

Champ de vitesse d'un filament tourbillon proche de la fibre

Le champ de vitesse créé par un filament tourbillon en un point \underline{x} de l'espace est donné par la formule :

$$\vec{V}(\underline{x}) = \frac{r}{4\pi} \int_0^s -\vec{\tau}(s') \wedge (\underline{x}(s') - \underline{x}_0) \frac{ds'}{|\underline{x}(s') - \underline{x}_0|^3}$$

Lorsqu'on est assez proche de la fibre pour utiliser les coordonnées locales à la fibre : $\underline{x} = \underline{x}(s) + r\vec{n}(\phi, s)$, l'expression précédente devient :

$$\begin{aligned} \vec{V}(r, \phi, s) &= \frac{r}{4\pi} \int_0^s -\vec{\tau}(s') \wedge (\underline{x}(s') - \underline{x}(s) - r\vec{n}(\phi, s)) \frac{ds'}{|\underline{x}(s') - \underline{x}(s) - r\vec{n}(\phi, s)|^3} \\ &= \frac{r}{4\pi} \int_{s^-}^{s^+} -\vec{\tau}(\bar{s} + s) \wedge (\underline{x}(\bar{s} + s) - \underline{x}(s) - r\vec{n}(\phi, s)) \frac{ds}{|\underline{x}(\bar{s} + s) - \underline{x}(s) - r\vec{n}(\phi, s)|^3} \\ &= \frac{r}{4\pi} \int_{s^-}^{s^+} K(\bar{s}, s, r, \phi) ds \end{aligned}$$

en posant $\bar{s} = s' - s$ et
 $s^+ = s^- = s/2$.

On voudrait obtenir un DA en fonction de r de ce champ de vitesse pour $r=0$.

Si on pose brutalement $r=0$ dans cette expression on obtient une intégrale divergente au voisinage de $\bar{s}=0$. Comme la perturbation est singulière, on va se servir de la DAR en faisant apparaître deux domaines distincts d'intégration (cf François).

On écrit donc $\vec{v}(r, \phi, z) = E_r^- + I_r^- + I_r^+ + E_r^+$ où :

$$E_r^- = \int_{-\eta}^{-\eta} K d\tilde{s} \quad I_r^- = \int_{-\eta}^0 K d\tilde{s} \quad I_r^+ = \int_0^\eta K d\tilde{s} \quad E_r^+ = \int_\eta^{s+} K d\tilde{s}$$

où l'on choisit de prendre

$$1 \gg \eta \gg r$$

On pose $E_r = E_r^- + E_r^+$ et $I_r = I_r^- + I_r^+$

on fait le changement de variable $\tilde{s} = \tilde{s}/r$ dans les intégrals intérieures I_r^- et I_r^+ :

$$\begin{aligned} I_r^+ &= r \int_0^{r/\eta} - \frac{\vec{e}^1(r\tilde{s}+r) \wedge (X(r\tilde{s}+r) - X(r) - r\vec{n}(\phi, r))}{|X(r\tilde{s}+r) - X(r) - r\vec{n}(\phi, r)|^3} d\tilde{s} \\ &= r \int_0^{r/\eta} L(r, \tilde{s}, r, \phi) d\tilde{s} \\ I_r^- &= r \int_{-\eta/r}^0 L(r, \tilde{s}, r, \phi) d\tilde{s}. \end{aligned}$$

I Le problème extérieur :

① limite extérieure de $K(\tilde{s}, s, r, \phi)$:

C'est la limite : $\lim_{r \rightarrow 0} \overline{K(\tilde{s}, s, r, \phi)} = K_p$

$$\begin{aligned} &\bullet [X(\tilde{s}+r) - X(r) - r\vec{n}(\phi, r)] \cdot [X(\tilde{s}+r) - X(r) - r\vec{n}(\phi, r)] \\ &= |X(\tilde{s}+r) - X(r)|^2 - 2r\vec{n} \cdot (X(\tilde{s}+r) - X(r)) + r^2 \\ &= \left\{ 1 - 2r\vec{n} \cdot \frac{(X(\tilde{s}+r) - X(r))}{|X(\tilde{s}+r) - X(r)|^2} + \frac{r^2}{|X(\tilde{s}+r) - X(r)|^2} \right\} |X - X(r)|^2 \end{aligned}$$

Donc

$$|X(\tilde{s}+r) - X(r) - r\vec{n}(\phi, r)|^{3/2} = \left\{ 1 + 3r\vec{n} \cdot \frac{(X(\tilde{s}+r) - X(r))}{|X(\tilde{s}+r) - X(r)|^2} \right\} |X(\tilde{s}+r) - X(r)|^3 + O(r^2)$$

$$-\vec{C}(\bar{s}+s) \wedge \left\{ (X(\bar{s}+s) - X(s)) - r \vec{n}(\psi, s) \right\}$$

$$= -\vec{C}(\bar{s}+s) \wedge (X(\bar{s}+s) - X(s)) - r \vec{n}(\psi, s) \wedge \vec{C}(\bar{s}+s)$$

D'où

$$\begin{aligned} K_\ell &= -\frac{\vec{C}(\bar{s}+s) \wedge (X(\bar{s}+s) - X(s))}{|X(\bar{s}+s) - X(s)|^3} + \frac{r \vec{n}(\psi, s) \wedge -\vec{C}(\bar{s}+s)}{|X(\bar{s}+s) - X(s)|^3} \\ &\quad - 3r \vec{n}(\psi, s) \cdot \frac{(X(\bar{s}+s) - X(s))}{|X(\bar{s}+s) - X(s)|^5} \vec{C}(\bar{s}+s) \wedge (X(\bar{s}+s) - X(s)) + O(s^2) \\ &= f(\bar{s}, s) + g(\bar{s}, s, r, \psi) + h(\bar{s}, s, r, \psi) + O(s^2). \end{aligned}$$

$$\begin{aligned} \text{ou } f(\bar{s}, s) &= -\frac{\vec{C}(\bar{s}+s) \wedge (X(\bar{s}+s) - X(s))}{|X(\bar{s}+s) - X(s)|^3} \quad g(\bar{s}, s, r, \psi) = \frac{r \vec{n} \wedge -\vec{C}(\bar{s}+s)}{|X(\bar{s}+s) - X(s)|^3} \\ &= r \vec{n}(s) \wedge g_1(\bar{s}, s) \\ h(\bar{s}, s, r, \psi) &= -3r \vec{n}(\psi, s) \cdot \frac{(X(\bar{s}+s) - X(s))}{|X(\bar{s}+s) - X(s)|^5} \vec{C}(\bar{s}+s) \wedge (X(\bar{s}+s) - X(s)). \end{aligned}$$

