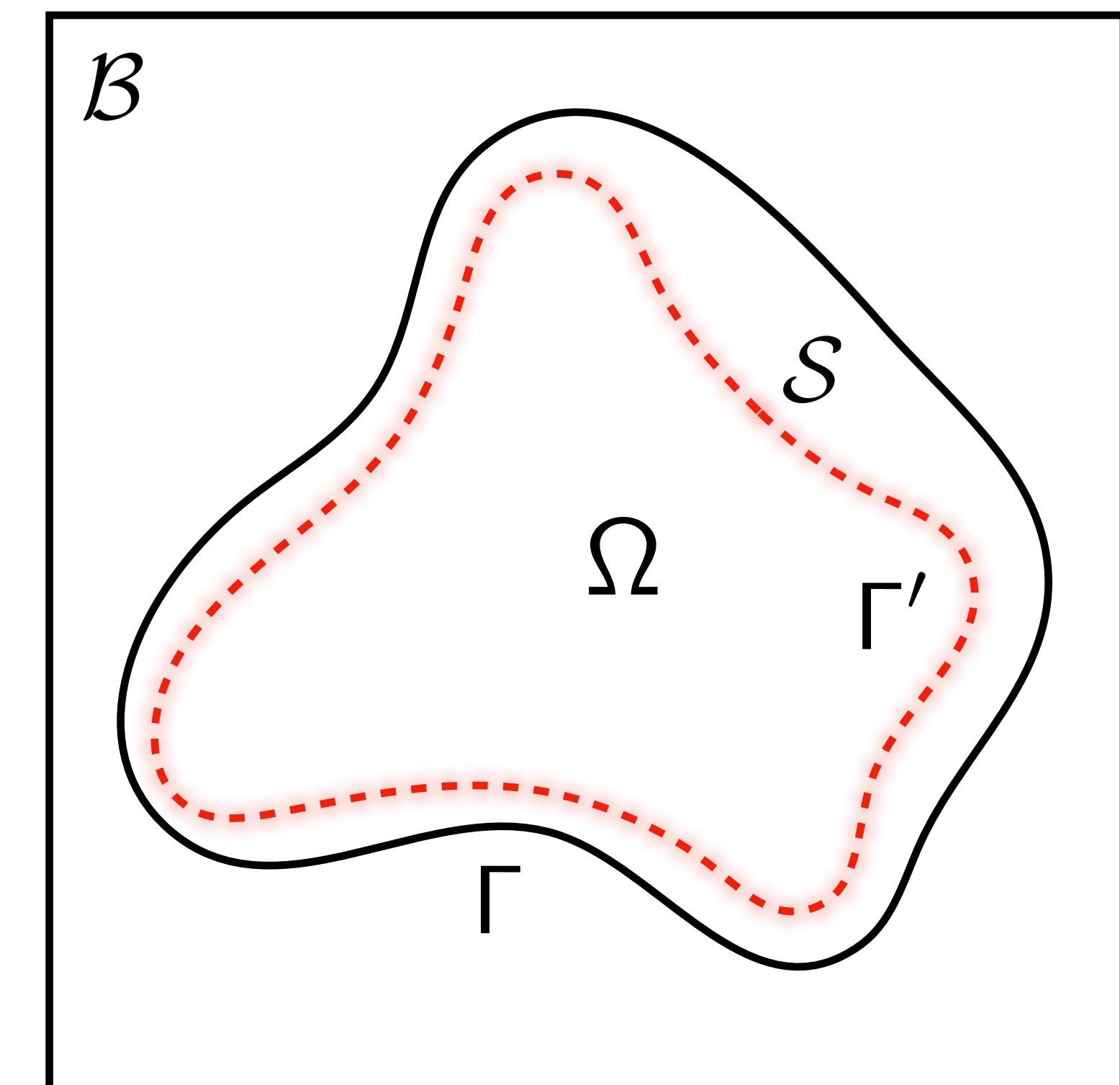
are harmonic and satisfy the jump relations

$$S[\sigma]|_{\Gamma'}^+ - S[\sigma]|_{\Gamma'}^- = 0$$

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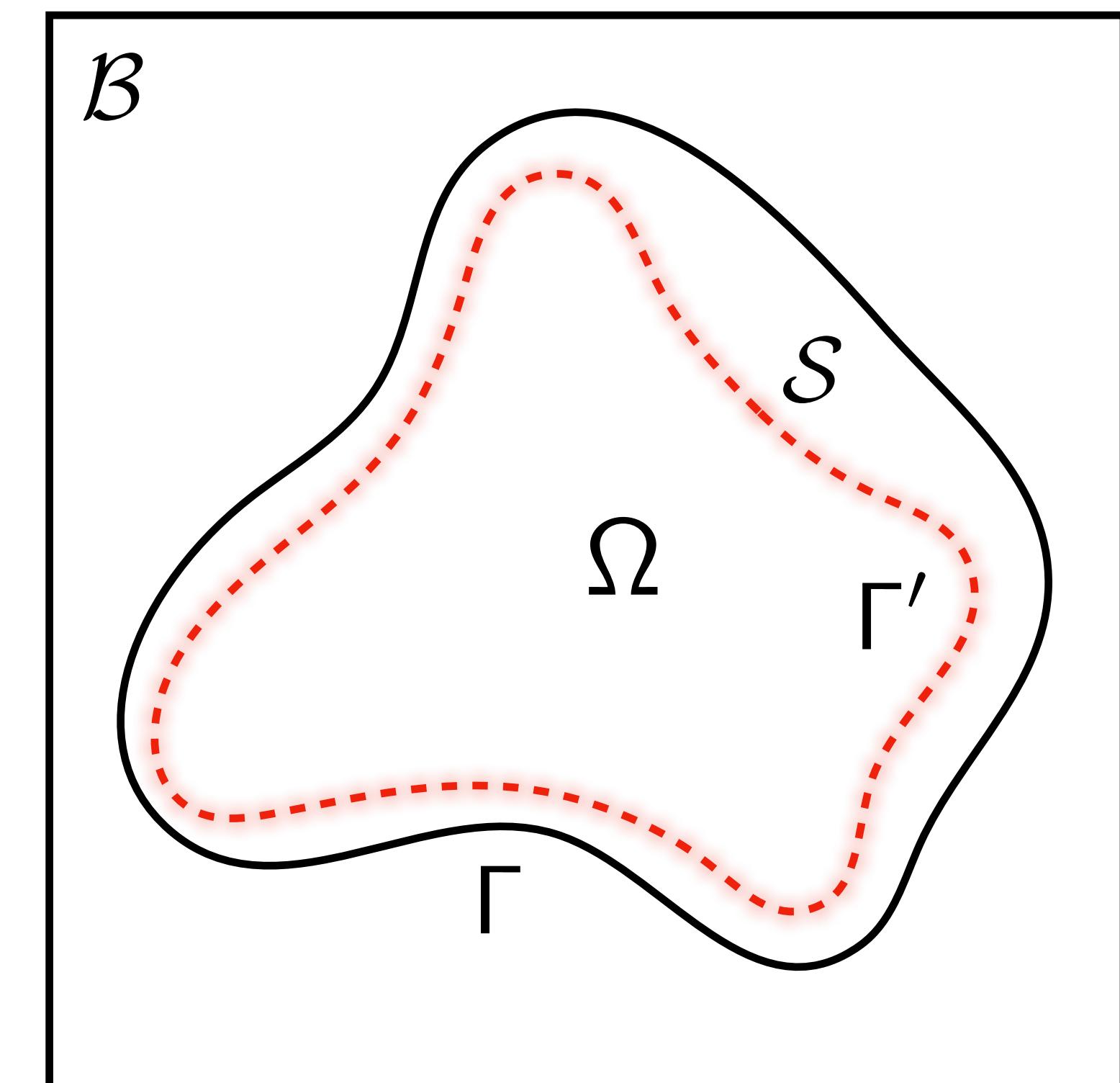
$$\partial_{\mathbf{n}} S[\sigma]|_{\Gamma'}^+ - \partial_{\mathbf{n}} S[\sigma]|_{\Gamma'}^- = -\sigma$$

