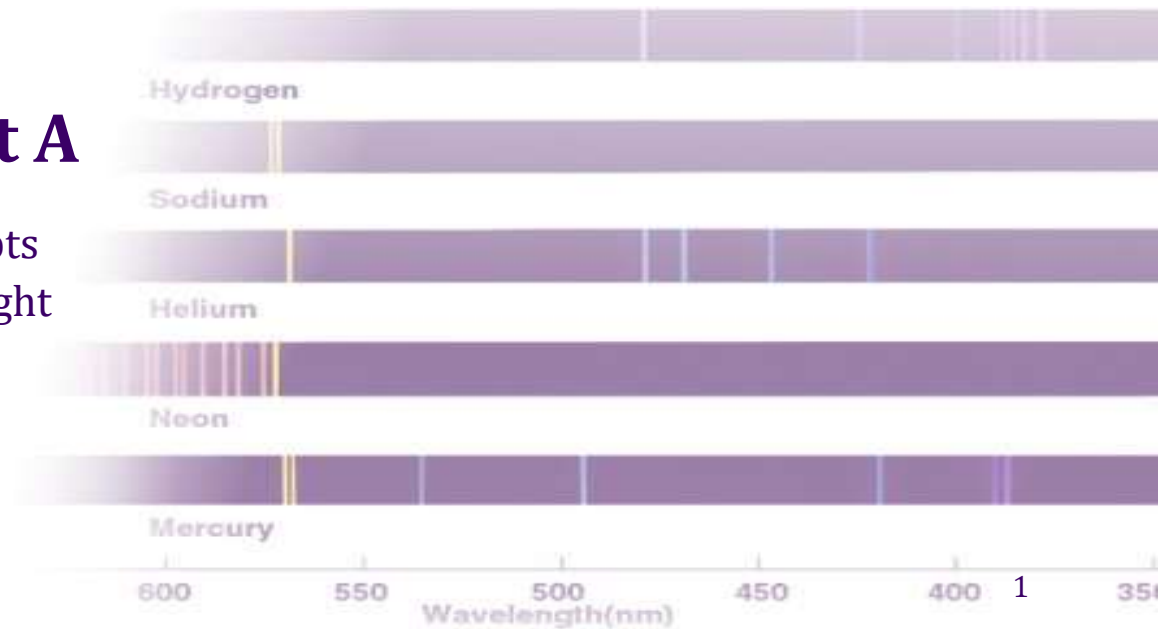


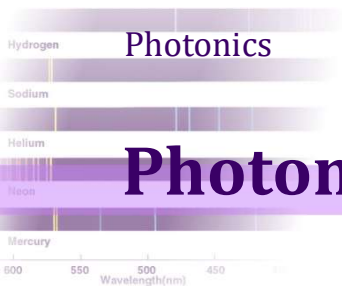
Photonics

R. Baets / E. Bente

Photons and Atoms Part A

Introduction quantum mechanics concepts
Atoms, molecules and interaction with light
Thermal light
Luminescence light





Photons and atoms

Generation/absorption of light \Rightarrow interaction of light with matter

- Matter= atoms
 - isolated (noble gas, gas vapour)
 - interacting (molecules, fluid, solid)

Structure of atoms

Interaction electromagnetic waves - atoms

described by **quantum mechanics**

(Quantum mechanics was discovered by studying this subject)

Light mainly interacts with the outer (valence) electrons in atoms/molecules/solid
these **electrons absorb or emit the energy**

Particles as waves

Louis de Broglie

	photon		particle
momentum	$\vec{p} = \hbar \vec{k}$	wavelength	$\vec{p} = \hbar \vec{k} = m \vec{v}$
energy	$E = \hbar \omega$	frequency	$E = \hbar \omega$

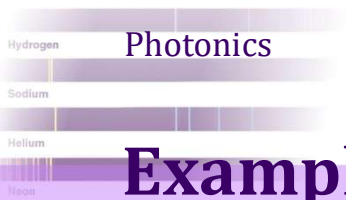
Assign a **wavelength** and frequency to a **particle**

probability finding a photon
 proportional to
 amplitude **electric field squared**
 (= power)

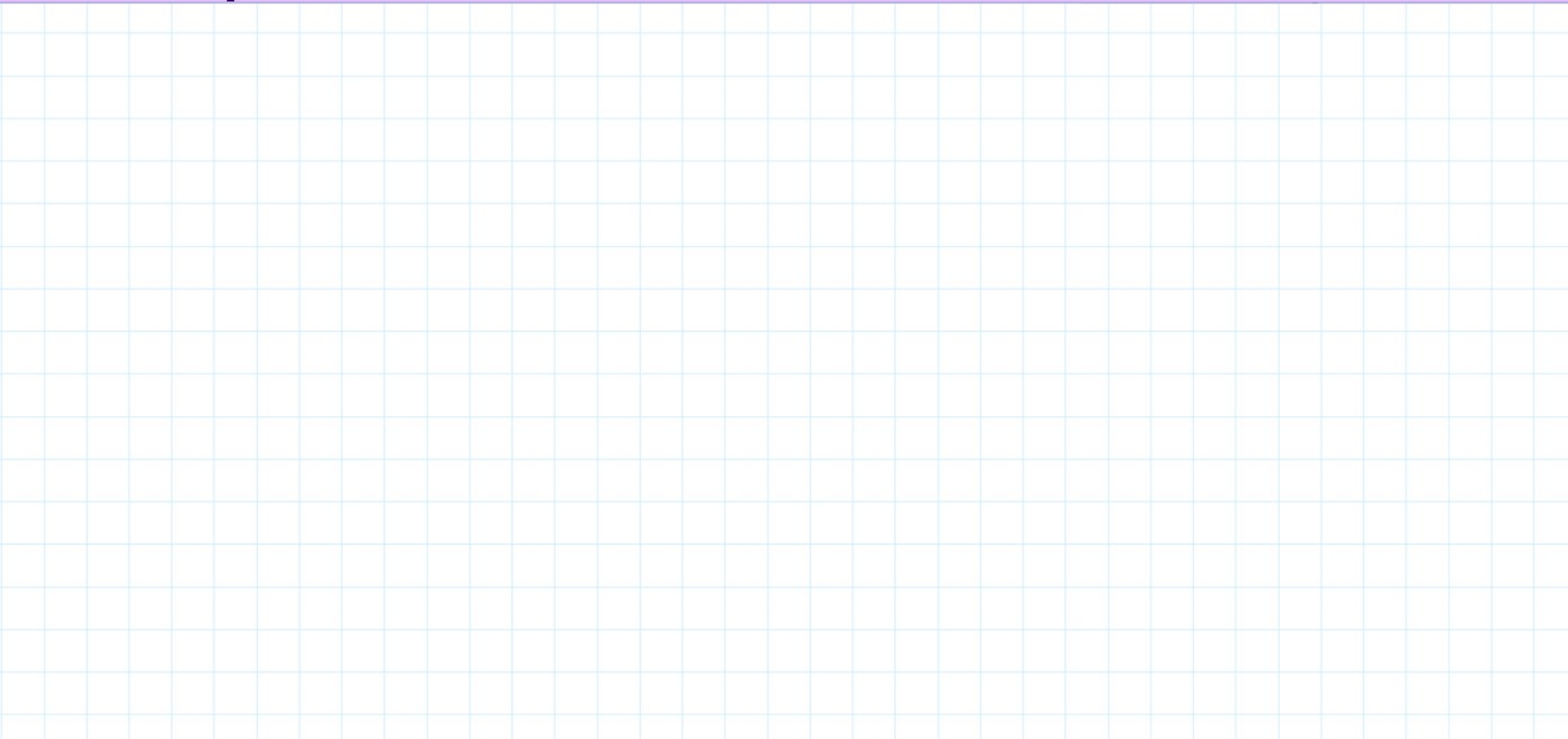
electric field is solution of
 the **classical wave equation**

probability finding a particle
 proportional to
 amplitude **wave function squared**

wave function is solution of
 the **Schrödinger equation**



Example 10.1



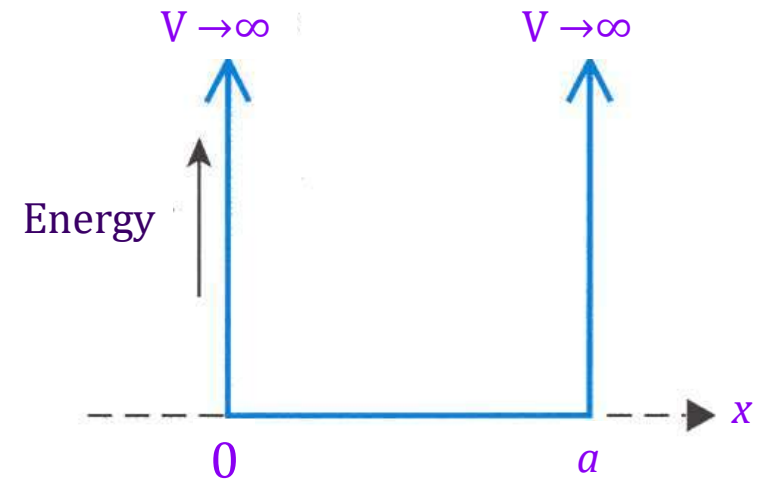
Particle in an infinite potential well 1dim

Particle mass m

Potential energy $V(x)$ (time independent)

$$V(x) = 0 \text{ for } 0 < x < a$$

$$V(x) \rightarrow \infty \text{ for } x \leq 0 \text{ and } x \geq a$$



This is a **bound** system

The particle (e.g. electron) cannot be outside $0 < x < a$

The particle has only **kinetic** energy

$$E_{\text{kin}} = \frac{p^2}{2m} = \frac{\hbar^2 k^2}{2m} = \frac{h^2}{\lambda^2 2m}$$

Particle in an infinite potential well 1dim

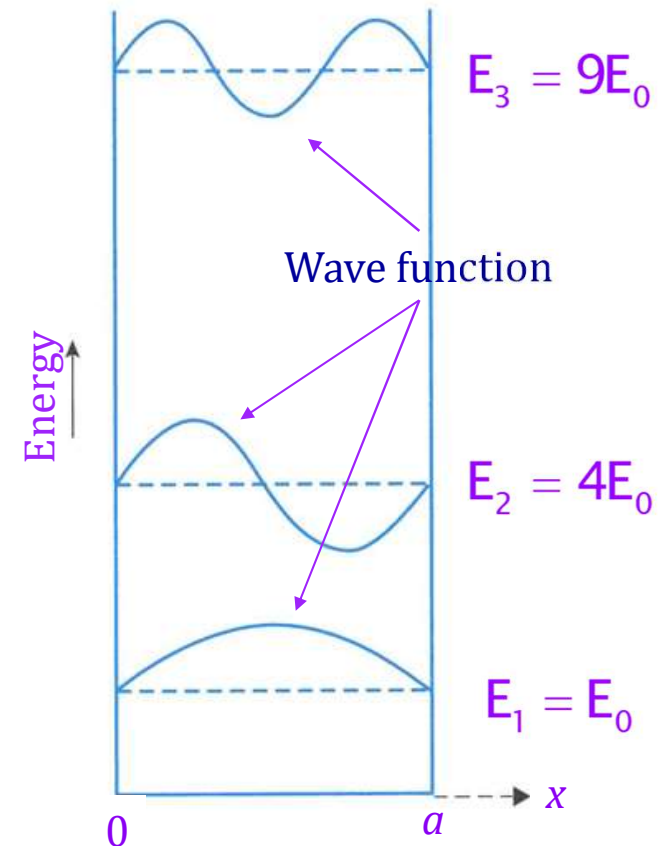
Solution: Particle wavelength has to fit in the potential well like a standing wave in a resonator

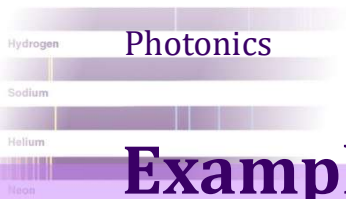
$$n \frac{\lambda}{2} = a \quad n = 1, 2, 3, \dots$$

$$E_n = \frac{p^2}{2m} = \frac{h^2}{\lambda^2 2m} = E_0 n^2 \quad E_0 \equiv \frac{\hbar^2 \pi^2}{2ma^2}$$

