Homework 2

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PROBLEM 1

Recall the φ^4 Lagrangian is given by

$$\mathcal{L} = -\frac{1}{2}\partial^{\mu}\varphi\partial_{\mu}\varphi - \frac{1}{2}m^{2}\varphi^{2} - \frac{1}{3!}g\varphi^{3} - \frac{1}{4!}\lambda\varphi^{4}$$

and has an energy-momentum tensor

$$T^{\mu\nu} = \partial^{\mu}\varphi \partial^{\nu}\varphi + q^{\mu\nu}\mathscr{L}$$

Part A

Problem. Derive the equation of motion for φ subject to the φ^4 Lagrangian.

To calculate the equation of motion for φ , we just have to find the stationary points of

$$S = \int d^4x \mathcal{L} = \int d^4x \left(\frac{-1}{2} \partial^{\mu} \varphi \partial_{\mu} \varphi - \frac{1}{2} m^2 \varphi^2 - \frac{1}{3!} g \varphi^3 - \frac{1}{4!} \lambda \varphi^4 \right)$$

That is, we find when $\delta S = 0$. To do so, we calculate

$$\begin{split} \delta S &= \int d^4x \delta \mathcal{L} \\ &= \int d^4x \left(\frac{-1}{2} \delta(\partial^\mu \varphi \partial_\mu \varphi) - \frac{1}{2} m^2 \delta(\varphi^2) - \frac{1}{3!} g \delta(\varphi^3) - \frac{1}{4!} \lambda \delta(\varphi^4) \right) \\ &= \int d^4x \left(\frac{-1}{2} (\partial^\mu \delta \varphi \partial_\mu \varphi + \partial^\mu \varphi \partial_\mu \delta \varphi) - m^2 \varphi \delta \varphi - \frac{1}{2} g \varphi^3 \delta \varphi - \frac{1}{3!} \lambda \varphi^3 \delta \varphi \right) \\ &= \int d^4x \left(\partial^2 \varphi \delta \varphi - m^2 \varphi \delta \varphi - \frac{1}{2} g \varphi^3 \delta \varphi - \frac{1}{3!} \lambda \varphi^3 \delta \varphi \right) \\ &= \int d^4x \left(\partial^2 \varphi - m^2 \varphi - \frac{1}{2} g \varphi^3 - \frac{1}{3!} \lambda \varphi^3 \right) \delta \varphi \end{split}$$

Which is zero for arbitrary variation if $\left(\partial^2 \varphi - m^2 \varphi - \frac{1}{2} g \varphi^3 - \frac{1}{3!} \lambda \varphi^3\right) = 0$. Thus, this is the equation of motion for φ .

Part B

Problem. Show that the energy-momentum tensor $T^{\mu\nu}$ satisfies $\partial_{\mu}T^{\mu\nu} = 0$.

This is just an exercise in direct calculation:

$$\begin{split} \partial_{\mu}T^{\mu\nu} &= \partial_{\mu} \left(\partial^{\mu}\varphi \partial^{\nu}\varphi \right) + \partial_{\mu}g^{\mu\nu}\mathcal{L} \\ &= \partial_{\mu}\partial^{\mu}\varphi \partial^{\nu}\varphi + \partial^{\mu}\varphi \partial_{\mu}\partial^{\nu}\varphi + \partial^{\nu} \left(\frac{-1}{2}\partial^{\mu}\varphi \partial_{\mu}\varphi - \frac{1}{2}m^{2}\varphi^{2} - \frac{1}{3!}g\varphi^{3} - \frac{1}{4!}\lambda\varphi^{4} \right) \\ &= \partial_{\mu}\partial^{\mu}\varphi \partial^{\nu}\varphi + \partial^{\mu}\varphi \partial_{\mu}\partial^{\nu}\varphi - \frac{1}{2}\partial^{\nu}\partial^{\mu}\varphi \partial_{\mu}\varphi - \frac{1}{2}\partial^{\mu}\varphi \partial^{\nu}\partial_{\mu}\varphi - m^{2}\varphi\partial^{\nu}\varphi - \frac{1}{2}g\varphi^{2}\partial^{\nu}\varphi - \frac{1}{3!}\lambda\varphi^{3}\partial^{\nu}\varphi \\ &= \partial_{\mu}\partial^{\mu}\varphi \partial^{\nu}\varphi - m^{2}\varphi\partial^{\nu}\varphi - \frac{1}{2}g\varphi^{2}\partial^{\nu}\varphi - \frac{1}{3!}\lambda\varphi^{3}\partial^{\nu}\varphi \\ &= \left(\partial^{2}\varphi - m^{2}\varphi - \frac{1}{2}g\varphi^{2} - \frac{1}{3!}\lambda\varphi^{3} \right)\partial^{\nu}\varphi \end{split}$$

Which is clearly zero if φ follows its equation of motion.