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So define $v_{\text{glue}} = S[\sigma] - D[\tau]$ with

$$\sigma = \partial_{\mathbf{n}} v_{\text{bulk}}|_{\Gamma'} - \partial_{\mathbf{n}} v_{\text{strip}}|_{\Gamma'}$$

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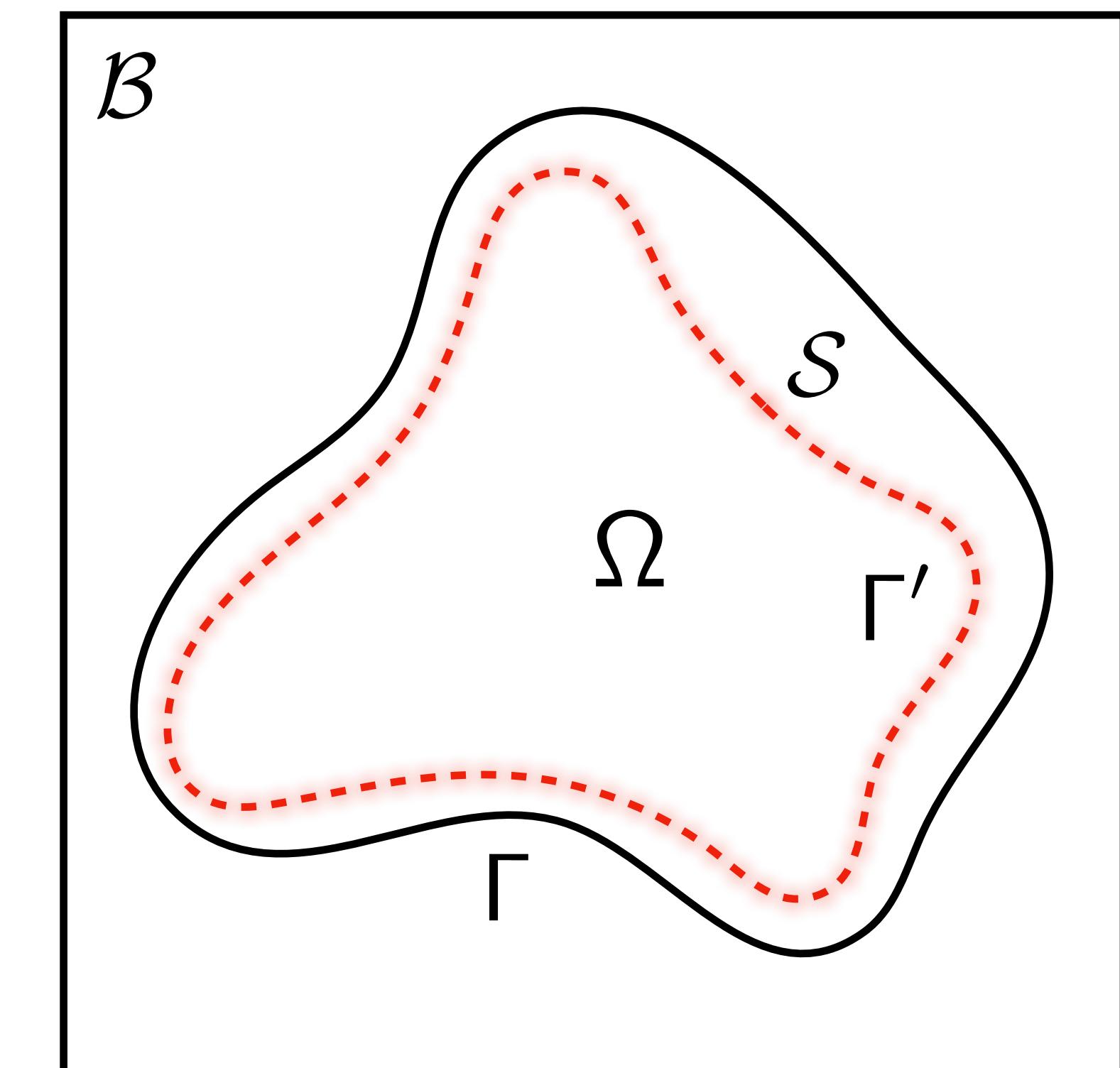
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Then

$$v = \begin{cases} v_{\text{bulk}} + v_{\text{glue}} & \text{in } \Omega \setminus \mathcal{S} \\ v_{\text{strip}} + v_{\text{glue}} & \text{in } \mathcal{S} \end{cases}$$

satisfies $\Delta v = f$ in Ω and is smooth across Γ' .



Correcting for boundary data

Homogeneous solution

Finally, the boundary conditions may not be satisfied.

A homogeneous correction may be computed using standard BIE techniques:

