

# Distribution of First Hitting Time

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## 1 Derivation for the Laplace transform of first hitting time for GBM

Let  $dS = \alpha S dt + \sigma S dW_t$  and  $\tau$  be a stopping time such that  $S_\tau = m$ . By Dynkin's formula,

$$\mathbb{E}[e^{-u\tau} f(S_\tau)] = f(S_0) + \mathbb{E} \left[ \int_0^\tau \left( f_S \alpha S + \frac{1}{2} f_{SS} \sigma^2 S^2 - u f \right) e^{-us} ds \right]$$

If  $f(x)$  satisfies

$$f_x \alpha x + \frac{1}{2} f_{xx} \sigma^2 x^2 - u f = -u \quad (1)$$

Then Dynkin's formula implies that

$$\mathbb{E}[e^{-u\tau} f(S_\tau)] = f(S_0) + \mathbb{E} [e^{-u\tau}] - 1$$

Which can be solved for  $\mathbb{E} [e^{-u\tau}]$ , the Laplace transform of the density of  $\tau$ .

One boundary condition is determined: if  $S > m$  then  $\lim_{S \rightarrow \infty} f(S) = 0$ , and if  $S < m$  then  $\lim_{S \rightarrow 0} f(S) = 0$ . We have some flexibility over other boundary conditions, but  $f(S_\tau) = f(m) = 0$  is an obvious choice. If  $f(m) = 0$  is chosen, then  $\mathbb{E}[e^{-u\tau} f(S_\tau)] = 0$ , and so

$$f(S_0) = 1 - \mathbb{E} [e^{-u\tau}]$$

Or, rearranging,

$$\mathbb{E} [e^{-u\tau}] = 1 - f(S_0) \quad (2)$$

### 1.1 Directly solving $f$

To solve equation 1, introduce  $\gamma_+$  and  $\gamma_-$  where

$$\gamma = \left( \frac{\sigma^2}{2} - \alpha \pm \sqrt{\left( \alpha - \frac{\sigma^2}{2} \right)^2 + 2u\sigma^2} \right) / \sigma^2$$

The general solution to the complementary homogeneous equation can be written as  $c_1 S^{\gamma_+} + c_2 S^{\gamma_-}$ . If  $S < m$  then  $c_2$  must be zero or  $f$  tends towards positive or negative infinity as  $S \rightarrow 0$ . On the other hand, if  $S > m$  then  $c_1$  must be zero or  $f$  tends towards positive or negative infinity as  $S \rightarrow \infty$ . A particular solution to equation 1 is  $f^p(x) = 1$ , so a full solution is  $f(x) = cS^\gamma + 1$ , where  $c$  is the remaining unknown constant, and  $\gamma = \gamma_-$  if  $S > m$  and  $\gamma = \gamma_+$  if  $S < m$ . Selecting  $f(S_\tau) = f(m) = 0$  as the other boundary condition,

$$cm^\gamma + 1 = 0 \implies c = -m^{-\gamma}$$

Therefore the solution to equation 1 is as follows (making explicit the dependence on  $u$ ):

$$f(u; S) = \begin{cases} 1 - \left(\frac{S}{m}\right)^{\gamma_+}, & \text{if } S < m \\ 1 - \left(\frac{S}{m}\right)^{\gamma_-}, & \text{if } S > m \end{cases} \quad (3)$$

This can then be re-arranged using equation 2 to solve for the Laplace transform.

## 1.2 Solving for $g$

However, notice that  $g(x) = \mathbb{E}[e^{-u\tau}] = 1 - f(x)$  also solves an ODE. Since  $g_x = -f_x$ ,  $g_{xx} = -f_{xx}$  and substituting into equation 1,

$$\begin{aligned} -g_x \alpha x - \frac{1}{2} g_{xx} \sigma^2 x^2 - u(1 - g) &= -u \\ \implies g_x \alpha x + \frac{1}{2} g_{xx} \sigma^2 x^2 - ug &= 0 \end{aligned}$$

With boundary condition  $g(m) = 1$ . This, being a homogenous differential equation, is easier to solve than the one for  $f$ , and solving it directly yields the Laplace transform.

Either way, the Laplace transform can be found as follows:

$$\mathcal{L}(u; S) = \begin{cases} \left(\frac{S}{m}\right)^{\gamma_+}, & \text{if } S < m \\ \left(\frac{S}{m}\right)^{\gamma_-}, & \text{if } S > m \end{cases} \quad (4)$$

## 2 Characteristic function of general diffusion

Consider a diffusion satisfying

$$dX = \alpha(X, t)dt + \sigma(X, t)dW_t$$

Then using the approach in subsection 1.2, the solution  $g$  to

$$g_x \alpha(X, t) + \frac{\sigma^2}{2} g_{xx} - uig = 0$$

is the characteristic function of  $\tau$ . The distribution can then be recovered using inversion techniques. However, the ODE tends not to have analytical solutions, being non-linear and second order.

