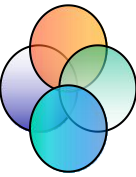


## § 5.3 Magnetic Dipole



Only a brief introduction and comparison

### Electric Dipole

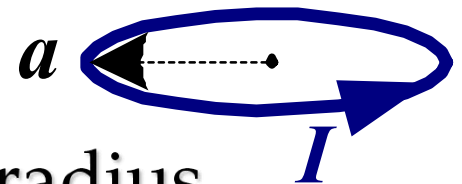
A pair of opposite charges very close to each other.

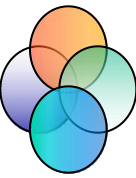
- Distance:  $l$
- Point charges:  $q_1 = q$ ,  $q_2 = -q$



### Magnetic Dipole

A circular current with a very small radius





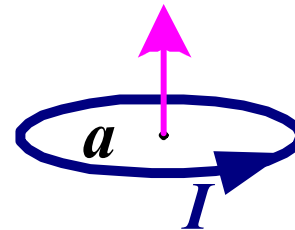
## *Electric Dipole Moment*

$$\vec{p} = q\vec{l}$$

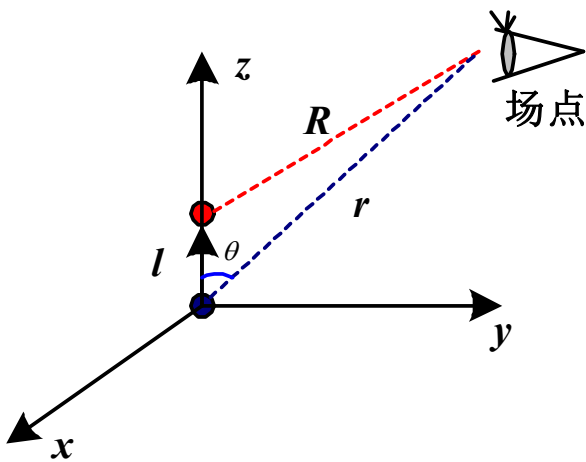
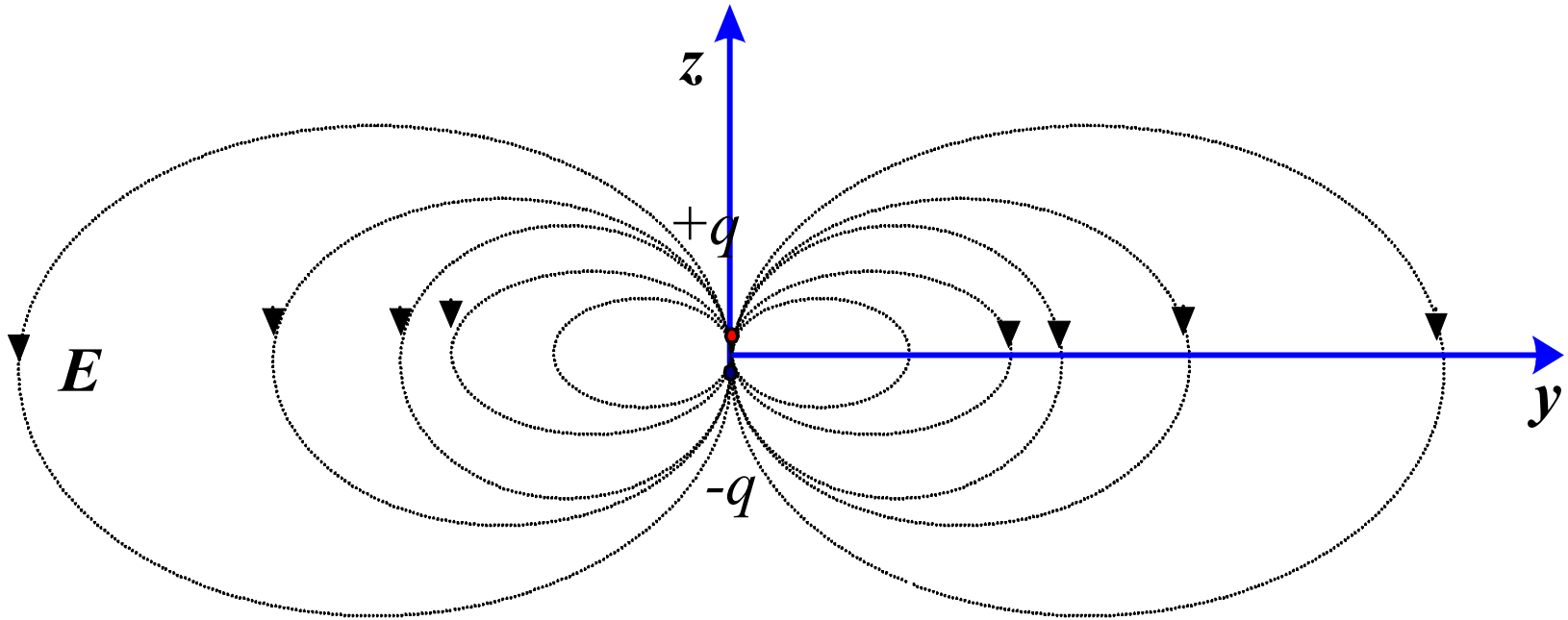
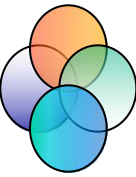


