

# Lecture : Eigenvalues and eigenvectors

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## Traces and Determinants

Two useful numbers associated with square matrices are the *trace* and the *determinant*. The trace is the sum of the diagonals

$$\text{Tr}(A) = \sum_i A_{ii} \quad (1)$$

Traces are very well behaved algebraic. One can check immediately the following identities.

$$\text{Tr}(A) = \text{Tr}(A^T), \quad \text{Tr}(A + B) = \text{Tr}(A) + \text{Tr}(B), \quad \text{Tr}(AB) = \text{Tr}(BA) \quad (2)$$

Formulas for the determinant are generally complicated but they compute how the volume of the unit cube changes under the transformation  $A$ .

$$\det(A) = \text{signed volume of the unit cube transformed by } A \quad (3)$$

The sign of the determinant flips if the unit cube is reflected across some axis.

Determinants have the properties

$$\det(A) = \det(A^T), \quad \det(A^{-1}) = \det(A)^{-1}, \quad \det(AB) = \det(BA) = \det(A)\det(B) \quad (4)$$

Both the trace and determinant have special relationships with the eigenvalues of  $A$  (see below for discussion of eigenvalues). If the eigenvalues of  $A$ ,  $\lambda_1, \dots, \lambda_n$  then we have that

$$\text{Tr}(A) = \sum_i \lambda_i, \quad \det(A) = \prod_i \lambda_i \quad (5)$$

## Eigenvectors, Eigenvalues, and Diagonalization

In general, multiplying a column vector  $x \in \mathbb{R}^n$  by a square matrix  $A \in \mathbb{R}^{n \times n}$  causes that vector to stretch and to rotate. However, some vectors in specific subspaces are *only stretched, not rotated*. Another way to say this is that those subspaces are *invariant* with respect to  $A$ . These invariant subspaces are called *right eigenspaces* and vectors within them are called *right eigenvectors*. The amount each eigenvector is stretched is called its *eigenvalue*. We can also consider a similar situation where left multiplying  $A$  by specific row vectors only causes them to stretch. These

row vectors are called *left eigenvectors* and they live in *left eigenspaces*. (The eigenvalues for left and right eigenvectors turn out to be the same, ie. left and right eigenspaces come in pairs.) Finding a linearly independent sets of eigenvectors (either left or right) for a square matrix  $A$  is one of the fundamental problems of linear algebra. **If we represent vectors as coordinates with respect to a basis of eigenvectors, then the action of the matrix simply becomes scaling each individual coordinate by the appropriate eigenvalue.** If a matrix has a linearly independent basis of eigenvectors then we say it is *diagonalizable*. Not all matrices are diagonalizable, but if we choose a matrix at random then it will be (with probability 1), ie. we have to specifically work to construct a matrix that is not diagonalizable. The reason for this is that non-diagonalizable matrices are a low dimensional subset of the space of all matrices. Many arguments in linear algebra are best understood by understanding them for diagonalizable matrices and then generalizing them to the non-diagonalizable case.

The right and left eigenvector equations are given by

$$\lambda v = Av, \quad \lambda w^T = w^T A \quad (6)$$

respectively. Suppose the columns of  $P \in \mathbb{R}^{n \times n}$  are a linearly independent set of right eigenvectors of  $A$  and with eigenvalues  $\lambda_1, \dots, \lambda_n$ . Let  $D \in \mathbb{R}^n$  be a diagonal matrix with the eigenvalues on the diagonal, ie.  $D = \text{diag}([\lambda_1, \dots, \lambda_n])$ . The columns of  $P$  being right eigenvectors is equivalent to the equation

$$AP = PD \quad (7)$$

$$\Rightarrow A = PDP^{-1} \quad (8)$$

We say that the matrix of eigenvectors  $P$  diagonalizes  $A$  because it relates  $A$  to a diagonal matrix  $D$  via a similarity transform. In other words if  $x = Pz$ ,  $z' = Px'$  and  $x' = Ax$ , then  $z' = Dz$ . Note that in the  $z$ -coordinates,  $D$  simply scales each coordinate by the appropriate eigenvalue.

