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Caputo fractional standard map: Scaling invariance analyses Daniel Borin

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ABSTRACT

In this paper, we investigate the scaling invariance of survival probability in the Caputo fractional standard map of the order $1 < \alpha < 2$ considered on a cylinder. We consider relatively large values of the nonlinearity parameter K for which the map is chaotic. The survival probability has a short plateau followed by an exponential decay and is scaling invariant for all considered values of α and K.

1. Introduction

Differential equations are a fundamental tool for studying dynamical systems, enabling the investigation of the evolution of observables over time and capturing the long-term behavior of these systems. These equations elucidate the rates of change in system variables and typically involve integer-order derivatives. A generalization of this theory has led to the rise of the field of fractional differential equations (FDEs) [1], extending the conventional framework of differential equations by incorporating fractional derivatives. This extension has countless applications in various scientific fields, such as quantum mechanics [2–7], electrodynamics [8–11], economics [12–15], engineering [16–18], biology [19,20], and geology [21].

Fractional derivatives come in various types, each with its unique characteristics and applications. The most common among them include Riemann–Liouville, Caputo, Grünwald–Letnikov, Hadamard type, etc. [1]. As initial conditions are crucial in physical dynamics, adopting the Caputo derivative in mechanical systems is convenient. This choice is particularly useful since the initial conditions for specific fractional dynamical systems align with those for usual dynamical systems, allowing for a clearer physical interpretation of the system.

Sometimes, differential equations can be complicated. A way to simplify the problem is to reduce it to the study of simple discrete maps. This approach provides a simplified yet valuable perspective on the dynamics of the system, enhancing its suitability for simulations and computational studies. Maps, derived from equations of motion involving fractional derivatives, exhibit a noteworthy characteristic: long-term memory effects [22,23]. In this context, the present state of the

system is uniquely determined by all past states, adding an intriguing layer of complexity and historical dependence to the analysis.

Memory is a significant property of human beings and is the subject of extensive biophysical and psychological research. The possibility of investigating the dynamics of memory in neurons and proteins [24] becomes attractive for the study of fractional maps.

An essential and insightful measure when exploring the statistical characteristics of transport phenomena is the concept of survival probability. This probability quantifies the likelihood of a diffusing particle remaining within a specific region at a given point in time. In the context of fully chaotic systems, the behavior of survival probability exhibits a distinct pattern over time, characterized by exponential decay [25]. These properties have not been explored for Caputo fractional maps.

Motivated by the above descriptions, this paper aims to explore the dynamic characteristics of a Caputo Fractional Map obtained by replacing the integer derivative with the Caputo derivative in the equation of motion that originates the famous standard map (Chirikov–Taylor Map). More specifically, we vary control parameters and observe how they affect the survival probability, seeking invariance of scaling.

The layout of the study is as follows: The Caputo Fractional Standard Map is introduced in Section 2. In Section 3, a study of survival probability is presented, varying all control parameters and demonstrating that this observable is scaling invariant concerning the same parameters. Finally, a discussion is provided in Section 4.