② DA de E_ℓ pour r proche de 0 :

$$\begin{aligned} \text{Posons : } F^+ &= \int_{-s^-}^{s^+} f(\bar{s}, s) ds \quad F^- = \int_{-s^-}^{-s^+} f(\bar{s}, s) ds \quad \text{et } F = F^+ + F^- \\ G^+ &= \int_{-s^-}^{s^+} g(\bar{s}, s, r, \psi) ds \quad G^- = \int_{-s^-}^{-s^+} g(\bar{s}, s, r, \psi) ds \quad \text{et } G = G^+ + G^- \\ H^+ &= \int_{-s^-}^{s^+} h(\bar{s}, s, r, \psi) ds \quad H^- = \int_{-s^-}^{-s^+} h(\bar{s}, s, r, \psi) ds \quad \text{et } H = H^+ + H^- \end{aligned}$$

Comme $s^+ = s^-$, on laissera toute partie impaire de côté lors

du calcul de F^+, G^+, H^+ .

$$\text{On cherche } E_{1\ell} = \int_{-s^-}^{-s^+} K_\ell d\bar{s} + \int_{-s^-}^{s^+} K_\ell d\bar{s} = F + G + H + O(s^2)$$

③

⑨1) Détermination de F :

⑨11) Développement limité de f en $\bar{s} = 0$:

- $X(\bar{s}+\gamma) - X(\gamma) = X_{\gamma}(s) \bar{s} + X_{ss}(s) \frac{\bar{s}^2}{2} + X_{sss}(s) \frac{\bar{s}^3}{6} + O(\bar{s}^4)$
 $= \vec{C}'(s) \bar{s} + K(s) \vec{m}'(s) \frac{\bar{s}^2}{2} + (K_s \bar{s} + K \vec{T}^b - K^2 \vec{C}') \frac{\bar{s}^3}{6} + O(\bar{s}^4)$
 $= \bar{s} \left\{ \vec{C}'(s) + K \frac{\vec{m}'}{2} \bar{s} + (K_s \vec{m} + K \vec{T}^b - K^2 \vec{C}') \frac{\bar{s}^2}{6} + O(\bar{s}^3) \right\}$
- $|X(\bar{s}+\gamma) - X(\gamma)|^2 = \left(\bar{s} - \frac{K^2}{3} \bar{s}^3 \right)^2 + \left(K \frac{\bar{s}^2}{2} + K_s \bar{s}^3 \right)^2 + \left(\frac{K \bar{s}^3}{6} \right)^2 + O(\bar{s}^5)$
 $= \bar{s}^2 - \frac{K^2}{3} \bar{s}^4 + \frac{K^2}{6} \bar{s}^4 + O(\bar{s}^5)$
 $= \bar{s}^2 \left(1 - \bar{s}^2 \frac{K^2}{12} + O(\bar{s}^3) \right)$
 $|X(\bar{s}+\gamma) - X(\gamma)|^3 = |\bar{s}|^3 \left(1 + \frac{3}{2} \bar{s}^2 \frac{K^2}{12} + O(\bar{s}^3) \right)$.
- $-\vec{C}'(\bar{s}+\gamma) = -\vec{C}'(s) - \vec{C}'_s \bar{s} - \vec{C}_{ss} \frac{\bar{s}^2}{2} + O(\bar{s}^3)$
 $= -\vec{C}'(s) - K(s) \vec{m}' \bar{s} - (K_s \vec{m} + K \vec{T}^b - K^2 \vec{C}') \frac{\bar{s}^2}{2} + O(\bar{s}^3)$
- $- \vec{C}' \wedge (X(\bar{s}+\gamma) - X(\gamma))$
 $= \bar{s} \left\{ K \vec{B}' \bar{s} - \frac{K}{2} \vec{B}' \bar{s} - \vec{C}' \wedge (K_s \vec{m} + K \vec{T}^b - K^2 \vec{C}') \frac{\bar{s}^2}{6} - (K_s \vec{m} + K \vec{T}^b - K^2 \vec{C}') \Lambda \vec{m}_0 \frac{\bar{s}^2}{2} \right. \\ \left. + O(\bar{s}^3) \right\}$
 $= \bar{s}^2 \left\{ \frac{K}{2} \vec{B}' + \frac{3}{6} \underbrace{[-K_s \vec{B}' + K \vec{T}^b + 3 \vec{B}' \cdot K_s - 3 K \vec{T} \vec{m}']}_{2(K_s \vec{B}' - K \vec{T} \vec{m}')} + O(\bar{s}^3) \right\}$
 $= \bar{s}^2 \left\{ \frac{K}{2} \vec{B}' + \frac{1}{2} [K_s \vec{B}' - K \vec{T} \vec{m}'] + O(\bar{s}^3) \right\}$

$$\text{d'où } \frac{-\bar{\epsilon}(\bar{s}) \wedge (X(\bar{s}+\eta) - X(\bar{s}))}{|X(\bar{s}+\eta) - X(\bar{s})|^3} = \frac{1}{|\bar{s}|^3} \left\{ \frac{K}{2} \bar{b} + \frac{5}{3} [K_s \bar{b} - K_T \bar{m}] + O(\eta^2) \right\}$$

$$\boxed{-\frac{\bar{\epsilon}(\bar{s}) \wedge (X(\bar{s}+\eta) - X(\bar{s}))}{|X(\bar{s}+\eta) - X(\bar{s})|^3} = \frac{K}{2} \frac{\bar{b}}{|\bar{s}|} + O(1)}$$

Q12 Détermination de F :

$$F^+ = \sum_{n=1}^{s+} f(\bar{s}_n, s_n) d\bar{s} = \sum_{n=1}^{s+} \left[f(\bar{s}_n, s_n) - \frac{K}{2} \frac{\bar{b}}{|\bar{s}_n|} \right] d\bar{s} + \sum_{n=1}^{s+} \frac{K}{2} \frac{\bar{b}}{|\bar{s}_n|} d\bar{s} \quad \text{fonction paire}$$

$$F^- = \sum_{n=s-}^{-1} \left[f(\bar{s}_{n+1}, s_n) - \frac{K}{2} \frac{\bar{b}}{|\bar{s}_n|} \right] d\bar{s} + \sum_{n=s-}^{-1} \left[f(\bar{s}_n, s_n) - \frac{K}{2} \frac{\bar{b}}{|\bar{s}_n|} \right] d\bar{s} + 2 \sum_{n=1}^{s+} \frac{K}{2} \frac{\bar{b}}{|\bar{s}_n|} d\bar{s}$$

$$2 \sum_{n=1}^{s+} \frac{K}{2} \frac{\bar{b}}{|\bar{s}_n|} d\bar{s} = K [\ln s^+ - \ln s^-] \bar{b}$$

$$\sum_{n=s-}^{s+} \left[f(\bar{s}_n, s_n) - \frac{K}{2} \frac{\bar{b}}{|\bar{s}_n|} \right] d\bar{s} = \sum_{n=s-}^{s+} \left[f(\bar{s}_n, s_n) - \frac{K}{2} \frac{\bar{b}}{|\bar{s}_n|} \right] d\bar{s} + \eta \overset{\text{impair}}{\cancel{\bar{b}}} + O(n^2)$$

$$F = \sum_{n=s-}^{s+} \left[f(\bar{s}_n, s_n) - \frac{K}{2} \frac{\bar{b}}{|\bar{s}_n|} \right] d\bar{s} = K(s) \bar{b}(s) \ln \frac{M}{s^+} + O(n^2)$$

$$\text{soit } \bar{A} = \sum_{n=s-}^{s+} \left[f(\bar{s}_n, s_n) - \frac{K}{2} \frac{\bar{b}}{|\bar{s}_n|} \right] d\bar{s}$$

$$\text{mais } \boxed{F = \bar{A} - K(s) \bar{b}(s) \ln \frac{M}{s^+} + O(n^2)}$$

$$\boxed{|\ln n| \gg 1 \gg n \gg s \gg \eta}$$

$n \gg n^2 \gg \sqrt{n} \gg n \rightarrow$ on prend n tel que $\boxed{\sqrt{n} \gg n \gg s}$