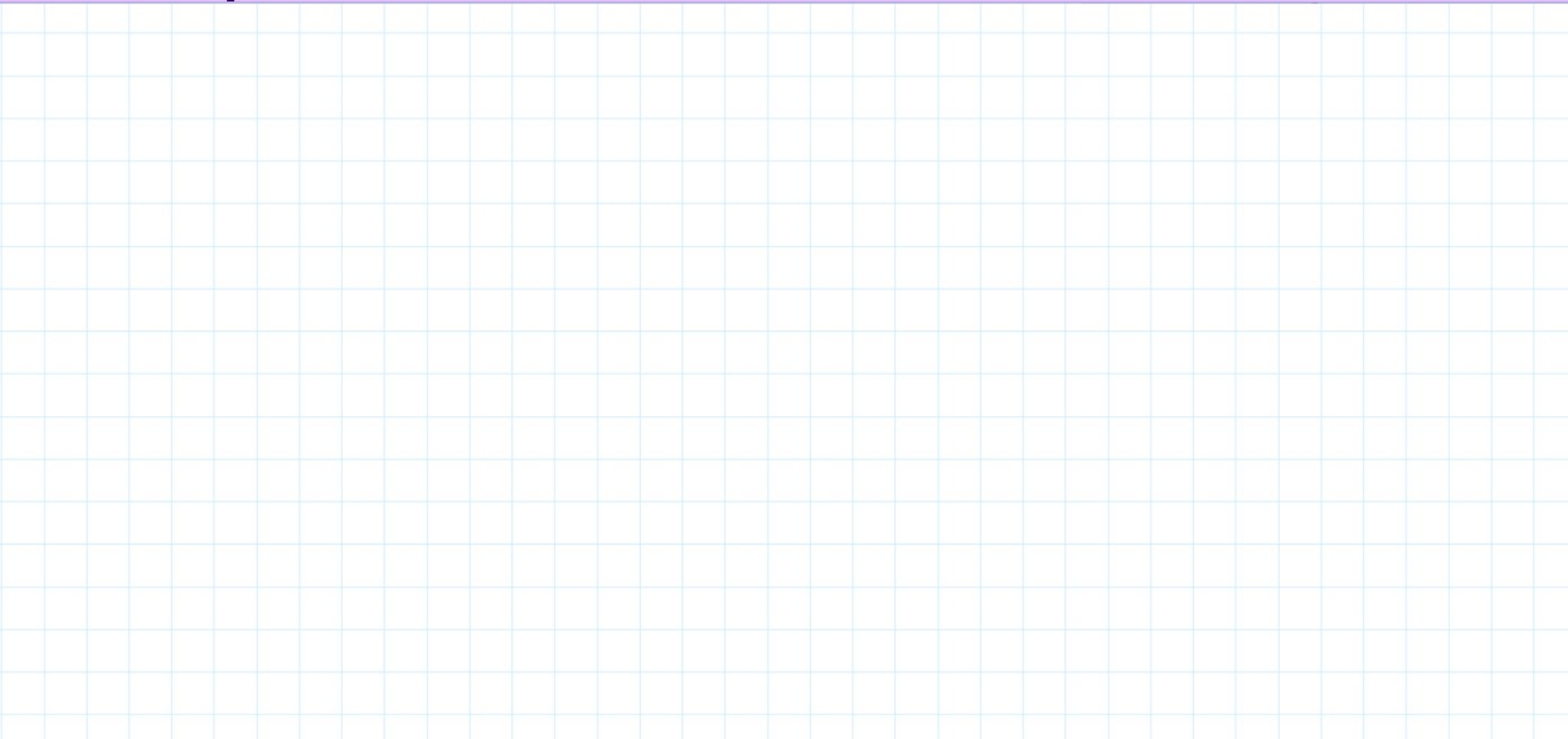
Bound system:

- Energy quantization (counting states)
- States identified by n : quantum number
- Lowest energy (ground state) $E_0 > 0$
- Values of the energy level depend on $V(x)$





Example 10.2



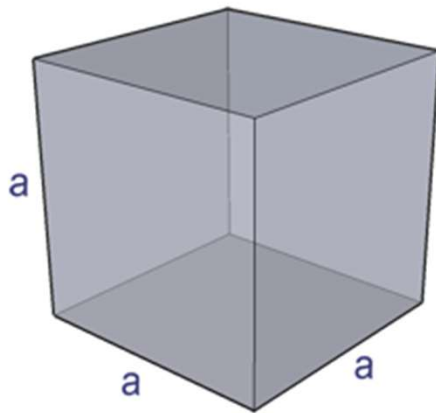
Particle in 3D potential well (a cube)

Solutions are like EM plane waves in a 3D resonator

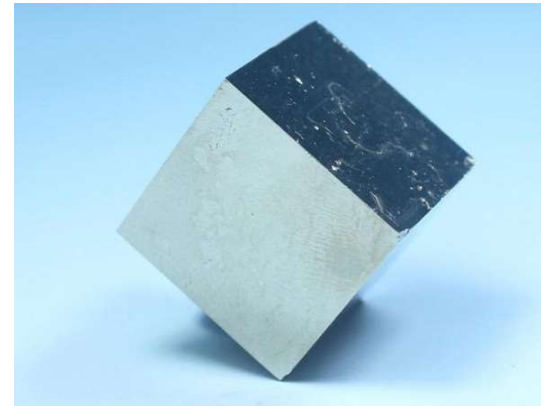
Energy levels of particle

$$E_{n_1, n_2, n_3} = E_0 (n_1^2 + n_2^2 + n_3^2) \quad E_0 \equiv \frac{\hbar^2 \pi^2}{2ma^2}$$

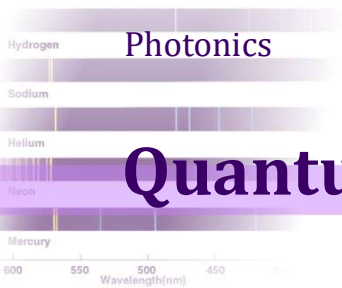
Three quantum numbers n_1, n_2, n_3 , characterise each state in 3D well
one quantum number per degree of freedom



http://www.onlineconversion.com/object_surfacearea_cube.htm



<https://m.dhgate.com/product/the-cube-natural-mineral-crystal-natural/373597893.html>



Quantum states – energy levels

Each quantum **state** of a system is characterized by its set of quantum numbers.

One quantum number per degree of freedom of the system

E.g. (n_1, n_2, n_3)

Each state of a system means that it has a specific energy, we call this:
the energy level of the system.

Multiple states of the system (multiple sets of quantum numbers) can have the same energy (energy levels).

The values of energies of the system in the different states depend on the specifics of the system (potential energy of particle as a function of its position)

Density of states 3D potential well (cube)

If the cube becomes large ($a \gg \lambda$ electron) (e.g. several cm³)
 E_0 becomes very small (but still discrete energy states)

In practice we then don't count the individual states any more but:

Use the concept of density of states: $g(E)$

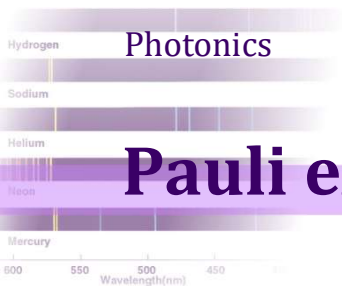
If N is the number of states at energy E in an interval ΔE

$$N \approx g(E) \cdot \Delta E$$

For the 3D potential well (cube) the density of states is

$$g(E) = C\sqrt{E} \quad C = \frac{\pi}{2E_0^{3/2}} = \frac{a^3(2m)^{3/2}}{2\hbar^3\pi^2} = Volume \cdot \frac{(2m)^{3/2}}{2\hbar^3\pi^2}$$

Density of states per unit volume is independent of size or shape!



Pauli exclusion principle

- Atoms, molecules and solid state crystal: electron occupy states.
- Question:
 - How many electrons in each state within one quantum system?

Pauli exclusion principle:

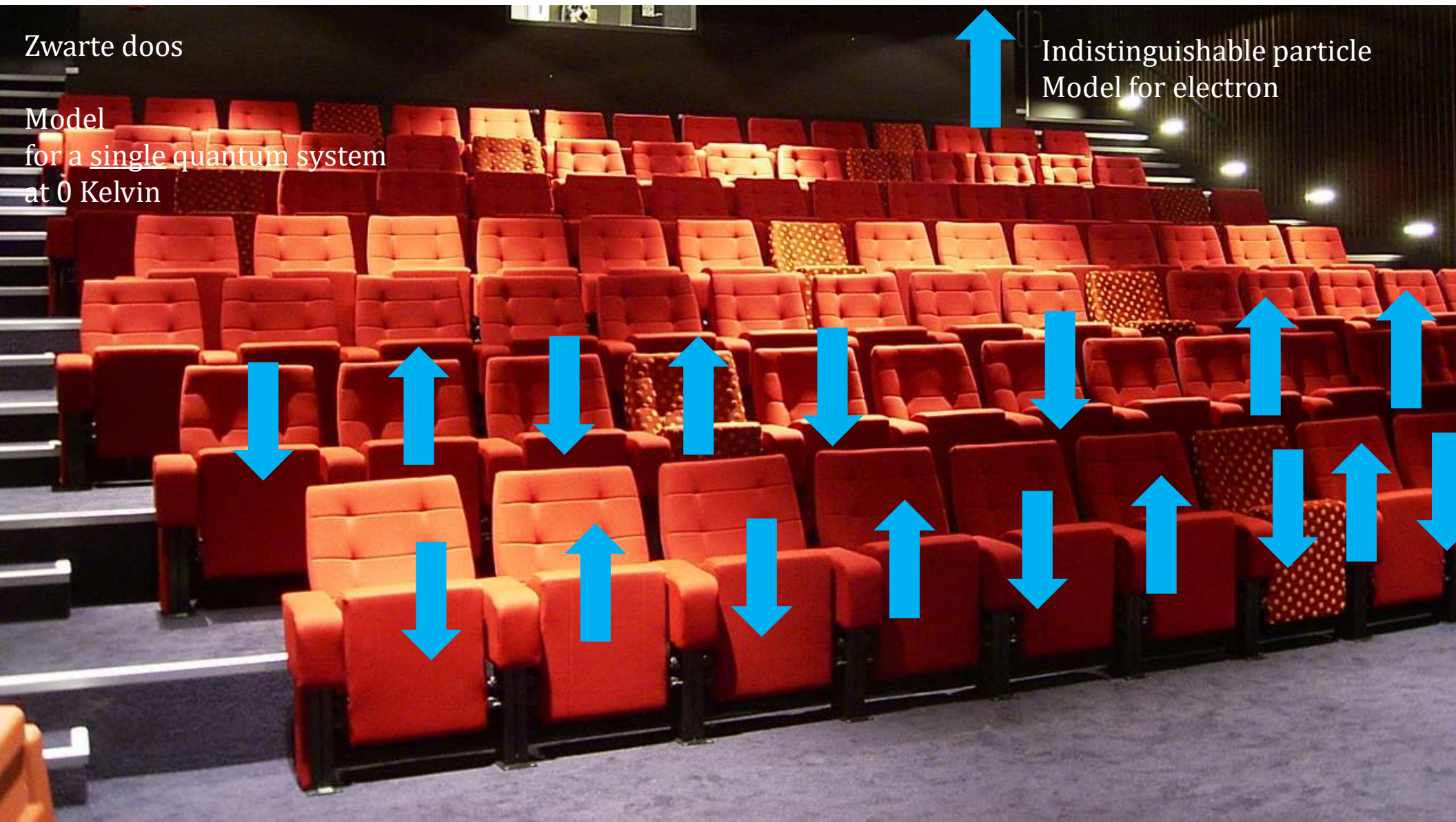
no two identical particles (of the category fermions) can be in the same **state** of a **quantum system**

- each **state** can be occupied by only **two electrons**
(one with spin up, spin down)
this makes each energy state is actually two states with (nearly) the same energy

Zwarte doos

Model
for a single quantum system
at 0 Kelvin

Indistinguishable particle
Model for electron



Energy levels of atoms

- Atom is a spherically symmetric system → potential energy of electron $V(\mathbf{r})$ only depends on \mathbf{r}
- Hydrogen atom** has simple structure of the energy levels:

