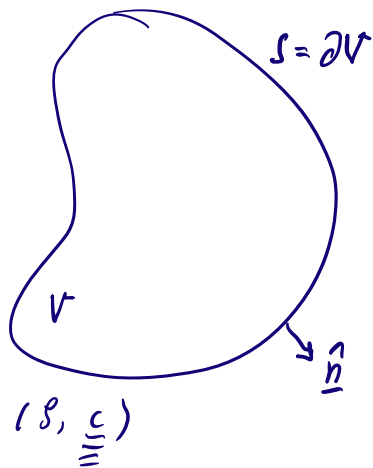


Basic theorems in dynamic elasticity



What is the displacement \underline{u} , given
 a source \underline{f} and material properties
 $(\rho, \underline{\underline{c}})$?

From Newton's (second) law $\underline{F} = m \underline{\underline{u}}''$, we can write

$$\underbrace{\int_V \underline{f} \, dx^3}_{\text{body forces}} + \underbrace{\int_{\partial V} \underline{\hat{n}} \cdot \underline{\underline{\tau}} \, dx^2}_{\text{surface forces}} = \int_V \rho \underline{\underline{u}}'' \, dx^3$$

which leads to the equations of motion

$$\boxed{\rho \underline{\underline{u}}'' - \underline{\nabla} \cdot \underline{\underline{\tau}} = \underline{f}}$$

ρ : density
 $\underline{\underline{c}}$: elastic
 tensor

The constitutive relationship relates stress with displacement.
 For elastic bodies, it is described by Hooke's law

$$\underline{\underline{\tau}} = \underline{\underline{c}} : \underline{\underline{\varepsilon}}$$

$$\underline{\underline{\varepsilon}} = \frac{1}{2} [\underline{\nabla} \underline{u} + (\underline{\nabla} \underline{u})^T]$$

strain

We can thus write the equations of motion for an elastic medium as

$$\rho \ddot{\underline{u}} - \nabla \cdot (\underline{\underline{c}} : \nabla \underline{u}) = \underline{f}$$

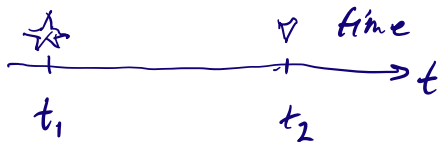
Note that the complexity of wave propagation relates to the complexity of the medium.

Betti's theorem : Determines a body's response to any forcing once its response to a specific forcing is found.

Betti states that displacement \underline{u} is a unique solution to a given initial stress $\underline{\underline{\tau}}$ and force \underline{f} .

That is, if $(\underline{\underline{\tau}}, \underline{f}) = (\underline{\underline{g}}, \underline{g})$ then we must have $\underline{u}(\underline{x}, t) = \underline{v}(\underline{x}, t)$ as well.

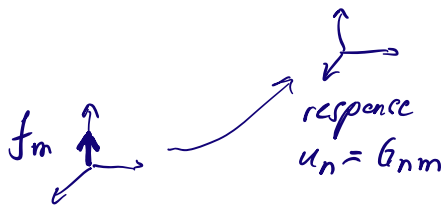
Green's function: Find the response to an excitation
impulsive in space & time.



reciprocity in time:

$$\underline{G}(\underline{x}, t_2; \underline{\xi}, t_1) = \underline{G}(\underline{x}, -t_1; \underline{\xi}, -t_2)$$

$\underbrace{\hspace{10em}}$
source at ξ
at time t_1
 $\underbrace{\hspace{10em}}$
source at ξ
at time $-t_2$



reciprocity in space:

$$G_{mn}(\underline{x}, t_2; \underline{\xi}, t_1) = G_{nm}(\underline{\xi}, t_2; \underline{x}, t_1)$$

$\underbrace{\hspace{2em}}$ $\underbrace{\hspace{2em}}$
 switching location
and component

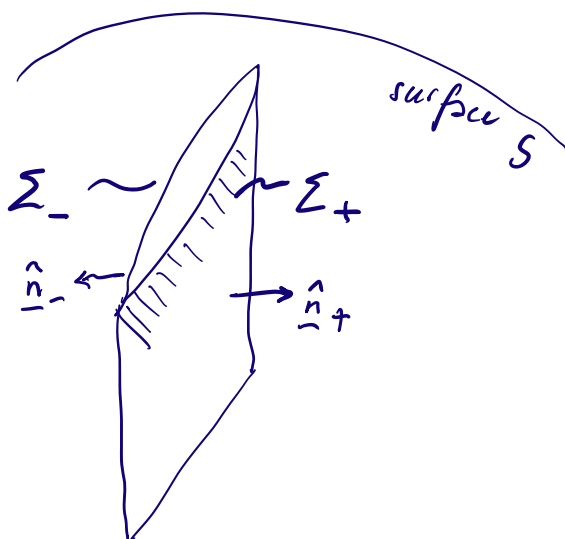
Representation theorem: Knowing displacement on the fault is enough to determine the displacement anywhere

$$\underline{u}(\underline{x}, t) = \int_{-\infty}^{+\infty} dT \int_{\Sigma} \underbrace{[\underline{u}(\underline{\xi}, T)]}_{\text{displacement at receiver away from fault}} \cdot \underbrace{[\underline{\underline{c}} : \underline{\underline{\nabla}} \underline{G}(\underline{x}, t-T; \underline{\xi}, 0)] \cdot \underline{\hat{n}}}_{\text{stress } \sigma \text{ on fault}} d\Sigma$$

displacement at receiver away from fault

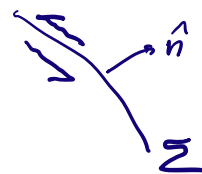
displacement on fault

stress σ



fault $\Sigma = \Sigma_- \cup \Sigma_+$

welded boundary,
no opening at fault:



$\underline{\underline{G}}$ and $\underline{\hat{n}} \cdot \underline{\underline{\tau}} = \underline{\hat{n}} \cdot (\underline{\underline{c}} : \underline{\underline{\nabla}} \underline{u})$
are continuous across Σ

Displacement written as convolutions,

$$f(t) * g(t) = \int_{-\infty}^{\infty} f(\tau) g(t-\tau) d\tau,$$

becomes

$$\underline{u} = \underbrace{S}_{\text{source}} * \underbrace{P}_{\substack{\text{propagation} \\ \triangleq \underline{G}}} * \underbrace{I}_{\text{instrument response}}$$

Index notation: Aki & Richards (3.2)

$$u_n(\underline{x}, t) = \int_{-\infty}^{\infty} dT \sum_i \underbrace{[u_i(\underline{\xi}, T)]}_{\text{slip function}} c_{ijpq} n_j \frac{\partial}{\partial \xi_q} G_{np}(\underline{x}, t-T; \underline{\xi}, 0) d\underline{\xi}$$