Q2) Détermination de G :

$$G^+ = \pi \vec{r}(\phi, s) \wedge \int_{\eta}^{s^+} g_1(\bar{s}, s) d\bar{s}$$

$$G^- = \pi \vec{r}(\phi, s) \wedge \int_{-s^-}^{-\eta} g_1(\bar{s}, s) d\bar{s}$$

$$G_1^+ = \int_{\eta}^{s^+} g_1(\bar{s}, s) d\bar{s} \quad G_1^- = \int_{-\eta}^{-s^-} g_1(\bar{s}, s) d\bar{s}.$$

$$G = G^- + G^+ = \pi \vec{r}(\phi, s) \wedge G_1$$

Q21) Développement limité de g_1 en $\bar{s} = 0$:

On a (cf Q21) :

$$-\vec{r}(\bar{s}+s) = -\vec{r}(s) - K(s)\vec{m}\bar{s} - (K_s \vec{m} + KT\vec{b} - K^2\vec{r})\frac{\bar{s}^2}{2} + O(\bar{s}^3)$$

$$|x(\bar{s}+s) - x(s)|^3 = |\bar{s}|^3 \left(1 + \frac{3}{2} \frac{\bar{s}^2 K^2}{72} + O(\bar{s}^3) \right)$$

D'où :

$$\frac{-\vec{r}(\bar{s}+s)}{|x(\bar{s}+s) - x(s)|^3} = \frac{1}{|\bar{s}|^3} \left[-\vec{r}(s) - K(s)\vec{m}\bar{s} - (K_s \vec{m} + KT\vec{b} - K^2\vec{r})\frac{\bar{s}^2}{2} - \frac{3}{2} \frac{K^2}{72} \vec{r}(s) \bar{s}^2 + O(\bar{s}^3) \right]$$

Q22) Détermination de G :

$$\begin{aligned} G_1^+ &= \int_{\eta}^{s^+} \left\{ -\frac{\vec{r}(s+\bar{s})}{|x(s+\bar{s}) - x(s)|^3} + \frac{1}{|\bar{s}|^3} \left[\vec{r}(s) + K(s)\vec{m}\bar{s} + (K_s \vec{m} + KT\vec{b} - \frac{3}{4} K^2\vec{r})\frac{\bar{s}^2}{2} \right] \right\} d\bar{s} \\ &\quad - \vec{r}(s) \left\{ \frac{1}{|\bar{s}|^3} \left[\frac{1}{B^3} - \frac{3}{8} K^2 \frac{1}{|\bar{s}|} \right] d\bar{s} - \vec{m} \int_{\eta}^{s^+} \underbrace{\left[K(s) \frac{\bar{s}}{|\bar{s}|^3} + \frac{K_s}{2} \frac{1}{|\bar{s}|} \right]}_{\text{impair}} d\bar{s} - \vec{b} \int_{\eta}^{s^+} \frac{K^2}{2} \frac{1}{|\bar{s}|} d\bar{s} \right\} \\ &= \int_{0}^{s^+} \left\{ -\frac{\vec{r}(s+\bar{s})}{|x(s+\bar{s}) - x(s)|^3} + \frac{1}{|\bar{s}|^3} \left[\vec{r}(s) + K(s)\vec{m}\bar{s} + (K_s \vec{m} + KT\vec{b} - \frac{3}{4} K^2\vec{r})\frac{\bar{s}^2}{2} \right] \right\} d\bar{s} + O(\eta) \\ &\quad - \vec{r}(s) \left\{ -\frac{1}{2} \left[\frac{1}{\bar{s}^2} \right]_{\eta}^{s^+} - \frac{3}{8} K^2 [\ln \bar{s}]_{\eta}^{s^+} \right\} - \vec{m} \frac{K_s}{2} [\ln \bar{s}]_{\eta}^{s^+} - \vec{b} \frac{K^2}{2} [\ln \bar{s}]_{\eta}^{s^+} \end{aligned}$$

$$G_1 = \int_{-\bar{s}}^{\bar{s}} \left\{ -\frac{\vec{v}'(s+\bar{s})}{|X(s+\bar{s}) - X(s)|^3} + \frac{1}{|\bar{s}|^3} [\vec{v}'(s) + k(s) \bar{s} \bar{s} + (k_s \bar{s} + k T \bar{b} - \frac{3}{4} k^2 \bar{v}') \frac{\bar{s}^2}{2}] \right\} ds$$

$$- \vec{v}'(s) \left\{ \frac{1}{\bar{s}^2} - \frac{1}{\bar{s}^2} + \frac{3}{4} k^2 \ln \frac{M}{\bar{s}^2} \right\} + \bar{s} k_s \ln \frac{M}{\bar{s}^2} + \bar{b}' k T \ln \frac{M}{\bar{s}^2} + O(\eta)$$

On a $\vec{v}' \wedge \bar{s} = \vec{\phi}'$ $\bar{s}' \wedge \bar{s} = \sin \phi \vec{v}'$ $\bar{b}' \wedge \bar{s} = -\vec{v}' \cos \phi$

puisque $\bar{s}' = \bar{s} \cos \phi + \bar{b}' \sin \phi$.

Soit $\boxed{\vec{B} = \bar{s}'(s) \wedge \int_{-\bar{s}}^{\bar{s}} \left\{ -\frac{\vec{v}'(s+\bar{s})}{|X(s+\bar{s}) - X(s)|^3} + \frac{1}{|\bar{s}|^3} [\vec{v}'(s) + k(s) \bar{s} \bar{s} + (k_s \bar{s} + k T \bar{b} - \frac{3}{4} k^2 \bar{v}') \frac{\bar{s}^2}{2}] \right\} ds}$

$$G = r \vec{B} + r \vec{\phi}' \left\{ \frac{1}{\bar{s}^2} - \frac{1}{\bar{s}^2} + \frac{3}{4} k^2 \ln \frac{M}{\bar{s}^2} \right\} - r \sin \phi \vec{v}' k_s \ln \frac{M}{\bar{s}^2} + r \cos \phi \vec{v}' k T \ln \frac{M}{\bar{s}^2} + O(\eta)$$

$$\boxed{G = r \vec{B} + r \vec{\phi}' \left\{ \frac{1}{\bar{s}^2} - \frac{1}{\bar{s}^2} + \frac{3}{4} k^2 \ln \frac{M}{\bar{s}^2} \right\} + r \vec{v}' \left\{ \sin \phi k_s + \cos \phi k T \right\} \ln \frac{M}{\bar{s}^2} + O(\eta)}$$