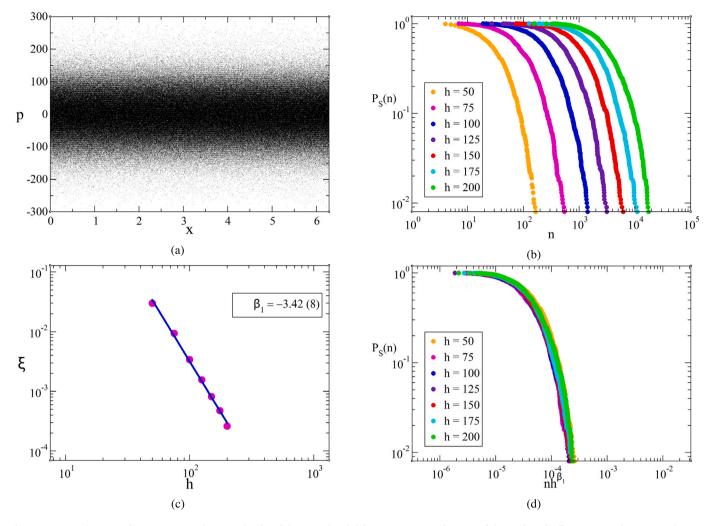


Fig. 1. (a) Space phase (x,p) for K=23.0602 and $\alpha=1.7$; (b) Plot of the survival probability curves $P_S(n)$ as functions of the number of collisions n, considering survival region height in the range $h \in [25,200]$ computed for $M=10^3$ orbits. The curves can be described by an exponential $P_S(n) \propto \exp(-\xi n)$; (c) The behavior of the critical exponent ξ as a function of h. As can be seen, the relationship between height h and critical exponent ξ is given by a power law, such that $\xi \propto h^{\beta_1}$, with an exponent $\beta_1 = -3.42(8)$; (d) We display the overlapping of the $P_S(n)$ curves in a single universal curve, indicating the scaling invariance of the survival probability for the variations of survival region height h.

2. The mapping

The standard map (also known as Chirikov-Taylor Map) [26,27]

$$\begin{cases} p_{n+1} = p_n - K \sin(x_n) \\ x_{n+1} = x_n + p_{n+1} \pmod{2\pi} \end{cases}$$
 (1)

is the solution of the equation

$$\ddot{x} = -K\sin(x)\sum_{n=0}^{\infty}\delta\left(\frac{t}{T} - n\right) \tag{2}$$

In this scenario, the perturbation involves a recurring sequence of delta-function-like pulses, or "kicks", with a periodicity of T and an amplitude represented by K.

We can generalize this model by substituting the second derivative with a Caputo fractional derivative. The left-sided Caputo fractional derivative [1,28] of order $\alpha > 0$ is defined by

$${}_{0}^{C}D_{t}^{\alpha}f(t) := \frac{1}{\Gamma(l-\alpha)} \int_{0}^{t} \frac{f^{(l)}(\tau)d\tau}{(t-\tau)^{\alpha-l+1}}.$$
 (3)

where $l \in \mathbb{N}$ such that $l-1 < \alpha \le l$. Note that when $\alpha = m \in \mathbb{N}$, the Caputo Fractional derivative of order α becomes an mth usual derivative, i.e.,

$${}_{0}^{C}D_{t}^{m}f(t) = \frac{d^{m}f(t)}{dt^{m}} \equiv f^{(m)}(t), \quad \text{for } m \in \mathbb{N}.$$

In this way, the fractional Caputo generalization of (2) is given by

$${}_{0}^{C}D_{t}^{\alpha}x(t) + K\sin(x(t))\sum_{t=0}^{\infty}\delta\left(\frac{t}{T} - n\right) = 0 \tag{4}$$

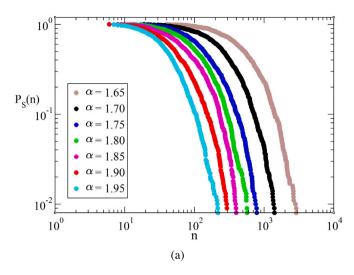
We can reduce the Caputo fractional differential equations to a Caputo Generalization nonlinear Volterra integral equations of second kind [29,30], such that x(t) is a solution of (4) if and only if it is a solution of the integral equation

$$x(t) = \sum_{k=0}^{l-1} \frac{x^{(k)}(0)}{k!} t^k + \frac{1}{\Gamma(\alpha)} \int_0^t (t-\tau)^{\alpha-1} \left[K \sin(x(\tau)) \sum_{j=0}^\infty \delta\left(\frac{\tau}{T} - j\right) \right] d\tau$$
(5)

where $l-1 < \alpha \le l$. Consider $1 < \alpha \le 2$, thus Eq. (5) is equivalent to the discrete map equations [22,23]

$$\begin{cases} p_{n+1} = p_0 - \frac{KT^{\alpha - 1}}{\Gamma(\alpha - 1)} \sum_{j=0}^{n} (n + 1 - j)^{\alpha - 2} \sin(x_j) \\ x_{n+1} = x_0 + p_0(n + 1)T - \frac{KT^{\alpha}}{\Gamma(\alpha)} \sum_{j=0}^{n} (n + 1 - j)^{\alpha - 1} \sin(x_j) \\ (\text{mod } 2\pi) \end{cases}$$
 (6)

for $1 < \alpha \le 2$. The Eq. (6) is named the Caputo Fractional Standard Map (CFSM). Note that if $\alpha = 2$, then Eq. (6) gives the usual map (1),



