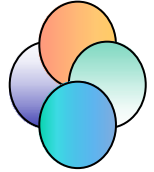


# Foreword



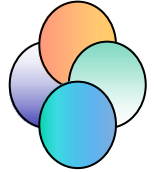
## ➤ In static field

- E-potential is a scalar.
- E-intensity is a vector.
- It is much easier to determine E-potential than to calculate E-intensity.
- Once potential is known, its negative grad. is intensity.

## ➤ How to determine E-potential?

- If the charge distribution is typical, we can write out the potential directly, as in former sections.
- However in most cases, we need to set up differential equations for E-potential.
- These **differential equations** are of 2nd order, and set up according to fundamental equ. of electrostatics.

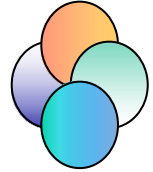
## § 3.11 Two Differential Equations



1. *Poisson's Equation* 泊松方程

2. *Laplace's Equation* 拉普拉斯方程

# Derivation of the Equations



From fundamental electrostatic equations

1. Electrostatic Conservation Law

$$\nabla \times \vec{E} = 0 \quad \Rightarrow \quad \vec{E} = -\nabla u$$

2. E-Gauss's Law in differential form

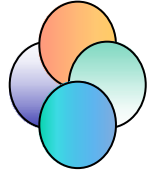
$$\nabla \cdot \vec{D} = \rho$$

3. Material Equation

$$\vec{D} = \epsilon \vec{E}$$

$$\nabla \cdot \vec{E} = \frac{\rho}{\epsilon}$$

# Poisson's Equation



$$\rightarrow \begin{cases} \nabla \cdot \vec{E} = \frac{\rho}{\epsilon} \\ \vec{E} = -\nabla u \end{cases} \Rightarrow \nabla \cdot \nabla u = \nabla^2 u = -\frac{\rho}{\epsilon}$$

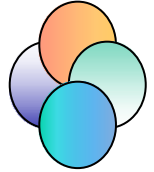
$$\nabla^2 u = -\frac{\rho}{\epsilon}$$

— *Poisson's Equation*

$\nabla^2$  *Laplacian* Or in Chinese 拉普拉斯算符

It's a **part differential function** of 2nd order.

# Laplace's Equation



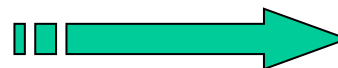
At the **source-free point** or in the **source-free region**, there is no charge scattered and the charge volume density is 0.

$$\nabla^2 u = -\rho / \epsilon$$



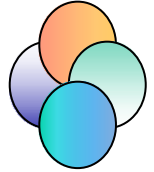
$$\nabla^2 u = 0$$

Poisson's Equation



Laplace's Equation

# About Laplacian



$\nabla^2$  refers to *div. of a grad.*

$$\nabla^2 = \nabla \bullet \nabla$$

$$\nabla^2(u) = \nabla \bullet (\nabla u)$$

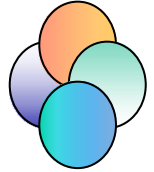
Laplacian  $u$  is a **scalar operator**.

➡ Laplacian in Cartesian Coordinates should be remembered.

$$\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$$

Its a second-order differential operator.

## In Cartesian Coordinates



$$\nabla \bullet \vec{X} = (\vec{a}_x \frac{\partial}{\partial x} + \vec{a}_y \frac{\partial}{\partial y} + \vec{a}_z \frac{\partial}{\partial z}) \bullet \vec{X}$$

$$\nabla \psi = \vec{a}_x \frac{\partial \psi}{\partial x} + \vec{a}_y \frac{\partial \psi}{\partial y} + \vec{a}_z \frac{\partial \psi}{\partial z}$$

$\nabla^2 \psi$  is just the **dot product of del with gradient**

$$\nabla^2 \psi = \nabla \bullet \nabla \psi = \frac{\partial}{\partial x} \psi_x + \frac{\partial}{\partial y} \psi_y + \frac{\partial}{\partial z} \psi_z$$

$$\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$$



- Expressions are more complicate in cylindrical and spherical coordinates.