③ Détermination de H :

③① Développement limité de h en $\bar{s}=0$:

On a (ch 21) :

$$- \vec{v}'(s+\bar{s}) \wedge (X(s+\bar{s}) - X(s)) = \bar{s}^2 \left\{ \frac{k}{2} \bar{b}' + \frac{2}{3} [k_s \bar{b}' - k T \bar{s}] + O(\bar{s}^3) \right\}$$

$$X(s+\bar{s}) - X(s) = \bar{s} \left\{ \vec{v}'(s) + K \frac{\bar{s}}{2} \bar{s} + (k_s \bar{s} + k T \bar{b} - k^2 \bar{v}') \frac{\bar{s}^2}{6} + O(\bar{s}^3) \right\}$$

A partir du développement de $|X(s+\bar{s}) - X(s)|^2$ il vient :

$$|X(s+\bar{s}) - X(s)|^{-5} = |\bar{s}|^{-5} \left(1 + \frac{5}{2} \bar{s}^2 \frac{k^2}{12} + O(\bar{s}^3) \right)$$

Alors :

$$\begin{aligned}\vec{\pi}^1(s) \cdot (\vec{x}(s+\tau) - \vec{x}(s)) &= \tau \left\{ \frac{K}{2} \cos \phi + (K \sin \phi + K T \sin \phi) \frac{\tau^2}{6} + O(\tau^3) \right\} \\ &= \tau^2 \left\{ \frac{K}{2} \cos \phi + (K \sin \phi + K T \sin \phi) \frac{\tau}{6} + O(\tau^2) \right\}\end{aligned}$$

$$\vec{n} \cdot \vec{\pi}^1 = 0 \quad \vec{n} \cdot \vec{m} = \cos \phi \quad \vec{n} \cdot \vec{b} = \sin \phi \text{ puisque } \vec{n} = \vec{m} \cos \phi + \vec{b} \sin \phi.$$

$$\frac{\tau^2 + \tau^2}{|\vec{s}|^5} = \frac{1}{|\vec{s}|}$$

D'où :

$$-\vec{\pi}(s) \cdot \frac{(\vec{x}(s+\tau) - \vec{x}(s)) [\vec{\pi}^1(s+\tau) \wedge (\vec{x}(s+\tau) - \vec{x}(s))]}{|\vec{x}(s+\tau) - \vec{x}(s)|^5} = \frac{1}{|\vec{s}|} \left\{ \frac{K}{2} \vec{b} \frac{K}{2} \cos \phi \right\} + O\left(\frac{\tau}{|\vec{s}|}\right)$$

③ Détermination de H :

$$\begin{aligned}H^+ &= 3 \pi \int_{-s}^{s} -\frac{\vec{\pi}(s) \cdot (\vec{x}(s+\tau) - \vec{x}(s)) [\vec{\pi}^1(s+\tau) \wedge (\vec{x}(s+\tau) - \vec{x}(s))]}{|\vec{x}(s+\tau) - \vec{x}(s)|^5} d\tau \\ &= 3 \pi \int_{-s}^{s} \left\{ -\frac{\vec{\pi}(s) \cdot (\vec{x}(s+\tau) - \vec{x}(s)) [\vec{\pi}^1(s+\tau) \wedge (\vec{x}(s+\tau) - \vec{x}(s))] - \left(\frac{K}{2}\right)^2 \frac{\vec{b}}{|\vec{s}|} \cos \phi}{|\vec{x}(s+\tau) - \vec{x}(s)|^5} d\tau \right. \\ &\quad \left. - 3 \pi \left(\frac{K}{2}\right)^2 \vec{b} \cos \phi \left[\ln \frac{M}{s} \right] \right\}\end{aligned}$$

Soit $\vec{C} = \int_{-s}^{s} \left\{ \frac{\vec{\pi}(s) \cdot (\vec{x}(s+\tau) - \vec{x}(s)) [\vec{\pi}^1(s+\tau) \wedge (\vec{x}(s+\tau) - \vec{x}(s))]}{|\vec{x}(s+\tau) - \vec{x}(s)|^5} + \left(\frac{K}{2}\right)^2 \frac{\vec{b}}{|\vec{s}|} \cos \phi \right\} d\tau$

On a :

$$H = -3 \pi \vec{C} - 6 \pi \left(\frac{K}{2}\right)^2 \vec{b} \cos \phi \left[\ln \frac{M}{s} \right] + O(\pi)$$

④ Expression de Er, DA extérieur de Er en $r=0$:

$$Er = F + G + H + O(r^2)$$

II Le problème intérieur:

(1) Limite intérieure de $L(\tilde{\pi}, s, \alpha, \phi)$:

C'est la limite : $\lim_{n \rightarrow 0} \stackrel{\pi \text{ fixé}}{L}(\tilde{\pi}, s, \alpha, \phi) = L_0$

$$\bullet -\vec{r}'(x\tilde{\pi}+s) = -\vec{r}'(s) - k(s)\vec{m} \cdot x\tilde{\pi} - (k_s \vec{m} + kT\vec{b} - k^2 \vec{r}') \frac{x^2 \tilde{\pi}^2}{2} + O(s^3)$$

en se servant du développement de $\vec{r}'(\tilde{\pi}+s)$ en $\tilde{\pi}=0$ (cf 211).

$$\bullet X(x\tilde{\pi}+s) - X(s) = x\tilde{\pi} \left\{ \vec{r}'(s) + \frac{k}{2} \vec{m} \cdot x\tilde{\pi} + (k_s \vec{m} + kT\vec{b} - k^2 \vec{r}') \frac{x^2 \tilde{\pi}^2}{6} + O(s^3) \right\} \quad (\text{cf 211})$$

$$X(x\tilde{\pi}+s) - X(s) - x\tilde{\pi}(\phi, s) = x \left\{ \frac{1}{2} \vec{r}' - \vec{r}' + \frac{k}{2} \vec{m} \cdot \tilde{\pi}^2 s + (k_s \vec{m} + kT\vec{b} - k^2 \vec{r}') \frac{\tilde{\pi}^3}{6} s^2 + O(s^3) \right\}$$

$$\bullet \vec{r}' \wedge \vec{r} = \vec{\phi}' \quad \vec{m} \wedge \vec{r} = \sin \phi \vec{b} \quad \vec{b} \wedge \vec{r} = -\cos \phi \vec{r}'$$

puisque $\vec{r} = \cos \phi \vec{m} + \sin \phi \vec{b}$ et $\vec{\phi}' = -\vec{m} \sin \phi + \vec{b} \cos \phi$

$$\bullet -\vec{r}'(x\tilde{\pi}+s) \wedge (X(x\tilde{\pi}+s) - X(s) - x\tilde{\pi}(\phi, s))$$