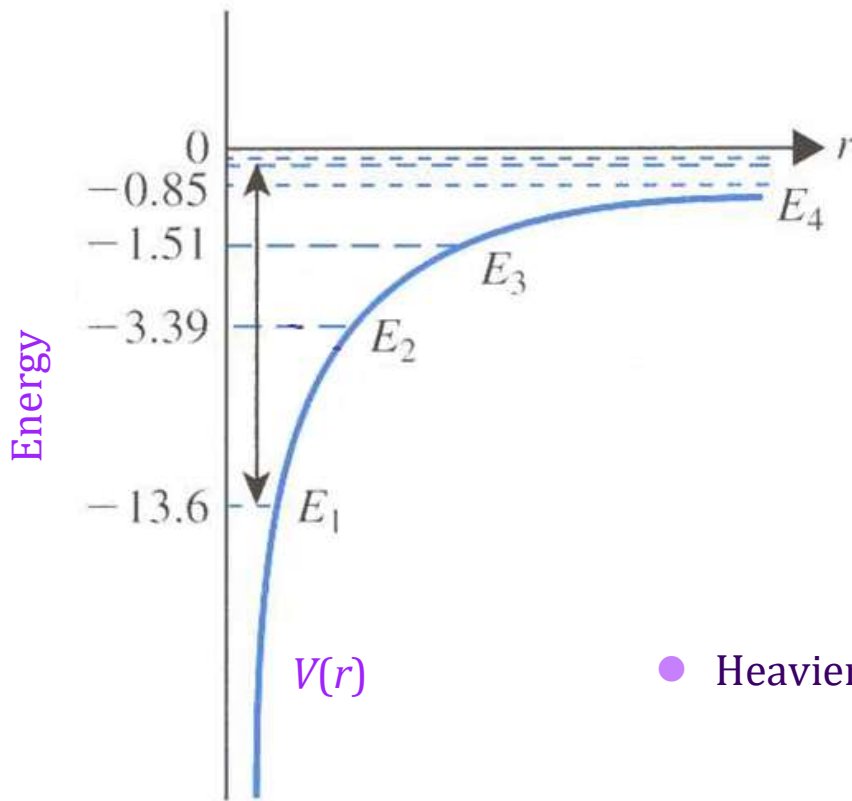
$$E_n = -\frac{m_r e^4}{2\hbar^2 n^2} = E_1 \frac{1}{n^2} \quad n = 1, 2, 3, \dots$$

- m_r is the reduced mass of the system.
 - e is the elementary charge.
 - Energy difference between lower states in the order of several eV.
→ visible light
- $$m_r = \frac{m_e \cdot M_N}{m_e + M_N} \approx m_e \quad (m_e \ll M_N)$$

Hydrogen atom – electronic states

$$E_n = -\frac{m_r e^4}{2\hbar^2 n^2} = E_1 \frac{1}{n^2} \quad n = 1, 2, 3, \dots$$

$$E_1 = -13.6 \text{ eV}$$

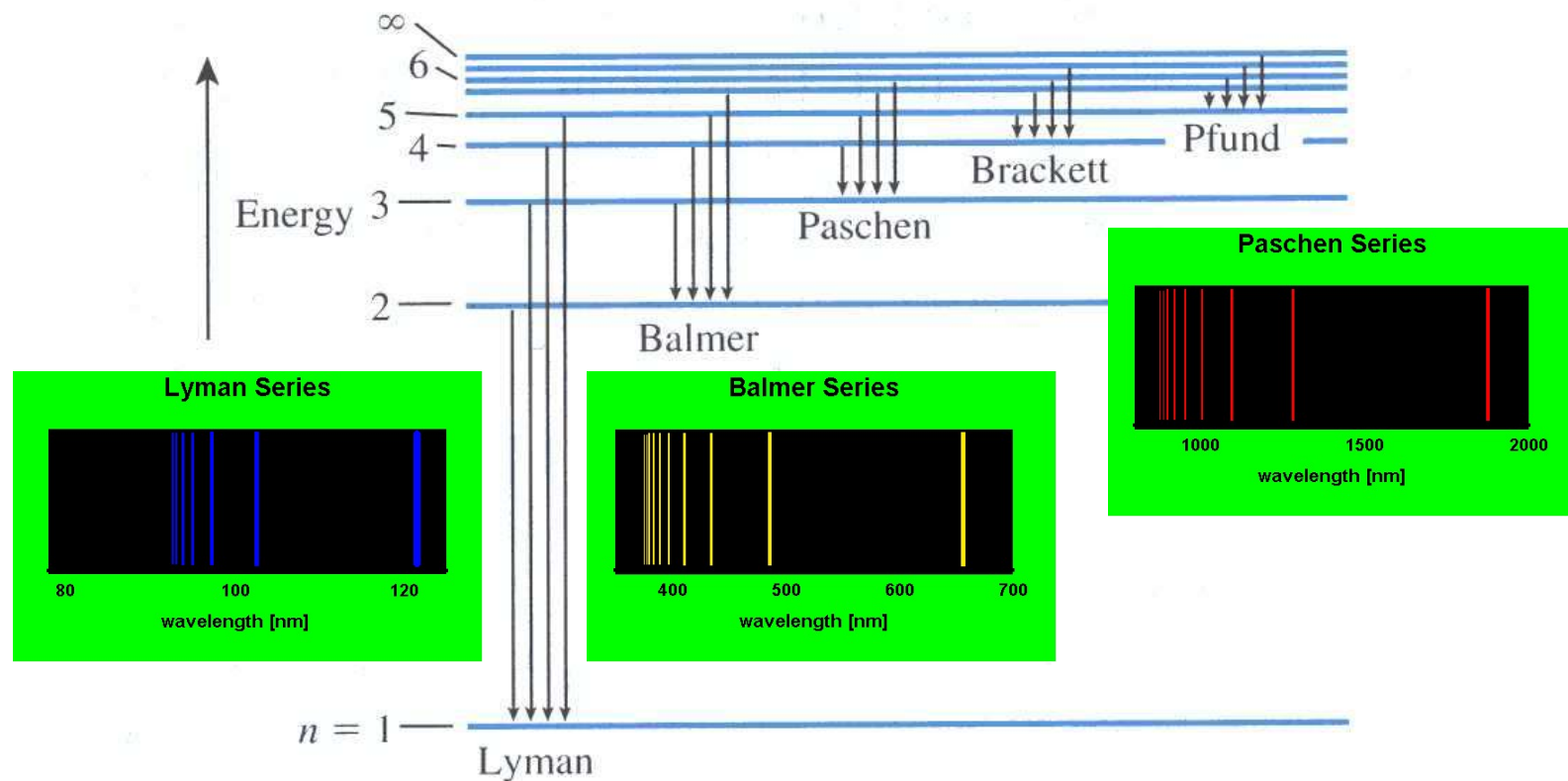


- Finite potential well – Coulomb attraction
- Infinite number of states
- Distances between lower levels
several 0.1 eV -10eV
interaction with photons
in the NIR, VIS, UV ranges
- Not all transitions allowed

- Heavier atoms have a more complex energy state diagram

Hydrogen emission lines

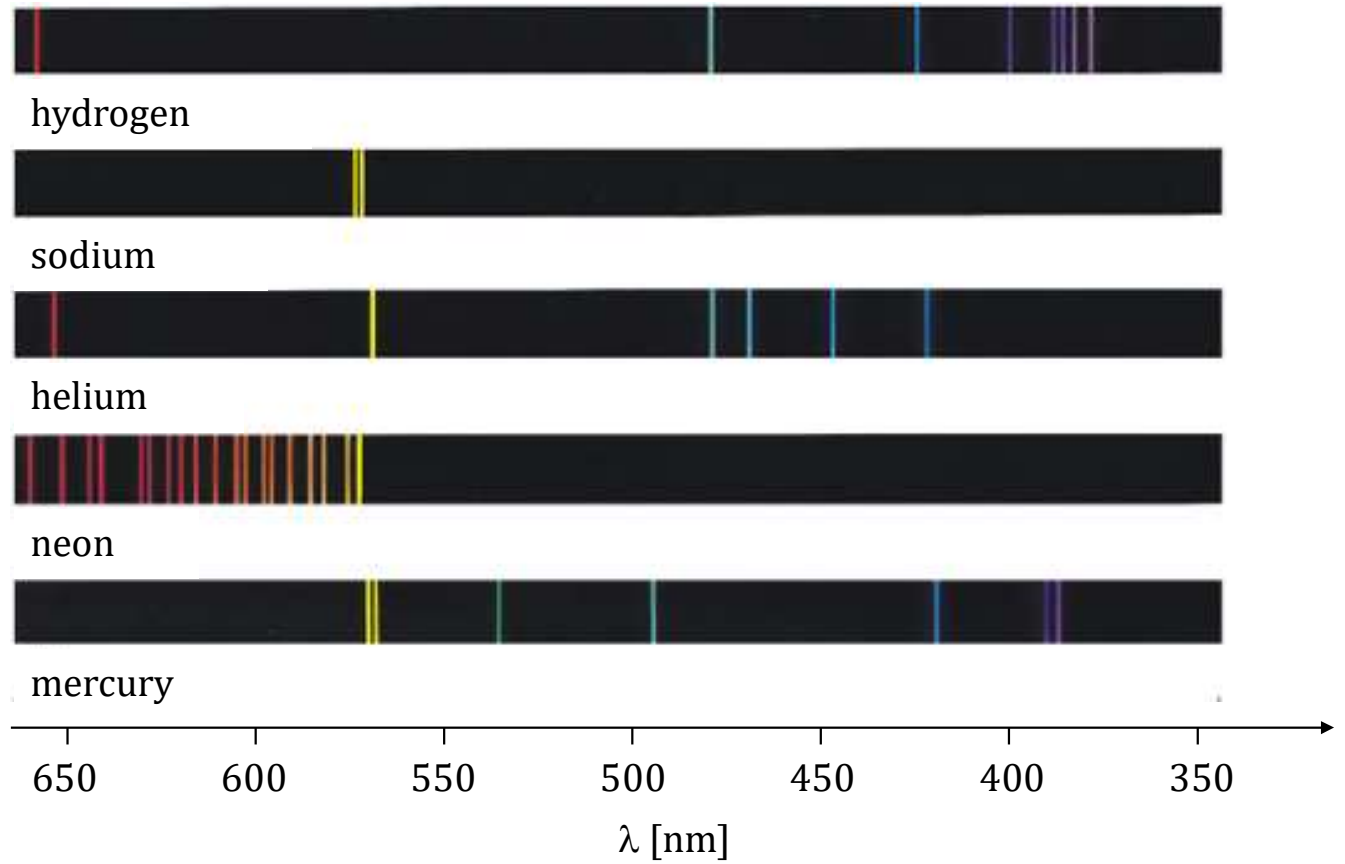
Transitions: change in energy in the atom = photon energy



Spectral lines

Every atom

- Different energy levels
- Fingerprint of atom



Energy levels in molecules (multiple atoms in one quantum system)

Electronic states

- Spacing: 1-10 eV

Vibrational states

- Atoms are vibrating
~Harmonic oscillator

$$E_v = \left(v + \frac{1}{2}\right) \hbar \omega_{res} \quad v = 1, 2, 3, \dots \quad \omega_{res} = \sqrt{\frac{k_s}{m_r}}$$

- Spacing: 0.05-0.5 eV

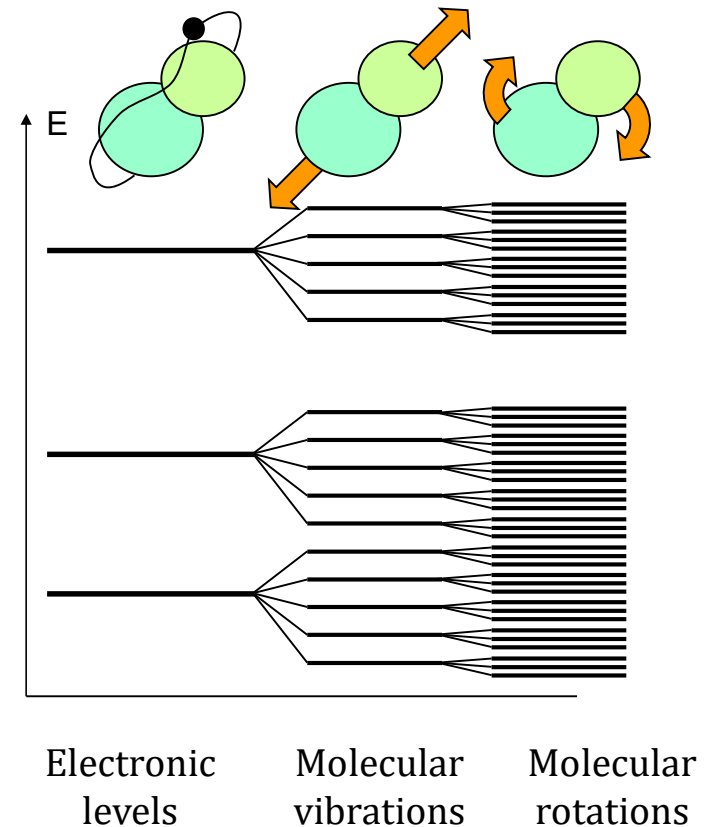
Rotational states

- Molecule rotating around axis

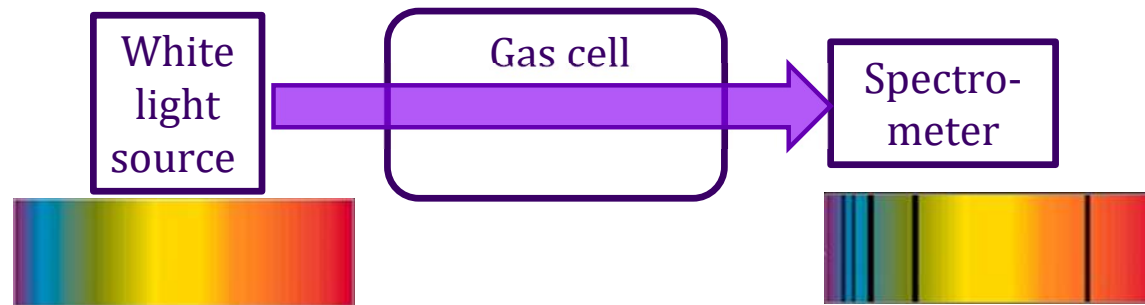
$$E_r = r \left(r + \frac{1}{2}\right) \frac{\hbar^2}{2I} \quad r = 1, 2, 3, \dots \quad I = \text{inertia}$$

- Spacing: 0.001-0.01 eV

- Not all transitions are allowed



Boltzman statistical distribution



Observation: Absorption by molecules from states above ground state.