## In Cylindrical Coordinates

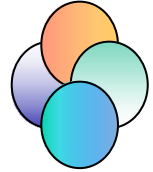
$$\nabla^2 u(r, \varphi, z) = \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial u}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 u}{\partial \varphi^2} + \frac{\partial^2 u}{\partial z^2}$$

## In Spherical Coordinates

$$\nabla^2 u(r, \theta, \varphi) = \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial u}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial u}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 u}{\partial \varphi^2}$$

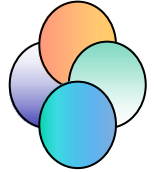


## Example 1.



- A *conductor* ball, with Radius of  $a$  & E-potential of  $U$ .
- Please determine E-potential outside the ball.
- Analysis:
  - Any Symmetry? — — Yes, **Spherical** symmetry.
  - How many approaches to determine E-potential?
    - ⊕ Via differential equations
    - ⊕ Via E-intensity
    - ⊕ Direct solution — — via integral or sum

# Solution 1, Laplace's Equation.



Because ??? we obtain  $\nabla^2 u = 0$

Because ??? we infer  $u = u(r)$

Express Laplace's Equ.  
in spherical coordinates  $\nabla^2 u = \frac{1}{r^2} \frac{d}{dr} r^2 \frac{du}{dr} = 0$

Through integral of above equ.  $u = -\frac{C_1}{r} + C_2$

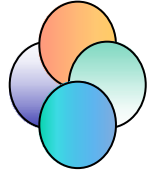
By boundary conditions

$$u \xrightarrow{r=a} U$$

$$u \xrightarrow{r=\infty} 0$$

$$u = \begin{cases} r > a & \frac{a}{r} \cdot U \\ r = a & U \\ r < a & U \end{cases}$$

## Solution 2, E-intensity



Via E-intensity

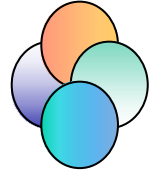
$$u(r) = \int_{\text{point A}}^{\infty} \vec{E} \cdot d\vec{l} = \int_r^{\infty} E_r \cdot dr$$

**Assuming** the ball is charged by  $Q \Rightarrow \vec{E} = \vec{a}_r E_r = ?$

Recall that we have calculated the E-intensity outside a conductor ball.

$$\vec{E} = \frac{1}{4\pi\epsilon_0 \cdot r^2} \cdot Q\vec{a}_R$$

## Recall:



Applying *E*-Gauss's Law

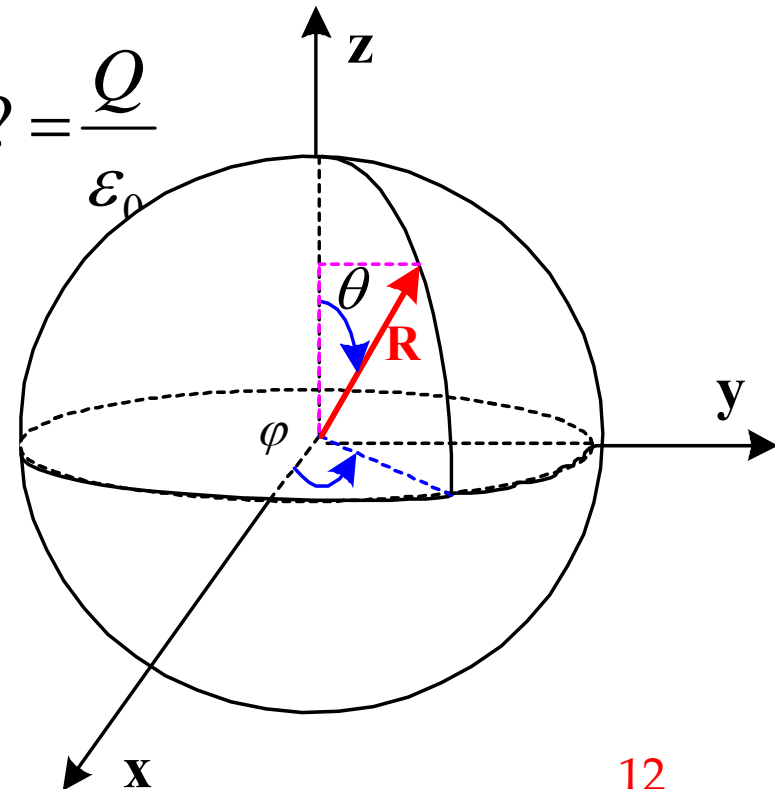
$$\oint_S \vec{E} \cdot d\vec{S} = \frac{1}{\epsilon_0} \int_V \rho dV = \frac{Q}{\epsilon_0}$$