$$= -x \left\{ -\vec{r}' \wedge \vec{r} + \underbrace{k \tilde{\pi}^2 \vec{m} \wedge \vec{r}'}_{} s - k \tilde{\pi} \vec{m} \wedge \vec{r}' s + \underbrace{\frac{k}{2} \tilde{\pi}^2 \vec{r}' \wedge \vec{m}}_{} s \right. \\ \left. + \vec{r}' \wedge (k_s \vec{m} + kT\vec{b}) \frac{x^3}{6} s^2 + (k_s \vec{m} + kT\vec{b} - k^2 \vec{r}') \frac{\tilde{\pi}^2}{2} \wedge (\frac{1}{2} \vec{r}' - \vec{r}) s^2 + O(s^3) \right\}$$

$$= -x \left\{ -\vec{\phi}' - \frac{k}{2} \tilde{\pi}^2 \vec{b} s - k \sin \phi \tilde{\pi} \vec{b} s + k_s \frac{\tilde{\pi}^3}{6} \vec{b} s^2 - \frac{kT}{6} \tilde{\pi}^3 \vec{m} s^2 + \frac{k^2 \tilde{\pi}^2}{2} s^2 \vec{\phi}' \right. \\ \left. - k_s \frac{\tilde{\pi}^3}{2} \vec{b} s^2 + \frac{kT}{2} \tilde{\pi}^3 \vec{m} s^2 - k_s \frac{\tilde{\pi}^2}{2} \sin \phi \vec{b} s^2 + kT \frac{\tilde{\pi}^2}{2} \cos \phi \vec{r}' s^2 + O(s^3) \right\}$$

$$= -x \left\{ -\vec{\phi}' - \left(\frac{k}{2} \tilde{\pi}^2 \vec{b} + k \sin \phi \tilde{\pi} \vec{b} \right) s + \left(-\frac{k_s \tilde{\pi}^3}{3} \vec{b} + \frac{1}{3} kT \tilde{\pi}^3 \vec{m} + \frac{k^2 \tilde{\pi}^2}{2} s \vec{\phi}' \right. \right. \\ \left. \left. + \frac{\tilde{\pi}^2}{2} [kT \cos \phi - k_s \sin \phi] \vec{b} \right) \vec{r}' s^2 + O(s^3) \right\}$$

$$\begin{aligned}
& \cdot (X(\alpha x + s) - X(s) - \alpha \bar{x}^1(\phi, s)) \circ (X(\alpha \bar{x} + s) - X(s) - \alpha \bar{x}^1(\phi, s)) \\
& = \alpha^2 \left\{ |\bar{s}\bar{C} - \bar{x}|^2 - K \bar{s}^2 \bar{x} \cdot \bar{m} x + \frac{\bar{s}^3}{3} [\bar{s}\bar{C}^2 - \bar{x}^2] \cdot [K_s \bar{m} + K_T \bar{b} - K^2 \bar{C}] x^2 + O(x^3) \right\} \\
& = \alpha^2 \left\{ 1 + \bar{s}^2 - K \bar{s}^2 \cos \phi x + \frac{\bar{s}^3}{3} [-K_s \cos \phi - K_T \sin \phi - \bar{s} \frac{K^2}{4}] x^2 + O(x^3) \right\} \\
& \quad \left| X(\alpha \bar{x} + s) - X(s) - \alpha \bar{x}^1(\phi, s) \right|^{-3} \\
& = \frac{1}{x^3} \frac{1}{(1 + \bar{s}^2)^{3/2}} \left\{ 1 + \frac{3}{2} \frac{K \bar{s}^2}{1 + \bar{s}^2} \cos \phi x - \frac{\bar{s}^3}{2} \frac{[-K_s \cos \phi - K_T \sin \phi - \bar{s} \frac{K^2}{4}]}{1 + \bar{s}^2} x^2 \right. \\
& \quad \left. + \frac{15}{8} \frac{K^2 \bar{s}^4 \cos^2 \phi}{(1 + \bar{s}^2)^2} x^2 + O(x^3) \right\} \cos(1 + \varepsilon)^{-3/2} = 1 - \frac{3}{2} \varepsilon + \frac{15}{4} \varepsilon^2
\end{aligned}$$

• On a donc :

$$\begin{aligned}
r^L e &= -\frac{1}{x} \frac{1}{(1 + \bar{s}^2)^{3/2}} \left\{ -\bar{\Phi} - \frac{3}{2} \frac{K \bar{s}^2}{1 + \bar{s}^2} \cos \phi \bar{\Phi} x - \left(\frac{K}{2} \bar{s}^2 \bar{b} + K \sin \phi \bar{s} \bar{C} \right) x \right. \\
& \quad + \frac{\bar{s}^3}{2} \frac{-K_s \cos \phi - K_T \sin \phi - \bar{s} \frac{K^2}{4}}{1 + \bar{s}^2} \bar{\Phi} x^2 - \frac{15}{8} \frac{K^2 \bar{s}^4 \cos^2 \phi}{(1 + \bar{s}^2)^2} \bar{\Phi} x^2 \\
& \quad \left. - \left(\frac{K}{2} \bar{s}^2 \bar{b} + K \sin \phi \bar{s} \bar{C} \right) \frac{3}{2} \frac{K \bar{s}^2 \cos \phi}{1 + \bar{s}^2} x^2 \right. \\
& \quad \left. + \left(-\frac{K_s \bar{s}^3}{3} \bar{b} + \frac{1}{3} K_T \bar{s}^3 \bar{m} + \frac{\bar{s}^2}{2} [K_T \cos \phi - K_s \sin \phi] \bar{C} + K^2 \frac{\bar{s}^2}{2} \bar{\Phi} \right) x^3 + O(x^3) \right\} \\
& = f' + g' + h' + O\left(\frac{x^2}{(1 + \bar{s}^2)^{3/2}}\right)
\end{aligned}$$

$$\text{or } f' = \frac{\bar{\Phi}}{x(1 + \bar{s}^2)^{3/2}} \quad g' = \frac{3}{2} \frac{K \bar{s}^2}{(1 + \bar{s}^2)^{3/2}} \cos \phi \bar{\Phi} + \frac{\frac{K}{2} \bar{s}^2 \bar{b} + K \sin \phi \bar{s} \bar{C}}{(1 + \bar{s}^2)^{3/2}}$$

$$\begin{aligned}
h' &= x \left\{ \frac{\bar{s}^4 \frac{K^2}{4} \bar{\Phi}}{(1 + \bar{s}^2)^{5/2}} + \frac{15}{8} K^2 \bar{s}^4 \frac{\cos^2 \phi}{(1 + \bar{s}^2)^{5/2}} \bar{\Phi} + \left(\frac{K}{2} \bar{s}^2 \bar{b} \right) \frac{3}{2} \frac{K \bar{s}^2 \cos \phi}{(1 + \bar{s}^2)^{5/2}} \right. \\
& \quad \left. + \frac{\bar{s}^2}{2} \frac{[K_T \cos \phi - K_s \sin \phi] \bar{C} + \frac{\bar{s}^2}{2} K^2 \bar{\Phi}}{(1 + \bar{s}^2)^{5/2}} \right\} \text{ en laissant les termes impairs.} \\
& \quad \text{10}
\end{aligned}$$

