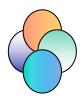
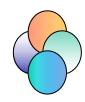
### § 2.9 Boundary Conditions for Electrostatics



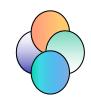
- → The Question
  - → In general, E-intensity or potential may be obtained through two equations:
    - In source free region --- Laplace's Equation
    - Otherwise --- Poisson's Equation
  - **→** Boundary conditions are useful to
    - determine the *undecided constants* in the solution;
    - turn the general solution into *specific solution*.
  - → Boundary Value Problem: solving partial differential equations under given boundary conditions.
    - Poisson's Equation subjected to boundary conditions
    - Laplace's Equation subjected to boundary conditions

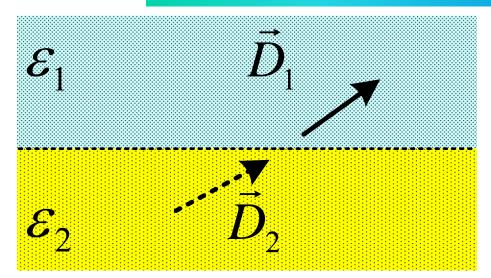
## What is Boundary Conditions?



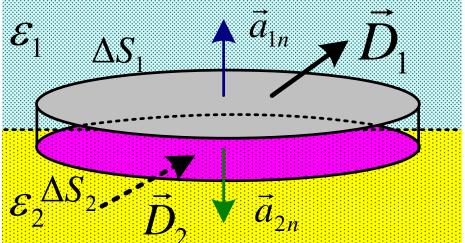
- **→** Equations governing the behavior of fields at the boundary (interface) between 2 mediums.
- Categories of Boundary Conditions:
  - → First: the field potential at boundary is given
    - ♥ Dirichlet Problems 秋理赫利问题
  - → Second: the normal component of the derivative of the field potential at boundary is known
    - \* Neumann Problems 纽曼问题
  - → Third: Hybrid Problems 混合问题

#### 1. Boundary Conditions in Normal Direction





Construct an auxiliary closed surface of a very flat box.



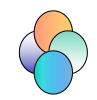
Applying Gauss's Law

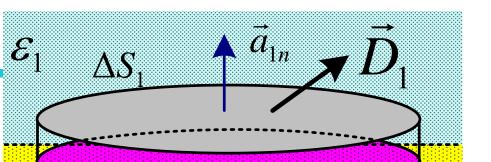
$$\oint \vec{D} \cdot d\vec{S} = \rho_s \cdot \Delta S$$

$$\vec{D}_1 \cdot \Delta \vec{S}_1 + \vec{D}_2 \cdot \Delta \vec{S}_2 = D_{1n} \Delta S - D_{2n} \Delta S = \rho_s \Delta S$$

$$\therefore D_{1n} - D_{2n} = \rho_s$$

---surface density of free charges





$$D_{1n} - D_{2n} = \rho_s$$

$$: \vec{E} = -\nabla \psi$$

$$\therefore \vec{D} \bullet \vec{a}_n = (\varepsilon \vec{E}) \bullet \vec{a}_n = \varepsilon (-\nabla \psi) \bullet \vec{a}_n = -\varepsilon \cdot \frac{\partial \psi}{\partial n}$$

$$\therefore \varepsilon_2 \cdot \frac{\partial \psi_2}{\partial n} - \varepsilon_1 \cdot \frac{\partial \psi_1}{\partial n} = \rho_s$$

#### **Discussions**



$$D_{1n} - D_{2n} = \rho_s$$

$$\varepsilon_2 \cdot \frac{\partial \psi_2}{\partial n} - \varepsilon_1 \cdot \frac{\partial \psi_1}{\partial n} = \rho_s$$

→ If there is no free charge at the boundary,

$$D_{1n} = D_{2n}$$

$$\varepsilon_1 \cdot \frac{\partial \psi_1}{\partial n} = \varepsilon_2 \cdot \frac{\partial \psi_2}{\partial n}$$

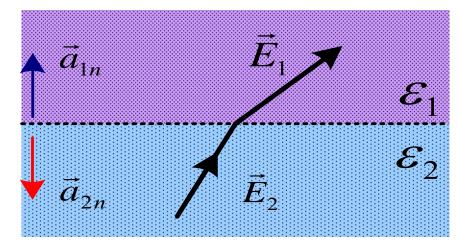
**◆** If the boundary is the interface of a conductor

$$D_{2n}=0$$

$$D_{1n}=\rho_s$$

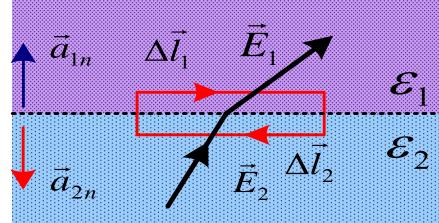
## 2. Boundary Conditions in Tangential Direction





Applying the conservative law

Make an auxiliary closed path in rectangular shape,  $\Delta h \rightarrow 0$ 



$$\oint_{c} \vec{E} \cdot d\vec{l} = 0$$

$$\therefore \vec{E}_{1} \cdot \Delta \vec{l}_{1} + \vec{E}_{2} \cdot \Delta \vec{l}_{2} = E_{1t} \cdot \Delta l - E_{2t} \cdot \Delta l = 0$$