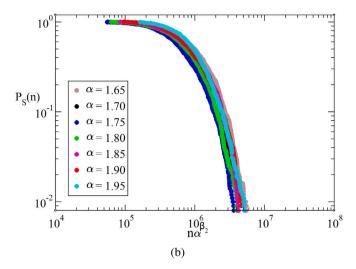


Fig. 2. (a) Plot of the survival probability curves $P_S(n)$ as functions of the number of collisions n, considering α in range the $\alpha \in [1.65, 1.95]$, K = 23.0602, h = 100 and computed for $M = 10^3$ orbits. The curves can be described by an exponential $P_S(n) \propto \exp(-\xi n)$; (b) We display the overlapping of the $P_S(n)$ curves in a single universal curve, indicating the scaling invariance of the survival probability concerning variations of α .

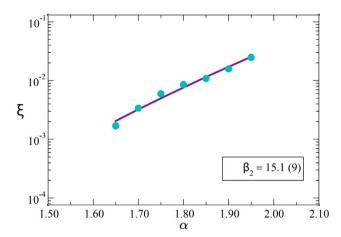


Fig. 3. The behavior of the critical exponent ξ as a function of α is depicted. As can be seen, the relationship between α and the critical exponent ξ is governed by a power law, such that $\xi \propto \alpha^{\beta_2}$, with an exponent $\beta_2 = 15.1(9)$.

making the CFSM a generalization of the standard map. The CFSM has no periodicity in p and can be considered on a cylinder.

The CFSM may generate periodic sinks, attracting slow-diverging trajectories, attracting accelerator mode trajectories, chaotic attractors, a cascade of bifurcations, and an inverse cascade of bifurcations type of attracting trajectories have been reported in several studies [31–37].

Self-similarity, distribution of exit times, and survival probability in the regular standard map were investigated in [38]. However, the scattering properties of the CFSM map have not been explored yet. Thus, in the next section, we undertake this task and characterize the escape of orbits from the phase space of the CFSM.

3. Scattering properties: scaling invariance of survival probability

The concept of scaling invariance is a well-known idea across various scientific and mathematical domains. It delves into the notion that certain phenomena can remain unchanged even when their parameters undergo rescaling or transformation. Essentially, if a system exhibits scaling invariance, its expected behaviors stay consistent regardless of scale. This characteristic is observed in diverse systems such as chaotic dynamics [25], human dynamics [39] and social networks [40].

The main goal of this section is the study of the survival probability of a trajectory remaining inside a specific region in the space phase, considering the CFSM Map proposed in the previous section. More specifically, our focus is on in investigating whether such behavior exhibits scale invariance when subjected to certain alterations of the mapping parameters.

To begin the analysis, it is essential to define the object of study: the measure of the probability of survival, denoted as $P_S(n)$ [41]. This probability refers to the possibility of a given particle remaining inside of a determined region after n interactions. Numerically, this measure is computed as [25]

$$P_S(n) = \frac{1}{M} \sum_{i=1}^n H_S(M, j), \quad \text{for } n \in \mathbb{N} < N,$$
 (7)

