Comparing Classical and Nonclassical Symmetries of Nonlinear Partial Differential Equations

William Helman and Daniel Sinderson

Southern Oregon University

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Project Goal

Our research objective for this project was to calculate the classical and nonclassical symmetry groups for the Born-Infeld equation and the reduced Gibbons-Tsarev equation and compare them.

Definition

A symmetry is a transformation that leaves an object invariant.

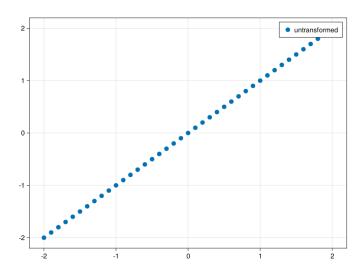
Definition

A symmetry is a transformation that leaves an object invariant.

Definition

A symmetry is a change that doesn't change anything.

Let's see this in action using the simple linear equation x - y = 0.



A Non-Example

Example (1)

- For our first transformation, let's define new variables $\bar{x} = x + 1$ and $\bar{y} = y$.
- ▶ Now let's rewrite our equation using these new variables.

$$\bar{x} - \bar{y} = 0$$
 by definition $x + 1 - y = 0$ by substitution $y = x + 1$ by rewriting in slope-intercept form

▶ This transformation is not a symmetry:

$$x - y + 1 \neq x - y$$

A Transformation that is a Symmetry

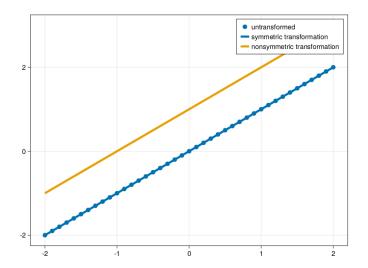
Example (2)

- Let's define some new variables again $\bar{x} = x + 1$ and $\bar{y} = y + 1$.
- ▶ Now let's rewrite our equation using these new variables.

▶ This transformation is a symmetry:

$$x - y = x - y$$

The graphs of our three equations.



Who cares?

➤ Symmetries help us to understand and to solve equations that we wouldn't normally be able to solve.

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- Symmetries help us to understand and to solve equations that we wouldn't normally be able to solve.
- ➤ Symmetries also encode physically meaningful aspects of equations, like conservation laws.

Definition

▶ A differential equation is an equation that contains both an unknown function and information about how that function relates to its rates of change.

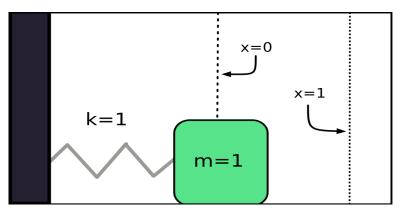
Definition

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- ➤ Differential equations show up everywhere we model something using information about how that thing changes. This includes everything from population dynamics to planetary orbits.

Definition

- ▶ A differential equation is an equation that contains both an unknown function and information about how that function relates to its rates of change.
- ▶ Differential equations show up everywhere we model something using information about how that thing changes. This includes everything from population dynamics to planetary orbits.
- ▶ Differential equations are different from algebraic equations, and they can't be solved in the same ways.

As an example, let's look a frictionless spring-mass system.

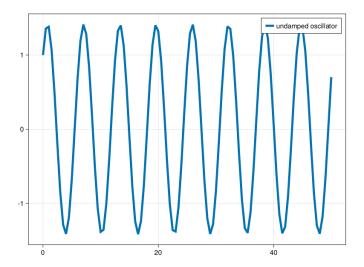


This system can be modeled by the following differential equation.

$$m\ddot{x} = -\kappa x$$

- ▶ Here m is the mass, \ddot{x} is the acceleration, κ is the spring constant, and x is the position.
- ▶ If we set both m and κ to 1, we get $\ddot{x} = -x$.
- This is a simple differential equation with the algebraic solution of $x(t) = c_1 \cos(t) + c_2 \sin(t)$.

This is the graph of our spring system where $c_1 = c_2 = 1$.



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- ➤ The Born-Infeld equation comes from a relativistic formulation of quantum electrodynamics (circa 1930).
- ▶ In classical electrodynamics the energy of a particle grows to infinity as its radius becomes very small.
- ► This is a problem if you want to study electrodynamics on the quantum scale.
- Max Born and Leopold Infeld's "fix" was to introduce a relativistic factor, similar to how $E=mc^2$ becomes $E=mc^2(1-\sqrt{1-v^2/c^2})$ in special relativity.