## *Magnetic Dipole Moment*

$$\vec{p}_m = I\vec{S}$$

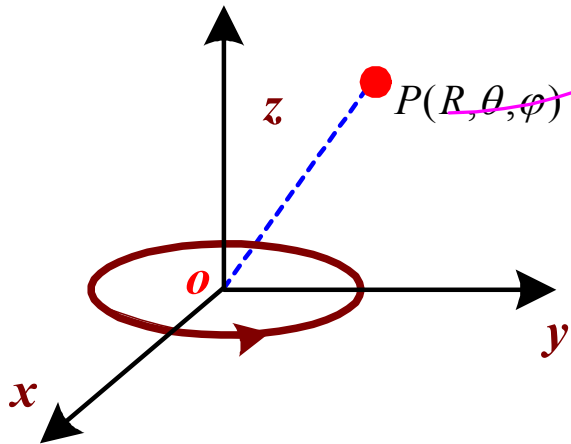
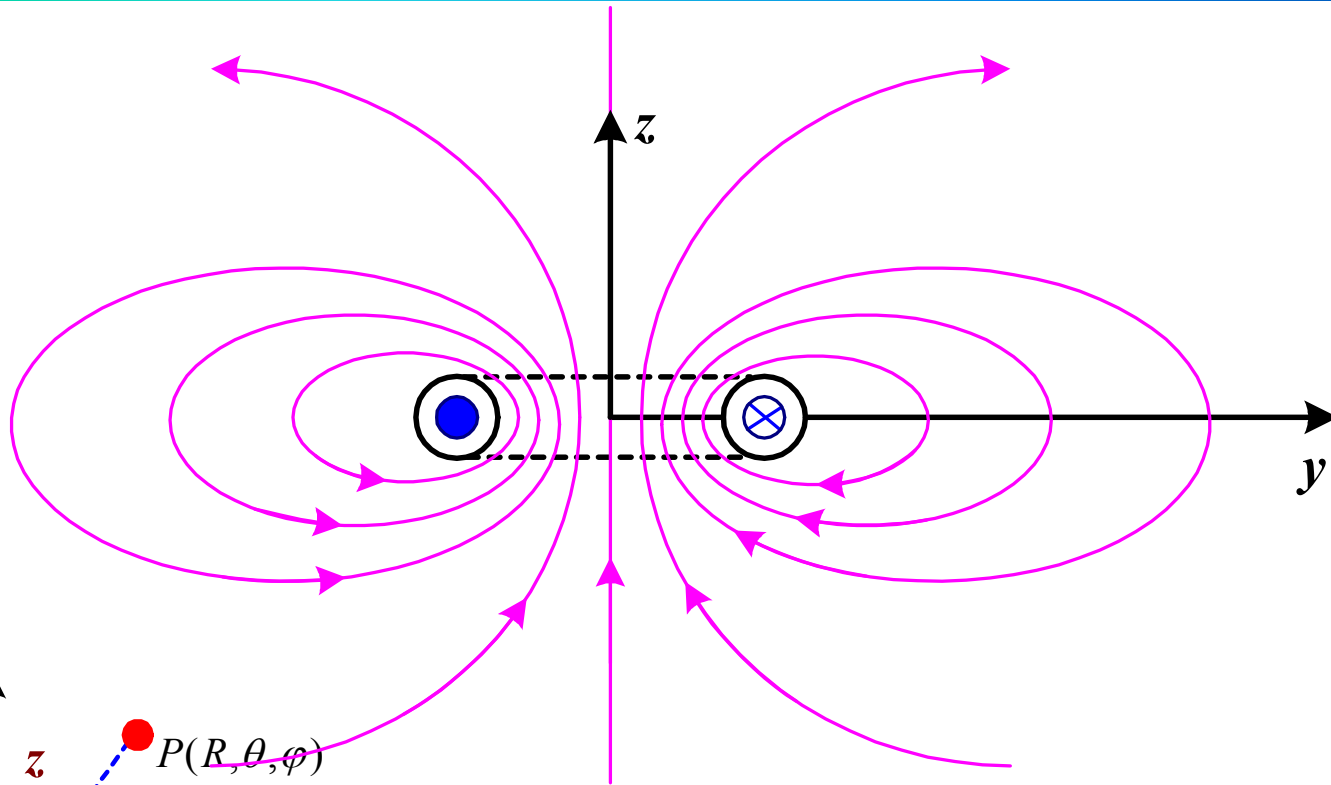
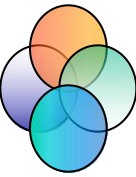