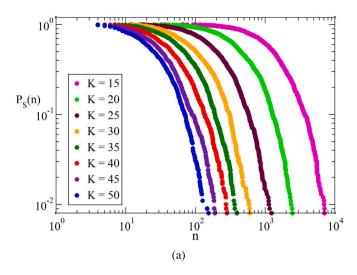
where $H_S(M,j)$ denotes the histogram associated with M distinct particles released in j collisions. To calculate this observable, an ensemble of M distinct orbits is considered, with initial conditions set as $x_0=0$ and random p_0 uniformly distributed in the interval $(0,2\pi/100)$. Each orbit evolves according to CFSM (6) until the trajectory escapes a rectangular region characterized by $p \leq |h|$. Thus the number of interactions up to that instant is recorded, and then a new particle in the system is initialized. This a procedure is repeated several times until the entire ensemble of initial conditions is exhausted. Here, it is important to mention that in our approach, each particles has the opportunity to evolve within the system for up to $N=10^6$ interactions if it has not escaped before.

With the incorporation of this new characterization of the dynamics, which is the survival region, three parameters, namely K, α , and h, influence the study of scattering properties. The goal from now on is to examine the influence of these parameters on the survival probability and determine whether it exhibits the property of scale invariance.

To avoid the cascade of bifurcation-type trajectories (CBTT) [42], which exists only in fractional dynamical systems, larger values of K will be considered, ensuring that the dynamics are essentially chaotic. The influence of these new types of attractors on the scattering properties will be investigated in future works.

Starting by keeping K=23.0602 and $\alpha=1.7$ constants, the phase space of (6) is illustrated in Fig. 1(a) and the survival probability $P_S(n)$ in the function of the number of interactions n considering different sizes for the survival region height h, as displayed in Fig. 1(b). One can note that for medium and long values of interactions P_S obeys an exponential decay as

$$P_S(n) \propto e^{-\xi n}$$
. (8)



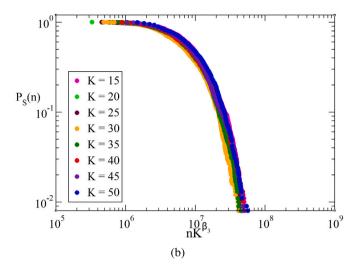


Fig. 4. (a) Plot of the survival probability curves $P_S(n)$ as functions of the number of collisions n, considering K in range $K \in [15,50]$, $\alpha = 1.7$, h = 100 and computed for $M = 10^3$ orbits. The curves can be described by an exponential $P_S(n) \propto \exp(-\xi n)$; (b) We display the overlapping of the $P_S(n)$ curves in a single universal curve, indicating the scaling invariance of the survival probability concerning variations of K.

Furthermore, the behavior of P_S depends on the value of h assumed such that, from a logical point of view, the higher the value of h, the slower the decay of P occurs, suggesting a relationship between h and ξ which is evident in Fig. 1(c), showing a power law function

$$\xi \propto h^{\beta_1}$$
,

 $\xi \propto \alpha^{\beta_2}$,

where $\beta_1 = -3.42(8)$, (·) denotes the error of the previous digit. The knowledge of β_1 allows us to re-scale the horizontal axis as $n \to nh^{\beta_1}$ producing an overlap of the curves plotted in Fig. 1(b) into a single and universal plot indicating that the exponential decay rate is scaling invariant with size of survival regions, as shown in Fig. 1(d).

Now, let us vary α , keeping K=23.0602 and h=10 as constants. As one can notice in Fig. 2(a), the shape presented for the decay of $P_S(n)$ for several values of α seems to be similar. This observation suggests a potential scaling invariance for this measure. To investigate this possibility, we examined $P_S(n)$ under the following scaling hypotheses:

where β_2 represents a critical exponent, which can be found through the analysis of the numerical fitting of ξ *vs.* α , as shown in Fig. 3, such that it is given by $\beta = 15.1(9)$. Thus, knowing the value of β , one can rescale the horizontal axis by transforming $n \to n\alpha^{\beta_2}$. This transformation results in overlapping curves as depicted in Fig. 2(a), converging into a unified and universal plot in Fig. 2(b). This outcome signifies the scaling invariance for $P_S(n)$ concerning the parameter α .