$$\Delta w = 0 \quad \text{in } \Omega$$

$$w = g - v_{\text{glue}}|_{\Gamma} \quad \text{on } \Gamma$$

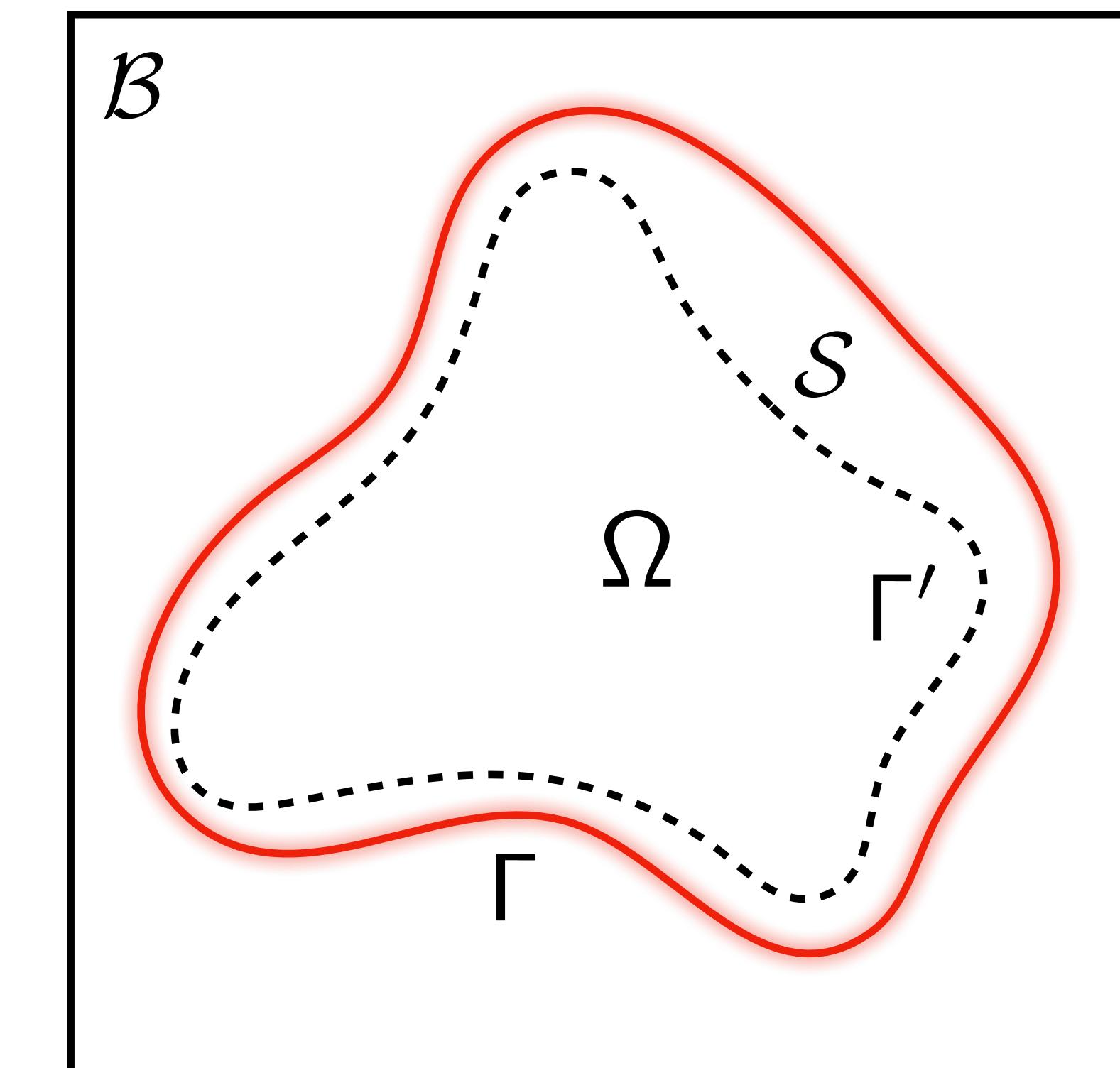
Then,

$$u = \begin{cases} v_{\text{bulk}} + v_{\text{glue}} + w & \text{in } \Omega \setminus \mathcal{S} \\ v_{\text{strip}} + v_{\text{glue}} + w & \text{in } \mathcal{S} \end{cases}$$