Question: Which atoms or molecules are in a state at energy E_n in equilibrium at a specific temperature T ?

Boltzmann distribution (classical) describes:

for a group of atoms or molecules that are separate systems (identical but distinguishable particles)

the probability of finding an atom or molecule in a particular state at a particular temperature.

Boltzmann distribution

- Gas (no interaction) of identical atoms: Probability an atom is occupying level E_m

$$P(E_m) = A e^{-\frac{E_m}{k_B T}}$$

■ k_B = Boltzmann constant

■ A = constant such that:

$$\sum_m P(E_m) = 1$$

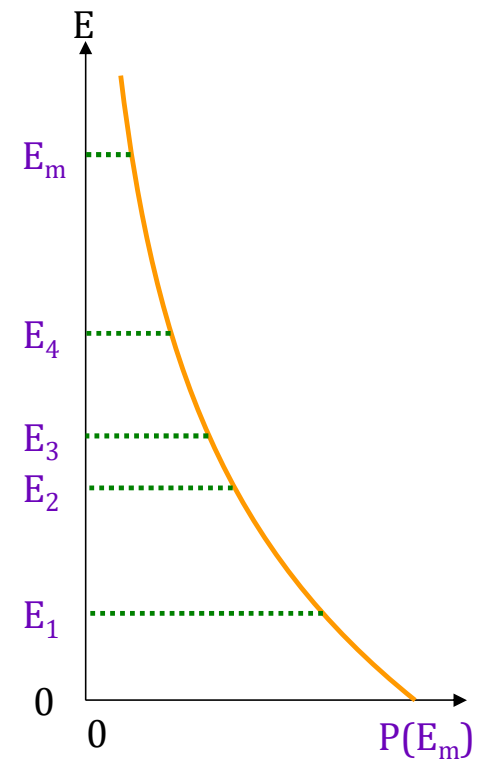
■ At temperature T in thermal equilibrium

- many (N) atoms, #atoms in state at E_m : $N_m = A N e^{-\frac{E_m}{k_B T}}$

- Population ratio: $\frac{N_i}{N_j} = \frac{g_i}{g_j} e^{-\frac{E_i - E_j}{k_B T}}$

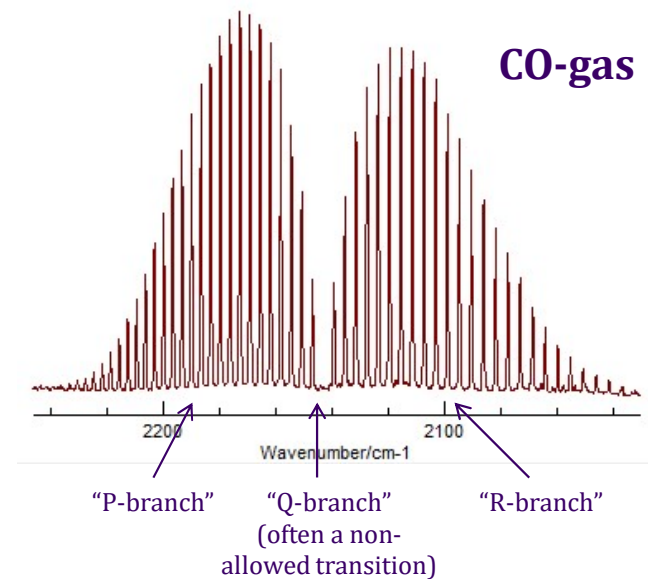
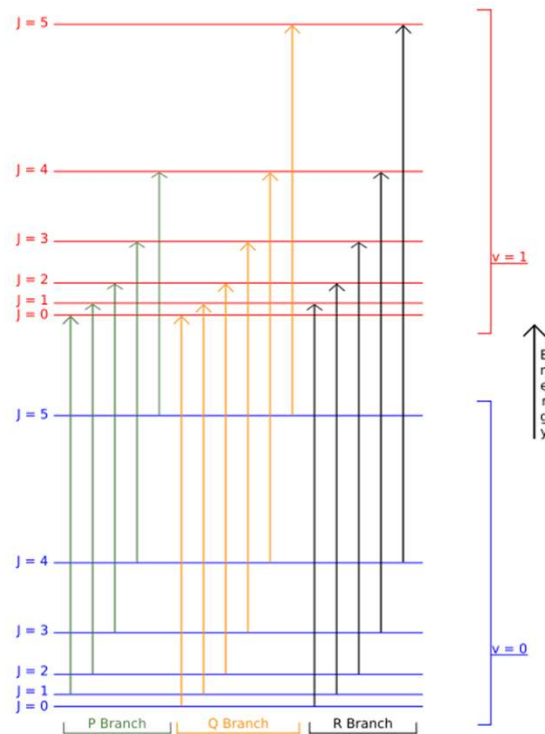
with degeneracy g_i and g_j .

$$E_1 < E_2 < E_3 < E_4 < \dots < E_m$$



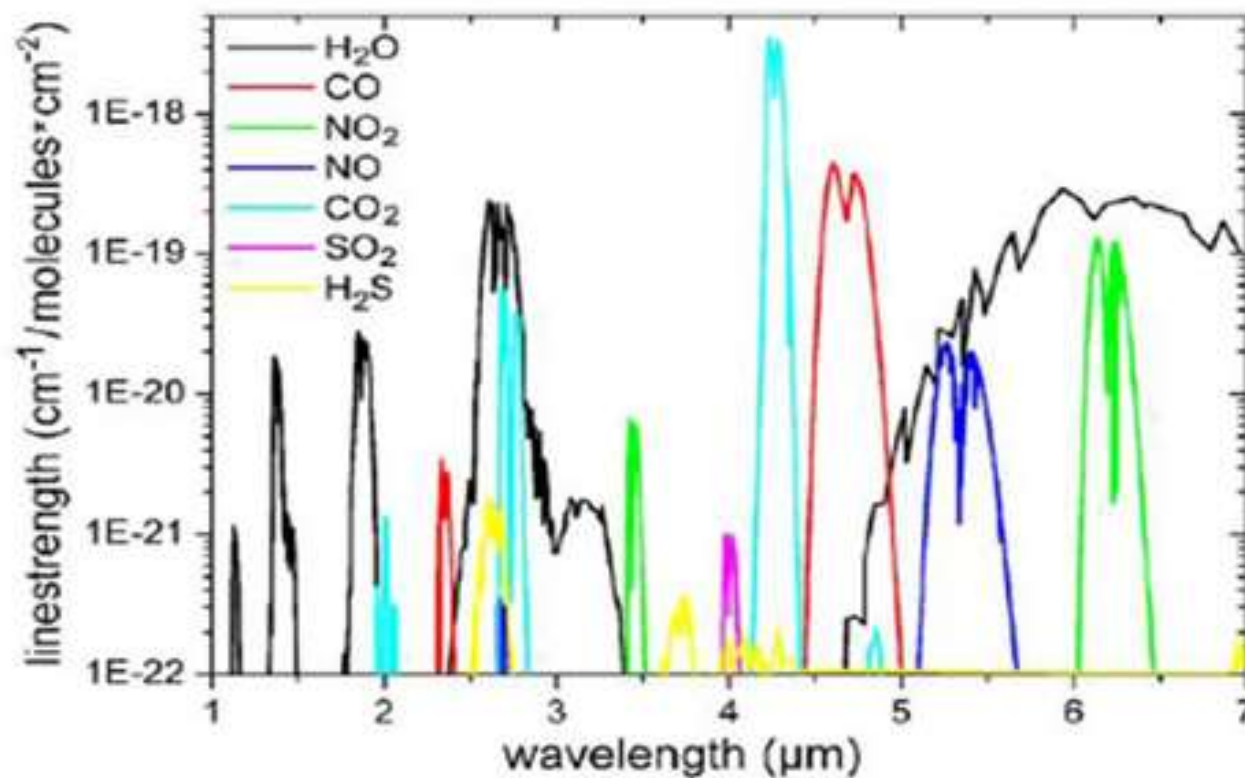
Rotational absorption spectrum

- Every vibrational transition consist of a fine structure of transitions between different rotational states



Vibrational absorption spectrum

- Infrared absorption lines caused by vibrational transitions form fingerprint





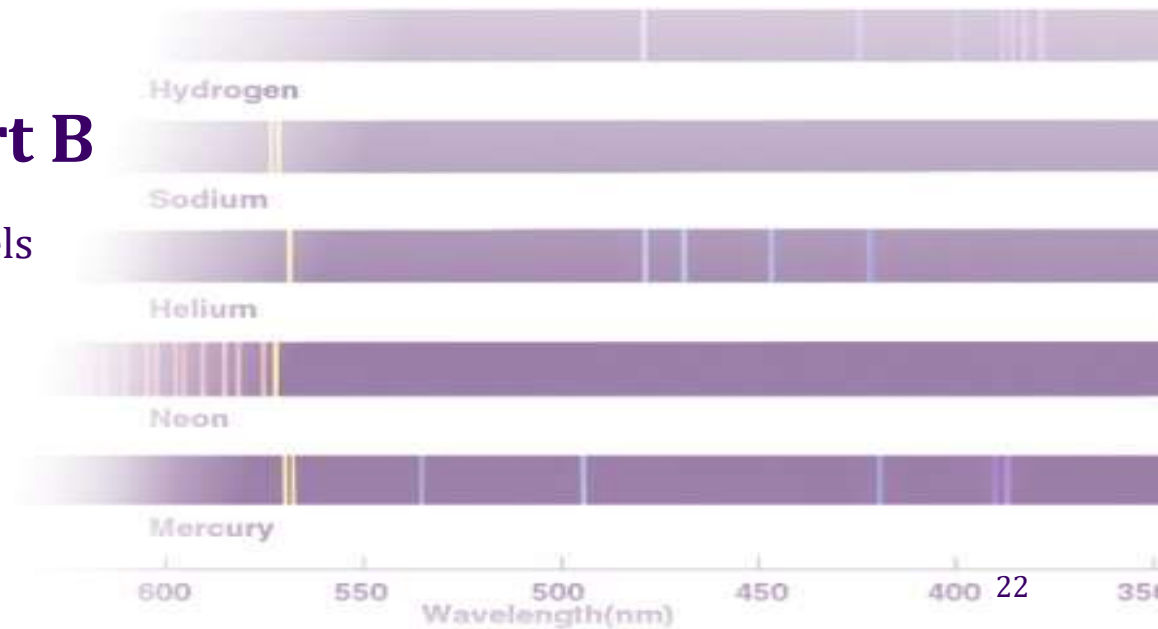
Photonics

R. Baets / E. Bente

Photons and Atoms Part B

Atoms, molecules, solid state energy levels

Interaction with light



Solid state: multiple atoms

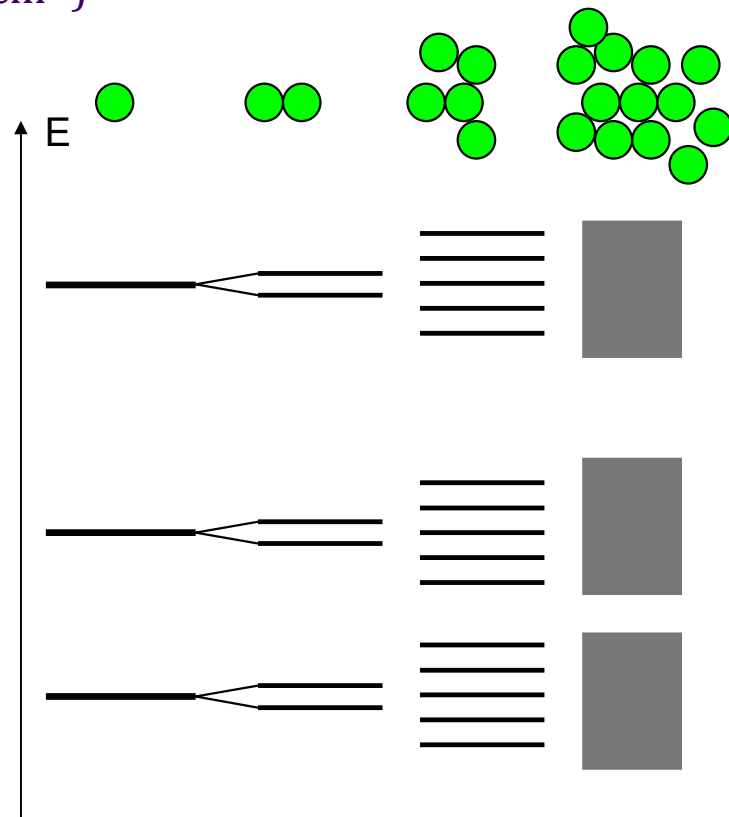
● System of N closely packed atoms (10^{23} atoms/cm³)

