

# Trotterization in Quantum Computing

## 1 Introduction

Trotterization is a key numerical technique used in quantum computing for approximating the time evolution operator of a quantum system. It is based on the Suzuki-Trotter expansion, which enables the decomposition of an exponential of sum of non-commuting operators into a sequence of exponentials of individual terms. This technique is particularly useful for simulating quantum many-body systems and Hamiltonian dynamics on quantum computers.

## 2 Trotter-Suzuki Decomposition

Consider a quantum system with a Hamiltonian that can be decomposed as a sum of non-commuting terms:

$$H = H_A + H_B. \quad (1)$$

The exact time evolution operator over a small time step  $t$  is given by:

$$U(t) = e^{-iHt} = e^{-i(H_A+H_B)t}. \quad (2)$$

However, due to the non-commutativity of  $H_A$  and  $H_B$ , we cannot simply split the exponentiation. Instead, using the first-order Trotter approximation:

$$e^{-i(H_A+H_B)t} \approx e^{-iH_A t} e^{-iH_B t} + \mathcal{O}(t^2). \quad (3)$$

This approximation becomes more accurate when we divide the time interval into smaller steps  $n$ :

$$U(t) = \left( e^{-iH_A t/n} e^{-iH_B t/n} \right)^n + \mathcal{O}(t^2/n). \quad (4)$$

In the limit  $n \rightarrow \infty$ , the error vanishes, and the decomposition becomes exact.

## 3 Why Non-Commuting Observables' Exponential Cannot Be Split

In quantum mechanics, the evolution of a system is often governed by an operator exponential of the form:

$$e^{(A+B)t}. \quad (5)$$

If the operators  $A$  and  $B$  commute, i.e.,

$$[A, B] = AB - BA = 0, \quad (6)$$

then we can factor the exponential as:

$$e^{(A+B)t} = e^{At} e^{Bt}. \quad (7)$$

However, when  $A$  and  $B$  do not commute, this factorization is no longer valid. The reason lies in the Baker-Campbell-Hausdorff (BCH) formula, which provides an expansion for the exponential of a sum of non-commuting operators:

$$e^{(A+B)t} = e^{At} e^{Bt} e^{-\frac{t^2}{2}[A,B]} e^{\mathcal{O}(t^3)}. \quad (8)$$

The presence of the additional exponential terms involving the commutator  $[A, B]$  and higher-order nested commutators means that the naive splitting of the exponentials introduces errors. Thus, for non-commuting operators, the direct separation:

$$e^{(A+B)t} \neq e^{At} e^{Bt} \quad (9)$$

is incorrect in general.

This property has important implications in quantum mechanics and quantum computing, particularly in time evolution and numerical approximations, where methods like Trotter-Suzuki decomposition are used to approximate such exponentials while managing the error introduced by non-commutativity.

## 4 Trotterization and Repeated Evolution Steps

In quantum mechanics, time evolution is given by the unitary operator:

$$U(t) = e^{-itH}, \quad (10)$$

where  $H$  is the Hamiltonian of the system. If the Hamiltonian consists of multiple non-commuting terms, such as:

$$H = H_X + H_Y + H_Z, \quad (11)$$

then the direct exponentiation of the sum is not equal to the product of exponentials due to non-commutativity:

$$e^{-it(H_X+H_Y+H_Z)} \neq e^{-itH_X} e^{-itH_Y} e^{-itH_Z}. \quad (12)$$

To approximate this evolution, we use Trotterization.

### 4.1 Small Time Step Approximation

Instead of evolving the system for a large time  $t$ , we break it into small steps:

$$t = m \cdot dt, \quad dt = \frac{t}{m}. \quad (13)$$

For a small time step  $dt$ , the first-order Trotter approximation is:

$$U(dt) \approx e^{-idtH_X} e^{-idtH_Y} e^{-idtH_Z}. \quad (14)$$

Since evolution is sequential, to achieve the full-time evolution, we apply this small-step evolution  $m$  times:

$$U(t) \approx (e^{-idtH_X} e^{-idtH_Y} e^{-idtH_Z})^m. \quad (15)$$

This effectively constructs the evolution over  $t$  by repeating small-step transformations.

### 4.2 Why Repeat Instead of Summation?

One might wonder why we do not use a summation instead of repeating the operation. The reason lies

in the nature of matrix exponentiation. A naive summation approach would be:

$$U(t) \approx \sum_{k=1}^m e^{-idtH}, \quad (16)$$

which is incorrect because exponentials of matrices do not sum linearly. Instead, the correct approach follows from:

$$U(t) = (U(dt))^m = (e^{-idtH})^m. \quad (17)$$

This is analogous to compound interest in finance. If an investment grows at rate  $r$ , instead of applying the rate once per year, it compounds in small steps:

$$(1 + rdt)^m. \quad (18)$$

As  $dt \rightarrow 0$ , this approaches the continuous exponential growth formula:

$$e^{rt}. \quad (19)$$

Similarly, breaking the quantum evolution into small steps and repeating them ensures the correct accumulation of time evolution effects.

### 4.3 Final Summary

- Instead of evolving for time  $t$ , we break it into small steps:  $t = m \cdot dt$ .
- We approximate  $U(dt)$  using first-order Trotterization.
- Since evolution is sequential, we repeat  $U(dt)$   $m$  times to get  $U(t)$ .
- This method ensures proper time evolution in quantum simulations and quantum computing.