satisfies

$$\Delta u = f \quad \text{in } \Omega$$

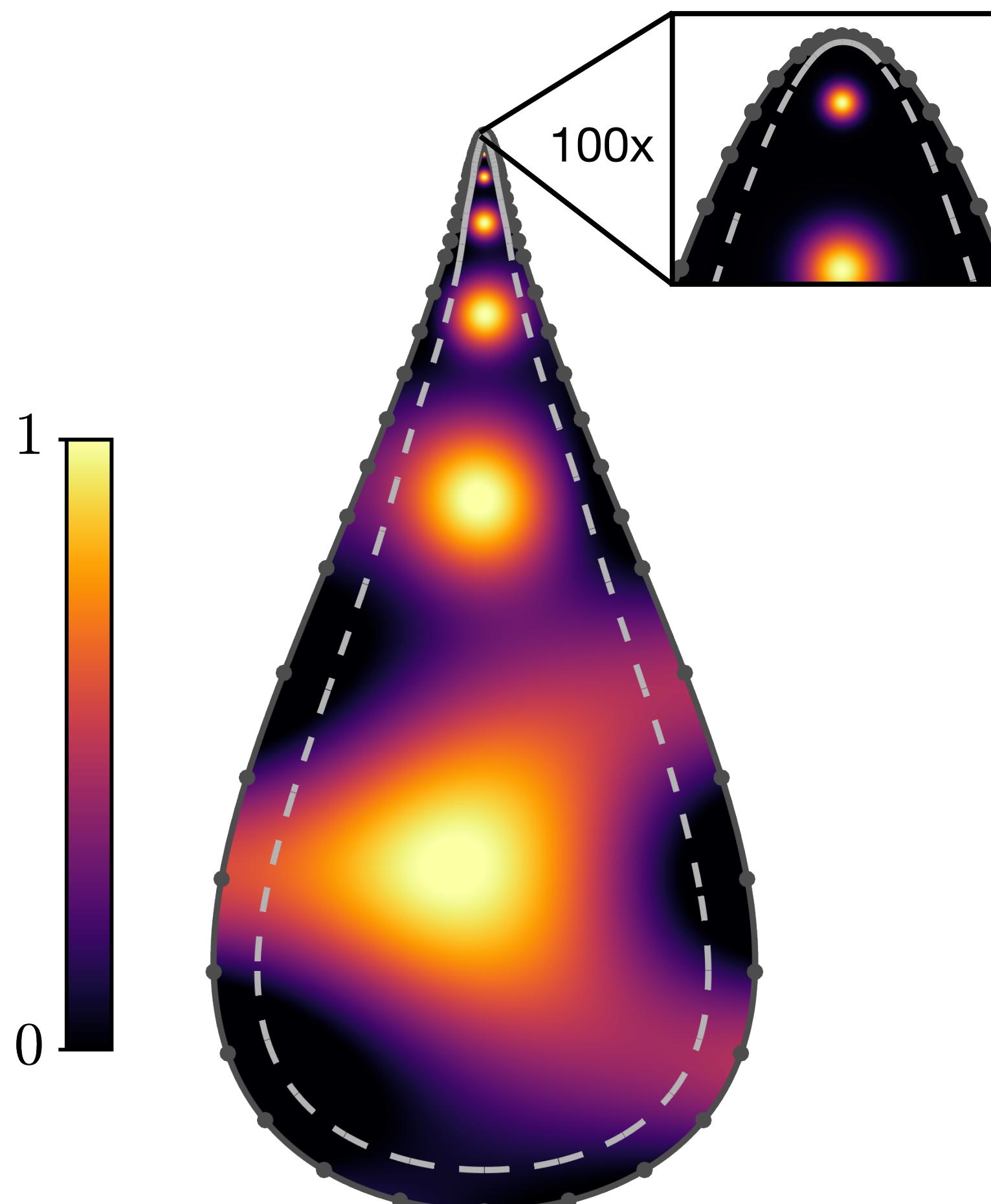
$$u = g \quad \text{on } \Gamma$$