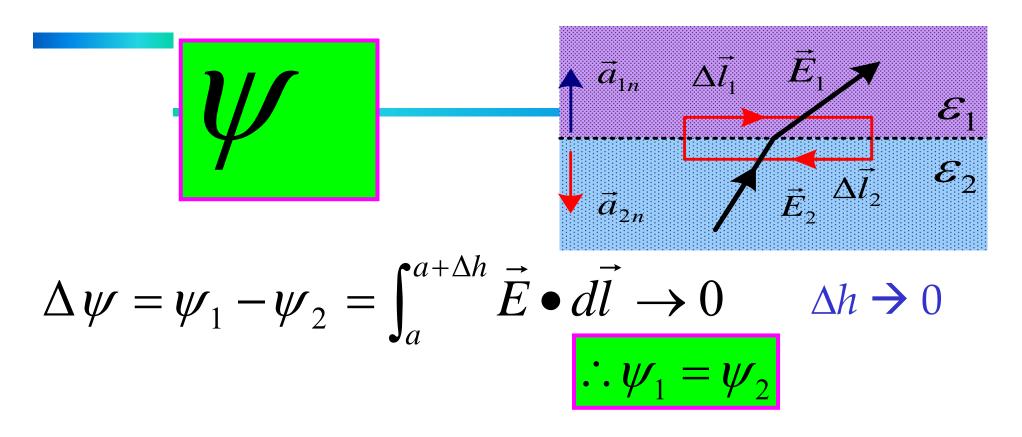
$$\therefore E_{1t} = E_{2t}$$

## 2. Boundary Conditions in Tangential Direction



→ The tangential components of E-intensity are always continuous across the boundary.

$$E_{1t} = E_{2t}$$



- **→** E-potential is continuous across the boundary;
- **→** In other word, E-potential is an integral function and all integral functions are continuous within their domains of definition.

## **Summary --- Boundary Conditions**



#### 1. In normal direction

$$D_{1n} - D_{2n} = \sigma_{fc}$$

$$\varepsilon_2 \cdot \frac{\partial \psi_2}{\partial n} - \varepsilon_1 \cdot \frac{\partial \psi_1}{\partial n} = \sigma_{fc}$$

#### 2. In tangential direction

$$E_{1t} = E_{2t}$$

$$\psi_1 = \psi_2$$

# 3. Applications of Boundary Conditions

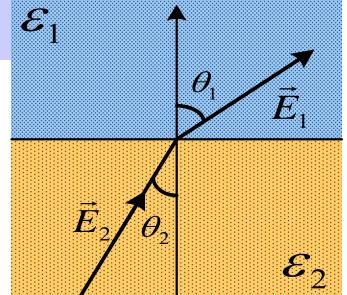


$$: D_{1n} = D_{2n} \Longrightarrow \varepsilon_1 \cdot E_{1n} = \varepsilon_2 \cdot E_{2n} \varepsilon_1$$

$$\varepsilon_1 \cdot E_{1n} = \varepsilon_2 \cdot E_{2n}$$

$$E_{1t} = E_{2t}$$

$$\therefore \begin{cases} E_1 \cdot \sin \theta_1 = E_2 \cdot \sin \theta_2 \\ \varepsilon_1 \cdot E_1 \cdot \cos \theta_1 = \varepsilon_2 \cdot E_2 \cdot \cos \theta_2 \end{cases}$$

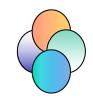


$$\therefore \frac{tg\,\theta_1}{tg\,\theta_2} = \frac{\varepsilon_1}{\varepsilon_2}$$

Similar to the refraction of light

Exercise 3.27 in textbook p127

## Example 2.



- → A conductor ball. Radius *a*. E-potential *U*.
- → Please determine the potential outside the ball.
- Analysis:
  - → Any Symmetry? Yes, point symmetry.
  - → How many approaches to determine E-potential?
    - Direct solution via integral or sum
    - Via E-intensity
    - Via differential equations

## Solution 1.



Because ??? we obtain

$$\nabla^2 \psi = 0$$

Because ??? we infer  $\psi = \psi(r)$ 

$$\psi = \psi(r)$$

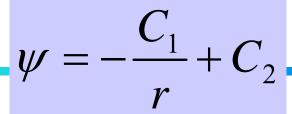
Express Laplace's Equ. in spherical coordinates

$$\nabla^2 \psi = \frac{1}{r^2} \frac{d}{dr} r^2 \frac{d\psi}{dr} = 0$$

Through integral of above equ.  $\psi = -\frac{C_1}{C_2} + C_2$ 

$$\psi = -\frac{C_1}{I} + C_2$$

Via boundary conditions 
$$\psi = \begin{cases} r > a & \frac{a}{r} \cdot U \\ r = a & U \\ r < a & U \end{cases}$$
 Field and Wave Electromagnetics





Determine the undecided constants via Boundary Conditions

If 
$$r \to \infty$$
, we know  $\psi = 0$ , and then  $C_2 = 0$ 

If 
$$r = a$$
 and  $\psi = U$ , we infer  $C_1 = -aU$ 

$$\psi = \begin{cases} r > a, & \frac{a}{r} \cdot U \\ r \le a, & U \end{cases}$$

▶ Please check that E-potential is continuous across the sphere of the conductor ball, i.e.  $\psi_1 = \psi_2$ 





Now, let's go on --->>>