● Energy levels:

- Valence electrons
- Rotations, Vibrations
- Splitting of each level into N energy levels

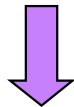


- Formation of **energy bands** if N is very large

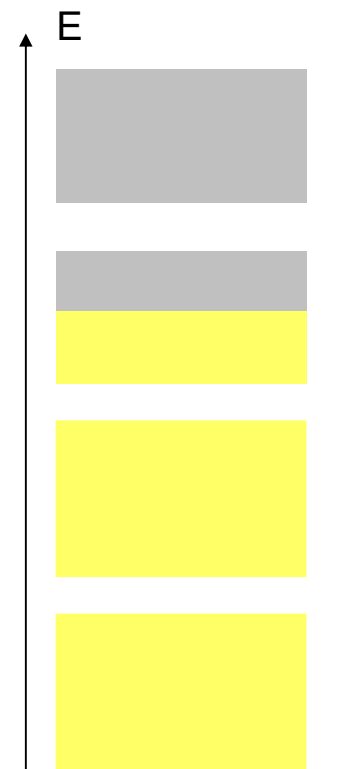


Band occupation: Metals

- Occupation of energy levels (bands)
 - Pauli: each level contains max 2 electrons
 - n electrons per atom
- single occupied levels:
 - upper band (= **valence band**) half filled: electrons can easily be excited from one level to another
 - thermal
 - by applying an electric field
 - ...



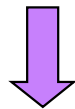
Good conduction: metals



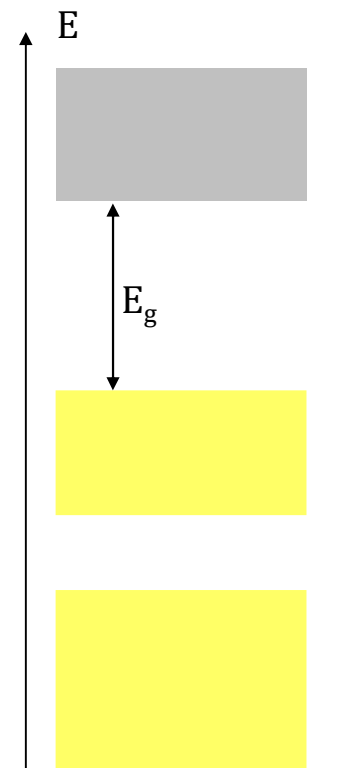
Band occupation: Isolator

- Occupation of energy levels (or bands)
 - Pauli: each level has max 2 electrons
 - n electrons per atom
- double occupied levels:
Valence band completely filled

If the distance E_g (= “Band Gap”) is large:
No excitation of electrons possible

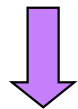


No conduction: ISOLATOR

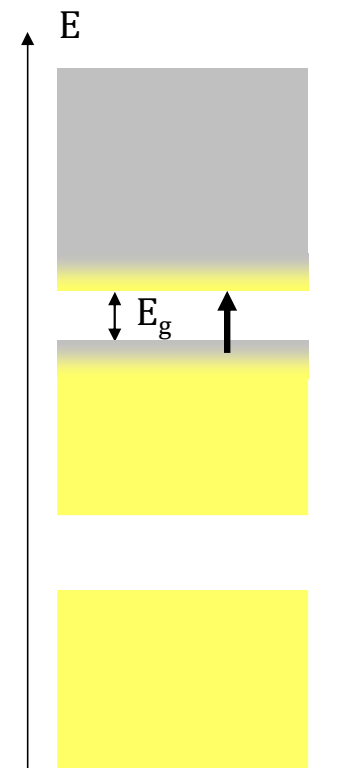


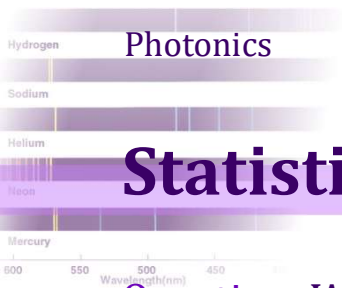
Band occupation: Semiconductor

- Occupation of energy levels (or bands)
 - Pauli: each level has max 2 electrons
 - n electrons per atom
- double occupied levels:
Valence band completely filled
- When the distance E_g (= "Band Gap") small:
Some electrons can be excited to the **conduction band**
where they can move freely
 - Thermal excitation
 - Excitation through electric field
 - ...



Some conduction: Semiconductor





Statistics electrons over states in a solid

Question: What is the probability to find an electron in an energy state at energy E ?

- Electrons in solid state:
the electrons are **fermions** and follow the **Pauli exclusion principle**
they have a strong interaction and form a single quantum system
the electrons are **indistinguishable**
- The electrons therefore **do not** distribute themselves over the states according to the **Boltzmann distribution**. That distribution allows more than one electron in each state.
- The electrons distribute themselves over the states according to the

Fermi-Dirac distribution.

Fermi-Dirac distribution

- Electrons in solid state: **indistinguishable** particles in a system

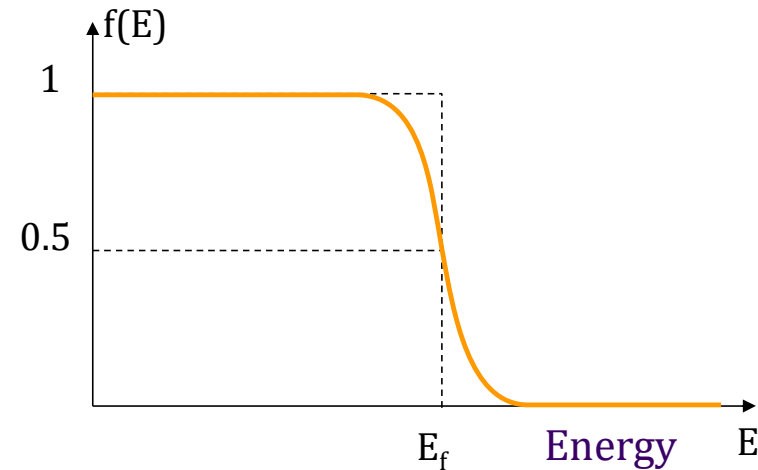
- Probability of the occupation of an energy level E :

$$f(E) = \frac{1}{e^{\frac{E-E_f}{k_B T}} + 1} \quad \text{where } E_f \text{ is the Fermi-energy } f(E_f) = 0.5$$

- For $E \gg E_f$:

$$f(E) \approx e^{-\frac{E-E_f}{k_B T}}$$

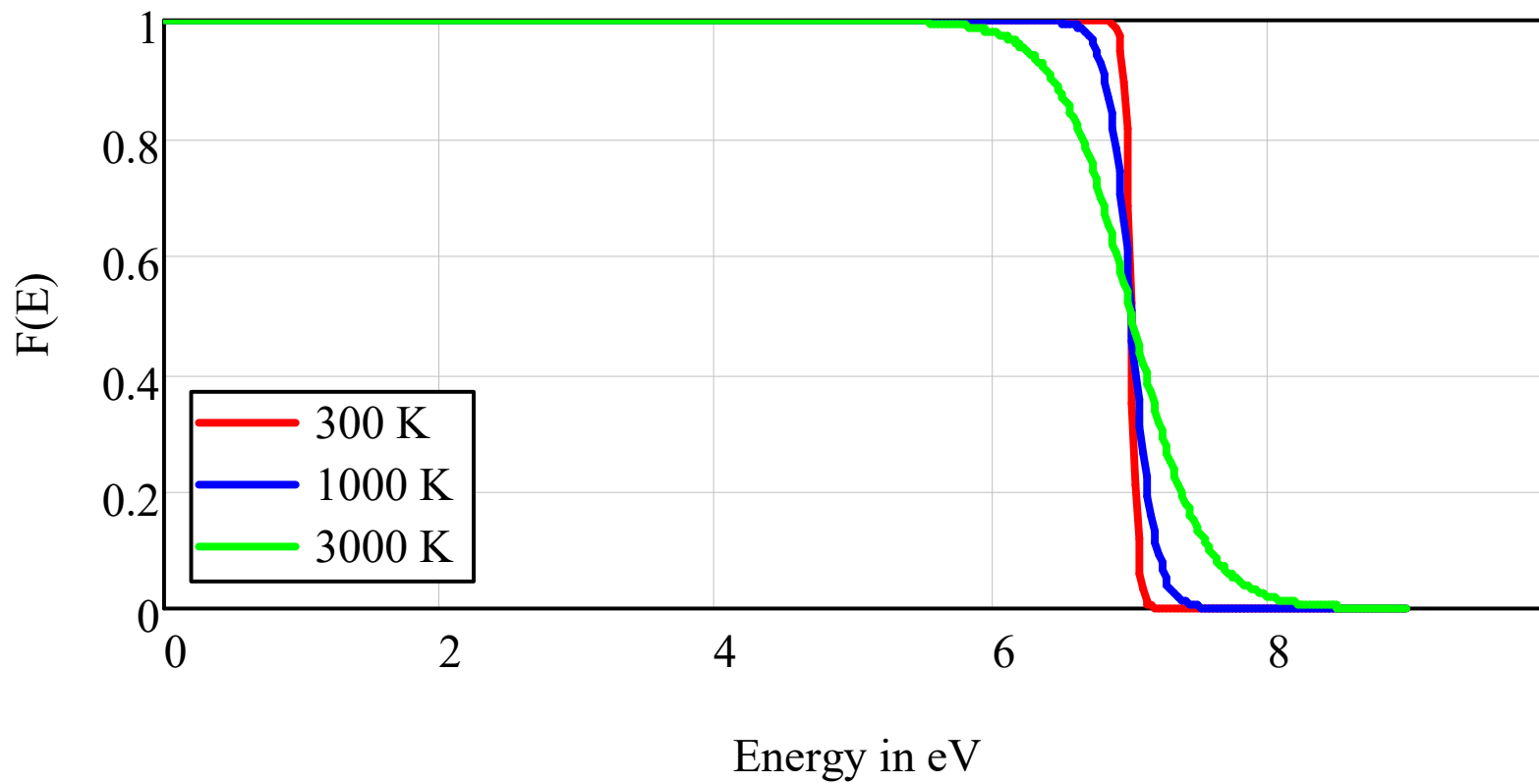
(Boltzmann approximation)