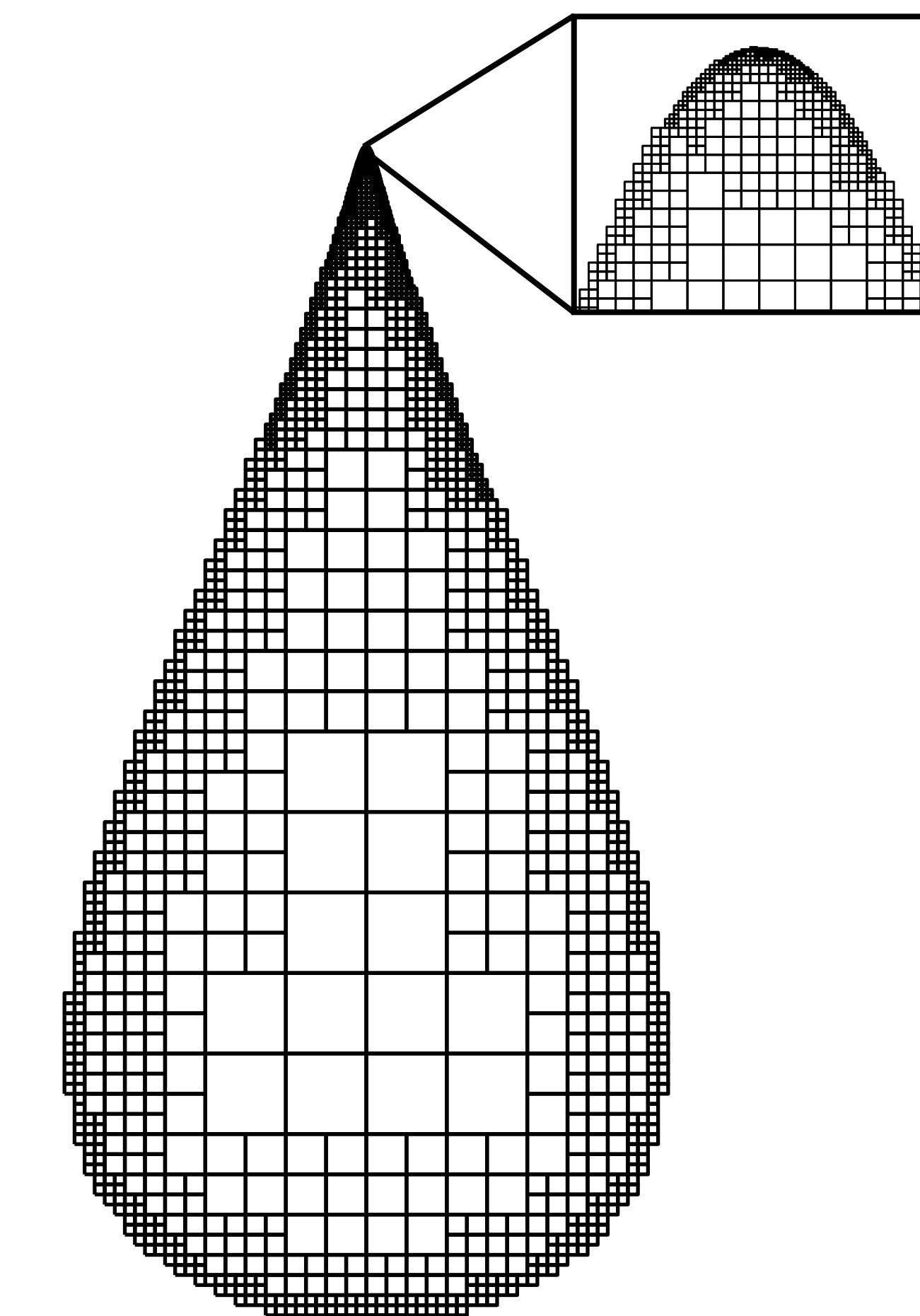
Example

A rounded cusp

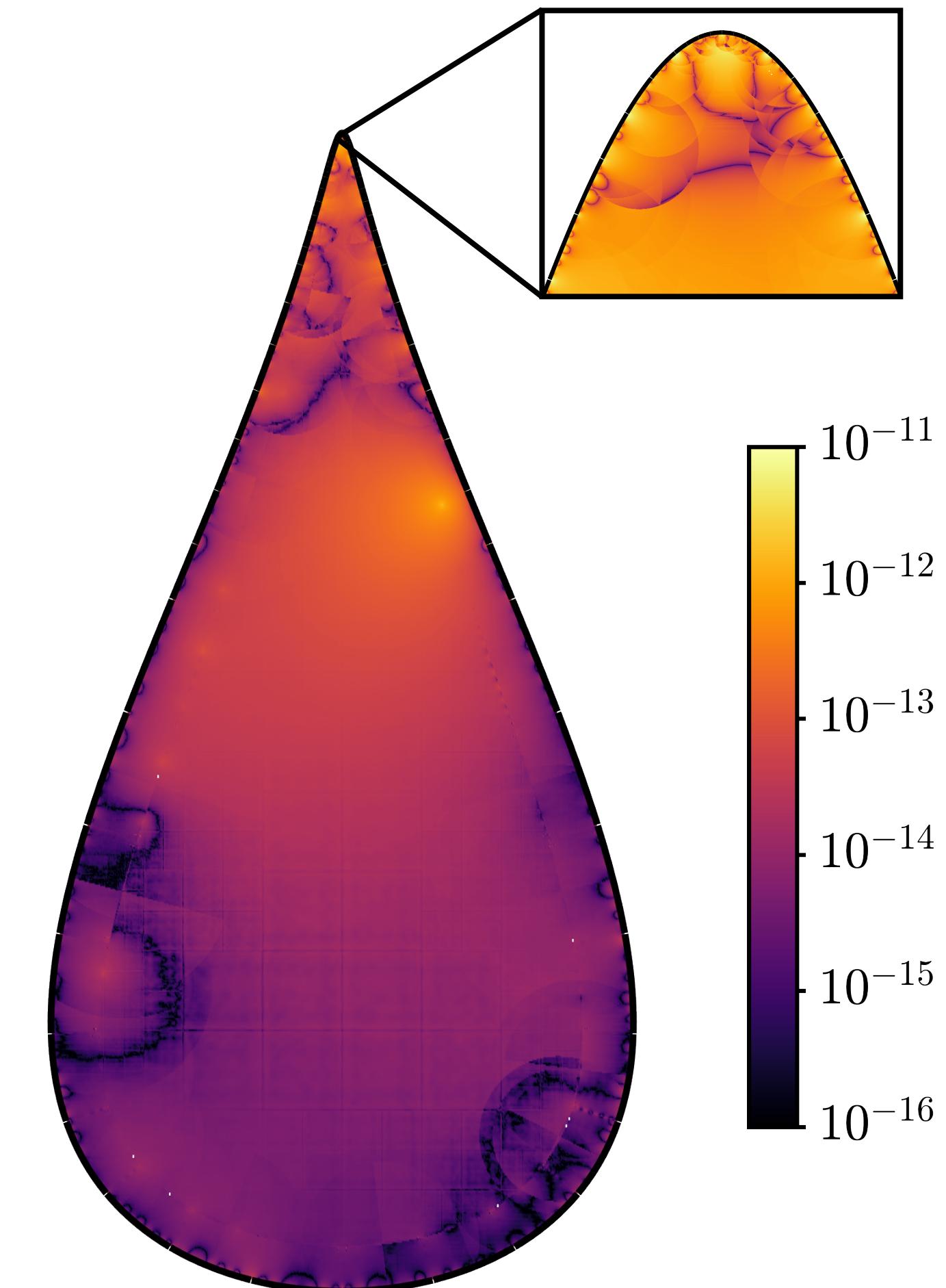
Inhomogeneity



Quadtree