Inside the ball ( $r < a$ ):  $\because \frac{1}{\epsilon_0} \int_V \rho dV = 0 \quad \therefore \vec{E} = 0$

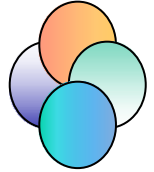
Outside the ball ( $r > a$ ):  $\because \frac{1}{\epsilon_0} \int_V \rho dV = ? = \frac{Q}{\epsilon_0}$

$$\oint_S \vec{E} \cdot d\vec{S} = E \cdot (4\pi \cdot r^2)$$

$$\therefore \vec{E} = \frac{1}{4\pi\epsilon_0 \cdot r^2} \cdot Q\vec{a}_R$$



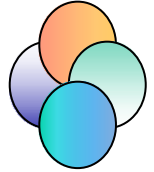

$$\vec{E} = \frac{1}{4\pi\epsilon_0 \cdot r^2} \cdot Q\vec{a}_R$$



$$\psi(r) = \int_r^\infty \vec{E} \bullet d\vec{l} = \int_r^\infty E_r \cdot dr = ?$$

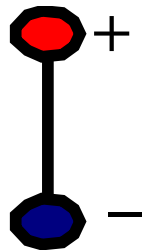
$$\because \psi(r) \big|_{r=a} = U \Rightarrow Q = ? \quad \psi(r) = ?$$

## § 3.6 Electric Dipole



A pair of equal charges of opposite signs that are very close together.

Two charges of equal charge but of opposite polarity and separated by a small distance.



Distance:  $l$

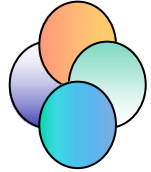


Point charges:  $q_1=q$ ,  $q_2=-q$

$$\vec{E} = \vec{E}_+ + \vec{E}_-$$

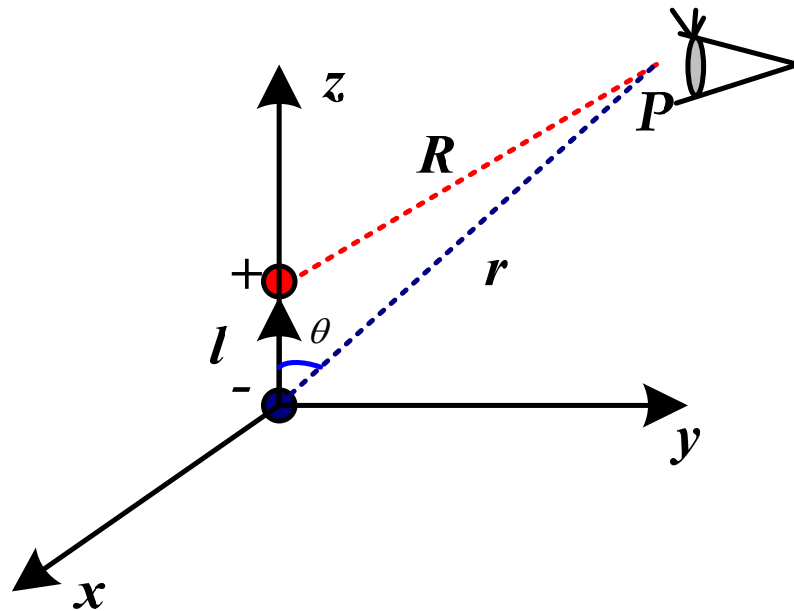
$$\vec{E}_- = -\frac{-q}{4\pi\epsilon_0} \nabla\left(\frac{1}{|\vec{r}|}\right)$$

$$\vec{E}_+ = -\frac{q}{4\pi\epsilon_0} \nabla\left(\frac{1}{|\vec{R}|}\right)$$

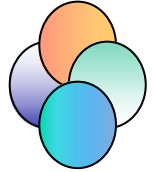


Cosine Theorem  $\frac{1}{|\vec{R}|} = \frac{1}{R} = \frac{1}{\sqrt{r^2 + l^2 - 2 \cdot r \cdot l \cos \theta}}$

Taylor Series ( $l \ll r$ )  $\frac{1}{R} = R^{-1} \approx \frac{1}{r} + \frac{1}{r^2} \cdot l \cdot \cos \theta$