## 2.1 Numerical algorithm

Discretizing the ODE and solving it numerically has linear complexity. Solving the ODE in this way yields the approximate solution for  $x \in (0, m)$ , but only for a single value of  $u$ . Therefore the ODE must be solved multiple times in order to obtain the characteristic function for each discretized  $u$ . Fortunately once the matrix of values is obtained the density function can be recovered for each discretized  $x \in (0, m)$ .

The total complexity for computing the matrix is  $O(l * n)$

## 3 Characteristic Function Algorithms

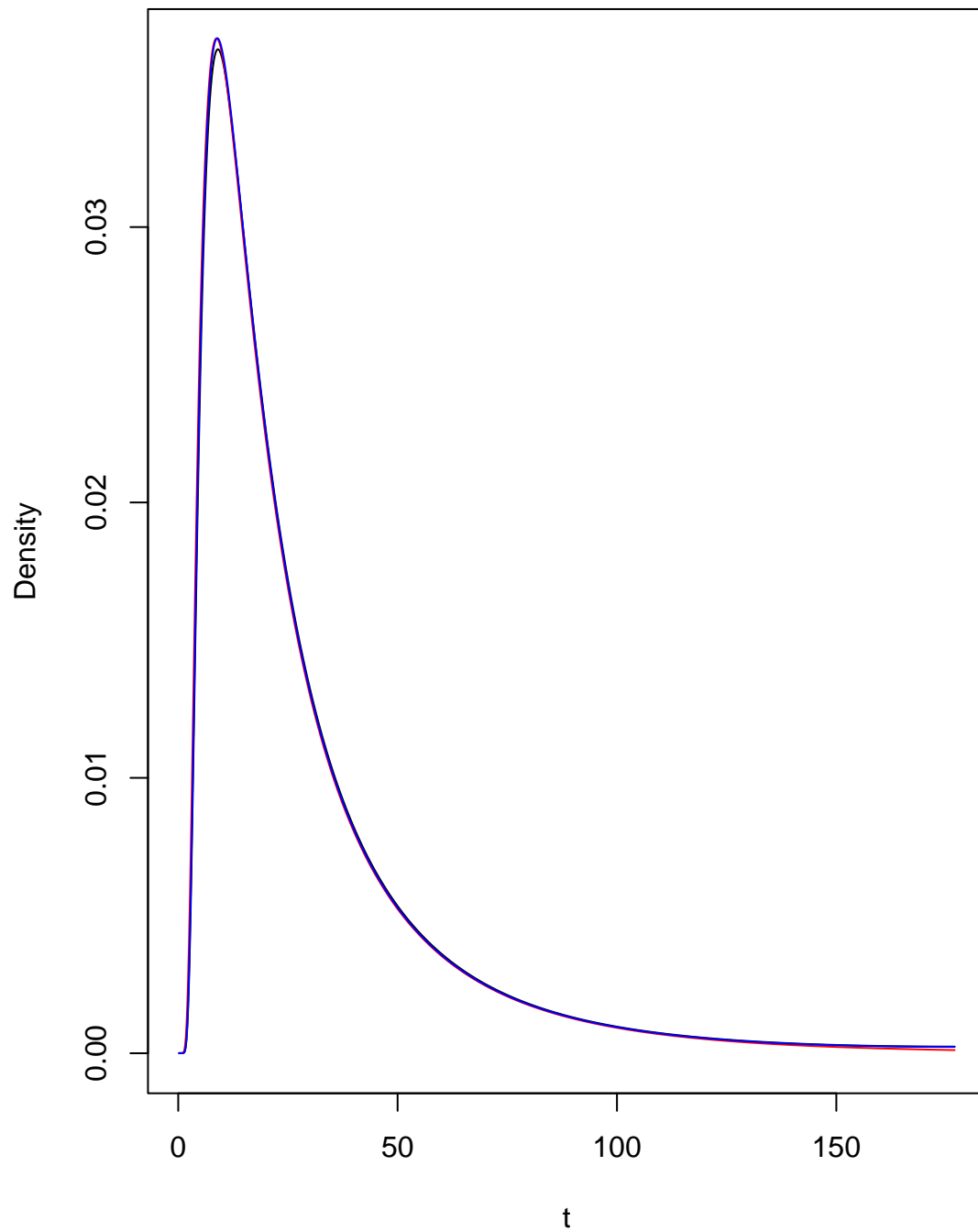
Let  $g(x)$  be the density function of  $\tau$ . By Fourier's theorem,