Example Fermi function

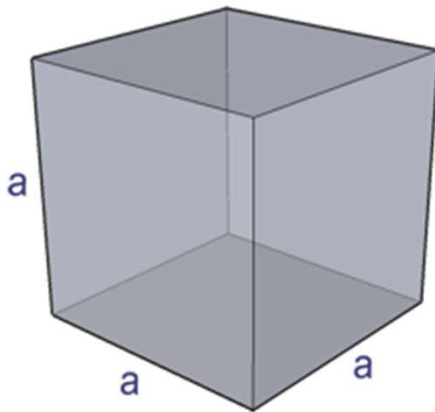
$k_B T = 0.0259$ eV at $T = 300$ K

$E_f = 7.0$ eV

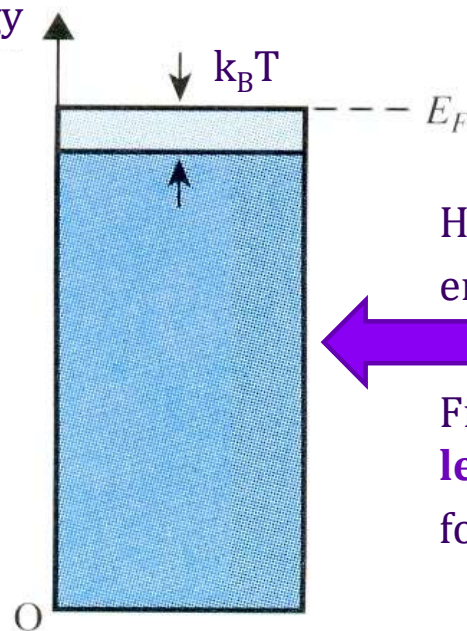


3D potential well – model for a metal

Macroscopic block of metal
e.g. $a = 1\text{cm}$



Energy



Huge number of
energy states

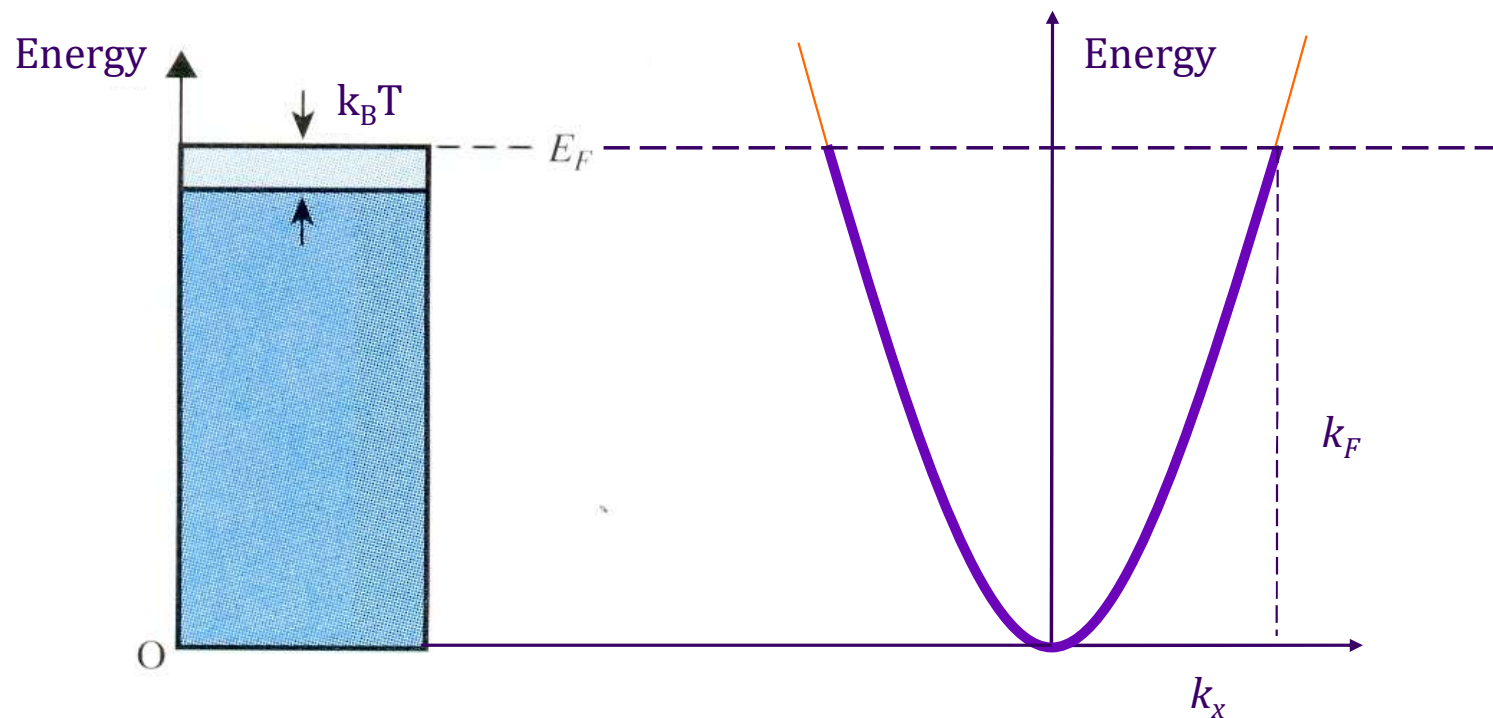
Filled with electrons up to the **Fermi level**
following the Fermi-Dirac distribution

- Density of electrons at energy E : $N(E)dE = g(E)f(E)dE$
- Such a model will be used in Chapters 14 and 15

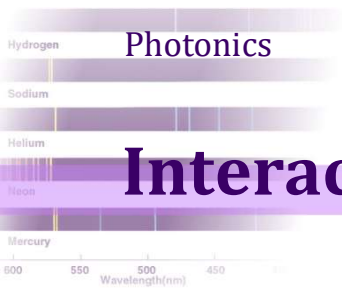
3D potential well – E-k relation

- All electron energy is kinetic energy

$$E = E_{kin} = \frac{p^2}{2m} = \frac{\hbar^2}{2m} k^2$$



- Such a model will be used in Chapters 14 and 15



Photonics

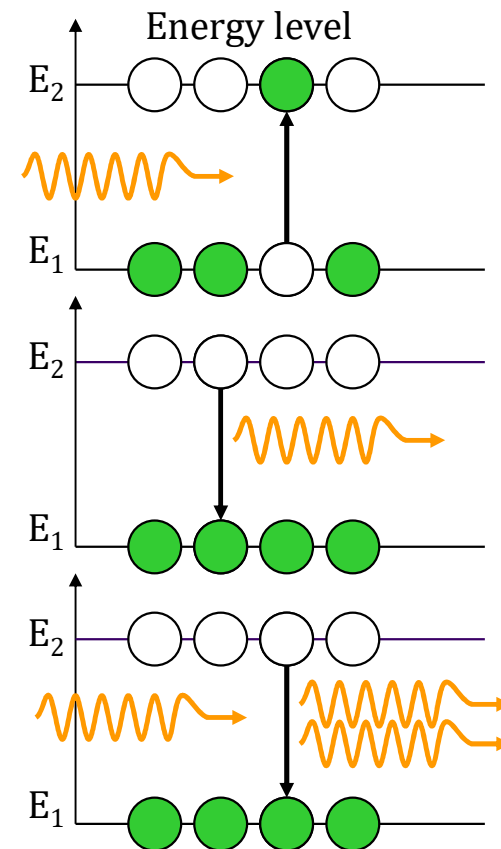
Photons and Atoms

Interaction between matter and light

- Basic processes two-energy level systems
- Relation between those processes

Emission and absorption

- Transition between atomic or molecular levels: emission or absorption of a photon
- Absorption:
 - Photon hits the system
 - System brought into excited state
- Spontaneous emission:
 - System goes back to lower state
 - Emission of photon
- Stimulated emission
 - Photon hits system
 - System goes back to lower state
 - Emission of identical photon



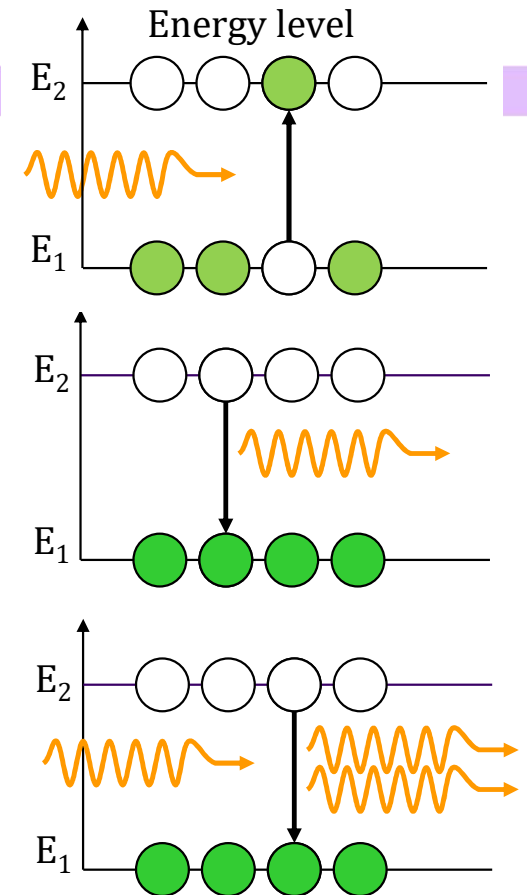
Einstein description

- Probability (per s) that an atom absorbs one photon from a specific EM mode and thereby goes to an excited state (p_{abs})
- =
- Probability (per s) that an atom spontaneously decays with the generation of one photon in a specific EM mode (p_{sp})
- =
- Probability (per s) that an atom falls back to the ground state under the influence of one photon in a specific EM mode (p_{st})

$$p_{sp} = p_{abs} = p_{st} \equiv \frac{c}{V} \sigma(\nu)$$

Cavity volume

transition cross-section



Spontaneous emission

Probability of transition with spontaneous emission of **one photon** in **one mode of cavity** with volume **V**

$$p_{sp} = \frac{c}{V} \sigma(\nu)$$

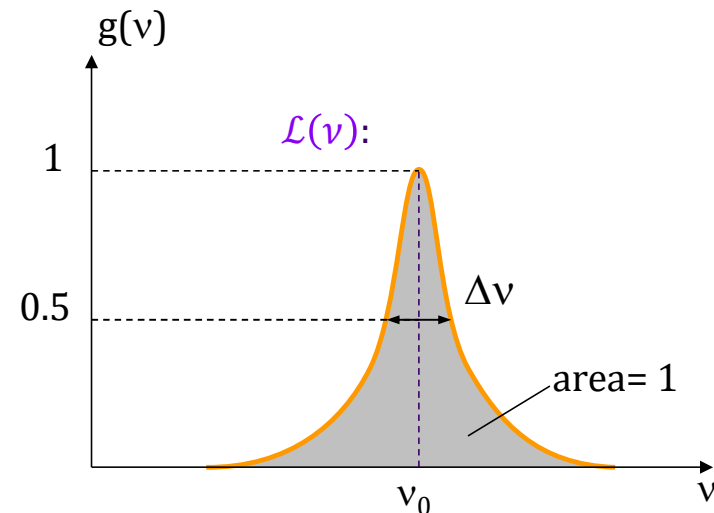
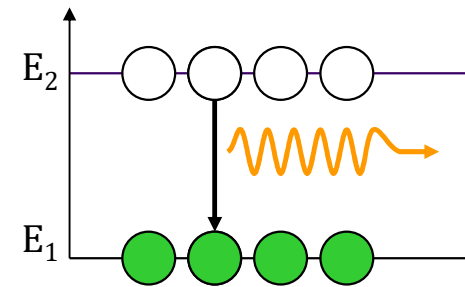
- With $\sigma(\nu)$ a function centered around

$$\nu_0 = \frac{E_2 - E_1}{h} \quad \leftarrow \text{Planck's constant } 6.626 \times 10^{-34} \text{ J}\cdot\text{s}$$

- After normalization:
line shape function $\mathcal{L}(\nu)$:

$$\mathcal{L}(\nu) = \frac{\sigma(\nu)}{\int \sigma(\nu) d\nu}$$

with 'linewidth' $\Delta\nu$. (Typically $\Delta\nu \ll \nu_0$)



Spontaneous emission

- Density of EM modes in a cavity
(= # EM modes in the cavity / volume freq.)
- Total emission probability of photon:

$$M(\nu) = \frac{8\pi\nu^2}{c^3} = \frac{8\pi}{\lambda^2 c}$$

$$P_{sp} = \int p_{sp} V M(\nu) d\nu = c \int \sigma(\nu) M(\nu) d\nu \approx c M(\nu_0) \int \sigma(\nu) d\nu$$

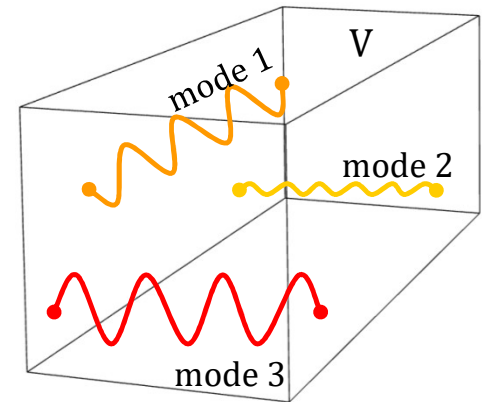
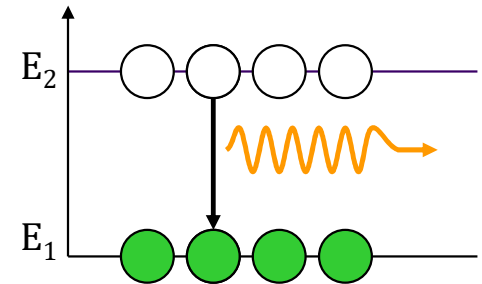
$$p_{sp} = \frac{c}{V} \sigma(\nu)$$

M constant over lineshape

$$P_{sp} \approx \frac{8\pi}{\lambda_0^2} \int \sigma(\nu) d\nu$$

= "A-coefficient of Einstein" A_{21} (independent of cavity volume V)

- Spontaneous lifetime: $\tau_{sp} = \frac{1}{P_{sp}} = \frac{1}{A_{21}}$



Stimulated emission

- Emission of an **IDENTICAL** photon triggered by an **incident photon** (basis of laser operation)
→ interaction with **only one mode**
- Probability for a transition with stimulated emission in a cavity with volume V :