## In Spherical Coordinates



$$\vec{E}_- = -\frac{-q}{4\pi\epsilon_0} \nabla\left(\frac{1}{r}\right) \quad \vec{E}_+ \approx -\frac{q}{4\pi\epsilon_0} \nabla\left(\frac{1}{r} + \frac{l}{r^2} \cdot \cos\theta\right)$$

$$\vec{E} = \vec{E}_+ + \vec{E}_- = -\frac{q}{4\pi\epsilon_0} [\nabla(?) - \nabla(?)]$$

$$\vec{E} = -\frac{q}{4\pi\epsilon_0} \left[ \nabla\left(\frac{1}{r} + ?\right) - \nabla\left(\frac{1}{r}\right) \right] = -\frac{q}{4\pi\epsilon_0} \nabla\left(\frac{l \cdot \cos\theta}{r^2}\right)$$



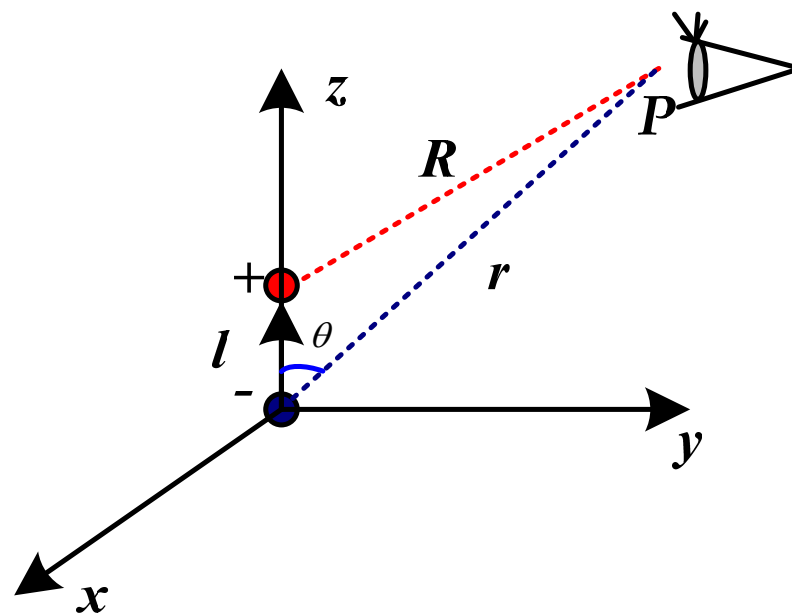
## Dipole Moment Vector (电偶极距)



$$\vec{E} = -\frac{q}{4\pi\epsilon_0} \left[ \nabla \left( \frac{1}{r} + ? \right) - \nabla \left( \frac{1}{r} \right) \right] = -\frac{q}{4\pi\epsilon_0} \nabla \left( \frac{l \cdot \cos \theta}{r^2} \right)$$

$$\vec{E} = -\frac{1}{4\pi\epsilon_0} \nabla \left( \frac{(ql) \cdot r \cdot \cos \theta}{r^3} \right)$$

$$\vec{E} = -\frac{1}{4\pi\epsilon_0} \nabla \left( \frac{(q\vec{l}) \bullet \vec{r}}{r^3} \right)$$



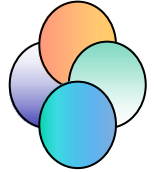
Let

$$\vec{p} = q\vec{l}$$

The quantity?

The direction ?

## Dipole Moment Vector




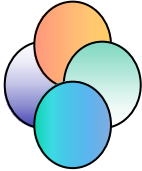
$$\vec{p} = q \vec{l}$$

Unit:  $C \cdot m$

$$\vec{E} = -\frac{1}{4\pi\epsilon_0} \nabla \left( (\vec{p} \bullet \vec{r}) \cdot \frac{1}{r^3} \right)$$

$$\nabla(u \cdot v) = ?$$

$$\nabla(u \cdot v) = u \nabla(v) + v \nabla(u) \quad \text{Vector}$$

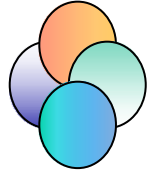



$$\vec{E} = -\frac{1}{4\pi\epsilon_0} \nabla \left( (\vec{p} \cdot \vec{r}) \cdot \frac{1}{r^3} \right) = ?$$