► From this new theory, we can derive an equation for electromagnetic waves in space. We get

$$\left(1-\frac{1}{b^2}\bigg(\frac{\partial u}{\partial y}\bigg)^2\right)\frac{\partial^2 u}{\partial x^2}+2\,\frac{1}{b^2}\frac{\partial u}{\partial x}\frac{\partial u}{\partial y}\frac{\partial^2 u}{\partial x\partial y}-\bigg(1-\frac{1}{b^2}\bigg(\frac{\partial u}{\partial x}\bigg)^2\bigg)\frac{\partial^2 u}{\partial y^2}=0,$$

where b is a constant.

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$$\left(1 - \frac{1}{b^2} \left(\frac{\partial u}{\partial y}\right)^2\right) \frac{\partial^2 u}{\partial x^2} + 2 \frac{1}{b^2} \frac{\partial u}{\partial x} \frac{\partial u}{\partial y} \frac{\partial^2 u}{\partial x \partial y} - \left(1 - \frac{1}{b^2} \left(\frac{\partial u}{\partial x}\right)^2\right) \frac{\partial^2 u}{\partial y^2} = 0,$$

where b is a constant.

Setting b = 1 for simplicity, we have our Born-Infeld equation

$$\left(1 - \left(\frac{\partial u}{\partial y}\right)^2\right) \frac{\partial^2 u}{\partial x^2} + 2 \frac{\partial u}{\partial x} \frac{\partial u}{\partial y} \frac{\partial^2 u}{\partial x \partial y} - \left(1 - \left(\frac{\partial u}{\partial x}\right)^2\right) \frac{\partial^2 u}{\partial y^2} = 0.$$

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- Note that this is now a differential equation of a function of two variables, where u = u(x, y).
- ▶ This is called a partial differential equation, because it makes use of these "partial derivatives" where $\partial u/\partial x$ reads as "the partial derivative of u with respect to x".

$$\left(1-\left(\frac{\partial u}{\partial y}\right)^2\right)\frac{\partial^2 u}{\partial x^2}+2\,\frac{\partial u}{\partial x}\frac{\partial u}{\partial y}\frac{\partial^2 u}{\partial x\partial y}-\left(1-\left(\frac{\partial u}{\partial x}\right)^2\right)\frac{\partial^2 u}{\partial y^2}=0$$

- Note that this is now a differential equation of a function of two variables, where u = u(x, y).
- ▶ This is called a partial differential equation, because it makes use of these "partial derivatives" where $\partial u/\partial x$ reads as "the partial derivative of u with respect to x".
- ▶ Not only is this a PDE, but it's also nonlinear. These equations are usually impossible to solve, and may not possess a general solution.

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- ➤ The Gibbons-Tsarev equation arose from special solutions to the nonlinear Schrödinger equation in 1996.
- ► Thus, it comes from studying some special behavior of light, gravitational waves, and plasma oscillations.
- Aside from its physical meaning, it has experienced a lot of research interest in a purely mathematical context.
- ▶ By studying the behavior and geometry of equations like Gibbons-Tsarev, we can learn more about how PDE work in general.

▶ Our equation is a reduced form of the Gibbons-Tsarev equation, dropping a +1, term as it arises naturally that way from its so-called "quotient construction". We studied

$$\frac{\partial^2 u}{\partial x^2} - \frac{\partial u}{\partial y} \frac{\partial^2 u}{\partial x \partial y} + \frac{\partial u}{\partial x} \frac{\partial^2 u}{\partial y^2} = 0.$$

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Input Equation \rightarrow Calculate Lie's invariance condition

 \rightarrow Solve linear system of PDEs

 \rightarrow Symmetries!

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- ▶ For a differential equation, this means that the function and any of its derivatives that are present in the equation must not change under the transformation.
- Lie's invariance condition gives us a way to generate transformations that meet all these requirements.
- ► If you can solve Lie's invariance condition for your equation, you can find its symmetries.

In general, Lie's invariance condition is the falsely unassuming

$$\Gamma^{(n)}(\Delta)\Big|_{\Delta=0} = 0$$

But there's a lot of complexity hiding under this notation.

Here Γ is the infinitesimal operator, Δ is your differential equation, and n is its order.