# Lines of E-Flux for E-Dipole

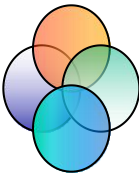


Unnecessary to be continuous.

# Lines of M-Flux for M-Dipole



They must be continuous.



# Key Parameters for M-Dipole

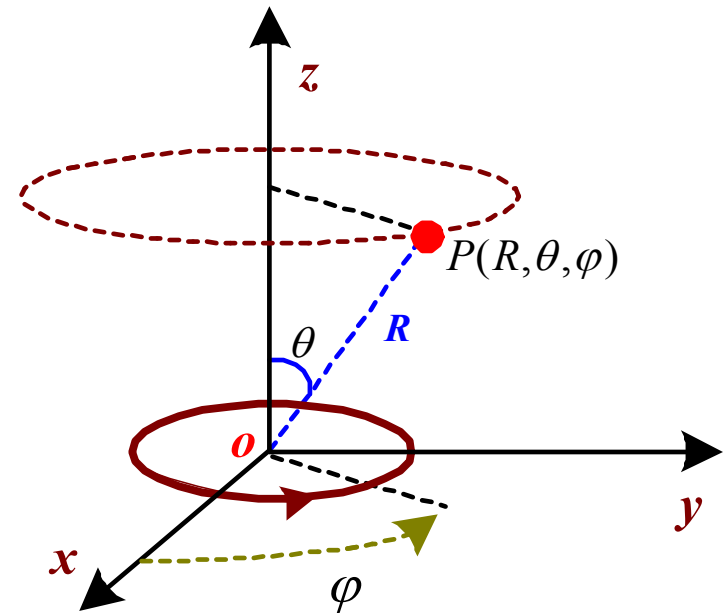
$$\vec{A} = \vec{a}_\varphi \left( \frac{\mu_0 I a^2 \cdot \sin \theta}{4R^2} \right)$$

$$\vec{A} = \frac{\mu_0 \vec{p}_m \times \vec{a}_R}{4\pi \cdot R^2} = -\frac{\mu_0}{4\pi} \vec{p}_m \times \nabla \left( \frac{1}{R} \right)$$