Finally, after the preliminary results indicated that the survivor probability is scaling invariant concerning h and α , a question is raised: Is the behavior of the survival probability $P_S(n)$ scale invariant due to variations of K?. To answer this, the analysis performed in Figs. 1(b) and 3 is repeated, considering the following scaling hypotheses:

$$\xi \propto K^{\beta_3}$$
.

In Fig. 5, it is found that the value of the critical exponent is $\beta_3 = 3.27(9)$. Thus, the curves in Fig. 4(a) overlap in a very good fashion after an appropriate scaling transformation, as shown in Fig. 4(b), confirming that the mapping presents scale invariance about K and answering the question proposed previously.

4. Discussion

In this paper, some dynamical and statistical properties of a fractional map were studied. Initially, a mapping was derived from a fractional differential equation by considering the Caputo derivative

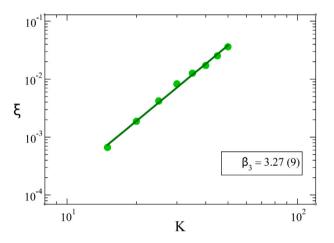


Fig. 5. The behavior of the critical exponent ξ as a function of K is depicted. As can be seen, the relationship between α and the critical exponent ξ is governed by a power law, such that $\xi \propto K^{\beta_3}$, with an exponent $\beta_3 = 3.27(9)$.

in the ODE that originates the Chirikov map. Thus, using an escape formalism, the behavior of orbits in the previously obtained mapping, referred to as the Caputo Fractional Standard Map (CFSM), was investigated. This map is parameterized by the control parameters K and α . In the study of survival probability, a new control parameter h is considered, and this observable now depends on (K, α, h) . Varying these parameters, $P_s(n)$ presents an initial plateau for short times, followed by an exponential decay for medium and long times, as illustrated in Figs. 1(b), 2(a), and 4(a). The behavior of survival probability allows us to propose that it is described by $P_S(n) \propto e^{-\gamma}$, where γ is a critical exponent that depends on the values of the assumed control parameters, as displayed in Figs. 1(c), 3, and 5.

The overlap of the curves shown in Figs. 1(d), 2(b) and 4(b) confirms a scaling invariance observed for Survival probability when studied in the context of Caputo Fractional Standard Mapping (CFSM), described in Eq. (6), for all control parameters (α , K and h).

CRediT authorship contribution statement

Daniel Borin: Writing – review & editing, Writing – original draft, Validation, Software, Methodology, Investigation, Formal analysis, Data curation, Conceptualization.

Declaration of competing interest

The authors declare the following financial interests/personal relationships which may be considered as potential competing interests: Daniel Borin reports financial support was provided by State of Sao Paulo Research Foundation. If there are other authors, they declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

Data availability

Data will be made available on request.