PROBLEM 2

Consider a complex scalar field φ governed by the Lagrangian

$$\mathcal{L} = -\partial^{\mu} \varphi^{\dagger} \partial_{\mu} \varphi - m^2 \varphi^{\dagger} \varphi + \Omega_0$$

Part A

Problem. Show φ obeys the Klein-Gordon equation.

We calculate the variation in $S = \int d^4x \mathcal{L}$ directly:

$$\begin{split} \delta S &= \int d^4x \delta \mathscr{L} \\ &= \int d^4x \left(-\delta(\partial^\mu \varphi^\dagger) \partial_\mu \varphi - \partial^\mu \phi^\dagger \delta(\partial_\mu \varphi) - m^2 \varphi \delta \varphi^\dagger - m^2 \varphi^\dagger \delta \varphi \right) \\ &= \int d^4x \left(-\partial^\mu \delta \varphi^\dagger \partial_\mu \varphi - \partial^\mu \phi^\dagger \partial_\mu \delta \varphi - m^2 \varphi \delta \varphi^\dagger - m^2 \varphi^\dagger \delta \varphi \right) \\ &= \int d^4x \left(\delta \varphi^\dagger \partial_2 \varphi + \partial^2 \phi^\dagger \delta \varphi - m^2 \varphi \delta \varphi^\dagger - m^2 \varphi^\dagger \delta \varphi \right) \\ &= \int d^4x \left(\left(\partial^2 \varphi - m^2 \varphi \right) \delta \varphi^\dagger + \left(\partial^2 \varphi^\dagger - m^2 \varphi^\dagger \right) \delta \varphi \right) \end{split}$$

Which is zero for arbitrary variations when both φ and φ^{\dagger} follow the Klein-Gordon equation.

Part B

Problem. Find the conjugate momenta for φ and φ^{\dagger} , and write down the Hamiltonian in terms of these.

We can read off the conjugate momenta easily:

$$\pi(x) = \frac{\partial \mathcal{L}}{\partial \dot{\varphi}}$$

$$= \frac{\partial}{\partial \dot{\varphi}} (\partial_{\mu} \varphi^{\dagger}(x) \partial_{\mu} \varphi(x) - m^{2} \varphi^{\dagger}(x) \varphi(x) + \Omega_{0})$$

$$= \dot{\varphi}^{\dagger}(x)$$

and similarly

$$\pi^{\dagger}(x) = \dot{\varphi}(x)$$

We write down the Hamiltonian density as

$$\mathscr{H} = \pi(x)\dot{\varphi}(x) + \pi^{\dagger}(x)\dot{\varphi}^{\dagger}(x) - \mathscr{L}$$

and calculate

$$\mathcal{H} = \pi(x)\dot{\varphi}(x) + \pi^{\dagger}(x)\dot{\varphi}^{\dagger}(x) + \left(\partial^{\mu}\varphi^{\dagger}(x)\partial_{\mu}\varphi(x) + m^{2}\varphi^{\dagger}(x)\varphi(x) - \Omega_{0}\right)$$

$$= \pi(x)\pi^{\dagger}(x) + \pi^{\dagger}(x)\pi(x) + \partial^{0}\varphi^{\dagger}(x)\partial_{0}\varphi(x) + \partial^{i}\varphi^{\dagger}(x)\partial_{i}\varphi(x) + m^{2}\varphi^{\dagger}(x)\varphi(x) - \Omega_{0}$$

$$= \pi(x)\pi^{\dagger}(x) + \pi^{\dagger}(x)\pi(x) - \pi(x)\pi^{\dagger}(x) + \partial^{i}\varphi^{\dagger}(x)\partial_{i}\varphi(x) + m^{2}\varphi^{\dagger}(x)\varphi(x) - \Omega_{0}$$

$$= \pi^{\dagger}(x)\pi(x) + \partial^{i}\varphi^{\dagger}(x)\partial_{i}\varphi(x) + m^{2}\varphi^{\dagger}(x)\varphi(x) - \Omega_{0}$$