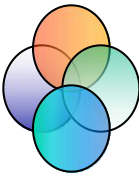
$$\vec{B} = \nabla \times \vec{A}$$

$$\vec{B} = \vec{a}_R \frac{\mu_0 P_m}{2\pi r^3} \cos \theta + \vec{a}_\theta \frac{\mu_0 P_m}{4\pi r^3} \sin \theta$$

$$\vec{E} = \vec{a}_R \frac{P_e}{2\pi \epsilon_0 r^3} \cos \theta + \vec{a}_\theta \frac{P_e}{4\pi \epsilon_0 r^3} \sin \theta$$



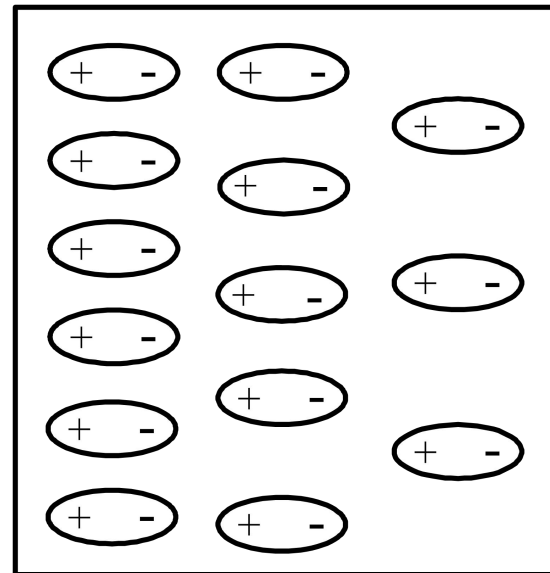
## § 5.4 Material in M-Field

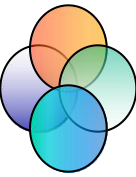


### *Magnetization*

Recall that

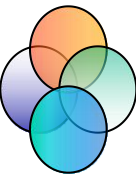
- Materials in E-field will be polarized. Subjected into an E-field, E-dipoles begin to queue orderly which induce bound charges on the surface.