Left multiplying (8) by  $P^{-1}$  gives  $P^{-1}A = DP^{-1}$ . Note that this means that the rows of  $P^{-1}$  are a set of linearly independent left-eigenvectors of  $A$ . Note that this also shows why the left and right eigenvectors come in pairs and share eigenvalues. To summarize, let

$$P = \begin{bmatrix} | & & | \\ v_1 & \cdots & v_n \\ | & & | \end{bmatrix}, \quad D = \begin{bmatrix} \lambda_1 & & 0 \\ \vdots & \ddots & \vdots \\ 0 & & \lambda_n \end{bmatrix}, \quad P^{-1} = \begin{bmatrix} - & w_1^* & - \\ & \vdots & \\ - & w_n^* & - \end{bmatrix}, \quad (9)$$

with  $v_i$  and  $w_j$  being right and left eigenvectors.  $A$  can be decomposed as

$$A = \begin{bmatrix} | & & | \\ v_1 & \cdots & v_n \\ | & & | \end{bmatrix} \begin{bmatrix} \lambda_1 & & 0 \\ \vdots & \ddots & \vdots \\ 0 & & \lambda_n \end{bmatrix} \begin{bmatrix} - & w_1^* & - \\ & \vdots & \\ - & w_n^* & - \end{bmatrix} = \sum_i \lambda_i v_i w_i^* \quad (10)$$

Note that real eigenvalues denote how much each eigenvectors get stretched when they are multiplied by the matrix.

## Computing Eigenvalues and Eigenvectors

As stated above the determinant of a matrix is equal to the product of its eigenvalues. This means that if a matrix has a zero eigenvalue then its determinant is zero. Any vector in the nullspace of a matrix is an eigenvector with an eigenvalue of 0. Note that if  $\lambda v = Av$  then  $(\lambda I - A)v = 0$ . In other words, if  $v$  is eigenvector of  $A$  with eigenvalue  $\lambda$ , then  $v$  is also an eigenvector of  $\lambda I - A$  with eigenvalue 0. We can find eigenvalues of  $A$  by finding values of  $\lambda$  such that  $(\lambda I - A)$  has a 0 eigenvalue. This leads us to characterize eigenvalues as solutions to the equation

$$\chi_A(s) = \det(sI - A) = 0 \quad (11)$$

$\chi_A(s)$  is called the *characteristic polynomial* of  $A$ .

$$\chi_A(s) = \det(sI - A) = s^n + \alpha_{n-1}s^{n-1} + \cdots + \alpha_1s + \alpha_0$$

Based on properties of determinants,  $\chi_A(s)$  will always have order  $n$  and the first term will always be  $s^n$ .

Once we find roots of  $\chi_A(s)$ ,  $\lambda_i$ , we find the corresponding right and left eigenvectors by finding vectors in the right and left nullspace of  $\lambda_i I - A$  respectively.

## Spectral Mapping Theorem

### Polynomial Functions

As stated above computing eigenvectors and eigenvalues simplifies matrix computations. In particular, note that given a diagonalization of  $A = PDP^{-1}$ , we can compute powers of  $A$  as

$$A^k = \underbrace{A \times \cdots \times A}_{\times k} = PD^k \underbrace{P^{-1} \times P}_I D^k P^{-1} \times \cdots \times PDP^{-1} = PD^k P^{-1} \quad (12)$$

This implies that if a function  $f : \mathbb{C}^{n \times n} \rightarrow \mathbb{C}^{n \times n}$  is a polynomial (or more generally analytic function) of  $A$ , then

$$f(A) = Pf(D)P^{-1} = P \begin{bmatrix} f(\lambda_1) & & 0 \\ \vdots & \ddots & \vdots \\ 0 & & f(\lambda_n) \end{bmatrix} P^{-1} \quad (13)$$

In other words, we can compute polynomial functions of  $A$  simply by applying that function to the eigenvalues of  $A$  and leaving the eigenvectors unchanged. This is known as the *spectral mapping theorem*. Note that this analysis applies to polynomials with an infinite number of terms such as Taylor expansions of functions such as  $e^{(\cdot)}$ ,  $\cos(\cdot)$ , and  $\sin(\cdot)$  as well.