② DA de I_n pour proche de 0 :

Posons

$$F'^+ = \int_0^{n/2} f' d\tilde{s} \quad F'^- = \int_{-n/2}^0 f' d\tilde{s} \quad \text{et } F' = F'^- + F'^+$$

$$G'^+ = \int_0^{n/2} g' d\tilde{s} \quad G'^- = \int_{-n/2}^0 g' d\tilde{s} \quad \text{et } G' = G'^- + G'^+$$

$$H'^+ = \int_0^{n/2} h' d\tilde{s} \quad H'^- = \int_{-n/2}^0 h' d\tilde{s} \quad \text{et } H' = H'^- + H'^+$$

On cherche $I_{n,\varepsilon} = \int_{-n/2}^{n/2} n L_\varepsilon d\tilde{s}$

$$= F' + G' + H' + O(\varepsilon^2)$$

③ Détermination de F' :

$$F' = \int_{-n/2}^0 f' d\tilde{s} + \int_0^{n/2} f' d\tilde{s} = \frac{2}{n} \vec{\Phi}'(1) \int_0^{n/2} \frac{d\tilde{s}}{(1+\tilde{s}^2)^{3/2}} = \frac{2}{n} \vec{\Phi}'(1) \left[\frac{\tilde{s}}{(1+\tilde{s}^2)^{1/2}} \right]_0^{n/2}$$

$$= 2 \frac{n}{\pi^2} \vec{\Phi}'(1) \left(1 + \left(\frac{n}{\pi} \right)^2 \right)^{-1/2} \quad (1+\varepsilon)^{-1/2} = 1 - \frac{1}{2}\varepsilon + \frac{3}{8}\varepsilon^2$$

$$= 2 \frac{\vec{\Phi}'(1)}{\pi} \left(1 - \frac{1}{2} \left(\frac{\pi}{n} \right)^2 + \frac{3}{4} \left(\frac{\pi}{n} \right)^4 + O \left(\left(\frac{\pi}{n} \right)^6 \right) \right)$$

$$F' = \frac{2\vec{\Phi}'(1)}{\pi} - \vec{\Phi}'(1) \frac{\pi}{n^2} + \frac{3}{4} \vec{\Phi}' \frac{\pi^3}{n^4} + O \left(\frac{\pi^5}{n^6} \right)$$

$$\pi \gg \frac{\pi^5}{n^6} \Rightarrow n \gg \pi^{2/3} \quad \text{Or } \pi^{1/2} \gg \pi^{2/3} \Rightarrow \pi \text{ car } 1/2 < 2/3 < 1$$

On prend n dans l'intervalle suivant: $\boxed{\pi^{1/2} \gg n \gg \pi^{2/3}}$

(22) Détermination de G' :

$$G' = \int_{-\eta/\lambda}^{\eta/\lambda} g' d\tilde{s} + \int_0^{\eta/\lambda} g' d\tilde{s} = 3K \cos \phi \vec{p} \cdot \int_0^{\eta/\lambda} \frac{\tilde{s}^2}{(1+\tilde{s}^2)^{3/2}} d\tilde{s} + K \vec{b} \cdot \int_0^{\eta/\lambda} \frac{\tilde{s}^2}{(1+\tilde{s}^2)^{3/2}} d\tilde{s}$$

$$\text{Or } \int_0^{\eta/\lambda} \frac{\tilde{s}^2}{(1+\tilde{s}^2)^{3/2}} d\tilde{s} = \frac{1}{3} \left[\frac{\tilde{s}^3}{(1+\tilde{s}^2)^{1/2}} \right]_0^{\eta/\lambda} = \frac{1}{3} \frac{(\eta/\lambda)^3}{[1+(\eta/\lambda)^2]^{1/2}} = \frac{1}{3} \frac{1}{[1+(\frac{\eta}{\lambda})^2]^{3/2}}$$

$$(1+(\frac{\eta}{\lambda})^2)^{-3/2} = 1 - \frac{3}{2} \left(\frac{\eta}{\lambda}\right)^2 + \frac{15}{8} \left(\frac{\eta}{\lambda}\right)^4 + O\left(\left(\frac{\eta}{\lambda}\right)^6\right)$$

$$\int_0^{\eta/\lambda} \frac{\tilde{s}^2}{(1+\tilde{s}^2)^{3/2}} d\tilde{s} = \frac{1}{3} \left[1 - \frac{3}{2} \left(\frac{\eta}{\lambda}\right)^2 + O\left(\left(\frac{\eta}{\lambda}\right)^4\right) \right].$$

$$\eta \gg (\frac{\eta}{\lambda})^4 \quad \text{et} \quad \eta \gg \lambda^3 \Leftrightarrow \eta \gg \lambda^{3/4} \quad \text{ou} \quad \eta \gg \lambda^{3/4} \quad \text{et} \quad \eta \gg \lambda^{2/3}.$$

$$\int_0^{\eta/\lambda} \frac{\tilde{s}^2}{(1+\tilde{s}^2)^{3/2}} d\tilde{s} = \left[-\frac{\tilde{s}}{(1+\tilde{s}^2)^{1/2}} + \ln(\tilde{s} + \sqrt{1+\tilde{s}^2}) \right]_0^{\eta/\lambda}$$

$$= -\frac{\eta/\lambda}{(1+(\eta/\lambda)^2)^{1/2}} + \ln\left(\frac{\eta}{\lambda} + \sqrt{1+(\eta/\lambda)^2}\right)$$

$$= -\frac{1}{(1+(\frac{\eta}{\lambda})^2)^{1/2}} + \ln\frac{\eta}{\lambda} + \ln\left(1 + \sqrt{1+(\eta/\lambda)^2}\right)$$

$$(1+(\frac{\eta}{\lambda})^2)^{-1/2} = 1 - \frac{1}{2} \left(\frac{\eta}{\lambda}\right)^2 + \frac{3}{8} \left(\frac{\eta}{\lambda}\right)^4 + O\left(\left(\frac{\eta}{\lambda}\right)^6\right) \quad (1+\varepsilon)^{-1/2} = 1 - \frac{\varepsilon}{2}$$

$$(1+(\frac{\eta}{\lambda})^2)^{1/4} = 1 + \frac{1}{2} \left(\frac{\eta}{\lambda}\right)^2 + O\left(\left(\frac{\eta}{\lambda}\right)^4\right)$$

$$\begin{aligned} \ln(1 + \sqrt{1+(\eta/\lambda)^2}) &= \ln(2 + \frac{1}{2}(\frac{\eta}{\lambda})^2 + O\left(\left(\frac{\eta}{\lambda}\right)^4\right)) = \ln 2 + \ln\left(1 + \frac{1}{4}(\frac{\eta}{\lambda})^2 + O\left(\left(\frac{\eta}{\lambda}\right)^4\right)\right) \\ &= \ln 2 + \frac{1}{4} \left(\frac{\eta}{\lambda}\right)^2 + O\left(\left(\frac{\eta}{\lambda}\right)^4\right) \quad \text{car} \quad \ln(1+\varepsilon) = \varepsilon + O(\varepsilon^2) \end{aligned}$$