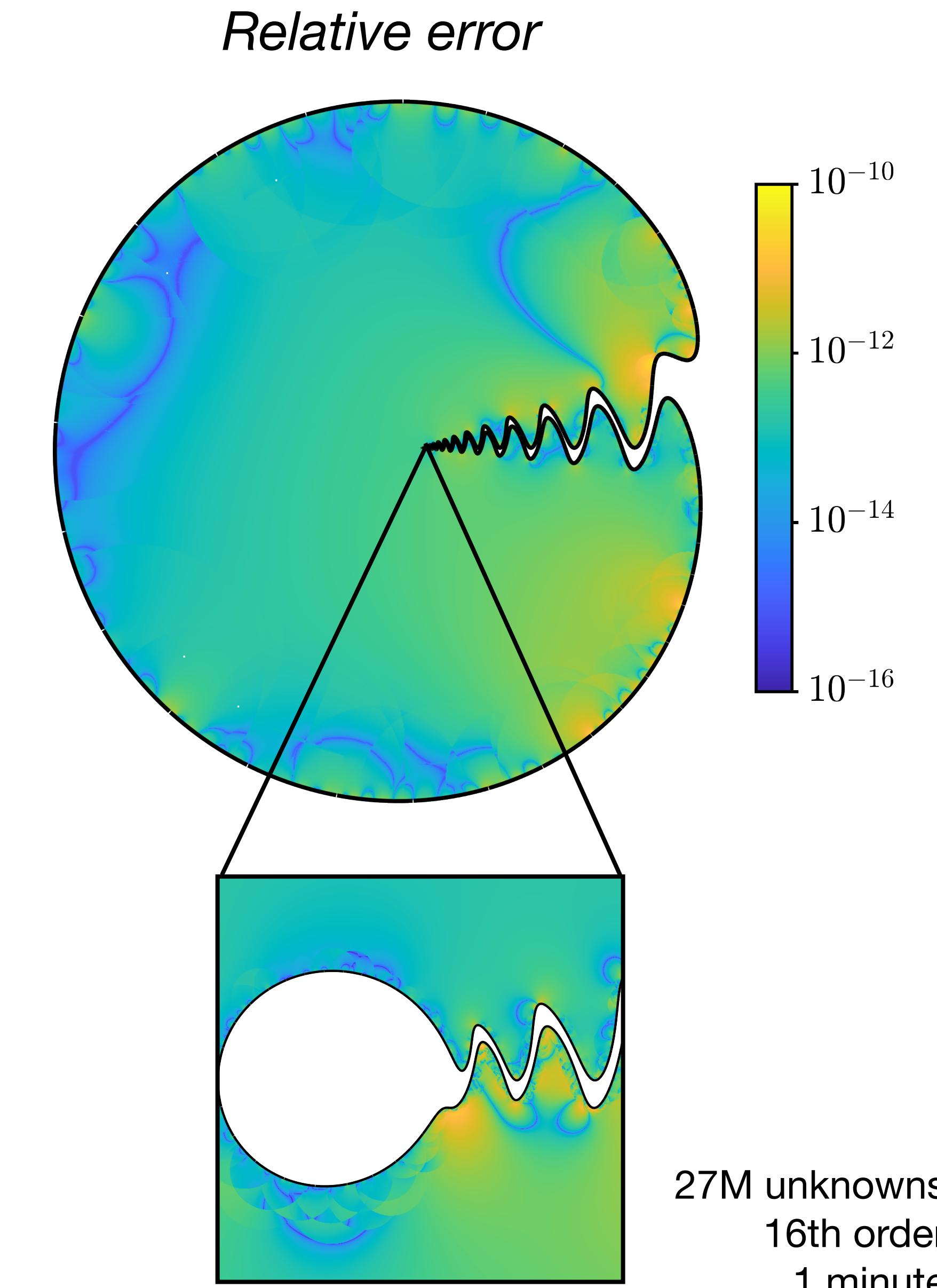
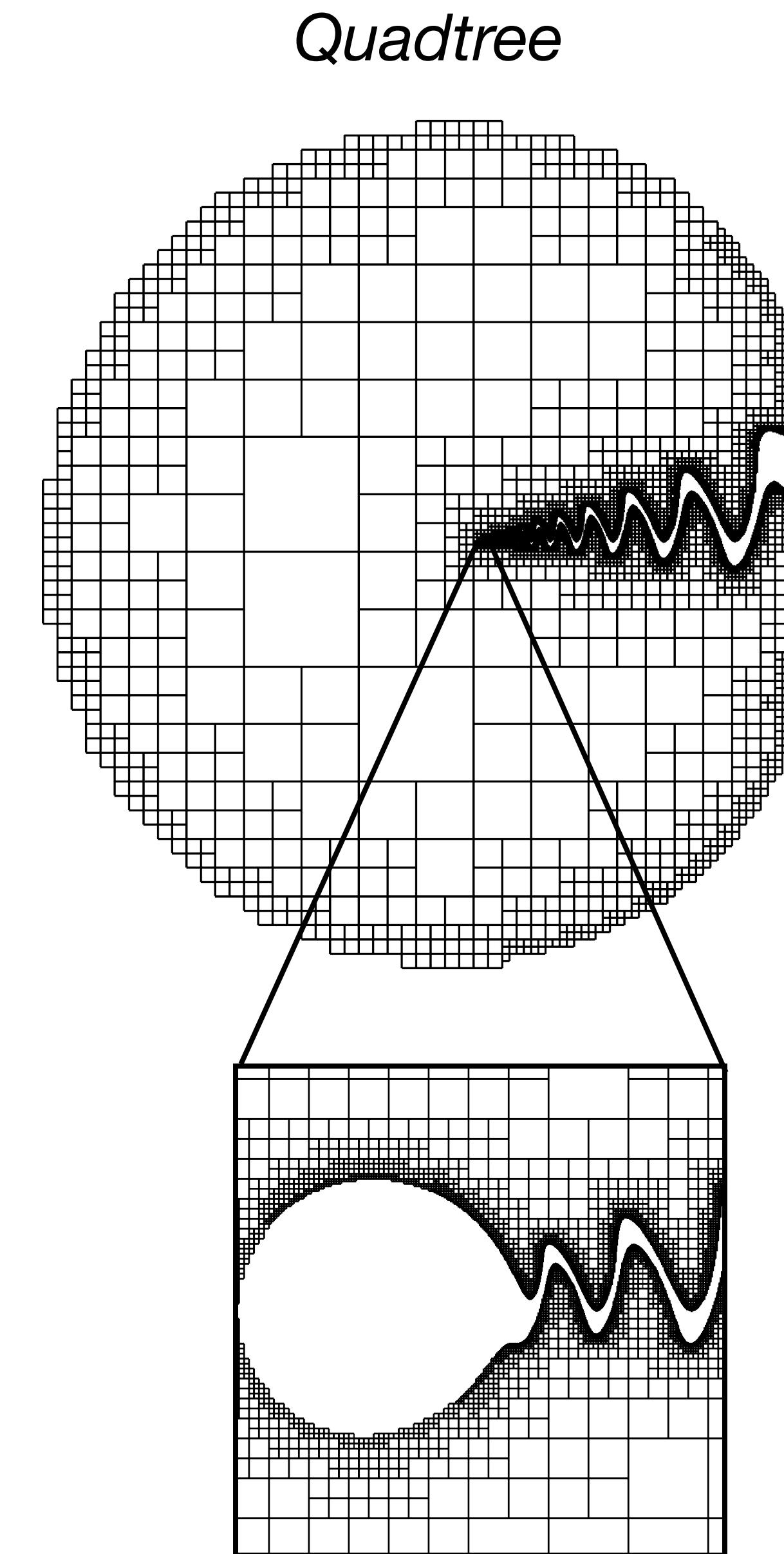
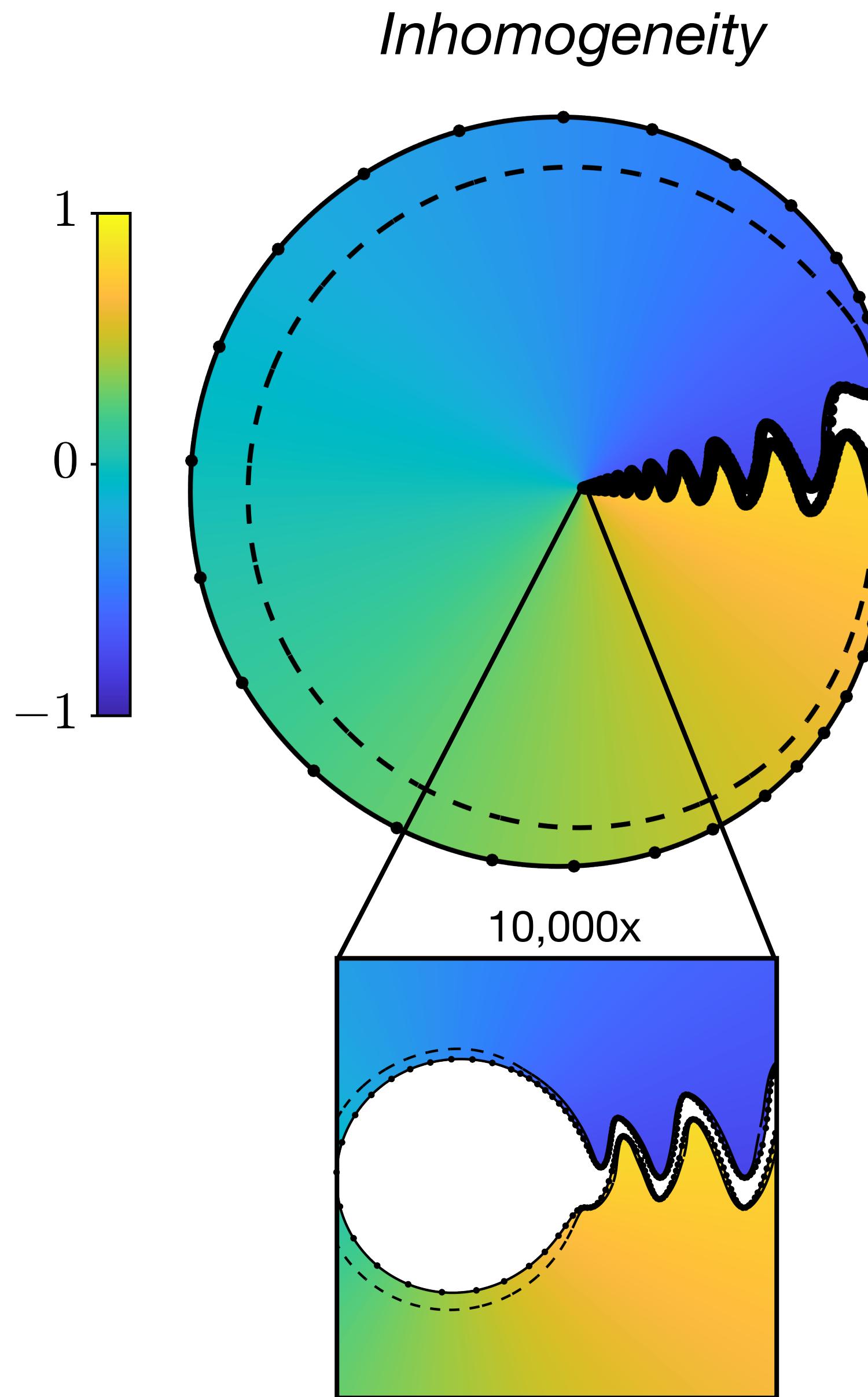
Relative error