Acknowledgments

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References

- Kilbas AA, Srivastava HM, Trujillo JJ. Theory and applications of fractional differential equations. Amsterdam, The Netherlands: Elsevier; 2006.
- [2] Liu ZY, Xia TC, Wang TT. Fractional-order quantum kicked top map and its discrete dynamic behaviors. Chaos 2023;33:013133. http://dx.doi.org/10.1063/ 5.0131776.
- [3] Liu ZY, Xia TC, Wang TT. Fractional quantum logistic map and its application in image encryption technology. Fractals 2021;29(07):2150196.
- [4] Atman KG, Şirin H. Nonlocal phenomena in quantum mechanics with fractional calculus. Rep Math Phys 2020;86(2):263–70. http://dx.doi.org/10.1016/S0034-4877(20)30075-6.
- [5] Laskin N. Time fractional quantum mechanics. Chaos Solitons Fractals 2017;102:16–28. http://dx.doi.org/10.1016/j.chaos.2017.04.010.
- [6] Lim S. Fractional derivative quantum fields at positive temperature. Phys A 2006;363(2). http://dx.doi.org/10.1016/j.physa.2005.08.005.
- [7] Laskin N. Fractional Schrodinger equation. Phys Rev E 2002;66:056108. http://dx.doi.org/10.1103/PhysRevE.66.056108.
- [8] Stefanski TP, Gulgowski J. Signal propagation in electromagnetic media described by fractional-order models. Commun Nonlinear Sci Numer Simul 2020;82. http://dx.doi.org/10.1016/j.cnsns.2019.105029.
- [9] Stefanski TP, Gulgowski J. Electromagnetic-based derivation of fractional-order circuit theory. Commun Nonlinear Sci Numer Simul 2019;79. http://dx.doi.org/ 10.1016/j.cnsns.2019.104897.
- [10] Sikora R, Pawlowski S. Fractional derivatives and the laws of electrical engineering. Compel 2018;37(4):1384–91. http://dx.doi.org/10.1108/COMPEL-08-2017-0347.
- [11] Moreles MA, Lainez R. Mathematical modelling of fractional order circuit elements and bioimpedance applications. Commun Nonlinear Sci Numer Simul 2017;46:81–8. http://dx.doi.org/10.1016/j.cnsns.2016.10.020.
- [12] Tarasov VE, Tarasova VV. Economic dynamics with memory: fractional calculus approach. Berlin: De Gruyter; 2021, http://dx.doi.org/10.1515/9783110627459.
- [13] Tarasov VE. Fractional econophysics: Market price dynamics with memory effects. Physica A 2020. http://dx.doi.org/10.1016/j.physa.2020.124865.
- [14] Tarasov VE. Mathematical economics: Application of fractional calculus. Mathematics 2020;8(5):660. http://dx.doi.org/10.3390/math8050660.
- [15] Tarasov VE. On history of mathematical economics: application of fractional calculus. Mathematics 2019;7(6):509. http://dx.doi.org/10.3390/math7060509.
- [16] Chen CK, Li YC. Machine chattering identification based on the fractionalorder chaotic synchronization dynamic error. Int J Adv Manuf Technol 2019;100(1):907–15. http://dx.doi.org/10.1007/s00170-018-2774-3.
- [17] Yau HT, Wang CC, Hsieh CT, Wu SY. Fractional order Sprott chaos synchronisation-based real-time extension power quality detection method. IET Gener Transm Distrib 2015;9(16):2775–81. http://dx.doi.org/10.1049/iet-gtd. 2015.0500.
- [18] Wang MH, Yau HT. New power quality analysis method based on chaos synchronization and extension neural network. Energies 2014;7(10):6340–57. http://dx.doi.org/10.3390/en7106340.