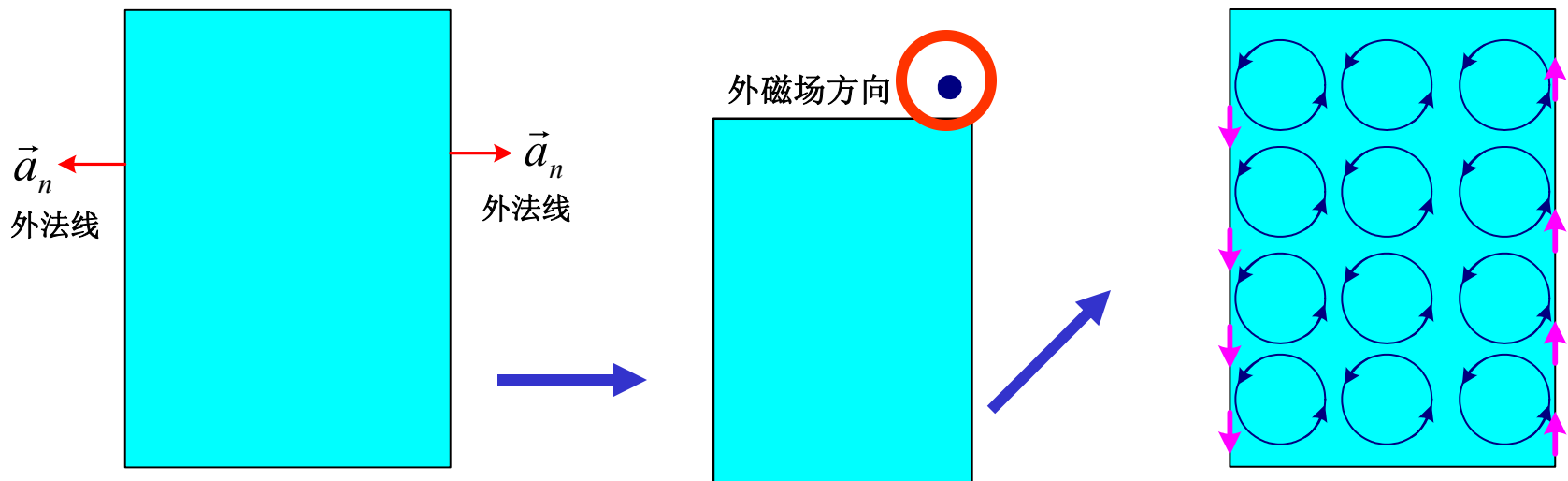


# Materials in M-field will be magnetized

- **Molecule currents, or atom currents are actually M-dipoles.**
- These M-dipoles oriented at random without external M-field.
- With external M-field, all M-dipoles point to the same direction, which is called magnetization.
  - Diamagnetic (反磁性体) : substance inside which the M-intensity is weaker than external M-intensity.
  - Paramagnetic (顺磁性的) : substance inside which the M-intensity is stronger than external M-intensity

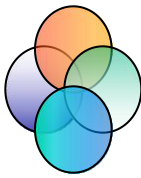


- Due to magnetization, all M-dipoles queue orderly and thus yield a kind of surface current, called **bound current**, or **magnetization current**.





# Magnetization Intensity 磁化强度 (Optional)



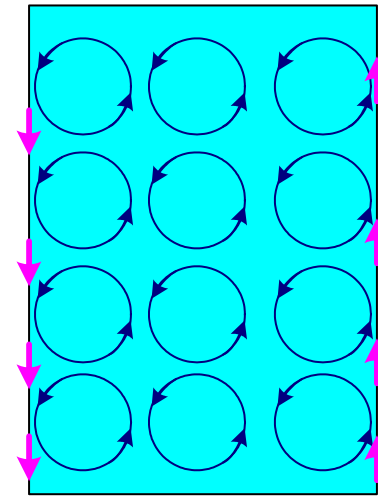
$$\vec{M} = \lim_{\Delta\tau \rightarrow 0} \frac{\sum \vec{p}_m}{\Delta\tau} \quad (A/m)$$

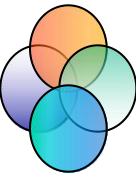
The magnetic dipole moment per unit volume