Which gives the Hamiltonian for the system as

$$H = \int d^4x \left(\pi^{\dagger}(x)\pi(x) + \partial^i \varphi^{\dagger}(x)\partial_i \varphi(x) + m^2 \varphi^{\dagger}(x)\varphi(x) - \Omega_0 \right)$$

Part C

Problem. Expanding φ as

$$\varphi(x) = \int d^{\tilde{3}}k \left(a(k) \exp(ikx) + b^{\dagger}(k) \exp(-ikx) \right)$$

solve for expressions of a(k) and b(k) in terms of $\varphi(x)$, $\varphi^{\dagger}(x)$ and their time derivatives.

We'll evaluate the integrals

$$\int d^3x \exp(-ikx)\varphi(x)$$
$$\int d^3x \exp(-ikx)\partial_t\varphi(x)$$
$$\int d^3x \exp(-ikx)\varphi^{\dagger}(x)$$
$$\int d^3x \exp(-ikx)\partial_t\varphi^{\dagger}(x)$$

So we first derive an expression for $\partial_t \varphi(x)$ and its conjugate.

$$\partial_t \varphi(x) = \int d^3k \partial_t \left(a(k) \exp(ikx) + b^{\dagger}(k) \exp(-ikx) \right)$$

$$= \int d^3k \left(a(k)(-i\omega) \exp(ikx) - b^{\dagger}(k)(-i\omega) \exp(-ikx) \right)$$

$$= \int d^3k (-i\omega) \left(a(k) \exp(ikx) - b^{\dagger}(k) \exp(-ikx) \right)$$

And similarly,

$$\partial_t \varphi^{\dagger}(x) = \int d^{\tilde{3}}k(-i\omega) \left(b(k) \exp(ikx) - a^{\dagger}(k) \exp(-ikx) \right)$$

Now, we can calculate those four integrals. One will be done explicitly, and the other three are done using the exact same calculation.

$$\int d^3x \exp(-ikx)\varphi(x) = \int d^3x \frac{d^3k'}{(2\pi)^3 2\omega} \exp(-ikx) \left(a(k') \exp(ik'x) + b^{\dagger}(k') \exp(-ik'x) \right)$$

$$= \int d^3x \frac{d^3k'}{(2\pi)^3 2\omega} \left(a(k') \exp(i(k'-k)x) + b^{\dagger}(k') \exp(-i(k'+k)x) \right)$$

$$= \int d^3x \frac{d^3k'}{(2\pi)^3 2\omega} (a(k') \exp(i(k'-k)^i x_i + i(\omega'-\omega)t)$$

$$+ b^{\dagger}(k') \exp(-i(k'+k)^i x_i + -i(\omega'+\omega_t))$$

$$= \int \frac{d^3k'}{(2\pi)^3 2\omega} (a(k')(2\pi)^3 \delta(k'-k) \exp(i(\omega'-\omega)t)$$

$$+ b^{\dagger}(k')(2\pi)^3 \delta(k'+k) \exp(-i(\omega'+\omega)t))$$

$$= \frac{1}{2\omega} (a(k) + b^{\dagger}(-k) \exp(-i(2\omega)t))$$

Following the same tactic, we find

$$\int d^3x \exp(-ikx)\varphi(x) = \frac{1}{2\omega}(a(k) + b^{\dagger}(-k)\exp(-i(2\omega)t))$$

$$\int d^3x \exp(-ikx)\partial_t\varphi(x) = \frac{-i}{2}(a(k) - b^{\dagger}(-k)\exp(-i(2\omega)t))$$

$$\int d^3x \exp(-ikx)\varphi^{\dagger}(x) = \frac{1}{2\omega}(b(k) + a^{\dagger}(-k)\exp(-i(2\omega)t))$$

$$\int d^3x \exp(-ikx)\partial_t\varphi^{\dagger}(x) = \frac{-i}{2}(b(k) - a^{\dagger}(-k)\exp(-i(2\omega)t))$$

So,

$$a(k) = \int d^3x \exp(-ikx) \left(\omega \varphi(x) + i\partial_t \varphi(x)\right)$$

and

$$b(k) = \int d^3x \exp(-ikx) \left(\omega \varphi^{\dagger}(x) + i\partial_t \varphi^{\dagger}(x)\right)$$

Part D

Problem. Derive the commutation relations for a(k) and b(k) and their conjugates.