- [19] Majee S, Barman S, Khatua A, Kar TK, Jana S. The impact of media awareness on a fractional-order SEIR epidemic model with optimal treatment and vaccination. Eur Phys J Spec Top 2023. http://dx.doi.org/10.1140/epjs/s11734-023-00910-9.
- [20] Jan A, Boulaaras S, Abdullah FA, Jan R. Dynamical analysis, infections in plants, and preventive policies utilizing the theory of fractional calculus. Eur Phys J Spec Top 2023. http://dx.doi.org/10.1140/epjs/s11734-023-00926-1.
- [21] Weiss CJ, van Bloemen Waanders BG, Antil H. Fractional operators applied to geophysical electromagnetics. Geophys J Int 2020;220(2):1242–59. http://dx.doi. org/10.1093/gji/ggz516.
- [22] Tarasov VE. Differential equations with fractional derivative and universal map with memory. J Phys A 2009;42:465102. http://dx.doi.org/10.1088/1751-8113/ 42/46/465102.
- [23] Tarasov VE, Zaslavsky GM. Fractional equations of kicked systems and discrete maps. J Phys A 2008;41(43):435101. http://dx.doi.org/10.1088/1751-8113/42/ 46/465102.
- [24] Min W, Luo G, Cherayil BJ, Kou SC, Xie XS. Observation of a power-law memory kernel for fluctuations within a single protein molecule. Phys Rev Lett 2005;94:198302.
- [25] Borin D, Livoratti ALP, Leonel ED. An investigation of the survival probability for chaotic diffusion in a family of discrete Hamiltonian mappings. Chaos Solitons Fractals 2023;175:113965. http://dx.doi.org/10.1016/j.chaos.2023.113965.
- [26] Chirikov BV. Research concerning the theory of nonlinear resonance and stochasticity. In: Preprint n 267. Novosibirsk: Institute of Nuclear Physics; 1969, (in Russian) [Engl. Transl. CERN Trans. 71-40, Geneva, (1971) Translated by A.T.Sanders].
- [27] Chirikov BV. A universal instability of many-dimensional oscillator systems. Phys Rep 1979;52(263). http://dx.doi.org/10.1016/0370-1573(79)90023-1.
- [28] Miller KS, Ross B. An introduction to the fractional calculus and fractional differential equations. New York: Wiley; 1993.
- [29] Kilbas AA, Marzan SA. Nonlinear differential equations with the Caputo fractional derivative in the space of continuously differentiable functions. Differential Equations 2005;41(1):84–9. http://dx.doi.org/10.1007/s10625-005-0137-y.
- [30] Kilbas AA, Marzan SA. Cauchy problem for differential equation with Caputo derivative. Fract Calc Appl Anal 2004;7(3):297–321.
- [31] Edelman M. Fractional standard map: Riemann–Liouville vs. Caputo. Commun Nonlinear Sci Numer Simul 2011;16:4573–80. http://dx.doi.org/10.1016/j.cnsns. 2011.02.007.
- [32] Edelman M. Fractional maps and fractional attractors. Part I: a-families of maps. Discontin Nonlinearity Complex 2013;1:305–24. http://dx.doi.org/10.5890/DNC. 2015.11.003.
- [33] Edelman M, Taieb LA. New types of solutions of non-linear fractional differential equations. In: Almeida A, Castro L, Speck F-O, editors. Advances in harmonic analysis and operator theory. Operator theory: advances and applications, Basel: Springer; 2013, p. 139–55, arXiv:1211.4012.
- [34] Edelman M. Fractional maps as maps with power-law memory. In: Afraimovich V, et al., editors. Nonlinear dynamics and complexity, nonlinear systems and complexity. vol. 8, New York: Springer; 2014, p. 79–120, arXiv: 1306.6361.
- [35] Edelman M. Universality in systems with power-law memory and fractional dynamics. In: Edelman M, Macau E, Sanjuan MAF, editors. Chaotic, fractional, and complex dynamics: new insights and perspectives. Understanding complex systems, Springer; 2018, p. 147–71, eBook.
- [36] Edelman M. Dynamics of nonlinear systems with power-law memory. In: Tarasov VE, editor. Handbook of fractional calculus with applications. Applications in physics, vol. 4, Berlin: De Gruyter; 2019, p. 103–32.
- [37] Edelman M, Helman AB. Asymptotic cycles in fractional maps of arbitrary positive orders. Fract Calc Appl Anal 2022. http://dx.doi.org/10.1007/s13540-021-00008-w.
- [38] Zaslavsky GM, Edelman M, Niyazov BA. Self-similarity, renormalization, and phase space nonuniformity of Hamiltonian chaotic dynamics. Chaos 1997;7:159–81. http://dx.doi.org/10.1063/1.166252.
- [39] Brockmann D, Hufnagel L, Geisel T. The scaling laws of human travel. Nature 2006;439:462–5. http://dx.doi.org/10.1038/nature04292.
- [40] Oliveira DFM, Chan KS, Leonel ED. Scaling invariance in a social network with limited attention and innovation. Phys Lett A 2018;382(47):3376–80. http: //dx.doi.org/10.1016/j.physleta.2018.09.034.
- [41] Mendez-Bermudez JA, Peralta-Martinez K, Sigarreta JM, Leonel ED. Leaking from the phase space of the Riemann–Liouville fractional standard map. Chaos Solitons Fractals 2023;172:113532.
- [42] Edelman M. Universal fractional map and cascade of bifurcations type attractors. Chaos 2013;23:033127. http://dx.doi.org/10.1063/1.4819165.