$$\begin{aligned} D' \text{ au } \int_0^{\eta/\lambda} \frac{\tilde{s}^2}{(1+\tilde{s}^2)^{3/2}} d\tilde{s} &= -1 + \frac{1}{2} \left(\frac{\eta}{\lambda}\right)^2 + O\left(\left(\frac{\eta}{\lambda}\right)^4\right) + \ln\left(\frac{\eta}{\lambda}\right) + \ln 2 + \frac{1}{4} \left(\frac{\eta}{\lambda}\right)^2 + O\left(\left(\frac{\eta}{\lambda}\right)^4\right) \\ &= -1 + \ln\frac{\eta}{\lambda} + \ln 2 + \frac{3}{4} \left(\frac{\eta}{\lambda}\right)^2 + O\left(\left(\frac{\eta}{\lambda}\right)^4\right) \end{aligned}$$

Il vient donc :

$$G' = k \cos \phi \vec{\phi} [1 - \frac{3}{2} (\frac{\pi}{n})^2] + k \vec{b} [-1 + \ln 2 + \ln \frac{n}{\pi} + \frac{3}{4} (\frac{\pi}{n})^2] + O((\frac{\pi}{n})^4)$$

③ Détermination de H' : $H' = \int_0^{n/2} h' d\tilde{x} + \int_{n/2}^{\pi/2} h' d\tilde{x}$

$$H' = -2\pi \left\{ \left[-\frac{k^2}{8} \vec{\phi} - \frac{3}{4} k^2 \cos \phi \vec{b} \right] \int_0^{n/2} \frac{\tilde{x}^4}{(1+\tilde{x}^2)^{5/2}} d\tilde{x} + \frac{k^2}{2} \vec{\phi} \int_{n/2}^{\pi/2} \frac{\tilde{x}^2}{(1+\tilde{x}^2)^{3/2}} d\tilde{x} \right.$$

$$\left. - \frac{15}{8} k^2 \cos^2 \phi \vec{\phi} \int_0^{n/2} \frac{\tilde{x}^4}{(1+\tilde{x}^2)^{7/2}} d\tilde{x} + [k T \cos \phi - k_s \sin \phi] \frac{\vec{\phi}}{2} \int_{n/2}^{\pi/2} \frac{\tilde{x}^2}{(1+\tilde{x}^2)^{3/2}} d\tilde{x} \right\}$$

$$\int_0^{n/2} \frac{\tilde{x}^4}{(1+\tilde{x}^2)^{5/2}} d\tilde{x} = \left[-\frac{\tilde{x}}{\sqrt{1+\tilde{x}^2}} - \frac{1}{3} \frac{\tilde{x}^3}{(1+\tilde{x}^2)^{3/2}} + \ln(\tilde{x} + \sqrt{1+\tilde{x}^2}) \right]_0^{n/2}$$

$$= -1 + O((\frac{\pi}{n})^2) - \frac{1}{3} + O((\frac{\pi}{n})^2) + \ln \frac{n}{\pi} + \ln 2 + O((\frac{\pi}{n})^2)$$

$$= -\frac{4}{3} + \ln 2 + \ln \frac{n}{\pi} + O((\frac{\pi}{n})^2) \quad \text{On a } n \gg \pi \quad \underline{(\frac{\pi}{n})^2}$$

$$\int_0^{n/2} \frac{\tilde{x}^2}{(1+\tilde{x}^2)^{3/2}} d\tilde{x} = -1 + \ln n/2 + \ln 2 + O((\frac{\pi}{n})^2) \quad \text{d'après 22.}$$

$$\int_0^{n/2} \frac{\tilde{x}^4}{(1+\tilde{x}^2)^{7/2}} d\tilde{x} = \left[\frac{1}{5} \frac{\tilde{x}^5}{(1+\tilde{x}^2)^{5/2}} \right]_0^{n/2}$$

$$= \frac{1}{5} \left(1 + (\frac{\pi}{n})^2 \right)^{-5/2} = \frac{1}{5} \left(1 + O((\frac{\pi}{n})^2) \right)$$

D'où

$$H' = -2\pi \left\{ \left[-\frac{k^2}{8} \vec{\phi} - \frac{3}{4} k^2 \cos \phi \vec{b} \right] \left(-\frac{4}{3} + \ln 2 + \ln \frac{n}{\pi} \right) + \frac{3}{8} k^2 \cos^2 \phi \vec{\phi} \right.$$

$$\left. + [k T \cos \phi - k_s \sin \phi] \frac{\vec{\phi}}{2} \left(-1 + \ln \frac{n}{\pi} + \ln 2 \right) + \frac{k^2}{2} \vec{\phi} (-1 + \ln \frac{n}{\pi} + \ln 2) \right\}$$

④ Expression de $I_{1\ell}$, DA ^{intérieur}_{extérieur} de I_2 en $\varepsilon = 0$:

$$I_{1\ell} = F' + G' + H' + O(\varepsilon^2)$$

III. Déf de $\vec{v}_r(r, \phi, \eta)$ en $r=0$:

$$\text{Onde } \vec{v}_r(r, \phi, \eta) = (E_{re} + I_{re}) \frac{r}{4\pi} \\ = (F + G + H + F' + G' + H') + O(r^2) \frac{r}{4\pi}$$