For our equations, the infinitesimal operator Γ we need is

$$\Gamma^{(2)} = Y \frac{\partial}{\partial y} + X \frac{\partial}{\partial x} + U \frac{\partial}{\partial u} + U_{[y]} \frac{\partial}{\partial u_y} + U_{[x]} \frac{\partial}{\partial u_x} + U_{[yy]} \frac{\partial}{\partial u_{yy}} + U_{[xy]} \frac{\partial}{\partial u_{xy}} + U_{[xx]} \frac{\partial}{\partial u_{xx}} = 0,$$

where

$$\begin{split} &U_{[y]} = D_y(U) - u_y D_y(Y) - u_x D_y(X), \\ &U_{[x]} = D_x(U) - u_y D_x(Y) - u_x D_x(X), \\ &U_{[yy]} = D_y(U_{[y]}) - u_{yy} D_y(Y) - u_{xy} D_y(X), \\ &U_{[xy]} = D_y(U_{[x]}) - u_{xy} D_y(Y) - u_{xx} D_y(X), \\ &U_{[xx]} = D_y(U_{[x]}) - u_{xy} D_x(Y) - u_{xx} D_x(X), \\ &D_y = \frac{\partial}{\partial y} + u_y \frac{\partial}{\partial u} + u_{yy} \frac{\partial}{\partial u_y} + u_{xy} \frac{\partial}{\partial u_x}, \\ &D_x = \frac{\partial}{\partial x} + u_x \frac{\partial}{\partial u} + u_{xy} \frac{\partial}{\partial u_y} + u_{xx} \frac{\partial}{\partial u_x}. \end{split}$$



➤ The result of putting our PDEs through Lie's invariance condition was a system of 24 linear PDEs for the Born-Infeld equation and 21 for the reduced Gibbons-Tsarev equation.

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- ► Though difficult, these systems were easier to solve than the original single equations because of their linearity.
- ▶ Solving them resulted in the following symmetries.

Classical Symmetries of the Born-Infeld Equation

- $X(x,y,u) = c_1x + c_2u + c_3y + c_4$
- $Y(x,y,u) = c_1y + c_3x + c_5u + c_6$
- $U(x, y, u) = c_1 u c_2 x + c_5 y + c_7$

Classical Symmetries of the reduced Gibbons-Tsarev Equation

- $X(x,y,u) = -2c_5y + (-c_2 + 2c_4)x + c_7$
- $Y(x,y,u) = -\frac{c_1x}{2} + c_4y + c_5u + c_6$
- $V(x,y,u) = c_1y + c_2u + c_3$

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- ▶ We need an equation called the invariant surface condition, which is a linear PDE, given by

$$\Delta_I = X(x, y, u) \frac{\partial u}{\partial x} + Y(x, y, u) \frac{\partial u}{\partial y} - U(x, y, u) = 0.$$

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$$\Delta_I = X(x, y, u) \frac{\partial u}{\partial x} + Y(x, y, u) \frac{\partial u}{\partial y} - U(x, y, u) = 0.$$

▶ We then seek the symmetries of the PDE system of our equation and the invariant surface condition. Our invariance condition then looks like

$$\Gamma^{(n)}(\Delta)\Big|_{\Delta=0, \Delta_I=0} = 0.$$

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- ➤ The result of running our PDE through the new invariance condition was a system of 28 nonlinear PDE for the Born-Infeld equation and 12 for the reduced Gibbons-Tsarev equation.
- ▶ Unlike the classical condition, this was a nonlinear system, and the equations were much harder to solve.
- ▶ But we did it!

Nonclassical Symmetries of the Born-Infeld Equation

- $X(x,y,u) = c_2x c_5u + c_7y + c_3$
- $Y(x, y, u) = c_1 u + c_2 y + c_7 x + c_4$
- $U(x, y, u) = c_1 y + c_2 u + c_5 x + c_6$

Nonclassical Symmetries of the reduced Gibbons-Tsarev Equation

- $X(x,y,u) = -2c_5y + (2c_4 c_2)x + c_7$
- $Y(x,y,u) = c_4y + c_5u \frac{c_1x}{2} + c_6$
- $V(x,y,u) = c_1y + c_2u + c_3$

Something Unexpected

Do you notice anything?

Classical Symmetries

$$X = c_1 x + c_2 u + c_3 y + c_4$$

$$Y = c_1 y + c_3 x + c_5 u + c_6$$

$$U = c_1 u - c_2 x + c_5 y + c_7$$

Nonclassical Symmetries

$$X = c_2 x - c_5 u + c_7 y + c_3$$

$$Y = c_1 u + c_2 y + c_7 x + c_4$$

$$U = c_1 y + c_2 u + c_5 x + c_6$$

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- ▶ Naturally, the next step is to find more examples of equations with this unique property, and try to understand why this is occurring.
- ► The Born-Infeld and Gibbons-Tsarev equations are both "Darboux integrable".
- ► It turns out, that every DI equation that we've checked also has this property.

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- ▶ Even further, what if we could prove the converse?
- ► Then we'd have a test for whether or not our equation is DI.
- ▶ This would be a big result, and great progress toward a test for whether or not a differential equation is solvable in general!

Thank you!