750k unknowns
16th order
2 seconds

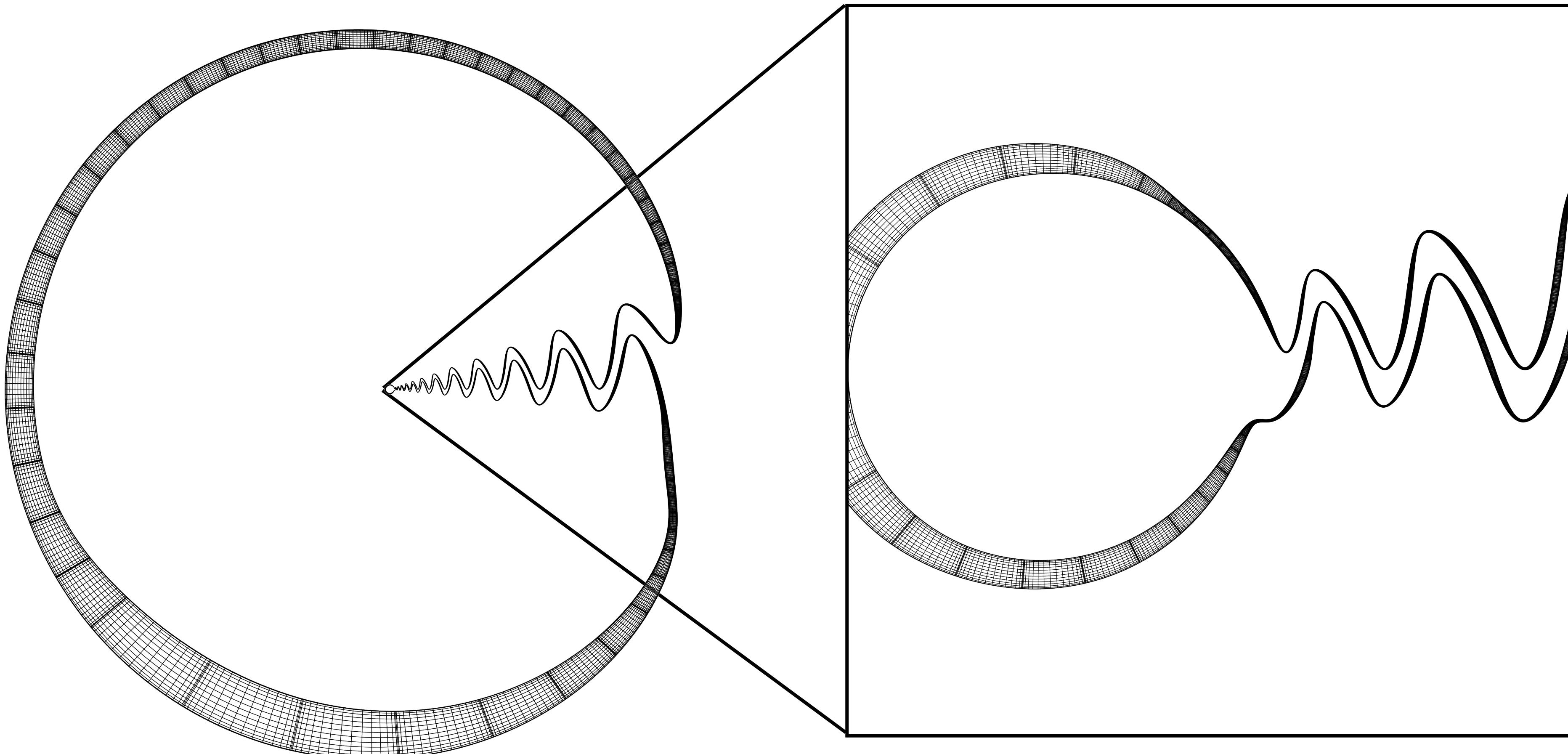
Example

A multiscale inhomogeneity and close-to-touching geometry



Example

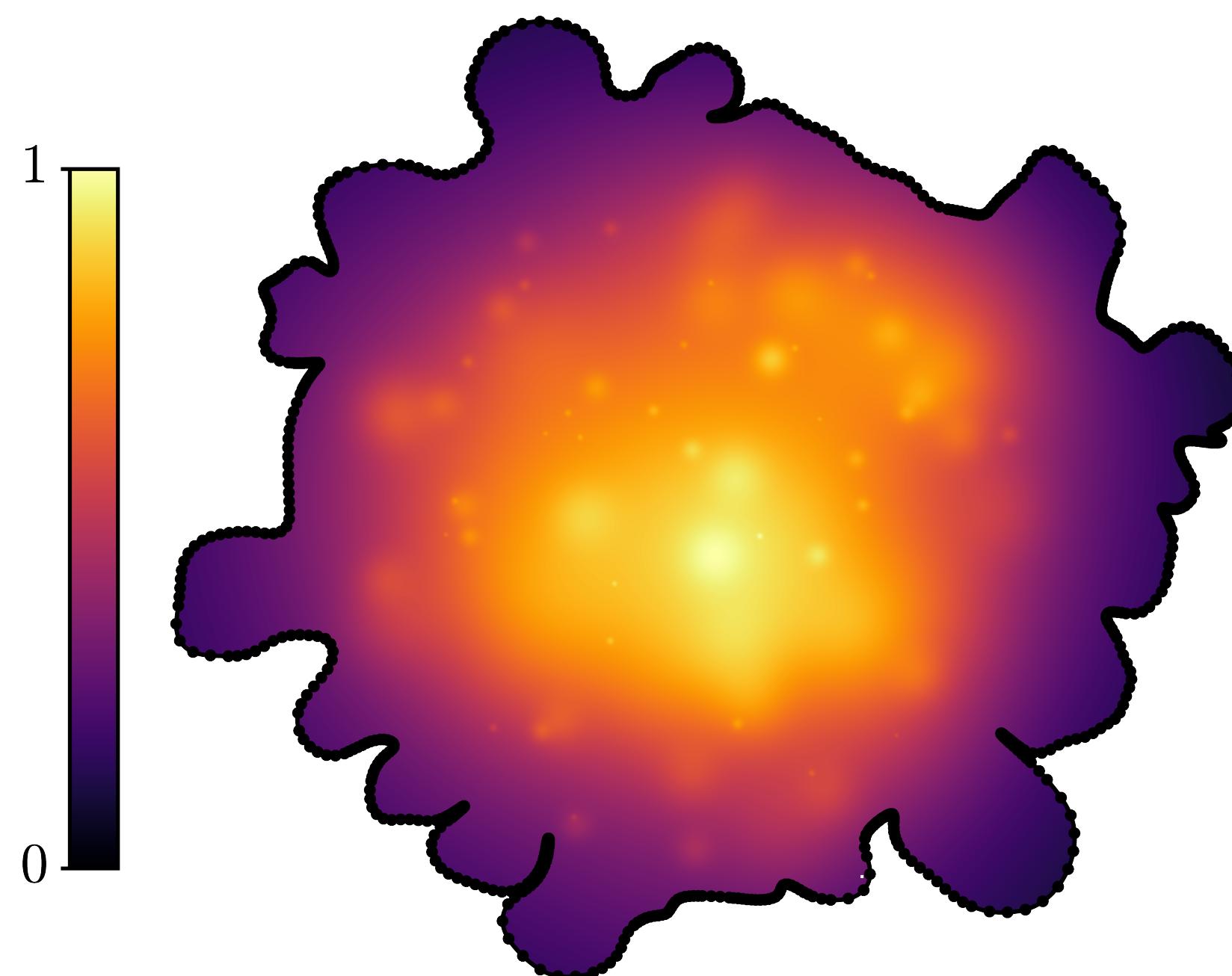
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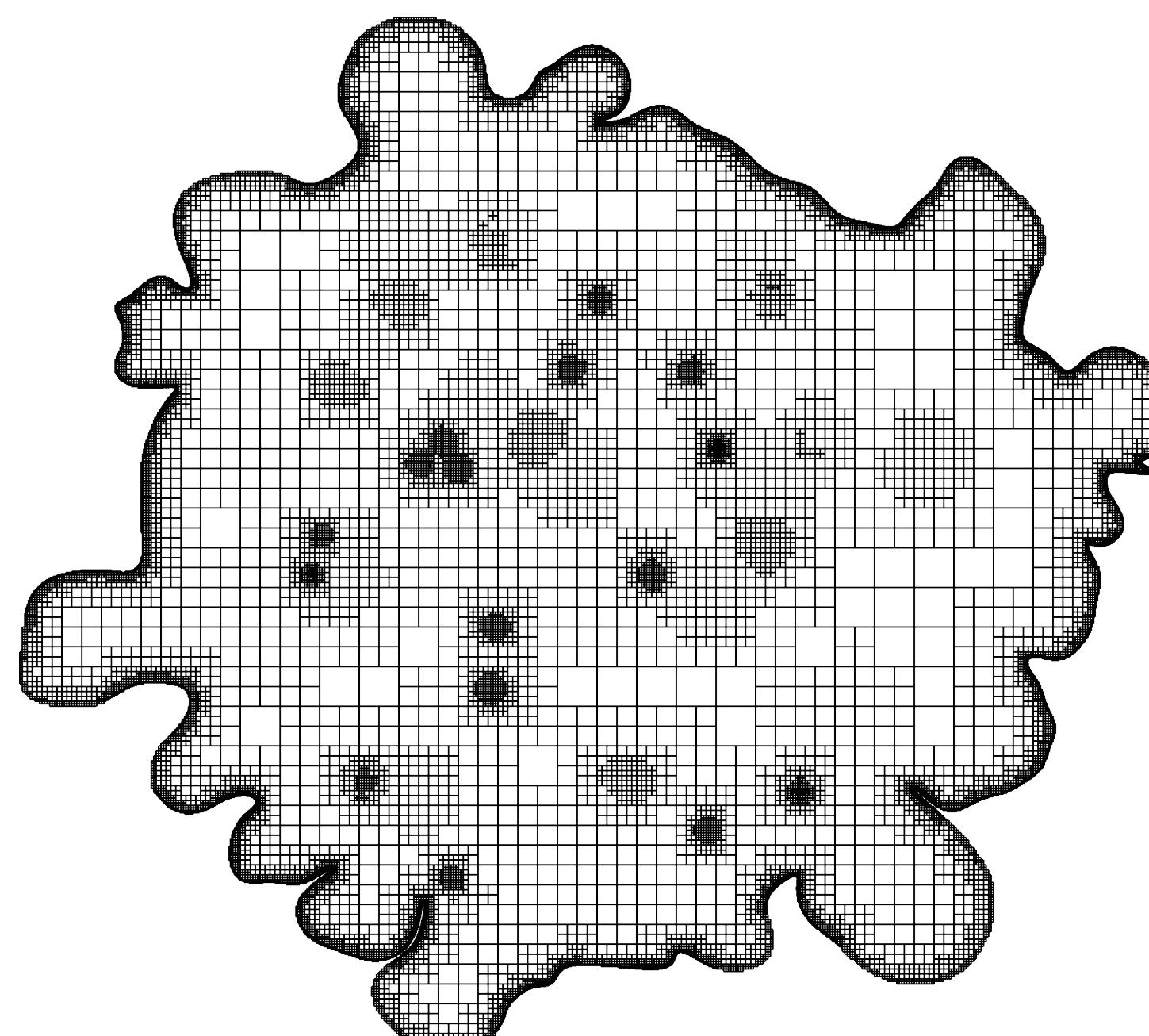
Example