$$g(x) = \frac{1}{2\pi} \int e^{iux} f(u) du$$

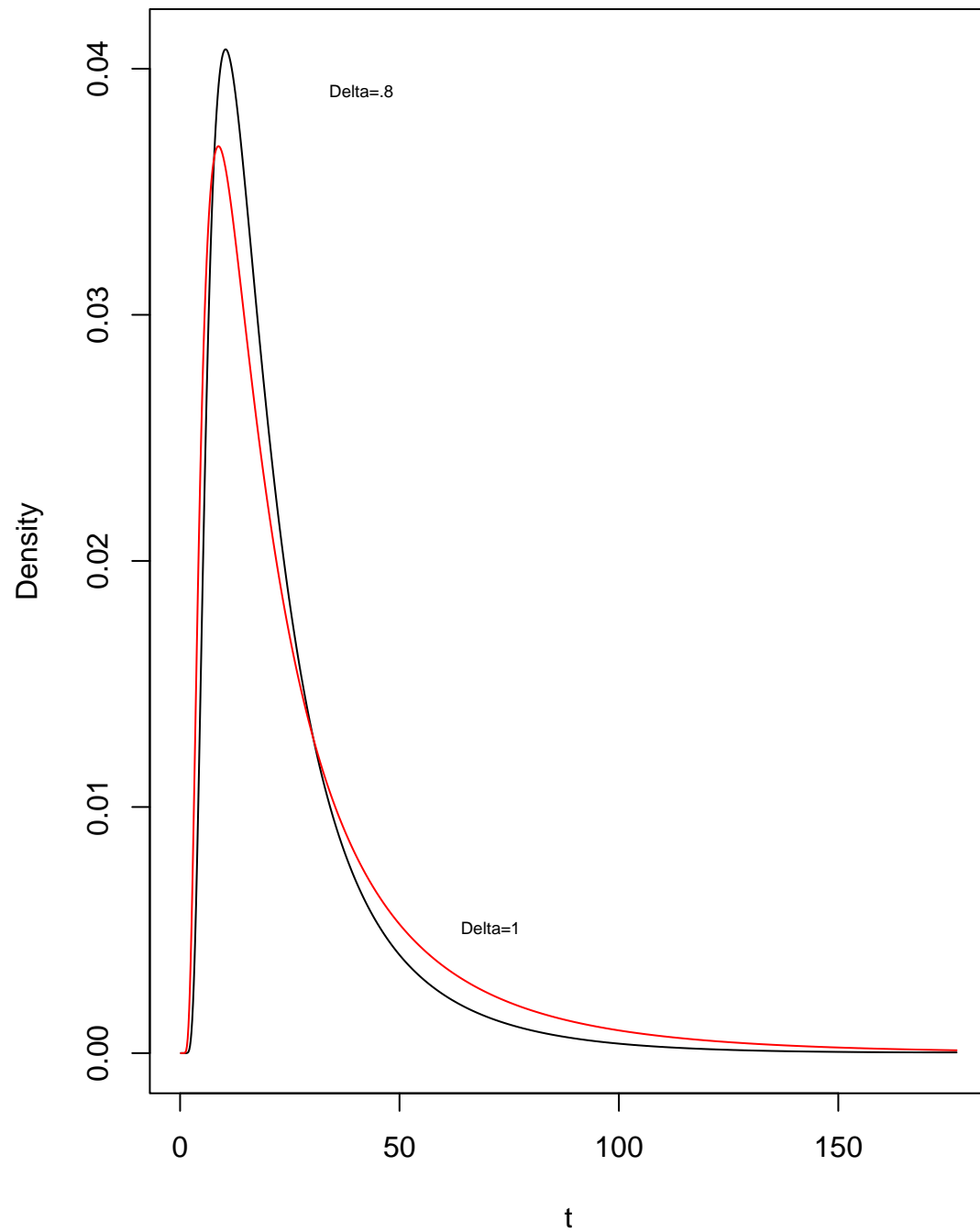
### 3.1 Cosine Method

Oosterlee and Fang (2008) proposed a novel method to price European options for a very diverse set of asset models. This method relies on inverting the characteristic function of the underlying log price; and prices options with many strikes simultaneously. Along the way they demonstrated the surprising accuracy of the method in computing the probability density from the characteristic function. Since the authors' main goal was to price options they spent minimal time on this result. Two key characteristics of this algorithm are the separate computations of the  $x$  and  $u$  domains and the exponential convergence of the algorithm. The method that the authors devised only requires  $O(nlk)$  operations to compute the density where  $k$  is the number of steps in the  $x$  domain and due to the excellent convergence properties  $m$  does not typically have to be larger than 128.

The following plot shows the actual pdf (red), the inverted analytical characteristic function (blue), and the numerically solved ODE for the characteristic function that is then inverted (black). The parameters are  $m = 5$ ,  $s = 1$ ,  $\sigma = .3$ ,  $\alpha = .1$ ,  $n = 512$ ,  $k = 1024$ ,  $l = 200$ .



The next plot is based off the CEV model with  $\delta = .8$ . Having no analytical solution to either the characteristic function or the density, only the numerical ODE solution is presented. The density of the distribution when  $\delta = 1$  is presented to show the effect that  $\delta$  has on the distribution of the first hitting time.



Since the algorithm computes the density for every node on the ODE, the following is the pdf for several values of  $S$ .