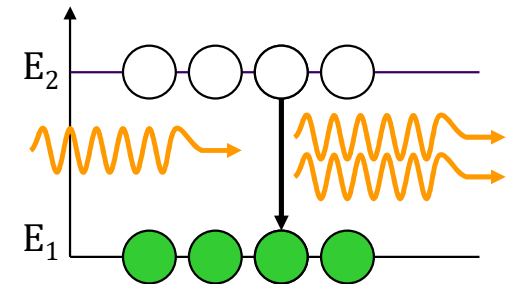
$$P_{st} = n p_{st} = n \frac{c}{V} \sigma(\nu)$$

with n the number of photons already present in that mode

- Total emission (stimulated plus spontaneous) probability:

$$p_{sp} + P_{st} = (n + 1) \frac{c}{V} \sigma(\nu)$$

└──────────┬──────────┐
└──────────┘ zero point energy of a mode
extra photons



Stimulated emission

$$P_{st} = n p_{st} = n \frac{c}{V} \sigma(\nu)$$

Interpretation:

- $n c/V$ is the FLUX DENSITY of the number of photons
(i.e. the number of incident photons per s and m²)
- $\sigma(\nu)$ is the cross section in m² for stimulated emission
- The product is the probability (per s) that stimulated emission occurs

Stimulated emission

● Cavity with spectral energy density $\rho(\nu)$ (per volume, per frequency).

photons in $[\nu, \nu + d\nu]$ in cavity $\frac{\rho(\nu)}{h\nu} V d\nu$

$$P_{st} = \int p_{st}(\nu) \frac{\rho(\nu)}{h\nu} V d\nu = c \int \frac{\rho(\nu)}{h\nu} \sigma(\nu) d\nu$$

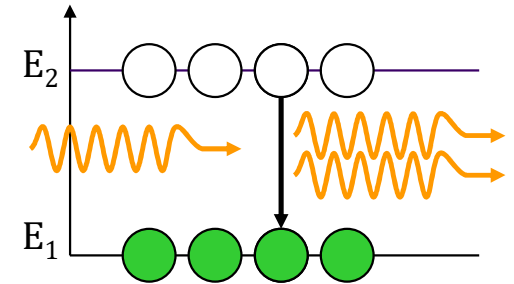
ρ const over lineshape

$$\approx c \frac{\rho(\nu_0) \lambda_0}{h} \int \sigma(\nu) d\nu = \frac{\lambda_0^3}{8\pi h \tau_{sp}} \rho(\nu_0)$$

$$P_{sp} = \frac{8\pi}{\lambda_0^2} \int \sigma(\nu) d\nu \equiv \frac{1}{\tau_{sp}} = A_{21}$$

$$= \frac{\bar{n}}{\tau_{sp}} = B_{21} \rho(\nu_0)$$

- with $\bar{n} = \frac{\lambda_0^3}{8\pi h} \rho(\nu_0)$ the average number of photons per mode $M(\nu) = \frac{8\pi \nu^2}{c^3} = \frac{8\pi}{\lambda^2 c}$
- With B_{21} the “Einstein B-coefficient”. $B_{21} = \frac{\lambda_0^3}{8\pi h \tau_{sp}} = \frac{\lambda_0^3 A_{21}}{8\pi h}$



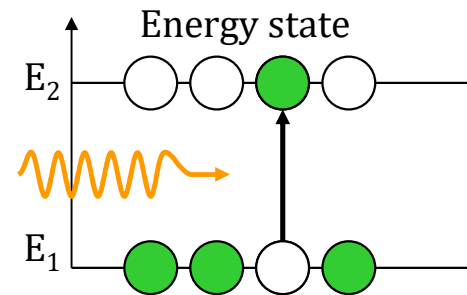
Absorption

- Probability of a transition due to absorption in a cavity with volume V :

$$P_{abs} = n p_{abs} = n \frac{c}{V} \sigma(\nu)$$

With n the number of photons present in the mode

$$\begin{aligned} P_{abs} &= P_{st} \\ &= \frac{\bar{n}}{\tau_{sp}} = B_{12} \rho(\nu_0) \\ &= \bar{n} P_{sp} \end{aligned}$$



$$B_{12} = B_{21}$$

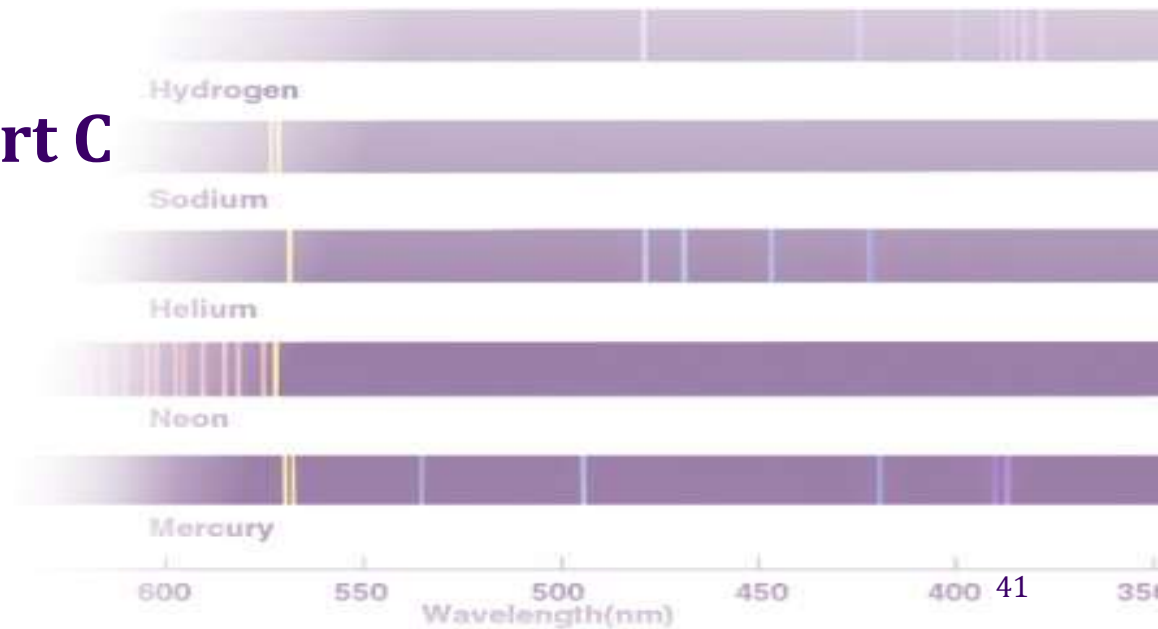
Photonics

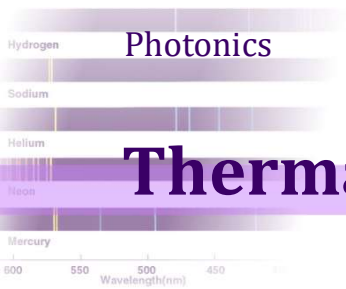
R. Baets / E. Bente

Photons and Atoms - Part C

Thermal light

Luminescence light





Photonics

Photons and Atoms

Thermal light

- Two-level atoms in thermal equilibrium with a cavity
- Black body radiation

Thermal equilibrium

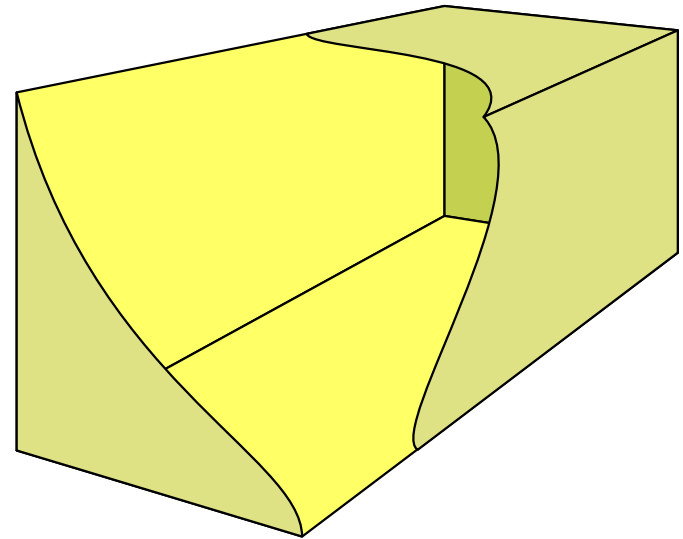
- Atoms, molecules of solids,... in thermal equilibrium: emission and absorption of photons
- Cavity with $V=1$, walls with many atoms per volume unit
 - N_1 atoms in state E_1 .
 - N_2 atoms in state $E_2 > E_1$.
 - \bar{n} photons in every EM-mode within linewidth of $\mathcal{L}(\nu)$:

$$P_{abs} = P_{st} \frac{\bar{n}}{\tau_{sp}}$$

- Thermal equilibrium:
 $\#(E_1 \Rightarrow E_2) = \#(E_2 \Rightarrow E_1)$



absorption = stimulated+ spontaneous emission



Equilibrium between material and photons

Spontaneous emission

- Proportional to N_2

$$\left. \frac{dN_2}{dt} \right|_{sp} = - \left. \frac{dN_1}{dt} \right|_{sp} = -N_2 P_{sp} = - \frac{N_2}{\tau_{sp}} = -A_{21} N_2$$

Stimulated emission

- Proportional to N_2
- Proportional to \bar{n}

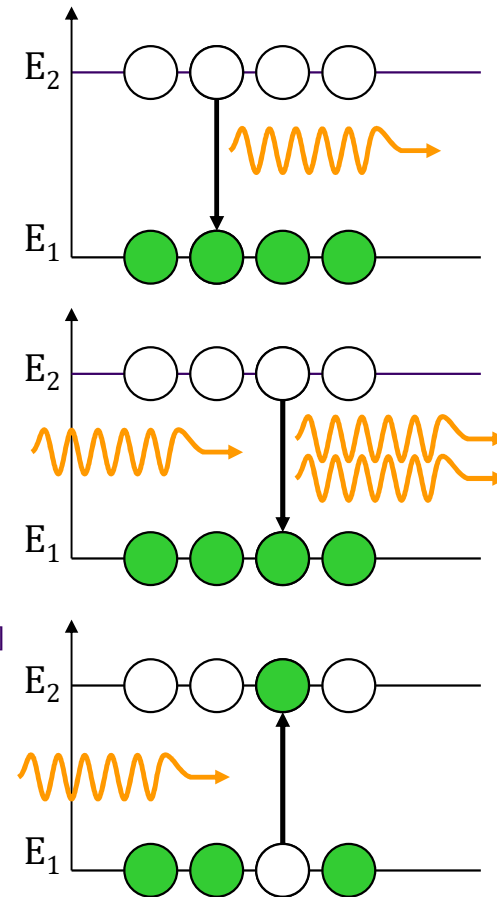
$$\left. \frac{dN_2}{dt} \right|_{st} = - \left. \frac{dN_1}{dt} \right|_{st} = -N_2 P_{st} = - \frac{\bar{n} N_2}{\tau_{sp}} = -B_{21} \rho(\nu_0) N_2$$

↑
spectral energy density [$\text{Jm}^{-3}\text{Hz}^{-1}$]

Absorption:

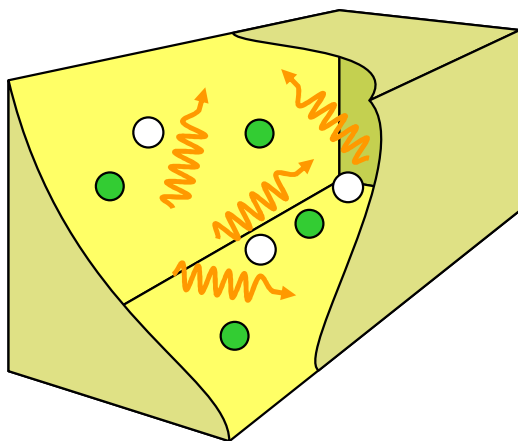
- Proportional to N_1
- Proportional to \bar{n} (photon flux)

$$\left. \frac{dN_2}{dt} \right|_{abs} = - \left. \frac{dN_1}{dt} \right|_{abs} = N_1 P_{abs} = \frac{\bar{n} N_1}{\tau_{sp}} = B_{12} \rho(\nu_0) N_1$$



Discuss

A box with 2-level atoms and photons:



\bar{n} photons in interacting mode

N_1 ground-state atoms

N_2 excited atoms

What fraction of atoms will be excited in thermal equilibrium?