$$\begin{aligned} \frac{4\pi}{r} \vec{v}_r &= \left\{ \vec{A} - [\vec{k}(r) \vec{b}(r) \ln \eta] + \vec{k}(r) \vec{b}(r) \ln St + O(r^2) \right\} \\ &+ \left\{ r \left(\vec{B} - \frac{\vec{\Phi}}{St} - \frac{3}{4} K^2 \ln St \vec{\Phi}' + \sin \phi \vec{C}^T \vec{k}_s \ln St - \cos \phi \vec{C}^T \vec{k}_T \ln St \right) \right. \\ &\quad + \boxed{\vec{\Phi}' \frac{\pi}{\eta^2}} + \boxed{\left[\frac{3}{4} K^2 \vec{\Phi}' - \sin \phi \vec{C}^T \vec{k}_s + \cos \phi \vec{C}^T \vec{k}_T \right] r \ln \eta} + O(r\eta) \\ &\quad - 3r \left(\vec{C} - 2 \left(\frac{K}{2} \right)^2 \cos \phi \vec{b} \ln St \right) - \boxed{6 \left(\frac{K}{2} \right)^2 \vec{b} \cos \phi r \ln \eta} + O(r\eta) \Big\} \\ &+ O(r^2) \\ &+ \left\{ 2 \frac{\vec{\Phi}}{r} - \boxed{\vec{\Phi} \frac{\pi}{\eta^2}} + \boxed{\frac{3}{2} \vec{\Phi}' \frac{\pi^3}{\eta^4}} + O\left(\frac{\pi^5}{\eta^6}\right) \right\} \\ &+ \left\{ K \cos \phi \vec{\Phi}' - \frac{3}{2} K \cos \phi \vec{\Phi} \vec{\Phi}' \left(\frac{\pi}{\eta} \right)^2 + O\left(\left(\frac{\pi}{\eta} \right)^4\right) \right. \\ &\quad + K \vec{b} \boxed{-1 + \ln 2} + \boxed{K \vec{b} \ln \eta} - K \vec{b} \ln r + \boxed{\frac{3}{4} K \vec{b} \left(\frac{\pi}{\eta} \right)^2} + O\left(\left(\frac{\pi}{\eta} \right)^4\right) \Big\} \\ &+ \left\{ r \left[\left(-\frac{3}{4} K^2 \vec{\Phi}' + \frac{3}{2} K^2 \cos \phi \vec{b} \right) \left(-\frac{4}{3} + \ln 2 \right) - \frac{1}{3} K^2 \vec{\Phi}' + \frac{3}{4} K^2 \cos^2 \phi \vec{\Phi}' \right. \right. \\ &\quad \left. \left. + (K \vec{C} \cos \phi - \vec{k}_s \sin \phi) \vec{\Phi}' (-1 + \ln 2) \right] \right. \\ &\quad + \boxed{\left(-\frac{3}{4} K^2 \vec{\Phi}' + \frac{3}{2} K^2 \cos \phi \vec{b} \right) r \ln \eta} - \left(-\frac{3}{4} K^2 \vec{\Phi}' + \frac{3}{2} K^2 \cos \phi \vec{b} \right) r \ln r + O\left(\left(\frac{\pi^3}{\eta^6} \right)\right) \\ &\quad \left. + \left(-K \vec{C} \cos \phi + \vec{k}_s \sin \phi \right) \vec{\Phi}' r \ln \eta - \left(-K \vec{C} \cos \phi + \vec{k}_s \sin \phi \right) \vec{\Phi}' r \ln r \right\} \\ &+ O(r^2). \end{aligned}$$

$$\begin{aligned} \text{Onde } \vec{v} \sin 2\phi + \vec{\Phi} \cos 2\phi &= 2 \vec{v} \sin \phi \cos \phi + 2 \vec{\Phi} \cos^2 \phi - \vec{\Phi}' \cos \phi \cos 2\phi = 2 \cos^2 \phi - 1 \\ &= 2 \cos \phi (\sin \phi \vec{v} + \cos \phi \vec{\Phi}') - \vec{\Phi}' \\ &= 2 \cos \phi \vec{b} - \vec{\Phi}'. \end{aligned}$$

$$-\frac{3}{4} K^2 \vec{\Phi}' + \frac{3}{2} K^2 \cos \phi \vec{b} = \frac{3}{4} K^2 [2 \cos \phi \vec{b} - \vec{\Phi}'] = \frac{3}{4} K^2 [\vec{v} \sin 2\phi + \vec{\Phi}' \cos 2\phi] = 0$$

Simplifions le terme en facteur de σ :

$$\begin{aligned}
 & \vec{B} - \frac{\vec{\Phi}}{s+2} - \frac{3}{4} k^2 \ln s^2 \vec{\phi}' + \sin \vec{\phi} \vec{C}^T K_s \ln s^2 - \cos \vec{\phi} \vec{C}^T k T \ln s^2 \\
 & - 3 \vec{C} + \frac{3}{2} k^2 b \vec{b} \cos \phi \ln s^2 \\
 & + g \left(-\frac{4}{3} + \ln 2 \right) + \frac{3}{4} k^2 \cos \phi \vec{\phi}' + \left[k T \cos \phi - K_s \sin \phi \right] \frac{T}{2} \left(-1 + \ln 2 \right) - \frac{1}{3} k^2 \vec{\phi}' \\
 = & \vec{B} - \frac{\vec{\Phi}}{s+2} - 3 \vec{C} + g \left(-\frac{4}{3} \right) + \ln s^2 \left(-\frac{3}{4} k^2 \vec{\phi}' + \frac{3}{2} k^2 b \vec{b} \cos \phi \right) + g \ln 2 + \frac{3}{4} k^2 \left[1 + \frac{\cos 2\phi}{2} \right] \vec{\phi}' \\
 & - \frac{1}{3} k^2 \vec{\phi}' + \vec{C} \left[(K_s \sin \phi - k T \cos \phi) \ln s^2 + \left[(K_s \sin \phi - k T \cos \phi) \left(1 - \frac{\ln 2}{2} \right) \right] \right] \\
 & \text{car } \cos^2 \alpha = 1 - \frac{\cos 2\alpha}{2} \\
 = & \vec{B} - 3 \vec{C} + g \left[\ln 2 s^2 - \frac{4}{3} \right] + \frac{3}{4} k^2 \left[1 + \frac{\cos 2\phi}{2} \right] \vec{\phi}' - \frac{\vec{\Phi}}{s+2} - \frac{1}{3} k^2 \vec{\phi}' \\
 & + \vec{C} \left[K_s \sin \phi - k T \cos \phi \right] \left[\ln s^2 + \frac{-1 - \frac{\ln 2}{2}}{1 + \frac{\ln 2}{2}} \right] \\
 = & \vec{B} - 3 \vec{C} + g \left[\ln 2 s^2 - \frac{4}{3} \right] + \frac{3}{8} k^2 \cos 2\phi \vec{\phi}' + k^2 \vec{\phi}' \left(\frac{3}{8} - \frac{1}{3} \right) - \frac{\vec{\Phi}}{s+2} \\
 & + \vec{C} \left[K_s \sin \phi - k T \cos \phi \right] \left[\ln s^2 + \frac{1 - \cancel{\frac{\ln 2}{2}}}{-1 + \cancel{\frac{\ln 2}{2}}} \right]
 \end{aligned}$$

D'où:

$$\begin{aligned}
 \vec{v}_d = & \frac{P}{q\pi r} \vec{\phi}' + \frac{rK}{4\pi} \left[\ln \frac{2s^2}{\pi} - 1 \right] \vec{b} + \frac{rK}{4\pi} \cos \phi \vec{\phi}' + \vec{A}' \\
 & + \frac{3}{16} \frac{rk^2 r}{\pi} \left\{ (\pi \sin 2\phi + \vec{\phi}' \cos 2\phi) \left[\ln \frac{2s^2}{\pi} - \frac{4}{3} \right] + \frac{1}{2} \cos 2\phi \vec{\phi}' + \frac{1}{18} \vec{\phi}' + \left(\vec{B} - 3\vec{C} - \frac{\vec{\Phi}}{s+2} \right) \frac{4}{3} \right\} \\
 & + \frac{r}{4\pi} \vec{C} \left[K_s \sin \phi - k T \cos \phi \right] \left[\ln \frac{2s^2}{\pi} + \frac{1 - \cancel{\frac{\ln 2}{2}}}{-1 + \cancel{\frac{\ln 2}{2}}} \right]
 \end{aligned}$$

(Voir FUKUMOTO 91)