Ray theory - High-frequency approximation

Let's consider the wave equation for an clastic isotropic medium

 $P\ddot{u} = (2+\mu) \nabla (D\cdot u) + \mu \nabla u + (D\cdot u) \nabla \lambda + (D\mu) (Du) + (Du) \cdot (D\mu)$

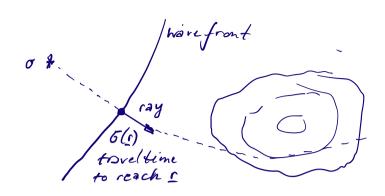
where the Lami parameter $\lambda = \lambda(\Sigma)$, $\mu = \mu(\Sigma)$ can vary in space.

We can substitute a monochromatic traveling wave $u(r,t) = U(r) e^{i\omega(f(r)-t)}$ $\omega: amgn/ar$ frequency factor amplitude

with spatial derivation $\nabla u(r,t) = \nabla u(r)e^{i\omega[6(r)-t]}$ $-i\omega u(r) \nabla 6(r)e^{i\omega[6(r)-t]}$

time derivation $\partial_t^2 u(r,t) = -U(r) \omega^2 e^{i\omega(\mathcal{L}(r)-t)}$

and find terms proportional to wo, wand we independent of co



The high-frequency approximation neglects we terms and sets the others to zero, as it looks for solutions independent of w U(1), 6(1)

$$\omega^{2} \text{ ferms}: \left(\frac{3+m}{g} \nabla G \cdot \nabla G + \frac{m}{g} \nabla G \cdot \nabla G + \frac{T}{2}\right) \cdot \mathcal{U} = 0$$

$$\text{has non-unital solutions for}$$

$$\text{Eikonal} \left(\nabla G(r) \cdot \nabla G(r) = \frac{g(r)}{A(r) + 2\mu(r)} \longrightarrow P\text{-wave rays}$$

$$\text{equations} \left(\nabla G(r) \cdot \nabla G(r) = \frac{g(r)}{\mu(r)} \longrightarrow S\text{-wave rays}\right)$$

w terms: set to zero - "transport" equations
for U(s)