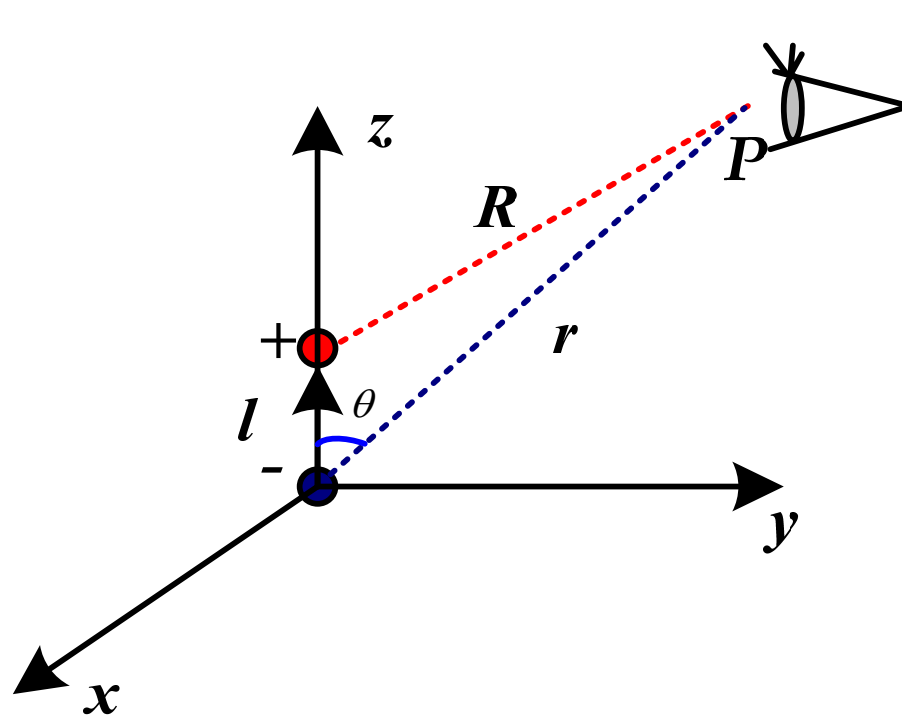
$$\vec{E} = ? \quad \nabla \left( (\vec{p} \cdot \vec{r}) \cdot \frac{1}{r^3} \right) = ? \left[ (\vec{p} \cdot \vec{r}) \nabla \left( \frac{1}{r^3} \right) + \frac{1}{r^3} \nabla (\vec{p} \cdot \vec{r}) \right]$$

$$= ? \left[ \left( \frac{-3 \cdot (\vec{p} \cdot \vec{r})}{r^5} \right) \vec{r} + \frac{1}{r^3} \vec{p} \right]$$

$\nabla(u \cdot v) = u \nabla(v) + v \nabla(u)$	<i>Vector</i>
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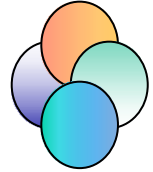
$$\vec{E}(\vec{p}, \vec{r}) = \frac{1}{4\pi\epsilon_0} \left[ \frac{3 \cdot (\vec{p} \cdot \vec{r})}{r^5} \vec{r} - \frac{1}{r^3} \vec{p} \right]$$