## Matrix Exponential

One important function of  $A$  that we want to compute is the *matrix exponential*  $e^A$  where which can be defined by its Taylor expansion.

$$e^A := I + A + \frac{1}{2!}A^2 + \frac{1}{3!}A^3 + \cdots = \sum_{k=0}^{\infty} \frac{1}{k!}A^k \quad (14)$$

Note that by the spectral mapping theorem we have that

$$e^A = Pe^DP^{-1} = P \begin{bmatrix} e^{\lambda_1} & & 0 \\ & \ddots & \\ 0 & & e^{\lambda_n} \end{bmatrix} P^{-1} \quad (15)$$

Exponential functions are interesting because they are functions who are equal to their own derivative (times some scaling), ie.  $\frac{d}{dt}e^{\lambda t} = \lambda e^{\lambda t}$ . (Note that  $e^{\lambda t}$  is actually an *eigenfunction* of the derivative operator  $\frac{d}{dt}$  with eigenvalue  $\lambda$ .)

## Cayley-Hamilton Theorem

The Cayley-Hamilton theorem says that a matrix satisfies its own characteristic polynomial, ie.  $\chi_A(A) = 0$ . For diagonalizable matrices, this is a direct application of the spectral mapping theorem.

$$\chi_A(A) = P\chi_A(D)P^{-1} = P \begin{bmatrix} \chi_A(\lambda_1) & \cdots & 0 \\ \vdots & & \vdots \\ 0 & \cdots & \chi_A(\lambda_n) \end{bmatrix} P^{-1} = 0$$

Consequently,

$$A^n = -\alpha_{n-1}A^{n-1} - \cdots - \alpha_1A - \alpha_0I$$

As a result of this, any polynomial function of  $A$  could be expressed in terms of powers of  $A$  only up through  $n-1$ . Higher powers of  $A$  can be reduced by iteratively plugging in the above equation.

Another application of Cayley-Hamilton gives a polynomial expression for a matrix inverse.

$$0 = (A^n + \alpha_{n-1}A^{n-1} + \cdots + \alpha_1A + \alpha_0I)A^{-1}$$

$$A^{-1} = -\frac{1}{\alpha_0}A^{n-1} - \frac{\alpha_{n-1}}{\alpha_0}A^{n-2} - \cdots - \frac{\alpha_1}{\alpha_0}I$$

## Jordan Form

To motivate a study of Jordan form, we consider the following matrix

$$J_i = \lambda_i I + N_i = \begin{bmatrix} \lambda_i & 1 & \cdots & \cdots & 0 \\ 0 & \lambda_i & & & \vdots \\ \vdots & & \ddots & & \vdots \\ \vdots & & & \lambda_i & 1 \\ 0 & \cdots & \cdots & 0 & \lambda_i \end{bmatrix}$$

where  $N_i$  is a matrix with 1's on the first super diagonal. This matrix  $N_i$  is an example of a *nilpotent matrix* since raising it to some power gives a matrix of 0's, ie. for example

$$\begin{bmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 0 & 0 & 0 \end{bmatrix}^3 = \begin{bmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

Note that any matrix similar to a nilpotent matrix is also nilpotent. If  $N_i^k = 0$ , then  $(PN_iP^{-1})^k = PN_i^kP^{-1} = 0$ . If  $J_i = \lambda_i I + N_i$ , then clearly,  $J_i - \lambda_i I$  is nilpotent, ie.  $J_i - \lambda_i I = N_i$ . Since the eigenvalues of a triangular matrix are just the diagonal values, we have that the only eigenvalue of  $N_i$  is simply 0. However,  $N_i$  clearly has  $n - 1$  linearly independent columns, ie. rank  $n - 1$ . Thus it only has a one dimensional nullspace. One can check that the characteristic polynomial of  $N_i$  is  $\chi_{N_i}(s) = s^n$  and the characteristic polynomial of  $J_i = \lambda_i I + N_i$  is  $\chi_{J_i}(s) = (s - \lambda_i)^n$ .

**A matrix is not diagonalizable when a full basis of eigenvectors does not exist.** For a matrix  $A \in \mathbb{R}^{n \times n}$  with  $n$  distinct eigenvalues, there must be a basis of  $n$  linearly independent eigenvectors since each eigenvalue  $\lambda_i$  is associated with the nullspace of  $\lambda_i I - A$ . We know these eigenvectors are linearly independent since if not

$$\begin{aligned} v_i &= \sum_{j \neq i} \alpha_j v_j \\ Av_i &= A \left( \sum_{j \neq i} \alpha_j v_j \right) \\ 0 &= \sum_{j \neq i} \alpha_j \lambda_j v_j - \lambda_i v_i \\ 0 &= \sum_{j \neq i} \alpha_j (\lambda_j - \lambda_i) v_j \end{aligned}$$

An inductive argument shows that  $\lambda_i = \lambda_j$  for some  $i$  and  $j$  which is a contradiction.