Cell blebbing geometry extract from microscopy

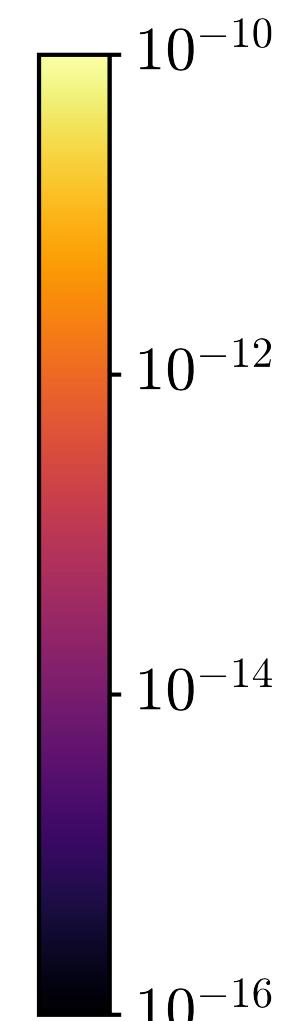
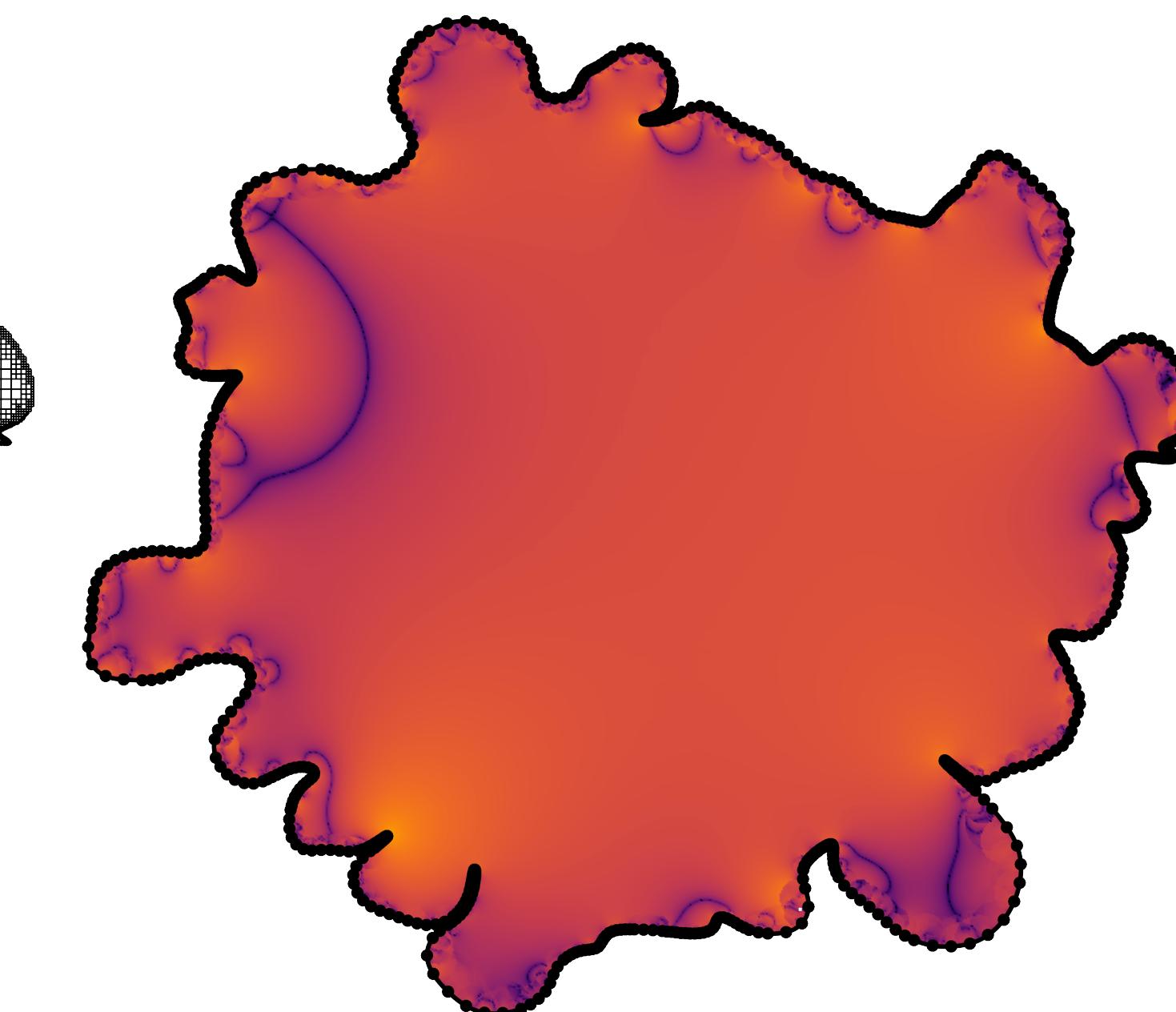
Inhomogeneity



Quadtree



Relative error



12M unknowns
16th order
30 seconds

Conclusion



David Stein



Alex Barnett

Conclusion

- Function extension can be challenging in multiscale geometries. Our scheme avoids extension.
- Particular solutions can be computed adaptively and efficiently using truncated volume potentials.
- Surface-conforming truncation idea designed to extend to 3D.



David Stein



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