$$\vec{p} = q \vec{l}$$

$$|\vec{E}| \propto \frac{1}{r^3}$$

## E-Flux Lines of Electric Dipole



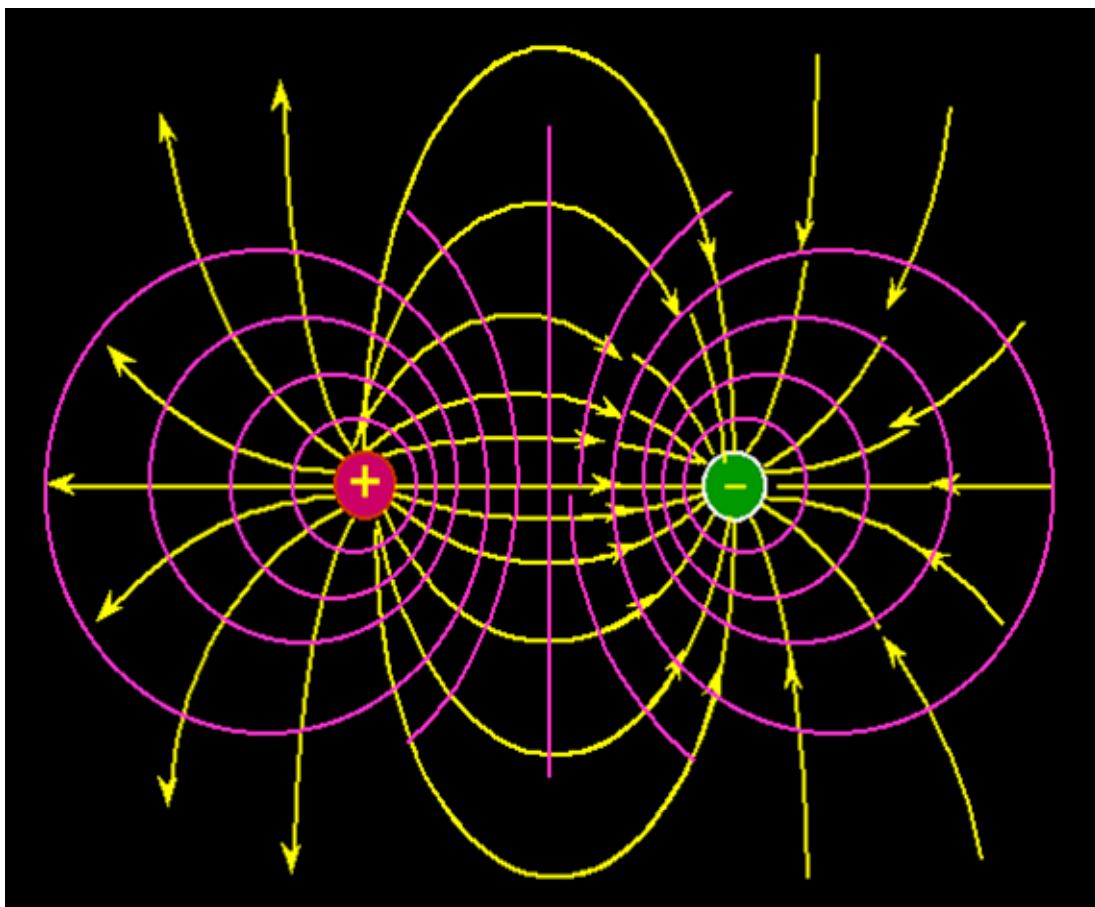
Equ. of E-Flux Lines: The lines shall be parallel to E-Field.

$$(1) \quad d\vec{l} \times \vec{E}(\vec{p}, \vec{r}) = 0 \quad (2) \quad kd\vec{l} = \vec{E}(\vec{p}, \vec{r})$$

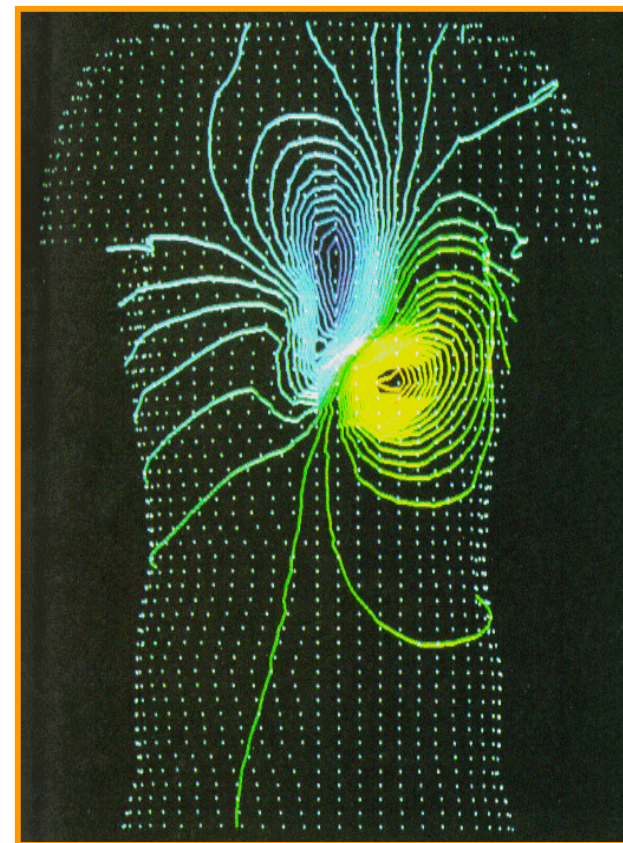
In spherical coordinates  $\vec{E}(\vec{p}, \vec{r}) = E_r \vec{a}_r + E_\theta \vec{a}_\theta + 0\vec{a}_\phi$

$$d\vec{l} \times \vec{E}(\vec{p}, \vec{r}) = ? + ? + ? = 0$$

Equ. of E-Field Lines  $r = C \cdot \sin^2 \theta$       The figure ?

