tor, this leads to discretization errors in the ratio  $f_{B_s}/f_{B_d}$  of  $\mathcal{O}(\alpha_s \times (am_s - am_d)) \sim 1.2\%$  plus  $\mathcal{O}(\alpha_s^2 a \Lambda_{\rm QCD} \times (\widetilde{m}_s - \widetilde{m}_d)/\Lambda_{\rm QCD}) \sim 0.6\%$  plus  $\mathcal{O}(a^2 \Lambda_{\rm QCD}^2 \times (\widetilde{m}_s - \widetilde{m}_d)/\Lambda_{\rm QCD}) \sim 1.7\%$ . Although we do not improve the heavy-light four fermion operator used to compute the B-mixing matrix element, the operator does not have any tree-level  $\mathcal{O}(a)$  errors [38]. Thus the leading discretization errors in the ratio  $\xi$  from the four-fermion operator are of  $\mathcal{O}(\alpha_s \times (am_s - am_d)) \sim 1.2\%$  plus  $\mathcal{O}(\alpha_s a \Lambda_{\rm QCD} \times (\widetilde{m}_s - \widetilde{m}_d)/\Lambda_{\rm QCD}) \sim 1.9\%$  plus  $\mathcal{O}(a^2 \Lambda_{\rm QCD}^2 \times (\widetilde{m}_s - \widetilde{m}_d)/\Lambda_{\rm QCD}) \sim 1.7\%$ .

Adding the contributions from light-quark and gluon discretization errors, heavy-quark discretization errors, and discretization errors in the heavy-light current or four-fermion operator in quadrature, we estimate the error in  $f_{B_s}/f_{B_d}$  to be  $\sim 3.2\%$  and the error in  $\xi$  to be  $\sim 3.7\%$ .

## D. Heavy-light current and four-fermion operator renormalization

We compute the renormalization factors needed to match the lattice axial current and four-fermion operator to the continuum using one-loop lattice perturbation theory. This leaves a residual error due to the omission of higher-order terms. Based on power-counting, we estimate the truncation error in the coefficients to be of  $\mathcal{O}(\alpha_s^2)$ , which is the size of the first neglected term in the series. As we noted earlier in Sec. III C, however, the matching coefficient  $Z_{\Phi}$  cancels in the ratio of decay constants  $f_{B_s}/f_{B_d}$ ; thus its contribution to the error in  $f_{B_s}/f_{B_d}$  is zero. Although such an exact cancellation does not occur for the ratio of mixing matrix elements  $\xi$ , the error in  $\xi$  due to the uncertainty in the ratio of matching coefficients  $Z_{SP}/Z_{VA}$  is suppressed by the SU(3)-breaking factor  $(\tilde{m}_s - \tilde{m}_d)/\Lambda_{\rm QCD}$ . This is because, in the SU(3) limit, the four-fermion operator matrix elements would be equal in the numerator and denominator, so the error in  $\xi$  from the renormalization factor uncertainty would be zero. We therefore expect the error in  $f_{B_s}/f_{B_d}$  to be 0% and the error in  $\xi$  to be of  $\mathcal{O}(\alpha_s^2 \times (\tilde{m}_s - \tilde{m}_d)/\Lambda_{\rm QCD}) \sim 2.2\%$ . This error will decrease with the inclusion of data at a finer lattice spacing because the smaller coupling constant will improve the convergence of the series.

to  $am_{res}$ . These effects, however, are expected to be sub-percent level in the matrix elements [75], and therefore negligible in the SU(3)-breaking ratios.