In this case, the characteristic polynomial is

$$\chi_A(s) = (s - \lambda_1)(s - \lambda_2) \cdots (s - \lambda_n)$$

In the general case with repeated eigenvalues, the characteristic polynomial is given by

$$\chi_A(s) = \prod_{i=1}^k (s - \lambda_i)^{k_i}$$

where  $k$  is the number of distinct eigenvalues and  $k_i$  is the number of times each eigenvalue is repeated. If  $\dim(\mathcal{N}(\lambda_i I - A)) = k_i$  for all  $i$ , then the matrix is diagonalizable. In this case,

$$\mathcal{N}(\lambda_i I - A) = \mathcal{N}((\lambda_i I - A)^2) = \mathcal{N}((\lambda_i I - A)^3) = \dots$$

and

$$\dim(\mathcal{N}(\lambda_i I - A)) = \dim(\mathcal{N}((\lambda_i I - A)^2)) = \dim(\mathcal{N}((\lambda_i I - A)^3)) = \dots = k_i$$

This happens when  $\mathcal{N}(\lambda_i I - A) \cap \mathcal{R}(\lambda_i I - A) = 0$  for all  $i$ .

It is also possible that  $\dim(\mathcal{N}(\lambda_i I - A)) < k_i$ . In this case

$$\mathcal{N}(\lambda_i I - A) \subset \mathcal{N}((\lambda_i I - A)^2) \subset \mathcal{N}((\lambda_i I - A)^3) \subset \dots$$

and

$$\dim(\mathcal{N}(\lambda_i I - A)) < \dim(\mathcal{N}((\lambda_i I - A)^2)) < \dim(\mathcal{N}((\lambda_i I - A)^3)) < \dots < k_i \quad (16)$$

ie.,  $\mathcal{N}(\lambda_i I - A) \cap \mathcal{R}(\lambda_i I - A) \neq 0$ . A regular eigenvector satisfies

$$(\lambda_i I - A)v_i = 0$$

If  $\dim(\mathcal{N}(\lambda_i I - A)) < \dim(\mathcal{N}((\lambda_i I - A)^2))$ , then we should be able to find generalized eigenvectors that satisfy

$$(\lambda_i I - A)w_i^2 \in \mathcal{N}(\lambda_i I - A), \quad (\lambda_i I - A)w_i^3 \in \mathcal{N}((\lambda_i I - A)^2), \quad \text{etc}$$

$w_i^2 \in \mathbb{C}^n$  is a 2nd order eigenvector,  $w_i^3 \in \mathbb{C}^n$  is a 3rd order eigenvector, etc.

Note that

$$(\lambda_i I - A)^2 w_i^2 = 0, \quad (\lambda_i I - A)^3 w_i^3 = 0, \quad \text{etc}$$

If we are careful in picking,  $v_i, w_i^2, w_i^3, \dots$  we can choose them so that

$$0 = (\lambda_i I - A)v_i, \quad v_i = (\lambda_i I - A)w_i^2, \quad w_i^2 = (\lambda_i I - A)w_i^3, \quad \text{etc} \quad (17)$$

A general organization of these equations is given by

$$AP = PJ = \underbrace{[V_1 \ \dots \ V_q]}_P \underbrace{\begin{bmatrix} J_1 & \dots & 0 \\ \vdots & & \vdots \\ 0 & \dots & J_q \end{bmatrix}}_J$$

where

$$V_i = \begin{bmatrix} | & | & | & \cdots \\ v_1 & w_1^2 & w_1^3 & \cdots \\ | & | & | & \end{bmatrix}, \quad J_i = \lambda_i I + N_i = \begin{bmatrix} \lambda_i & 1 & \cdots & \cdots & 0 \\ 0 & \lambda_i & & & \vdots \\ \vdots & & \ddots & & \vdots \\ \vdots & & & \lambda_i & 1 \\ 0 & \cdots & \cdots & 0 & \lambda_i \end{bmatrix}$$

$$N_i = \begin{bmatrix} 0 & 1 & \cdots & \cdots & 0 \\ 0 & 0 & & & \vdots \\ \vdots & & \ddots & & \vdots \\ \vdots & & & 0 & 1 \\ 0 & \cdots & \cdots & 0 & 0 \end{bmatrix}$$