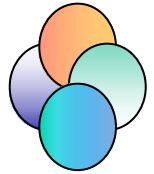


电偶极子的电场线和等势面



作心电图时人体的等势面分布

# Force of -Dipole

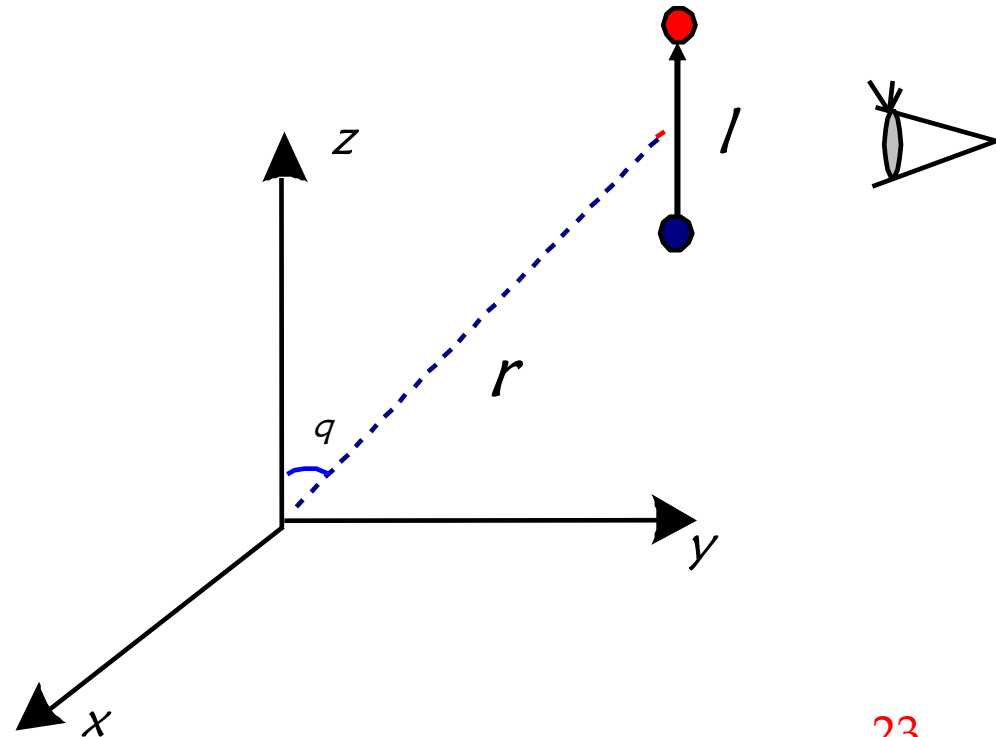


$$\vec{f}_{\vec{p}}(\vec{r}) = q\vec{E}(\vec{r} + \frac{\vec{l}}{2}) + (-q)\vec{E}(\vec{r} - \frac{\vec{l}}{2})$$

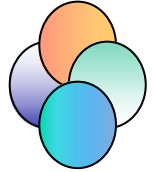
In static E-field

$$\vec{E} = -\frac{1}{4\pi\epsilon_0} \nabla \left( (\vec{p} \cdot \vec{r}) \cdot \frac{1}{r^3} \right) = ?$$

$$\vec{f}_{\vec{p}}(\vec{r}) = \nabla [\vec{p} \cdot \vec{E}(\vec{r})]$$



# Moment of Force of E-Dipole



$$\begin{aligned}\vec{T}_{\vec{p}}(\vec{r}) &= \frac{\vec{l}}{2} \times \left[ q\vec{E}(\vec{r} + \frac{\vec{l}}{2}) \right] - \frac{\vec{l}}{2} \times \left[ (-q)\vec{E}(\vec{r} - \frac{\vec{l}}{2}) \right] \\ &= \frac{q\vec{l}}{2} \times \left[ \vec{E}(\vec{r} + \frac{\vec{l}}{2}) + \vec{E}(\vec{r} - \frac{\vec{l}}{2}) \right] = \vec{p} \times \vec{E}(\vec{r})\end{aligned}$$



- Distance:  $l$
- Point Charges:  $q_1=q, q_2=-q$