Equilibrium between material and photons

- Total flux:

$$\frac{dN_2}{dt} = -\frac{dN_1}{dt} = \left. \frac{dN_2}{dt} \right|_{sp} + \left. \frac{dN_2}{dt} \right|_{ab} + \left. \frac{dN_2}{dt} \right|_{st} = \frac{1}{\tau_{sp}} (\bar{n}N_1 - N_2(\bar{n} + 1))$$

Does not take into account

- Other energy levels
- Other (non-radiative) processes

- Equilibrium: $\frac{dN_2}{dt} = -\frac{dN_1}{dt} = 0$ (absorption = emission)

$$\Rightarrow \frac{N_2}{N_1} = \frac{\bar{n}}{\bar{n} + 1}$$

Equilibrium between material and photons

- Equilibrium: $\frac{N_2}{N_1} = \frac{\bar{n}}{\bar{n} + 1}$
- Boltzmann distrib.: $\frac{N_2}{N_1} = e^{-\frac{E_2 - E_1}{k_B T}} = e^{-\frac{h\nu}{k_B T}}$
- Average number of photons with frequency ν :

$$\bar{n} = \frac{1}{e^{\frac{h\nu}{k_B T}} - 1}$$

It turns out **this expression is also valid** for

- Multiple energy levels
- continuum

Spectrum of a blackbody radiator

- Energy of a radiating mode with frequency ν :

$$\bar{E} = \bar{n}h\nu = \frac{h\nu}{e^{\frac{h\nu}{k_B T}} - 1}$$

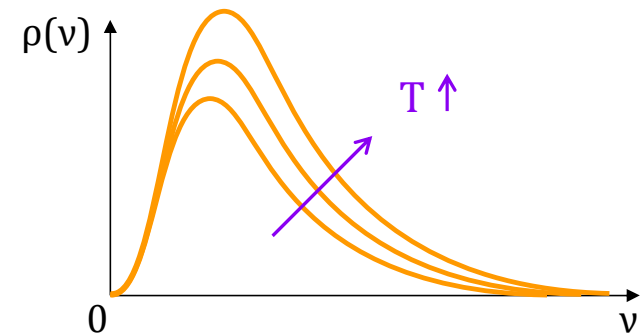
For low photon energies: $h\nu \ll k_B T \Rightarrow \bar{E} \approx k_B T$

- Spectral energy density = Mode-energy x mode density

$$\rho(\nu) = M(\nu)\bar{E} = \frac{8\pi h\nu^3}{c^3} \frac{1}{e^{\frac{h\nu}{k_B T}} - 1}$$

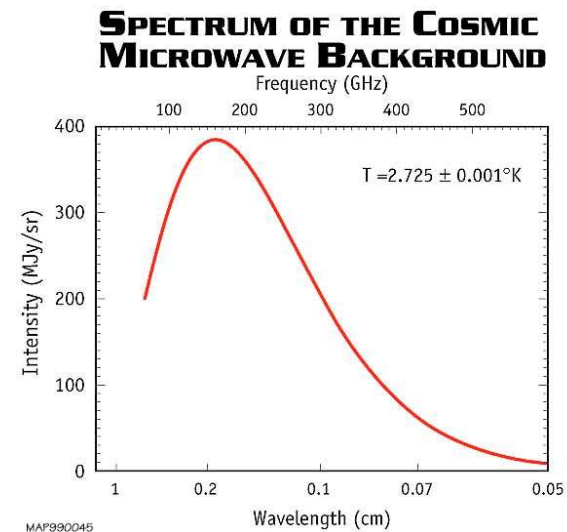
= Black body spectrum

- A black body radiator emits an EM spectrum only characterized by its temperature.



Blackbody radiator

- A black body radiator emits a EM spectrum only characterized by its temperature.
- It is an idealized concept. Real objects radiate with lower efficiency that can be wavelength dependent (emissivity)
- But real devices/situation can be close to it.
- Sun's surface 5778 K
- Earth approximately 300K
- Background in the universe 2.725 K



https://map.gsfc.nasa.gov/universe/bb_tests_cmb.html

Properties of a blackbody

The most important properties only depend on temperature

- Wien's law

$$\nu_{max} = 2.821436 \frac{k_B T}{h}$$

- Law of Stephan-Boltzmann

j total emitted power per unit area

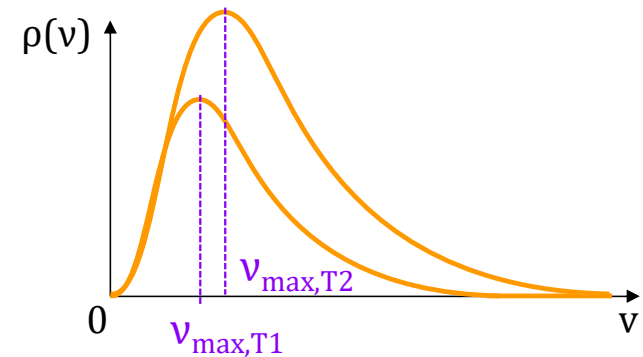
$$j = \sigma T^4 \quad \text{with} \quad \sigma = \frac{2\pi^5 k_B^4}{15c^2 h^3} = 5.6704 \cdot 10^{-8} \frac{W}{m^2 K^4}$$

- Example: temperature of sun $T = 5778 \text{ K}$

$$\rightarrow \nu_{max} = 3.378 \cdot 10^{14} \text{ Hz} \Rightarrow 887.4 \text{ nm}$$

(= where the maximum spectral energy density per unit frequency is)

<https://phet.colorado.edu/en/simulation/blackbody-spectrum>



Black body radiation energy density $\rho(\lambda)$

- Black body radiation energy density as a function of wavelength.

$$\rho(\lambda) = \frac{8\pi hc}{\lambda^5} \frac{1}{e^{\frac{hc}{\lambda k_B T}} - 1}$$

Energy per unit volume,
per unit of wavelength

- Wien's law in wavelength

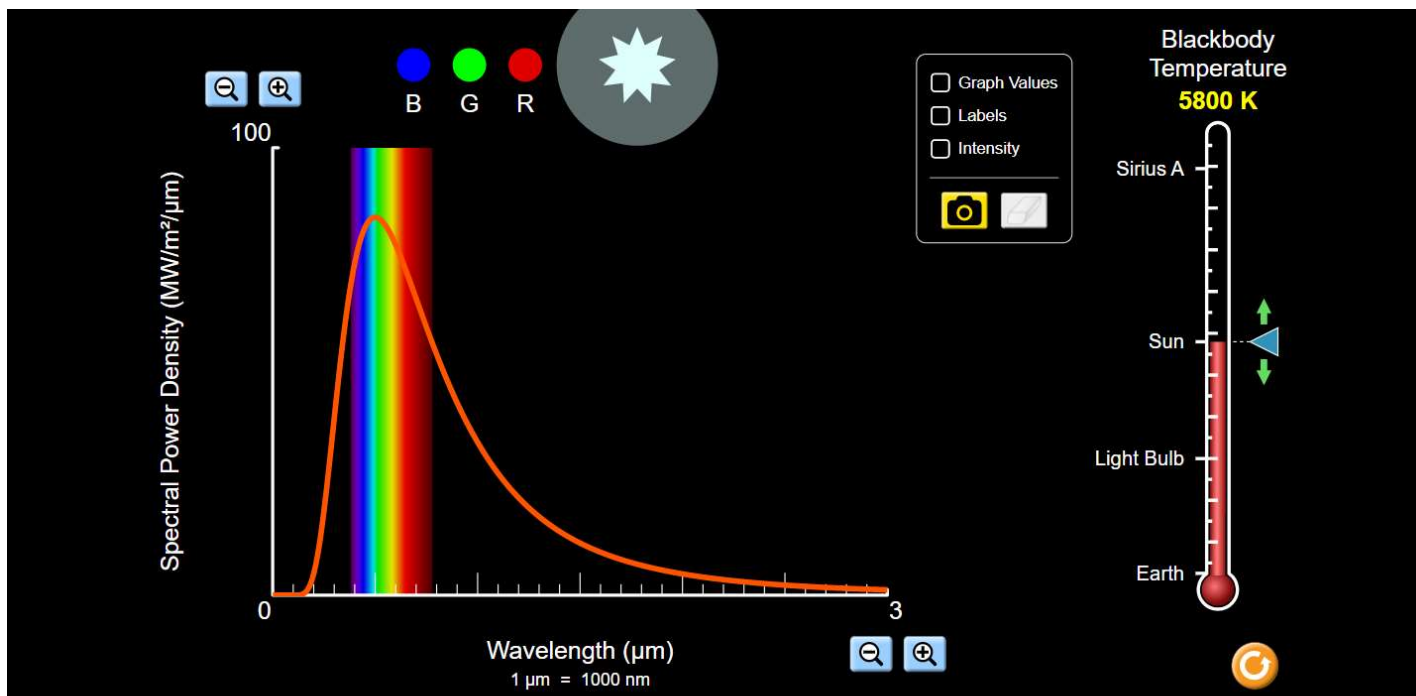
$$\lambda_{max} = 0.0028978 \frac{1}{T} \text{ m}$$

Note that $\lambda_{max} \neq \frac{c}{\nu_{max}}$ Since they stem from different entities!

ν_{max} as defined on previous slide:
the frequency at which the spectral power
density per unit of frequency is maximum

<https://phet.colorado.edu/en/simulations/>

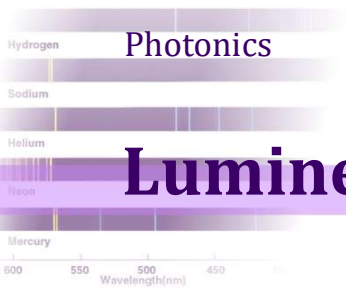
https://phet.colorado.edu/sims/html/blackbody-spectrum/latest/blackbody-spectrum_en.html



Blackbody Spectrum

Black body radiation spectral radiance

- In literature also expressions for the **spectral radiance** of the black body radiation are presented (without explicit mentioning)
- Spectral luminescence $L(\nu) = \frac{c \cdot \rho(\nu)}{4\pi}$ is Watt per m² per sr
- Spectral luminescence as a function of frequency
$$L(\nu) = \frac{2h\nu^3}{c^2} \left(\frac{1}{e^{\frac{h\nu}{k_B T}} - 1} \right)$$
- Spectral luminescence as a function of wavelength
$$L(\lambda) = \frac{2hc^2}{\lambda^5} \left(\frac{1}{e^{\frac{hc}{\lambda k_B T}} - 1} \right)$$



Photonics

Photons and Atoms

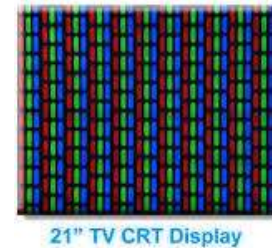
Luminescent light

- Common terminology for processes in which light is generated

Luminescent light

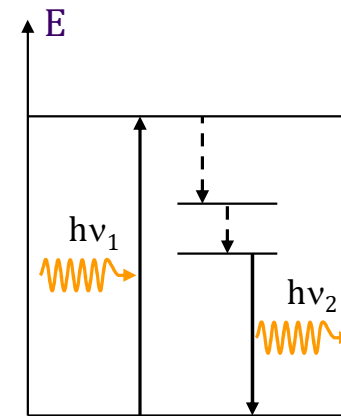
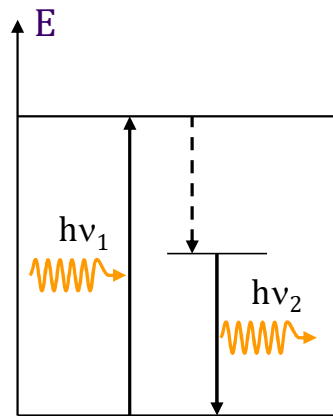
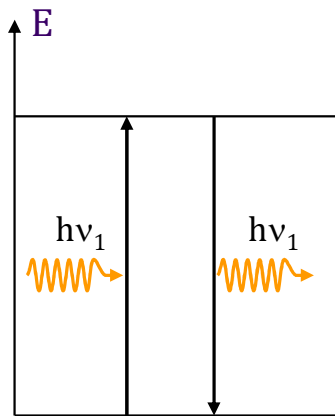
Excitation of system by external energy source

- Cathodoluminescence: accelerated electrons
e.g. old fashioned TV screen (CRT) – northern lights (auroras)
- photoluminescence: optical photons
e.g. White T-shirt illuminated by UV light, white LEDs
- Chemiluminescence: chemical reaction
e.g. Oxidizing phosphor
- Bioluminescence: biochemical reaction
e.g. Fireflies
- Electroluminescence: an external electrical field
e.g. LED
- Sonoluminescence: under the influence of sound, shock wave



Photoluminescence

- Atomic system is excited by a photon
- Falls back to ground state by emitting a photon
- Intermediate (non-radiative) transitions possible



Fluorescence and phosphorescence

- Fluorescence:
Singlet-Singlet transition
= easy
- Phosphorescence:
Triplet-Singlet overgang
= 'not allowed'
- Singlet: as many *spin-up* as *spin-down* electrons
- Triplet: Two electrons more *spin-down* than *spin-up* or vice versa
- Ground state= Singlet