Recall that the polarization intensity is E-moment per unit volume

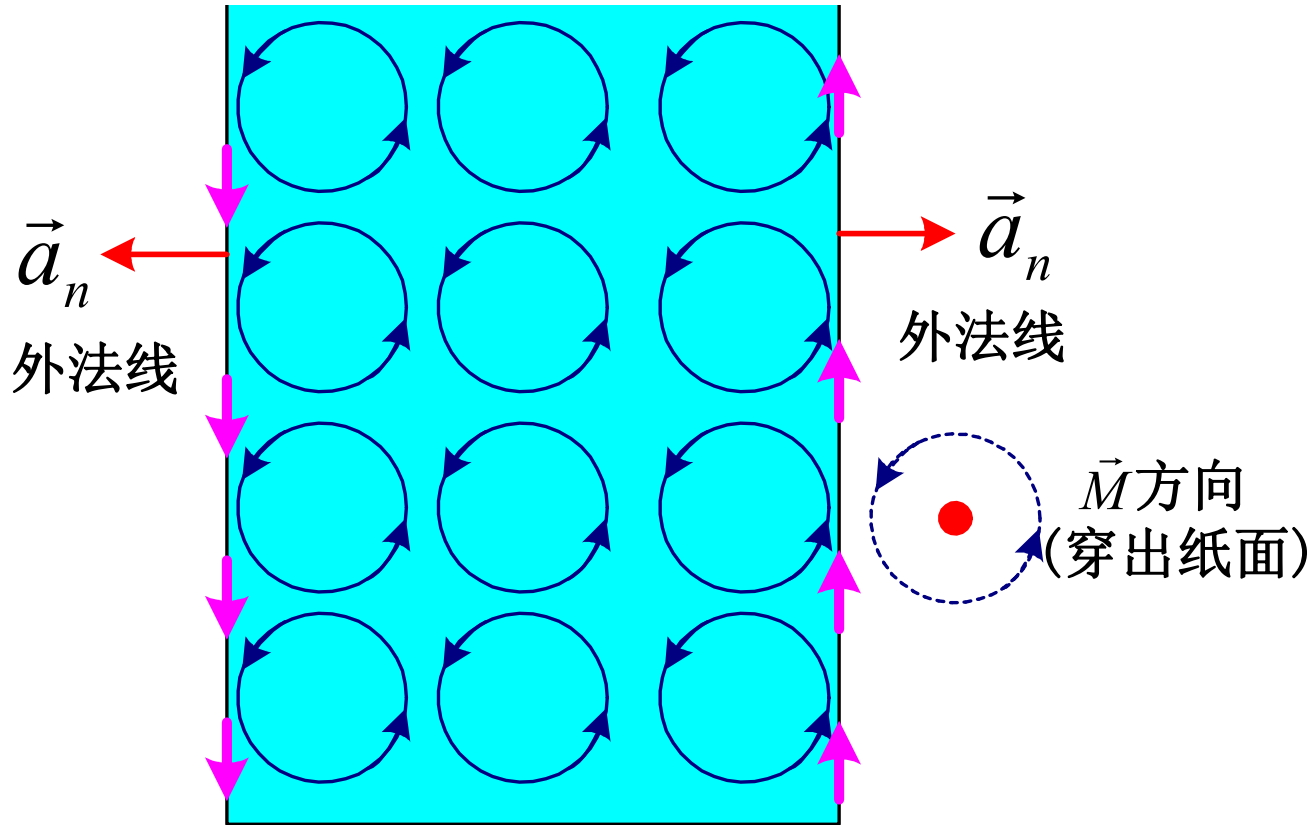
Density of Magnetization current

$$\begin{cases} \vec{J}_m = \nabla \times \vec{M} & (A/m^2) \\ \vec{J}_{ms} = \vec{M} \times \vec{a}_n & (A/m) \end{cases}$$





If a homogeneous substance is magnetized uniformly, the net current inside must be 0.

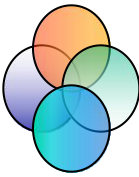


$$\vec{J}_m \begin{cases} \vec{J}_{mV} \\ \vec{J}_{mS} \end{cases}$$

Adjacent M-dipoles will  
**counteract** each other

$$\begin{cases} \vec{J}_{mV} = \nabla \times \vec{M} = 0 \\ \vec{J}_{mS} = \vec{M} \times \vec{a}_n = ? \end{cases}$$

# M-Intensity & Relative Permeability(磁导率)



## Question:

External M-field + Magnetized Substance  $\rightarrow$  New M-field  
How to describe new M-field inside the magnetized substance?