$J_i$  is called a Jordan block and  $q$  is the number of Jordan blocks. Each Jordan block corresponds to one true eigenvector and a chain of generalized eigenvectors as in (17). Note that if each distinct eigenvalue has only one Jordan block (and only one true eigenvector), then  $q = k$ , the number of distinct eigenvalues. It is possible that a distinct eigenvalue has more than one Jordan block. In this case,  $q > k$ . Most matrices are diagonalizable, but every matrix can be put in *Jordan form*. Note that

$$\begin{aligned} A - \lambda_1 I &= PJP^{-1} - \lambda_1 PP^{-1} \\ &= P(J - \lambda_1 I)P^{-1} \\ &= P \begin{bmatrix} N_1 & \cdots & 0 \\ \vdots & & \vdots \\ 0 & \cdots & J_q \end{bmatrix} \end{aligned}$$

and that

$$(A - \lambda_1 I)^\ell = P \begin{bmatrix} N_1^\ell & \cdots & 0 \\ \vdots & & \vdots \\ 0 & \cdots & J_q^\ell \end{bmatrix}$$

Since  $N_1$  is nilpotent, as  $\ell$  increases the nullspace of  $(A - \lambda_1 I)^\ell$  grows as in (16).

We now perform several manipulations with a simple non-diagonalizable matrix to illustrate

some simple properties of Jordan form. Consider

$$\begin{aligned}
A &= \begin{bmatrix} | & | & | \\ v & w^2 & w^3 \\ | & | & | \end{bmatrix} \begin{bmatrix} \lambda & 1 & 0 \\ 0 & \lambda & 1 \\ 0 & 0 & \lambda \end{bmatrix} \begin{bmatrix} | & | & | \\ v & w^2 & w^3 \\ | & | & | \end{bmatrix}^{-1} \\
&= \begin{bmatrix} | & | & | \\ v & w^2 & w^3 \\ | & | & | \end{bmatrix} \begin{bmatrix} \lambda & 1 & 0 \\ 0 & \lambda & 1 \\ 0 & 0 & \lambda \end{bmatrix} \begin{bmatrix} - & (q^3)^T & - \\ - & (q^2)^T & - \\ - & p^T & - \end{bmatrix} \\
&= \lambda v(q^3)^T + (v + \lambda w^2)(q^2)^T + (w^2 + \lambda w^3)p^T \\
&= \lambda v(q^3)^T + \lambda w^2(q^2)^T + \lambda w^3 p^T + v(q^2)^T + w^2 p^T
\end{aligned}$$

Note that

- The first order right eigenvector  $v$  matches up with the third order left generalized eigenvector  $(q^3)^T$
- The second order right eigenvector  $w^2$  matches up with the second order left generalized eigenvector  $(q^2)^T$
- The third order right eigenvector  $w^3$  matches up with the first order left eigenvector  $p^T$

We note that we could also write  $A$  in other ways related to Jordan form (These are just a sample of how the Jordan block and eigenvectors could be shuffled.)

$$\begin{aligned}
A &= \begin{bmatrix} | & | & | \\ w^2 & v & w^3 \\ | & | & | \end{bmatrix} \begin{bmatrix} \lambda & 0 & 1 \\ 1 & \lambda & 0 \\ 0 & 0 & \lambda \end{bmatrix} \begin{bmatrix} - & (q^2)^T & - \\ - & (q^3)^T & - \\ - & p^T & - \end{bmatrix} \\
&= \begin{bmatrix} | & | & | \\ w^3 & w^2 & v \\ | & | & | \end{bmatrix} \begin{bmatrix} \lambda & 0 & 0 \\ 1 & \lambda & 0 \\ 0 & 1 & \lambda \end{bmatrix} \begin{bmatrix} - & p^T & - \\ - & (q^2)^T & - \\ - & (q^3)^T & - \end{bmatrix} \\
&= \text{etc...}
\end{aligned}$$