## Solution:

Recall that for magnetostatics **in free space** we have

$$\begin{cases} \nabla \cdot \vec{B} = 0 \\ \frac{1}{\mu_0} \nabla \times \vec{B} = \vec{J} \end{cases}$$

For new M-field inside the magnetized substance

$$\frac{1}{\mu_0} \nabla \times \vec{B} = \vec{J} + \vec{J}_M$$

Corresponding  
to free current

Corresponding to  
magnetization current

$$\frac{1}{\mu_0} \nabla \times \vec{B} = \vec{J} + \vec{J}_M = \vec{J} + \nabla \times \vec{M} \quad \therefore \nabla \times \left( \frac{\vec{B}}{\mu_0} - \vec{M} \right) = \vec{J}$$

## Magnetic Field Intensity

$$\vec{H} = \frac{\vec{B}}{\mu_0} - \vec{M} \quad (A / m)$$

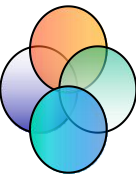
$$\therefore \nabla \times \vec{H} = \vec{J} \quad (\text{volume density of free current})$$

**In comparison with electrostatics:**

$$\oint_S \vec{E} \cdot d\vec{S} = \frac{\sum q_{fc} + \sum q_{pc}}{\epsilon_0} \quad \nabla \cdot (\epsilon_0 \vec{E} + \vec{P}) = \rho_{fc}$$

$$\nabla \cdot \vec{D} = \rho_{fc}$$

$$\nabla \times \vec{H} = \vec{J} \text{ (volume density of free current)}$$



so

$$\int_S (\nabla \times \vec{H}) \cdot d\vec{S} = \int_S (\vec{J}) \cdot d\vec{S}$$

Applying Stokes's Law, we have

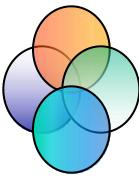
$$\oint_C \vec{H} \cdot d\vec{l} = I$$

*What is the Unit of “Magnetic Field Intensity” ?*

$$\vec{H}: (A / m)$$

*How about “Electric Field Intensity” ?*

$$\vec{E}: (V / m)$$



# In **Linear** & **Isotropic** Materials

$$\vec{M} = \chi_m \vec{H}$$

$\chi_m$  : magnetic susceptibility (磁化率 无量纲)

$$\therefore \vec{H} = \frac{\vec{B}}{\mu_0} - \vec{M}$$

$$\vec{B} = \dots = \mu_0 (1 + \chi_m) \vec{H} = \mu_0 \mu_r \vec{H} = \mu \vec{H}$$

$$\vec{B}: (Wb / m^2)$$

$$\vec{H}: (A / m)$$

$\mu_r$  : relative permeability (相对磁导率)

$\mu$  : absolute permeability (绝对磁导率)

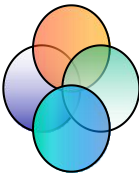
$$\mu_r = 1 + \chi_m = \frac{\mu}{\mu_0}$$

**In comparison with electrostatics:**

$$\vec{P} = \varepsilon_0 \chi_e \vec{E}$$

$$\vec{D} = \varepsilon_0 (1 + \chi_e) \vec{E} = \varepsilon \vec{E}$$

## A Discussion on Relative Permeability

 $\mu_r$ 

1. *diamagnetic* 抗磁性材料  $\mu_r \leq 1$   $\chi_m \approx -0$

Copper, lead, gold, silver, etc..

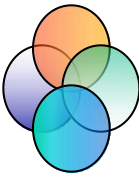
2. *paramagnetic* 顺磁性材料  $\mu_r \geq 1$   $\chi_m \approx +0$

Aluminum, tungsten (钨), etc..

3. *ferromagnetic* 铁磁性材料  $\mu_r \gg 1$   $\chi_m \gg 0$

Cobalt (钴), iron, etc..

# Summary on Material Parameters



真空中磁导率 (*Permeability*):

$$\mu_0 = 4\pi \cdot 10^{-7} (H / m)$$

真空中介电常数 (*Dielectric Constant*):

$$\varepsilon_0 = \frac{1}{4\pi \cdot 9 \times 10^9} = 8.85 \times 10^{-12} (F / m)$$

$$\frac{1}{\sqrt{\varepsilon_0 \cdot \mu_0}} = c$$

简单媒质——线性、均匀、各向同性

磁化率  $\chi_m$ : 无单位、**常数**

相对磁导率  $\mu_r$ : 无单位、**常数**